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A GENERAL STREAMLINE MODELLING TECHNIQUE FOR
HOMOGENEOUS AND HETEROGENEOUS POROUS MEDIA, WITH
APPLICATION TO STEAMFLOOD PREDICTION

The University of Oklahoma

PH.D. 1982

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A GENERAL STREAMLINE MODELLING TECHNIQUE FOR
HOMOGENEOUS AND HETEROGENEOUS POROUS MEDIA,
WITH APPLICATION TO STEAMFLOOD PREDICTION

A DISSERTATION
SUBMITTED TO THE GRADUATE FACULTY
in partial fulfillment of the requirements for the
degree of
DOCTOR OF PHILOSOPHY

By
DAOPU THOMPSON NUMBERE
Norman, Oklahoma
1982

A GENERAL STREAMLINE MODELLING TECHNIQUE FOR
HOMOGENEOUS AND HETEROGENEOUS POROUS MEDIA,
WITH APPLICATION TO STEAMFLOOD PREDICTION

A DISSERTATION

APPROVED FOR THE DEPARTMENT OF PETROLEUM AND
GEOLOGICAL ENGINEERING

APPROVED BY

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ABSTRACT

A general streamline/streamtube simulation method applicable to homogeneous, anisotropic, as well as piecewise homogeneous porous media is described with its application to steamflood recovery prediction. The boundaries of the media can be of any arbitrary but smooth shape with boundary conditions that can be any of or combination of either Neuman, Dirichlet, or mixed boundary conditions.

The (Laplace's) equation describing the distribution of potential was solved by the Boundary Element Method (BEM) also called the Boundary Integral Element Method (BIEM). This method utilizes the superposition of singular solutions whose distributions are determined by the boundary conditions of the problem. The streamlines are generated assuming single phase flow (unit mobility assumption). The streamtubes are assumed to be linear tubes of rectangular cross section. Streamlines and streamtubes were generated for a homogeneous reservoir assuming it was bounded on all sides by sealing faults. Then, it was assumed to have an oil-water contact (constant pressure) boundary on one side, while the remainder of the boundary remained sealed. Oil recovery calculations were made for a steam drive process

in each streamtube generated and added together to give the recovery in the entire field or pilot area. Finally, the reservoir was treated as a piecewise homogeneous reservoir having two regions of unequal permeabilities. The streamlines and streamtubes were again generated and oil recovery calculated for each of the two types of boundary arrangements mentioned above.

Steamflood prediction was made assuming each streamtube existed in isolation. Thus, there was no heat exchange between streamtubes. For homogeneous reservoirs, recovery was calculated using the Marx and Langenheim equations. For the piecewise homogeneous reservoir, new equations were derived for the location and rate of advance of the steam front while it is in the second permeability region of a streamtube containing two permeability regions. Comparison of the potentials obtained by the BEM and those obtained analytically for simple domains confirm the validity of the Boundary Element Method. This method provides a general technique to properly model different reservoirs with arbitrary boundaries and boundary conditions, particularly, piecewise homogeneous reservoirs which had not been possible using the method of images.

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A GENERAL STREAMLINE MODELLING TECHNIQUE FOR
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WITH APPLICATION TO STEAMFLOOD PREDICTION

CHAPTER I

INTRODUCTION

Steamflood recovery performance can be predicted either by the use of analytical equations derived for linear and radial homogeneous porous media, or by the use of numerical simulators. The numerical simulators have the benefits of giving more detailed and more reliable results. They also have the advantage of being applicable to a wide variety of types of porous media. Unfortunately, these simulators require a great deal of detailed reservoir data, making them expensive to run. Quite often, the available geological and production data are not sufficient, and economic considerations prohibit the collection of such extensive data. In such situations, simulators are often run with some of the data being assumed. This reduces the advantages of the numerical simulator. Improving the accuracy by obtaining more data may make the cost of running one unjustifiable.

The analytical methods have the advantage of simplicity and therefore are relatively cheap. But, in order to predict recovery using analytic equations for an entire field with multiple sources and sinks, the field must be reduced to an equivalent linear or radial system. One way to achieve this is by the so-called streamline/streamtube technique. Previous methods of determining the streamlines and streamtubes for reservoirs with regular as well as irregular boundaries were based on the technique of image wells. According to the well-image theory, hydrogeologic boundaries may be replaced, for analytical purposes, by imaginary wells which produce the same disturbance effects as the boundaries. Boundary problems are thereby simplified to the consideration of an infinite reservoir in which real and image wells operate simultaneously and their effects are superposed to give the potentials and streamlines at desired locations. This image-well method is inadequate in the following respects:

1. Theoretically, the number of image wells and their locations extend to infinity. The actual number of image wells used in the solution of a particular problem becomes a matter of experience or judgement.
2. It has only been applied to homogeneous reservoirs. Extension to heterogeneous systems is theoretically possible but the calculations very quickly become tedious.

3. Irregular boundaries are generally modelled by arbitrary locations of image wells whose rates are determined as those necessary to confine as much of the streamlines as desired. Hence, the derived potentials from which the streamlines and streamtubes are calculated do not represent the true potential solutions to the problem with its boundary conditions.

It is therefore easy to see that a new modelling technique is needed for cases where current streamline methods would either be over-simplifications or incorrect, and numerical simulators uneconomical. These would include:

- (i) Homogeneous reservoirs with irregularly shaped boundaries having a variety of boundary conditions.
- (ii) Heterogeneous reservoirs whose properties are such that they could be naturally divided into a number of subregions with homogeneous properties.

The research work described here develops such a model. A heterogeneous reservoir is assumed to be made up of a few subregions with constant material properties such as transmissibility, porosity, etc. The model exploits the simplicity of the analytical models by utilizing as its basis the point source solution in infinite medium, from which it determines other solutions by superposition. At the same time, it uses a numerical technique where the boundaries of the region are discretized, called the Boundary Element

Method (BEM) or the Boundary Integral Element Method (BIEM). This technique is used to generate the streamtubes (or stream channels) for heterogeneous reservoirs of any arbitrary shape or boundary conditions. These streamtubes form the conduits through which the steam displacement processes take place from injectors to producers.

CHAPTER II

LITERATURE SURVEY

In the year 1899, Slichter¹ used the concept of steady state streamlines in his theoretical investigation of the motion of ground waters. In 1962, Ferris et al² published a comprehensive study of ground water flow in bounded systems with various kinds of boundary conditions. The effects of the boundaries were simulated by the use of image wells. This image-well technique is now widely used by ground water hydrologists as well as petroleum engineers.^{2,3,4,5}

Petroleum engineers have extended the streamline concept to include so-called streamtubes or streamchannels whose dimensions are determined from the knowledge of the streamlines.^{6,7} The streamtubes act as the channels, inside which oil displacement processes are assumed to be taking place. The use of image-wells to simulate the boundary conditions has generally been limited to homogeneous domains with boundaries of simple geometrical shapes such as squares, rectangles, circles, etc. Even though, in theory, the image well method can be applied to heterogeneous media⁴ as well as to arbitrary boundary shapes, in

practice, the calculations quickly become tedious. This is partly due to the fact that the image wells and their locations extend to infinity even for the simplest closed systems. This problem was partially addressed by Chan,⁸ who in 1975 presented an improved image-well method for straight-line boundaries that involves a rearrangement of the image pattern depending on the boundary configurations. However, even with this improvement, the number of image wells required still prohibits its use in all but the simplest geometrical shapes. Secondly, the improved method does not account for heterogeneous media. The limitation to simple geometrical shapes was removed by Leblanc⁹ and Lin.¹⁰ Leblanc used the image-well technique to bound the streamlines in a homogeneous reservoir of arbitrarily shaped boundary by using a trial and error procedure to determine the image well rates required to adequately confine the streamlines. Lin extended the method to include reservoirs which are not sealed but have some restricted flow as in the case of partial water encroachment. Each method is unsatisfactory because locations of the image wells (and therefore the potentials) are determined by how well the streamlines are confined rather than by the true boundary conditions of the physical system. Thus, the image-well technique is inadequate for reservoirs that are:

(a) heterogeneous and (b) have arbitrarily shaped boundaries.

The Boundary Element Method (BEM), also called the Boundary Integral Element Method (BIEM), is a technique that can handle several kinds of shapes and heterogeneities. It is an integral method that utilizes the superposition of singular solutions of the partial differential equations of the system. The method has recently emerged as an important numerical technique for the solution of linear elliptic equations such as occur in potential flow and elastostatics.^{11,14} During the last decade, the BEM has been extended to solve linear parabolic equations as well.^{15,16} The extension of the method to piece-wise non-homogeneous bodies of arbitrary shape has been attempted by Banerjee,¹⁷ and Jaworski¹⁸ while Butterfield and Tomlin¹⁹ have used the method for solving zoned anisotropic continuum problems. The literature on the BEM is large and growing rapidly as more researchers develop new methods of application. Reference 17 lists more than 70 references on this subject and contains a good review of the literature on BEM up to the year 1976. The Boundary Element Method has been used in this research work to generate the steady state streamlines and associated streamtubes for heterogeneous reservoirs. The streamtubes serve as the channels inside which any displacement process (steam drive in this case) takes place. In what follows, a brief review of the literature on steam drive processes is given.

The benefits of applying heat to an oil reservoir to aid oil recovery were foreseen as early as 1917,⁵⁹ and

in the 1920's and 1930's, steam was used to remove paraffin from the wellbore.²² Since Shell Oil Company's successful stimulation of a California oil well in the early 1960's, steam injection has steadily gained prominence as an oil recovery method. In fact, steam injection is currently the most popular of all enhanced oil recovery methods--producing 405 million barrels per day (MMBPD) worldwide as compared with 221 MMBPD from all other enhanced recovery methods combined.²²

Even before Shell Oil Company's successful stimulation of the early 1960's, technical papers began to be published describing the transport of injected heat in porous media. One of these papers was by Lauwerier,²³ whose model for the injection of hot water into an oil bearing formation was published in 1955. He assumed a linear homogeneous reservoir of constant thickness, sandwiched between two oil sands. The thermal conductivity of the oil sand in the vertical direction was assumed equal to that of the cap and base rock which were identical. The horizontal thermal conductivity in the oil sand was assumed to be zero. The vertical thermal conductivity in the water sand was assumed infinite so that the temperature in the vertical direction for the water sand was always uniform. Instantaneous equilibrium was assumed between sand grains and reservoir fluid so that the sand grains had the same temperature as the reservoir fluid throughout. Finally,

he formulated his problem by assuming that at time t , the temperature at the boundaries between the two zones were elevated and kept at a constant temperature due to the injection of hot water. By making separate heat balances within the oil and water sands, two equations were obtained. They were solved using the Laplace transform technique to obtain the temperature distributions in the two sands. In 1959, Rubenstein²⁴ improved Lauwerier's model by assuming constant, isotropic thermal conductivities in the reservoir as well as in the confining cap and base rock.

In 1959, Marx and Langenheim²⁵ presented a model for calculating the heated area and its rate of advance based on idealized step temperature at the steam front. They also presented an equation for the temperature gradient for a linear temperature distribution in the steam zone. Finally, they proposed a method of predicting the oil recovery. In 1961, Willman, et al²⁶ presented data on oil recovery mechanisms for hot water and steam injection for a variety of oils and sand types. They presented a modified Buckley-Leverett method to predict the oil recovery for a radial system. Both the Marx and Langenheim, and the Willman, et al models have been further explained and expanded by several authors.²⁷⁻³⁰ Baker³¹ reported on an experimental study of heat flow in steam flooding in a radial system. Among some of Baker's results were: (a) the fraction of injected heat lost when expressed as a function

of time did not depend on injection rate. (b) A significant portion of injected heat was contained in the hot water zone ahead of and underlying the steam zone. (c) The division of heat between the steam and hot water zones did depend on injection rate. In the same year, Mandle and Volek³² made a rigorous calculation of the transport of heat during steamflooding. They found that earlier theory, based on neglecting the heat transport from the steam zone into the oil/water region was consistent with the physical model up to a critical time t_c . t_c marks the time at which the mode of heat transfer from the steam zone into the oil/water region changes from being predominantly conductive to being predominantly convective. The equation describing the growth of the steam zone changes accordingly at $t = t_c$. They presented approximate equations to describe the steam zone growth after $t = t_c$.

In 1970, Shutler and Boberg³³ presented a graphical method for calculating oil recovery during steamflooding. The method was a modification of the Buckley-Leverett method for isothermal, two phase flow in porous media. The 1970's saw continued and more diverse interest in steam drive mechanisms. Previous efforts had been directed more on obtaining analytical solutions to simplified models. But in the 1970's the analytical, experimental, as well as numerical models, became more sophisticated. Some of the experimental and analytical research reported were:

Neuman³⁴ presented a theoretical analysis for a steam flood model that included the effects of gravity override. Van Lookeren³⁵ derived equations to estimate the approximate shape of the steam/liquid interface. Huygen³⁶ who fitted curves to recovery and sweep data for a three dimensional model of a five-spot pattern. Gomaa³⁷ used a numerical method to determine the effects of various parameters on steamflood recovery. He reported the following: (a) above 150 ft., the reservoir thickness had little effect on vertical heat loss. (b) Pattern shape and spacing had insignificant effect on steamflood oil recovery if injection rate per unit reservoir volume was fixed and well productivity was not a limiting factor. Chu³⁸ studied the effects of well patterns on steamflood performance. He reported that the alteration of one pattern resulted in substantial loss of production in the surrounding patterns. He also reported that a 5-spot pattern produced more oil, the production getting progressively less in the inverted 7-spot, the inverted 13-spot and the inverted 9-spot, respectively. Atkinson and Ramey³⁹ presented a general relationship between the various mathematical models of heat transfer in porous media. They also presented new models for fractured and non-fractured porous media. Rhee and Doscher⁴⁰ presented a semi-analytical method for calculating oil recovery by steamflooding that combined the effects of steam distillation and gravity override. Myhill and

Stegemeier⁴¹ published correlations for the effects, of thermal, reservoir petrophysical, and steam properties on the equivalent oil/steam ratio.

CHAPTER III

RESEARCH OBJECTIVE AND PROPOSED MODEL

The objective of this research work is to develop a streamline model to predict oil recovery during steam-flooding in heterogeneous porous media of arbitrary shapes and boundary conditions. The prediction model is desired to be able to handle the following types of heterogeneities:

1. Piece-wise homogeneous reservoirs. That is, heterogeneous reservoirs that can be split into subregions such that rock properties such as permeability, porosity, etc. can be considered as uniform within each subregion, but only be different from region to region.
2. Anisotropic reservoirs.
3. Homogeneous reservoirs with impermeable shale inclusions.

The external boundaries of these reservoirs can be

1. completely closed
2. partially closed, with the remainder at infinity.

The boundary conditions can be any of Neuman, Dirichlet, mixed, or any combination of the aforementioned.

In this research work, only the piecewise homogeneous reservoirs will be treated in detail, the others

are handled by minor modifications. The proposed model uses the Boundary Element Method to generate steady state streamlines for the above mentioned heterogeneous reservoirs. As each streamline is traced, the dimensions of a hypothetical streamtube surrounding it are calculated. Thus, for each streamline, an associated streamtube is calculated. When this is done for all the streamlines, the heterogeneous system is now assumed to be replaced by a series of straight conduits connecting injectors to producers. Each of these conduits is considered to be an isolated system inside which the steam displacement processes take place. Lateral heat losses from the sides of the streamtubes are ignored. For piecewise homogeneous reservoirs, some, but not all, of the streamtubes will consist of regions of differing permeabilities. For these cases, a new generalized derivation of the equations of Marx and Langenheim or Mandle and Volek was made. The rate of advance of the steam front and recovery in each of the tubes is accumulated for any time to complete the desired model.

CHAPTER IV

STATEMENT OF THE PROBLEM

Consider a two-dimensional reservoir domain (D), of arbitrary shape enclosed by the boundary surface (S). The domain D is made up of a number of homogeneous sub-domains, D_i for $i = 1, \dots$, number of sub-domains (Figure 1). Thus, $D = D_1 + D_2 + \dots + D_n$. Each subdomain has its own value of material properties such as permeability, porosity etc., that may or may not be different from those of the adjacent region or regions.

The surface S is assumed to be sufficiently smooth and can be such that part or all of it is at infinity. S can be made up of any combination of three kinds of surfaces S_ϕ , S_n , and $S_{\phi,n}$. Each specifies a different kind of boundary condition. $S = S_\phi + S_n + S_{\phi,n}$. Over S_ϕ , the potentials are specified. Over S_n , the gradients $\left(\frac{\partial \phi}{\partial n}\right)$ are specified; and over $S_{\phi,n}$, a mixed boundary condition is specified.

Inside any or all of the domains are arbitrarily located sources and sinks. The first problem is to find the steady state potential, streamline and streamtube distributions for the whole system subject to any combination of the three types of boundary conditions:

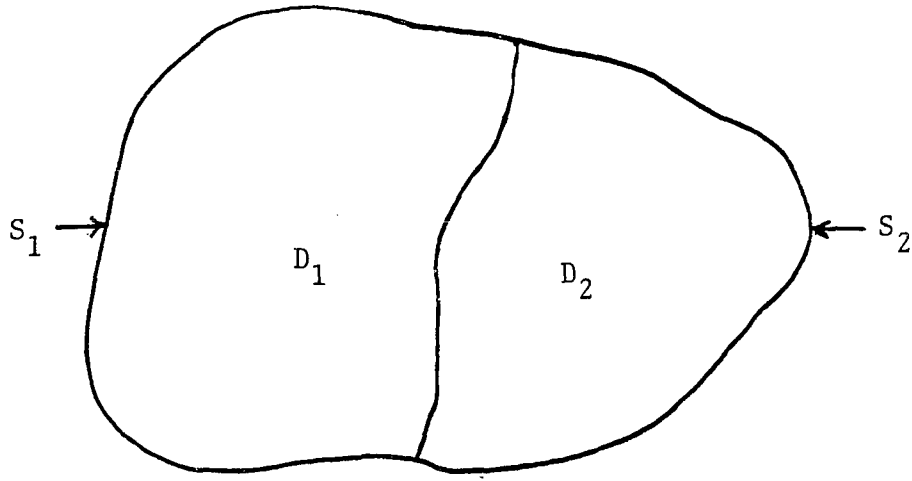


FIGURE 1: Piecewise-Homogeneous Domain

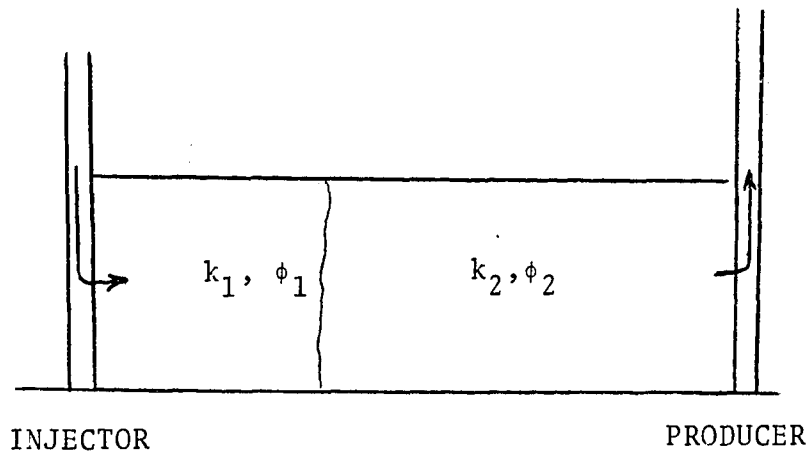


FIGURE 2

Linear one-dimensional model of piecewise-homogeneous reservoir

- a. The Dirichlet condition prescribes the potential on the S_ϕ boundary. That is,

$$\phi = \bar{\phi}(x,y) \text{ on } S = S_\phi \quad (1)$$

- b. The Neumann condition prescribes the normal derivative of the potential on the S_n boundary. That is,

$$\frac{\partial \phi}{\partial n} = \frac{\partial \bar{\phi}}{\partial n} \text{ on } S = S_n \quad (2)$$

- c. The mixed boundary condition prescribes a relation between ϕ and $\frac{\partial \phi}{\partial n}$ on the $S_{\phi,n}$ boundary. That is,

$$\phi = f\left(\frac{\partial \phi}{\partial n}\right) \text{ on } S = S_{\phi,n} \quad (3)$$

The third condition is found on the free-water surface in an aquifer and is not generally of concern to Petroleum Engineers.

The second half of the problem to be solved to complete the model is to derive equations that describe the rate of advance of the steam zone in a linear system made up of regions of different permeability (Fig. 2). These equations will be used to predict the oil recovery inside the generated streamtube.

CHAPTER V

MATHEMATICAL DEVELOPMENT

A common method of formulating the equations that describe the characteristics of a physical system is by considering some elemental portion of its domain and applying some physical law to it, such as the conservation principle. This method of formulation generally results in differential equations which can be solved analytically or numerically. The most widely-used numerical technique is the method of finite differences. In this method, the domain is divided into elements or cells and the differential operators are approximated by difference operators resulting in a

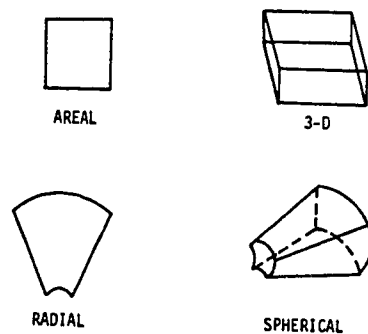


FIGURE 3. Elemental volumes⁴⁵

system of algebraic equations that are valid at the "nodes" of the cells within the domain. Because the domain of the problem is discretized, the method will be termed a domain method. Another way of solving the differential equations is by the Boundary Element Method. In this method, the differential equations in the domain are first transformed into integral equations which relate the unknown function and possibly certain of its derivatives to the given values on the boundary. The method of solution utilizes the superposition of known basic solutions to give the complete solution. After the differential equations that describe the system have been derived, they are first transformed into integral equations on the boundary before being solved by the boundary element technique.

The equations describing the potential distribution and flow of a single phase, incompressible fluid in a homogeneous porous medium is developed by taking an elemental volume of the reservoir and applying the law of conservation of mass to it. This is a procedure that can be found in standard texts.^{44,45,46} The result is called the continuity equation and can be expressed in general form as:

$$\nabla \cdot (\rho u) = - \frac{\partial}{\partial t} (\rho \phi) \quad (V.1.1)$$

Equation (V.1.1) is called the continuity equation where

$$\nabla \cdot (\rho u) = \text{divergence of } (\rho u)$$

$$\rho = \text{density}$$

u = volume flux per unit area

t = time

ϕ = porosity

Darcy's law relates the volume flux (u) to the potential gradient.

$$u = -\nabla \left(\frac{k}{\mu} \phi \right) \quad (V.1.2)$$

where

k = permeability

ϕ = potential

μ = fluid viscosity

Substituting Equation (V.1.2) into Equation (V.1.1) gives:

$$\nabla \cdot \left[\frac{\rho}{\mu} \nabla (k\phi) \right] = \frac{\partial}{\partial t} (\rho\phi) \quad (V.1.3)$$

Equation (V.1.3) is completely general. No assumptions have been made about the nature of the medium. It can be applied to any coordinate system, and is valid at any point of a three dimensional domain except at locations where there are sources or sinks.

The following simplifying assumptions are made in order to formulate the equations describing the physical system.

1. Incompressible, homogeneous, single phase fluid is assumed to be flowing in the system.
2. The medium is homogeneous and isotropic.
3. Two-dimensional flow exists. This means that flow

in the vertical direction is negligible. Therefore, flow is only in the horizontal x and y directions.

4. Gravitational effects are neglected.

Introducing these simplifying assumptions gives, in two dimensions:

$$\frac{\rho}{\mu} \frac{\partial}{\partial x} \left(k_x \frac{\partial \phi}{\partial x} \right) + \frac{\rho}{\mu} \frac{\partial}{\partial y} \left(k_y \frac{\partial \phi}{\partial y} \right) = - \frac{\partial}{\partial t} (\rho \phi) \quad (\text{V.1.4})$$

Making the further assumption of steady state flow, Equation (V.1.4) reduces to:

$$\frac{\partial}{\partial x} \left(k_x \frac{\partial \phi}{\partial x} \right) + \frac{\partial}{\partial y} \left(k_y \frac{\partial \phi}{\partial y} \right) = 0 \quad (\text{V.1.5})$$

For a homogeneous reservoir, $k_x = k_y = k$. Therefore, Equation (V.1.5) becomes:

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0 \quad (\text{V.1.6})$$

Equation (V.1.6) is known as Laplace's equation. It applies everywhere in the domain except where sources and sinks exist. When sources and sinks are present in the domain, Equation (V.1.6) changes to:

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\mu}{k} \sum_{j=1}^N q_j \delta(x - x_j, y - y_j) = 0 \quad (\text{V.1.7})$$

which is known as Poisson's equation, where

q_j = the volumetric flow rate per unit reservoir

thickness of the jth source or sink

N = the number of sources and sinks

δ = the Dirac delta function

x_j, y_j = the coordinates of the j 'th source or
sink location

For a heterogeneous reservoir made up of two or more homogeneous subregions, Darcy's law applies in each subregion. This means that Equation (V.1.6) as well as Equation (V.1.7) are applicable depending on whether there are sources or not. To obtain the potential distribution for the heterogeneous system, either Equation (V.1.6) or Equation (V.1.7) (depending upon the presence or absence of sources and sinks) is applied to each region. They are combined by taking care of the compatibility and continuity conditions on the various interfaces between adjacent regions.

V.2 TRANSFORMATION TO AN INTEGRAL FORMULATION

As stated earlier, when the domain is a heterogeneous one, made up of homogeneous subdomains, the method is to apply Laplace's or Poisson's equation to each subdomain. They are then combined (making sure to satisfy compatibility and continuity conditions on the common boundaries between the regions). Thus, the homogeneous subregions form the basic building blocks for the heterogeneous system. For this reason, the theory will be developed for a homogeneous domain.

Extensions to heterogeneous systems will be discussed in a later section. The objective, therefore, is to find

potential distributions that satisfy the differential equation

$$L\phi \equiv \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = -\frac{\mu}{k} \sum_{j=1}^N q_j \delta(x - x_j, y - y_j) \quad (V.2.1)$$

subject to:

$$\phi|_{S_\phi} = \bar{\phi}_{S_\phi} \quad (V.2.2)$$

$$\left. \frac{\partial \phi}{\partial n} \right|_{S_n} = (\bar{\phi}_n)_{S_n} \quad (V.2.3)$$

where L is the differential operator $\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$.

N is the number of sources and sinks

$\bar{\phi}_{S_\phi}$ is the specified potential on the S_ϕ boundary

$(\bar{\phi}_n)_{S_n}$ is the specified normal gradient $\left(\frac{\partial \phi}{\partial n} \right)$ on the S_n boundary

The first step in the integral formulation of the problem is to form the product of a function W with both sides of Equation (V.2.1) and integrate over the domain of interest. This is called the inner, or scalar, or dot product of W with $L\phi$ and is denoted $\langle W, L\phi \rangle$.⁴⁷ By definition,

$$\langle W, L\phi \rangle = \int_D W \{L\phi\} dD \quad (V.2.4)$$

L is the linear differential operator $\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$ on ϕ .

W is an arbitrary function which is sufficiently differentiable for $L^*\{W\}$ to exist, where L^* is the formal adjoint

differential operator associated with L . Integration by parts is performed on (V.2.4) until the operator on ϕ disappears (i.e., integrating twice by parts for two-dimensional problems). The result is of the form

$$\int_D W\{L\phi\} dD = \{\text{boundary terms}\} + \int_D \phi\{L^*W\} dD \quad (\text{V.2.5})$$

In two dimensions, the boundary terms are line integrals; and in three dimensions, they are surface integrals. Forming the inner product $\langle L\phi, W \rangle$ for the left hand side of Equation (V.1.2) gives

$$\langle L\phi, W \rangle = \int_D \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) W dD \quad (\text{V.2.6})$$

where D represents the domain of integration. Integration of (V.2.6) by parts is done by first rewriting it as a double integral:

$$\begin{aligned} \langle L\phi, W \rangle = & \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial x^2} W dx dy \\ & + \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial y^2} W dx dy \end{aligned} \quad (\text{V.2.7})$$

Integrate each term by parts to give, for the first term,

$$\begin{aligned} \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial x^2} W dx dy = & \int_{y_1}^{y_2} \left[W \frac{\partial \phi}{\partial x} \right]_{x_1}^{x_2} dy \\ & - \int_{x_1(y)}^{x_2(y)} \frac{\partial \phi}{\partial x} \frac{\partial W}{\partial x} dx dy \end{aligned} \quad (\text{V.2.8})$$

Integrate the last term on the R.H.S. of (V.2.8) by parts and substitute back into (V.2.8) to give

$$\begin{aligned}
 \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial x^2} W \, dx \, dy &= \int_{y_1}^{y_2} \left[W \left(\frac{\partial \phi}{\partial x} \right) \Big|_{x_1(y)}^{x_2(y)} \right. \\
 &\quad \left. - \frac{\partial W}{\partial x} (\phi) \Big|_{x_1(y)}^{x_2(y)} + \int_{x_1(y)}^{x_2(y)} \phi \frac{\partial^2 W}{\partial x^2} \, dx \right] dy \\
 &= \int_{y_1}^{y_2} W \left(\frac{\partial \phi}{\partial x} \right) \Big|_{x_1(y)}^{x_2(y)} dy - \int_{y_1}^{y_2} \frac{\partial W}{\partial x} \left(\phi \right) \Big|_{x_1(y)}^{x_2(y)} dy \\
 &\quad + \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \phi \frac{\partial^2 W}{\partial x^2} \, dx \, dy \tag{V.2.9}
 \end{aligned}$$

Now

$$dy = ds \cos \theta$$

and

$$\begin{aligned}
 \cos \theta &= \hat{i} \cdot \hat{n} \\
 \therefore dy &= \hat{i} \cdot \hat{n} \, ds \tag{V.2.10}
 \end{aligned}$$

where $\hat{n}$ is the outward unit normal and ds is the differential element of arc length along the boundary S .

$\hat{i}$ is the unit vector in the x direction.

Substituting for dy into (V.2.9) gives

$$\begin{aligned}
 \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial x^2} W \, dx \, dy &= \int_S \left(W \frac{d\phi}{dx} \right) \hat{i} \cdot \hat{n} \, ds - \int_S \left(\frac{\partial W}{\partial x} \phi \right) \hat{i} \cdot \hat{n} \, ds \\
 &\quad + \int_D \phi \frac{\partial^2 W}{\partial x^2} \, dD \tag{V.2.11}
 \end{aligned}$$

By a similar procedure

$$\begin{aligned} \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \frac{\partial^2 \phi}{\partial x^2} W \, dx \, dy &= \int_S \left(W \frac{\partial \phi}{\partial y} \right) \hat{j} \cdot \hat{n} \, ds - \int_S \left(\frac{\partial W}{\partial y} \phi \right) \hat{j} \cdot \hat{n} \, ds \\ &+ \int_D \frac{\partial^2 W}{\partial y^2} \, dD \end{aligned} \quad (V.2.12)$$

Combining (V.2.11) and (V.2.12) gives

$$\begin{aligned} \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \left(\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) W \, dx \, dy &= \int_S W \left(\frac{\partial \phi}{\partial x} \hat{i} + \frac{\partial \phi}{\partial y} \hat{j} \right) \cdot \hat{n} \, ds \\ &- \int_S \phi \left(\frac{\partial W}{\partial x} \hat{i} + \frac{\partial W}{\partial y} \hat{j} \right) \cdot \hat{n} \, ds + \int_D \phi \left(\frac{\partial^2 W}{\partial x^2} + \frac{\partial^2 W}{\partial y^2} \right) \, dD \end{aligned} \quad (V.2.13)$$

where $\hat{j}$ is the unit vector in the y-direction.

Therefore, the inner product of the left hand side of Equation (V.2.1) gives:

$$\langle L\phi, W \rangle = \int_S W \frac{\partial \phi}{\partial n} \, ds - \int_S \phi \frac{\partial W}{\partial n} \, ds + \int_D \left(\frac{\partial^2 W}{\partial x^2} + \frac{\partial^2 W}{\partial y^2} \right) \, dD \quad (V.2.14)$$

The inner product of the right hand side of Equation (V.2.1) is:

$$\begin{aligned} \langle \sum_{j=1}^N q_j \delta(x - x_j, y - y_j), W \rangle &= \frac{\mu}{k} \int_{y_1}^{y_2} \int_{x_1(y)}^{x_2(y)} \sum_{j=1}^N q_j \delta(x - x_j, y - y_j) W \, dx \, dy \end{aligned} \quad (V.2.15)$$

which is simply,

$$\frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) W(x_j, y_j) \quad (\text{V.2.16})$$

Therefore, from Equations (V.2.14) and (V.2.16), the inner product of both sides of Equation (V.2.1) is:

$$\begin{aligned} \int_S W \frac{\partial \Phi}{\partial n} ds - \int_S \Phi \frac{\partial W}{\partial n} ds + \int_D \Phi \nabla^2 W dD \\ = - \frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) W(x_j, y_j) \end{aligned} \quad (\text{V.2.17})$$

where ∇^2 is the two-dimensional Laplacian differential operator denoting the vector dot product $\nabla \cdot \nabla$. ∇ is the gradient operator defined in cartesian coordinates as

$$\nabla = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} .$$

Thus,

$$\nabla^2 = \nabla \cdot \nabla = \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} \right) \cdot \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} \right) = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$\hat{i}$ and $\hat{j}$ are unit vectors in the x and y directions, respectively. Equation (V.2.17) can be rewritten as:

$$\begin{aligned} \int_D \Phi \nabla^2 W dD = \int_S \Phi \frac{\partial W}{\partial n} ds - \int_S W \frac{\partial \Phi}{\partial n} ds \\ - \frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) W(x_j, y_j) \end{aligned} \quad (\text{V.2.18})$$

Now, by a clever choice of W , the left part of Equation (V.2.18), that is, the domain integral, can be reduced to

the potential ϕ only. When this is done, the original differential equation would have been transformed into an integral equation involving just boundary integrals only. The value of W to achieve this is that W which satisfies the equation

$$\nabla^2 W = -\delta(x - x_i, y - y_i) \quad (V.2.19)$$

without satisfying any boundary conditions. Equation (V.2.19) happens to be the equation describing the potential distribution that would occur due to the application of a unit charge at (x_i, y_i) , and the solution to it is variously called the fundamental solution, the free-space Green's function, or the principal solution. Using the right side of Equation (V.2.19) to substitute for $\nabla^2 W$ in Equation (V.2.18) gives:

$$\begin{aligned} -\int_D \phi \delta(x - x_i, y - y_i) dD = \int_S W \frac{\partial \phi}{\partial n} ds - \int_S \phi \frac{\partial W}{\partial n} ds \\ + \frac{\mu}{k} \sum_{j=1}^N q_j(x_i, y_j) W(x_j, y_j) \end{aligned} \quad (V.2.20)$$

The left side of Equation (V.2.20) is simply $-\phi(x_i, y_i)$; thus

$$\begin{aligned} \phi(x_i, y_i) = \int_S W \frac{\partial \phi}{\partial n} ds - \int_S \phi \frac{\partial W}{\partial n} ds \\ + \frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) W(x_j, y_j) \end{aligned} \quad (V.2.21)$$

Equation (V.2.21) relates the potential at an interior point (x_i, y_i) to both its value on the boundary and that

of its normal derivative to the boundary. It contains the function W which is the solution to Equation (V.2.19) and for this work will be called the fundamental solution given as,⁴⁸

$$W = \frac{1}{2\pi} \ln \frac{1}{r} \quad (\text{V.2.22})$$

where r is the distance between an observation point (x,y) and the point of application of the unit charge (x_i, y_i) . If the observation point is anywhere within the domain (except at the source/sink locations),

$$r = r_i = \sqrt{(x - x_i)^2 + (y - y_i)^2} \quad (\text{V.2.23})$$

and

$$W = \frac{1}{2\pi} \ln \frac{1}{\sqrt{(x - x_i)^2 + (y - y_i)^2}} \quad (\text{V.2.24})$$

Let

$$r_{i,j} = \sqrt{(x_j - x_i)^2 + (y_j - y_i)^2} \quad (\text{V.2.25})$$

be the distance between the unit charge point (x_i, y_i) and the source point (x_j, y_j) ; then

$$W(x_j, y_j) = \frac{1}{2\pi} \ln \frac{1}{\sqrt{(x_j - x_i)^2 + (y_j - y_i)^2}} \quad (\text{V.2.26})$$

Equation (V.2.21) can be rewritten as:

$$\begin{aligned} \phi(x_i, y_i) = & \frac{1}{2\pi} \int_S \frac{\partial \phi}{\partial n} \left(\ln \frac{1}{r_i} \right) ds - \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_i} \right) ds \\ & + \frac{u}{2\pi k} \sum_{j=1}^N q_j(x_j, y_j) \ln \frac{1}{r_{i,j}} \end{aligned} \quad (\text{V.2.27})$$

It is of interest to note that Equation (V.2.27) could have been obtained by different methods, such as application of Green's second identity directly to Equation (V.1.7), or by the method of weighted residuals.¹²

V.3 REPRESENTATION AS A BOUNDARY INTEGRAL EQUATION

Equations (V.2.21) or (V.2.27) are integral formulations of the differential equation (V.2.1). In order to apply the boundary element solution method, Equation (V.2.22) or (V.2.27) must be transformed to a boundary integral equation. This means that the unknown $\phi(x_i, y_i)$ on the left of Equation (V.2.22) or (V.2.27) should be in terms of boundary coordinates. Therefore, the point of application of the unit charge (Equation V.2.20) has to be on the boundary. Let the boundary point of application of the unit charge be denoted (x_b, y_b) , then the fundamental solution is, then:

$$W(x_b, y_b) = \frac{1}{2\pi} \ln \frac{1}{r_b} \quad (V.3.1)$$

where r_b is the distance from any point (x, y) to the boundary unit charge point (x_b, y_b) and is given as

$$r_b = \sqrt{(x - x_b)^2 + (y - y_b)^2} \quad (V.3.2)$$

The distance from the source/sink locations (x_j, y_j) to the boundary unit charge point is:

$$r_{b,j} = \sqrt{(x_b - x_j)^2 + (y_b - y_j)^2} \quad (V.3.3)$$

By taking the point of application of the unit charge to the boundary, the boundary integral formulation of the problem becomes:

$$\begin{aligned} \phi(x_b, y_b) = & \int_S W \frac{\partial \phi}{\partial n} ds - \int_S \frac{\partial W}{\partial n} ds \\ & + \frac{\mu}{k} \sum_{j=1}^{N_j} q_j(x_j, y_j) W(x_j, y_j) \end{aligned} \quad (V.3.4)$$

Substituting for $W(x_b, y_b)$ and $W(x_j, y_j)$ using Equations (V.3.2) and (V.3.3) gives:

$$\begin{aligned} \phi(x_b, y_b) = & \frac{1}{2\pi} \int_S \frac{\partial \phi}{\partial n} \ln \frac{1}{r_b} ds - \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds \\ & + \frac{\mu}{2\pi k} \sum_{j=1}^N q_j(x_j, y_j) \ln \frac{1}{r_{b,j}} \end{aligned} \quad (V.3.5)$$

where r_b and $r_{b,j}$ are given by Equations (V.3.2) and (V.3.3).

The evaluation of the right hand side of Equation (V.3.5) gives the potential at any boundary point with coordinates (x_b, y_b) .

V.3.1 EVALUATION OF THE IMPROPER INTEGRALS

All the integrals involve the term $\ln \frac{1}{r_b}$ which has a singularity at the point $r_b = 0$. These kinds of integrals whose integrands become infinite at some point in the domain of integration are known as improper integrals. The usual method of evaluating such integrals is as follows:^{12,48}

- a. For a two-dimensional problem, take a small circle or semi-circle of radius ϵ around the point of singularity. This divides the domain of integration (the boundary) into two parts; one part contains the singularity and the other part does not (Figure 4).
- b. Evaluate both integrals and take the limit as ϵ tends to zero. The result is the required value of the integral.

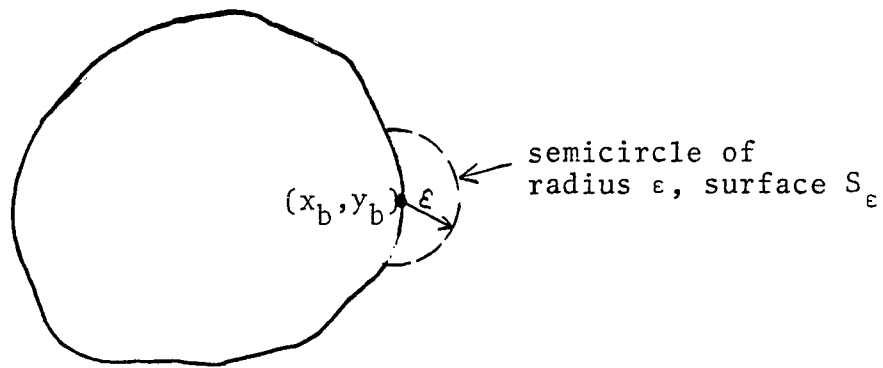


FIGURE 4. Method of evaluating improper integral.

On the part of the boundary $S - 2\epsilon$, $\ln \frac{1}{r_b}$ is continuous. Following the procedure outlined above, the integrals in Equation (V.3.5) can be evaluated as follows: For the second integral on the right of Equation (V.3.5),

$$\begin{aligned} \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds &= \frac{1}{2\pi} \int_{S-S_\epsilon} \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds \\ &+ \frac{1}{2\pi} \int_{S_\epsilon} \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds \end{aligned} \quad (V.3.6)$$

For the last integral of Equation (V.3.6), the normal direction is the same as the radial direction. That is,

$$\left. \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) \right|_{r_b=\epsilon} = \frac{\partial}{\partial r_b} \left(\ln \frac{1}{r_b} \right) \bigg|_{r_b=\epsilon} = -\frac{1}{\epsilon} \quad (V.3.7)$$

The last term of Equation (V.3.6) now becomes:

$$\frac{1}{2\pi} \int_{S_\epsilon} \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds = -\frac{1}{2\pi\epsilon} \int_{S_\epsilon} \phi ds \quad (V.3.8)$$

For the case where the surface S_ϵ is a semi-circle, Equation (V.3.8) becomes:

$$\begin{aligned} \frac{1}{2\pi\epsilon} \int_{S_\epsilon} \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds &= -\frac{1}{2\pi\epsilon} \int_{\frac{2\pi\epsilon}{2}} \phi ds = -\frac{1}{2\pi\epsilon} \left(\frac{2\pi\epsilon}{2} \phi^* \right) \\ &= -\frac{1}{2} \phi^* \end{aligned} \quad (V.3.9)$$

where ϕ^* is an average value of $\phi(x,y)$ on S_ϵ , and in the limit as ϵ tends to zero

$$\lim_{\epsilon \rightarrow 0} \left(-\frac{1}{2} \phi^* \right) = -\frac{1}{2} \phi(x_b, y_b) \quad (V.3.10)$$

The first integral on the right of Equation (V.3.5) is split into two integrals as:

$$\begin{aligned} \frac{1}{2\pi} \int_S \ln \frac{1}{r_b} \frac{\partial \phi}{\partial n} ds &= \frac{1}{2\pi} \int_{S-S_\epsilon} \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds \\ &+ \frac{1}{2\pi} \int_{S_\epsilon} \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds \end{aligned} \quad (V.3.11)$$

Following the procedure leading to Equation (V.3.9) gives

$$\begin{aligned} \frac{1}{2\pi} \int_{S_\epsilon} \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds &= - \frac{1}{2\pi\epsilon} \left(\frac{2\pi\phi}{2} \right) \left(\frac{\partial \phi}{\partial n} \right)^* \\ &= - \frac{1}{2} \frac{\partial \phi^*}{\partial n} \end{aligned} \quad (V.3.12)$$

where $\frac{\partial \phi^*}{\partial n}$ is an average value of $\frac{\partial \phi}{\partial n}$ on S_ϵ and

$$\lim_{\epsilon \rightarrow 0} \left(- \frac{1}{2} \frac{\partial \phi^*}{\partial n} \right) = 0 \quad (V.3.13)$$

From (V.3.6)

$$\lim_{\epsilon \rightarrow 0} \frac{1}{2\pi} \int_{S-S_\epsilon} \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds = \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds \quad (V.3.14)$$

From (V.3.11)

$$\lim_{\epsilon \rightarrow 0} \frac{1}{2\pi} \int_{S-S_\epsilon} \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds = \frac{1}{2\pi} \int_S \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds \quad (V.3.14a)$$

Substituting (V.3.10), (V.3.13), (V.3.14), and (V.3.14a) into (V.3.5) gives the potential at any boundary point (x_b, y_b) as:

$$\begin{aligned}
\phi(x_b, y_b) &= \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds - \frac{1}{2\pi} \int_S \left(\ln \frac{1}{r_b} \right) \frac{\partial \phi}{\partial n} ds \\
&+ \frac{1}{2} \phi(x_b, y_b) + \sum q_j(x_j, y_j) \frac{1}{2\pi} \ln \frac{1}{r_{b,j}} \quad (V.3.15)
\end{aligned}$$

Therefore,

$$\begin{aligned}
\frac{1}{2} \phi(x_b, y_b) &= \frac{1}{2\pi} \int_S \ln \frac{1}{r_b} \left(\frac{\partial \phi}{\partial n} \right) ds - \frac{1}{2\pi} \int_S \phi \frac{\partial}{\partial n} \left(\ln \frac{1}{r_b} \right) ds \\
&+ \frac{\mu}{k} \sum q_j(x_j, y_j) \frac{1}{2\pi} \ln \frac{1}{r_{b,j}} \quad (V.3.16)
\end{aligned}$$

Equation (V.3.16) is the boundary integral formulation of the problem. The boundary S is made up of two parts, S_1 and S_2 such that $S = S_1 + S_2$. On S_1 , ϕ is prescribed and $\frac{\partial \phi}{\partial n}$ is unknown. On S_2 , $\frac{\partial \phi}{\partial n}$ is prescribed and ϕ is unknown.

CHAPTER VI

SOLUTION BY DISCRETIZATION OF THE BOUNDARY

The integral equations (V.2.26) and V.3.16) are solved numerically by discretizing the boundaries of the domain into elements (hence the name: Boundary Element Method). A homogeneous medium will be assumed. Later, the method will be modified to handle heterogeneous media made up of homogeneous regions. The boundary of the domain is divided into M straight line segments as shown in Figure 5.

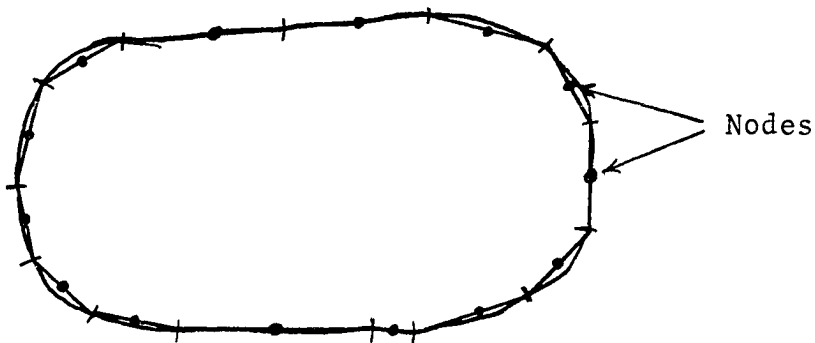


FIGURE 5

Straight-line boundary elements or segments

The unknown values of ϕ and $\frac{\partial \phi}{\partial n}$ are evaluated at the mid-points of each segment. These points are called the node points. In the present work, the node points will always be at the mid-points of straight line boundary segments. The cases where the boundary segments are curved and cases where the node points are at the intersections of the boundary elements are treated in Brebbia.¹² Of the M boundary segments, M_1 belong to the S_ϕ type boundary where ϕ is specified as constant ($\bar{\phi}$) on each segment while $\frac{\partial \phi}{\partial n}$ is unknown. The remaining M_2 segments belong to the S_n type boundary where $\frac{\partial \phi}{\partial n}$ is specified as constant $\left(\frac{\partial \bar{\phi}}{\partial n}\right)$ on each segment and ϕ is unknown. Equation (V.3.16), when put in discrete form over the two types of boundaries, is:

$$\begin{aligned}
 \frac{1}{2} \phi(x_b, y_b) &= \frac{1}{2\pi} \sum_{L=1}^{M_1} \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \\
 &+ \frac{1}{2\pi} \sum_{L=1}^{M_2} \left(\frac{\partial \bar{\phi}}{\partial n} \right)_L \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \\
 &- \frac{1}{2\pi} \sum_{L=1}^{M_1} \bar{\phi}_L \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) dS_L \\
 &- \frac{1}{2\pi} \sum_{L=1}^{M_2} \phi_L \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) dS_L \\
 &+ \frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) \frac{1}{2\pi} \ln \frac{1}{r_{b,j}}
 \end{aligned} \tag{VI.0.1}$$

which can be combined as:

$$\begin{aligned}
 \frac{1}{2} \phi(x_b, y_b) &= \frac{1}{2\pi} \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \\
 &\quad - \frac{1}{2\pi} \sum_{L=1}^M \phi_L \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) dS_L \\
 &\quad + \frac{\mu}{k} \sum_{j=1}^N q_j(x_j, y_j) \frac{1}{2\pi} \ln \frac{1}{r_{b,j}}
 \end{aligned} \tag{VI.0.2}$$

where

$$M = M_1 + M_2$$

$$S = S_\phi + S_n$$

$\bar{\phi}_L$ = given boundary condition of ϕ on S_ϕ of the Lth segment

$\left(\frac{\partial \bar{\phi}}{\partial n} \right)_L$ = given boundary condition of $\frac{\partial \phi}{\partial n}$ on S_n of the Lth segment

Let

$$G_{b,L} = \frac{1}{2\pi} \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \tag{VI.0.3}$$

$$H_{b,L} = \frac{1}{2\pi} \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) dS_L \tag{VI.0.4}$$

Equation (VI.0.2) then becomes:

$$\begin{aligned}
 \frac{1}{2} \phi(x_b, y_b) &= \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_L [G_{b,L}] - \sum_{L=1}^M \phi_L [H_{b,L}] \\
 &\quad + \frac{\mu}{2\pi k} \sum_{j=1}^N q_j(x_j, y_j) \ln \frac{1}{r_{b,j}}
 \end{aligned} \tag{VI.0.5}$$

$G_{b,L}$ and $H_{b,L}$ are evaluated for all the segments all around the boundary, including the segment containing the charge point (x_b, y_b) . Therefore, there are two kinds of integrals to be evaluated, namely:

1. Integration over the segment b that contains the unit charge point. This segment contains a singularity point at (x_b, y_b) . Integration over this segment happens when $L = b$.
2. Integration over the rest of the segments. These are the condition when $L \neq b$.

Because of this need to evaluate two kinds of integrals, the boundary integral equation (VI.0.5) can be rewritten to reflect these two kinds of integrals by defining the following:

$$G_{b,L} = \begin{cases} \check{G}_{b,L} & \text{when } L \neq b \\ G_{b,b} & \text{when } L = b \end{cases}$$

$$H_{b,L} = \begin{cases} \check{H}_{b,L} & \text{when } L \neq b \\ H_{b,b} & \text{when } L = b \end{cases}$$

Equation (VI.0.5) now becomes:

$$\begin{aligned} \frac{1}{2} \phi(x_b, y_b) = & \left[\left(\frac{\partial \phi}{\partial n} \right)_{L=b} G_{b,b} + \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_{L \neq b} \check{G}_{b,L} \right] \\ & - \left[\phi(x_b, y_b) H_{b,b} + \sum_{L=1}^M (\phi)_{L \neq b} \check{H}_{b,L} \right] \\ & + \frac{\mu}{2\pi k} \sum_{j=1}^N q_j(x_j, y_j) \ln \frac{1}{r_{b,j}} \end{aligned} \quad (\text{VI.0.6})$$

which can be rearranged to give (where $\phi(x_b, y_b) = \phi_{L=b}$)

$$\begin{aligned}
 & \left[\phi_{L=b} \left(H_{b,b} + \frac{1}{2} \right) + \sum_{L=1}^M (\phi)_{L \neq b} \check{H}_{b,L} \right] \\
 & - \left[\left(\frac{\partial \phi}{\partial n} \right)_{L=b} G_{b,b} + \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_{L \neq b} \check{G}_{b,L} \right] \\
 & = \frac{\mu}{2\pi k} \sum_{j=1}^N q_j(x_j, y_j) \ln \frac{1}{r_{bj}} \quad (VI.0.7)
 \end{aligned}$$

V.1 EVALUATION OF THE INTEGRALS

Equation (VI.0.7) requires the evaluation of the following integrals:

$$\begin{aligned}
 G_{b,L} &= \frac{1}{2\pi} \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \begin{cases} \text{for all } L \neq b, \text{ denoted as } \check{G}_{b,L} \\ \text{for all } L = b, \text{ denoted as } G_{b,b} \end{cases} \\
 H_{b,L} &= \frac{1}{2\pi} \int_{S_L} \frac{\partial}{\partial n} \ln \frac{1}{r_{b,L}} dS_L \begin{cases} \text{for all } L \neq b, \text{ denoted as } \check{H}_{b,L} \\ \text{for } L = b, \text{ denoted as } H_{b,b} \end{cases}
 \end{aligned}$$

These integrals can easily be evaluated analytically for the present case where the segments S_L are straight lines. However, for generality, they are expressed in dimensionless form and evaluated numerically using a four-point Gaussian quadrature method with the formula:

$$\int_{-1}^1 f(\alpha) d\alpha = \sum_{k=1}^n W_k f(\alpha_k) + E_n \quad (VI.1.1)$$

where

W_k = weighting factor

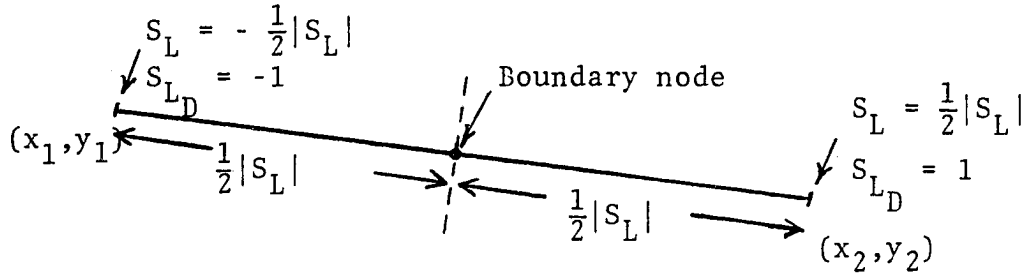
α_k = coordinate of the k'th integration point

n = total number of integration points ($n = 4$)

E_n = error term

In order to apply the quadrature formula to $G_{i,L}$ and $H_{i,L}$, they must be transformed to integrals with limits from -1 to +1. This is achieved by expressing the variable of integration dS in dimensionless form as:

For any element where $L \neq b$:



define a dimensionless variable S_{L_D} as:

$$S_{L_D} = \frac{S_L}{\frac{1}{2}|S_L|}, \quad \text{then} \quad dS_L = \frac{1}{2}|S_L| dS_{L_D}$$

where

$$S_{L_D} = 0 \text{ when } S_L = 0$$

$$S_{L_D} = -1 \text{ when } S_L = -\frac{1}{2}|S_L|$$

$$S_{L_D} = 1 \text{ when } S_L = \frac{1}{2}|S_L|$$

and

$$S_L = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}$$

Therefore, from Equation (A.1.10) of Appendix A,

$$\begin{aligned}
 H_{b,L} &= \frac{1}{2\pi} \int_{S_L} \frac{\partial}{\partial n_L} \left(\ln \frac{1}{r_{b,L}} \right) ds \\
 &= - \frac{1}{4\pi} \int_{-1}^1 \frac{(x_L - x_b)(y_2 - y_1) + (y_L - y_b)(x_2 - x_1)}{(x_L - x_b)^2 + (y_L - y_b)^2} dS_D
 \end{aligned} \tag{VI.1.2}$$

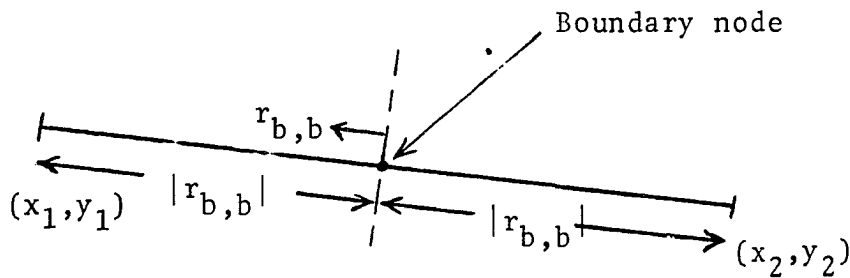
From Equation (A.1.11) of Appendix A

$$\begin{aligned}
 G_{b,L} &= \frac{1}{2\pi} \int_{S_L} \ln \frac{1}{r_{b,L}} ds \\
 &= \frac{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}}{8\pi} \int_{-1}^1 \ln \frac{1}{\sqrt{(x_L - x_b)^2 + (y_L - y_b)^2}} dS_D
 \end{aligned} \tag{VI.1.3}$$

For the element where $L = b$:

For this case, the vector r_b is along the length of the segment and is equal to one-half the length of the segment.

Therefore, $ds = dr$.



Since $r_{b,b}$ is along the segment,

$$\frac{\partial}{\partial n_b} = \frac{\partial}{\partial r_{b,b}} \cdot \frac{\partial r_{b,b}}{\partial n_b}$$

but $\frac{\partial r_{b,b}}{\partial n_b} = 0$ because $r_{b,b}$ is perpendicular to n_b . Define the dimensionless variable

$$r_D = \frac{r_{b,b}}{|r_{b,b}|} \quad dr = |r_{b,b}| dr_D$$

From Equations (A.1.12) and (A.1.15) of Appendix A,

$$H_{b,b} = \frac{1}{2\pi} \int_{S_b} \frac{\partial}{\partial n_b} \left(\ln \frac{1}{r_{b,b}} \right) ds_b = 0 \quad (\text{VI.1.4})$$

$$G_{b,b} = \frac{1}{2\pi} \int_{S_b} \ln \frac{1}{r_{b,b}} ds_b = \frac{|r_{b,b}|}{\pi} \left(\ln \frac{1}{|r_{b,b}|} - 1 \right) \quad (\text{VI.1.5})$$

$$G_{b,b} = \frac{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}}{2\pi} \left(\ln \frac{1}{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}} - 1.0 \right) \quad (\text{VI.1.6})$$

No integration is therefore required for the case when $L = b$. Applying the quadrature formula to Equations (VI.1.2) and (VI.1.3) result in:

$$H_{b,L} = - \frac{1}{4\pi} \sum_{k=1}^4 w_k \left(\frac{(x_{L_k} - x_b)(y_2 - y_1) + (y_{L_k} - y_b)(x_2 - x_1)}{(x_{L_k} - x_b)^2 + (y_{L_k} - y_b)^2} \right) \quad (\text{VI.1.7})$$

$$G_{b,L} = \frac{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}}{4\pi} \sum_{k=1}^4 w_k \left[\ln \left(\frac{1}{\sqrt{(x_{L_k} - x_b)^2 + (y_{L_k} - y_b)^2}} \right) \right] \quad (\text{VI.1.8})$$

where (x_{L_k}, y_{L_k}) are the quadrature points. The values of $H_{b,L}$ and $G_{b,L}$ for all the boundary segments can now be calculated using Equations (VI.1.4) through (VI.1.8). Since $H_{b,b} = 0$, Equation (VI.0.10) becomes:

$$\begin{aligned} \frac{1}{2} \phi_b &= \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_L G_{b,L} - \sum_{L=1}^M \phi_L H_{b,L} + \frac{1}{2\pi} \sum_{j=1}^n q_j(x_j, y_j) \\ &\quad \times \ln \frac{1}{r_{b,j}} \end{aligned} \quad (\text{VI.1.9})$$

which is put in the form

$$\begin{aligned} \sum_{L=1}^M \phi_L H_{b,L} - \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_L G_{b,L} \\ = \frac{\mu}{2\pi k} \sum_{j=1}^n q_j(x_j, y_j) \ln \frac{1}{r_{b,j}} \end{aligned} \quad (\text{VI.1.10})$$

Note that when $L = b$, the first term on the left of Equation (VI.1.10) is $\phi_b(\frac{1}{2} + H_{b,b})$. But, since $H_{b,b} = 0$ when $L = b$, the first term on the left of Equation (VI.1.10) becomes $\frac{1}{2} \phi_b$.

CHAPTER VII

DETERMINATION OF THE POTENTIAL AND VELOCITY AT INTERIOR POINTS

Equation (VI.1.10) is written for every node in the system; that is, for $b = 1, 2, \dots, m$ (number of nodes). This will result in m equations in m unknowns which can be put in matrix notation (Brebbia¹²) as:

$$[H]\phi = [G]\phi_n + B \quad (\text{VII.0.1})$$

where $\phi_n = \frac{\partial \phi}{\partial n}$ = the derivative normal to the boundary

$$B = \frac{u}{2\pi k} \sum_{j=1}^n q_j(x_j, y_j) \ln \frac{1}{r_{b,j}}$$

B is a column vector, H and G are m by m matrices.

Let $m = m_1 + m_2$ where m = number of elements. In the L.H.S. of (VII.0.1), m_1 values of ϕ are specified as boundary conditions, leaving m_2 unknown values of ϕ . Within the R.H.S. of (VII.0.1), m_2 values of $\left(\frac{\partial \phi}{\partial n}\right)$ are specified in the boundary conditions, leaving m_1 unknown values of $\left(\frac{\partial \phi}{\partial n}\right)$. All the values of the B vector are known. Therefore, Equation (VII.0.1) has m equations and a total of m unknowns which can be rearranged as:

$$Ax = F \quad (\text{VII.0.2})$$

where x now is a vector containing the m unknowns, m_1 of which are the unknown $\frac{\partial \phi}{\partial n}$'s and m_2 of which are the unknown ϕ 's on the boundaries. $m_1 + m_2 = m$. The matrix system of Equation (VII.0.2) can be solved by Gauss elimination method to give the m_1 values of $\frac{\partial \phi}{\partial n}$ and m_2 values of ϕ on the boundaries.

VII.1 POTENTIAL

After the solution, all the boundary values of ϕ and $\frac{\partial \phi}{\partial n}$ are now known and are used to determine the potentials and velocities at any desired interior point.

Interior points where the potentials are desired are denoted with the coordinates (x_i, y_i) . Therefore, in this case,

$$r_{i,L} = \sqrt{(x_L - x_i)^2 + (y_L - y_i)^2} \quad (\text{VII.1.1})$$

$$r_{i,j} = \sqrt{(x_j - x_i)^2 + (y_j - y_i)^2} \quad (\text{VII.1.2})$$

The potential at any interior point can be determined by Equation (V.2.27) which in discrete form is:

$$\begin{aligned} \phi(x_i, y_i) = & \sum_{L=1}^M \left(\frac{\partial \phi}{\partial n} \right)_L G_{i,L} - \sum_{L=1}^M \phi_L H_{i,L} \\ & + \frac{\mu}{2\pi k} \sum_{j=1}^n q_j(x_j, y_j) \ln \frac{1}{r_{i,j}} \end{aligned} \quad (\text{VII.1.3})$$

where:

$$G_{i,L} = \frac{1}{2\pi} \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{i,L}} \right) dS \quad (\text{VII.1.4})$$

$$H_{i,L} = \frac{1}{2\pi} \int_{S_L} \ln \frac{1}{r_{i,L}} dS \quad (\text{VII.1.5})$$

$r_{i,L}$ is the distance from the interior point (x_i, y_i) to the mid-point of the L 'th boundary point (x_L, y_L) . $r_{i,L}$ and r_{ij} are given by Equations (VII.1.1) and (VII.1.2) respectively.

After the application of the quadrature formula, $H_{i,L}$ and $G_{i,L}$ can be evaluated as:

$$H_{i,L} = - \frac{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}}{4\pi} d_{i,L} \times \sum_{k=1}^4 W_k \frac{1}{\sqrt{(x_i - x_{L_k})^2 + (y_i - y_{L_k})^2}} \quad (\text{VII.1.6})$$

$$G_{i,L} = \frac{\sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}}{4\pi} \sum_{k=1}^4 W_k \ln \frac{1}{\sqrt{(x_{L_k} - x_i)^2 + (y_{L_k} - y_i)^2}} \quad (\text{VII.1.7})$$

VII.2 VELOCITY

The velocities in the x and y directions at any point in the interior of the domain are obtained by Darcy's law:

$$v_x = - \frac{k}{\mu\phi} \frac{\partial}{\partial x} \phi(x, y) \quad (\text{VII.2.1})$$

$$v_y = - \frac{k}{\mu\phi} \frac{\partial \phi}{\partial y} \quad (\text{VII.2.2})$$

Utilizing (VII.1.4) and (VII.1.5), the potential at an interior point (x_i, y_i) can be written as:

$$\begin{aligned} \phi(x_i, y_i) = & \frac{1}{2\pi} \sum_{L=1}^m \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} \ln \left(\frac{1}{r_{i,L}} \right) dS_L \\ & - \frac{1}{2\pi} \sum_{L=1}^m \phi_L \int_{S_L} \frac{\partial}{\partial n} \left(\ln \frac{1}{r_{i,L}} \right) dS_L + \frac{\mu}{k} \frac{1}{2\pi} \sum_{j=1}^n q_j \ln \left(\frac{1}{r_{i,j}} \right) \end{aligned} \quad (\text{VII.2.3})$$

From Appendix A,

$$\frac{\partial}{\partial n} \left(\ln \frac{1}{r_{i,L}} \right) = - \frac{d_{i,L}}{r_{i,L}^2} \quad (\text{VII.2.4})$$

Substituting (VII.2.4) into (VII.2.3) and differentiating with respect to x and y gives:

$$\begin{aligned} \frac{\partial \phi(x_i, y_i)}{\partial x_i} = & \frac{1}{2\pi} \sum_{L=1}^m \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} \frac{\partial}{\partial x_i} \left(\ln \frac{1}{r_{i,L}} \right) dS_L \\ & - \frac{1}{2\pi} \sum_{L=1}^m \phi_L \int_{S_L} \frac{\partial}{\partial x_i} \left(\frac{-d_{i,L}}{r_{i,L}^2} \right) dS_L + \frac{\mu}{k} \frac{1}{2\pi} \sum_{j=1}^n q_j \frac{\partial}{\partial x_i} \ln \left(\frac{1}{r_{i,j}} \right) \end{aligned} \quad (\text{VII.2.5})$$

Similarly,

$$\begin{aligned} \frac{\partial \phi(x_i, y_i)}{\partial y_i} = & \frac{1}{2\pi} \sum_{L=1}^m \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} \frac{\partial}{\partial y_i} \left(\ln \frac{1}{r_{i,L}} \right) dS_L \\ & - \frac{1}{2\pi} \sum_{L=1}^m \phi_L \int_{S_L} \frac{\partial}{\partial y_i} \left(\frac{-d_{i,L}}{r_{i,L}^2} \right) dS_L + \frac{\mu}{k} \frac{1}{2\pi} \sum_{j=1}^n q_j \frac{\partial}{\partial y_i} \ln \left(\frac{1}{r_{i,j}} \right) \end{aligned} \quad (\text{VII.2.6})$$

where:

$$\begin{aligned}
 r_{i,L} &= \sqrt{(x_i - x_L)^2 + (y_i - y_L)^2} \\
 r_{i,j} &= \sqrt{(x_i - x_j)^2 + (y_i - y_j)^2} \\
 d_{i,L} &= \frac{\left(\frac{y_2 - y_1}{x_2 - x_1}\right)x_i - y_i + y_1 - \left(\frac{y_2 - y_1}{x_2 - x_1}\right)x_1}{\pm \sqrt{\left(\frac{y_2 - y_1}{x_2 - x_1}\right)^2 + 1}} \quad (\text{VII.2.7})
 \end{aligned}$$

The sign is chosen so that $d_{i,L}$ is always positive. The terms

$$\frac{\partial}{\partial x_i} \left(-\frac{d_{i,L}}{r_{i,L}^2} \right), \text{ and } \frac{\partial}{\partial y_i} \left(-\frac{d_{i,L}}{r_{i,L}^2} \right)$$

are given in Appendix A as:

$$\frac{\partial}{\partial x_i} \left(-\frac{d_{i,L}}{r_{i,L}^2} \right) = \frac{2(x_i - x_L)d_{i,L}}{r_{i,L}^4} - \left(\frac{1}{r_{i,L}^2} \right) \left(\frac{A}{\pm \sqrt{A^2 + 1}} \right) \quad (\text{VII.2.8})$$

$$\frac{\partial}{\partial y_i} \left(-\frac{d_{i,L}}{r_{i,L}^2} \right) = \frac{2(y_i - y_L)d_{i,L}}{r_{i,L}^4} - \left(\frac{1}{r_{i,L}^2} \right) \left(\frac{A}{\pm \sqrt{A^2 + 1}} \right) \quad (\text{VII.2.9})$$

where

$$A = \frac{y_2 - y_1}{x_2 - x_1}$$

Since

$$\frac{\partial}{\partial x_i} \left(\ln \frac{1}{r_{i,L}} \right) = -\frac{\partial}{\partial x_i} (\ln r_{i,L}) = -\frac{1}{r_{i,L}} \left(\frac{dr_{i,L}}{dx_i} \right) = -\frac{x_i - x_L}{r_{i,L}^2} \quad (\text{VII.2.10})$$

$$\frac{\partial}{\partial y} \left(\ln \frac{1}{r_{i,L}} \right) = - \frac{y_i - y_L}{r_{i,L}^2} \quad (\text{VII.2.11})$$

Equations (VII.2.8) to (VII.2.11) are substituted into (VII.2.5) and (VII.2.6) to give:

$$\begin{aligned} \frac{\partial \phi(x_i, y_i)}{\partial x_i} &= \frac{1}{2\pi} \sum_{L=1}^m \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} - \frac{x_i - x_L}{r_{i,L}^2} dS_L \\ &- \frac{1}{2\pi} \sum_{L=1}^m \phi_L \int_{S_L} \left[\frac{2(x_i - x_L)d_{i,L}}{r_{i,L}^4} - \left(\frac{1}{r_{i,L}^2} \right) \left(\frac{A}{\pm \sqrt{A^2 + 1}} \right) \right] dS_L \\ &+ \frac{\mu}{2\pi k} \sum_{j=1}^n q_j \left(\frac{x_i - x_j}{r_{i,j}^2} \right) \end{aligned} \quad (\text{VII.2.12})$$

$$\begin{aligned} \frac{\partial \phi(x_i, y_i)}{\partial y_i} &= \frac{1}{2\pi} \sum_{L=1}^m \left(\frac{\partial \phi}{\partial n} \right)_L \int_{S_L} - \frac{y_i - y_L}{r_{i,L}^2} dS_L \\ &- \frac{1}{2\pi} \sum_{L=1}^m \phi_L \int_{S_L} \left[\frac{2(y_i - y_L)d_{i,L}}{r_{i,L}^4} - \left(\frac{1}{r_{i,L}^2} \right) \left(\frac{1}{\pm \sqrt{A^2 + 1}} \right) \right] dS_L \\ &+ \frac{\mu}{2\pi k} \sum_{j=1}^n q_j \left(\frac{y_i - y_j}{r_{i,j}^2} \right) \end{aligned} \quad (\text{VII.2.13})$$

When Equations (VII.2.12) and (VII.2.13) are evaluated and substituted into Equations (VII.2.1) and (VII.2.2), the velocities in the x and y directions are obtained at any given locations in the interior of the domain. The integrals in Equations (VII.2.12) and (VII.2.13) are evaluated numerically using a four point Gaussian quadrature formula as given earlier in Section 5.1.

In matrix representation, Equations (VII.2.12) and (VII.2.13) can be expressed as:

$$\frac{\partial \phi(x_i, y_i)}{\partial x_i} = \sum_{L=1}^m [G'_x]_{i,L} \left(\frac{\partial \phi}{\partial n} \right)_L - \sum_{L=1}^m [H'_x]_{i,L} \phi_L + B' \quad (\text{VII.2.14})$$

$$\frac{\partial \phi(x_i, y_i)}{\partial y_i} = \sum_{L=1}^m [G'_y]_{i,L} \left(\frac{\partial \phi}{\partial n} \right)_L - \sum_{L=1}^m [H'_y]_{i,L} \phi_L + B' \quad (\text{VII.2.15})$$

After applying the quadrature formula,

$$H'_{x \ i, L} = -2 \sum_{k=1}^4 w_k \left[\frac{d_{i,L}(x_i - x_{L_k})}{r_{i,L_k}^4} + \frac{A}{r_{i,L_k}^2 \pm \sqrt{A^2 + 1}} \right] \left(\frac{1}{2} |S_L| \right) \quad (\text{VII.2.16})$$

$$H'_{y \ i, L} = -2 \sum_{k=1}^4 w_k \left[\frac{d_{i,L}(y_i - y_{L_k})}{r_{i,L_k}^4} + \frac{A}{r_{i,L_k}^2 \pm \sqrt{A^2 + 1}} \right] \left(\frac{1}{2} |S_L| \right) \quad (\text{VII.2.17})$$

$$G'_{x \ i, L} = -\sum_{k=1}^4 \frac{x_i - x_{L_k}}{r_{i,L_k}} \left(\frac{1}{2} |S_L| \right) \quad (\text{VII.2.18})$$

$$G'_{y \ i, L} = -\sum_{k=1}^4 \frac{y_i - y_{L_k}}{r_{i,L_k}} \left(\frac{1}{2} |S_L| \right) \quad (\text{VII.2.19})$$

where (x_{L_k}, y_{L_k}) are the coordinates of the k th quadrature point on the L th boundary segment.

r_{i,L_k} = the distance from an interior point (x_i, y_i) to the k th quadrature point (x_{L_k}, y_{L_k}) on the L th boundary segment

$$\begin{aligned}
 |S_L| &= \text{the length of the } L\text{th segment} \\
 &= \sqrt{(x_{2_L} - x_{1_L})^2 + (y_{2_L} - y_{1_L})^2}
 \end{aligned}$$

$d_{i,L}$ = the perpendicular distance from an interior
 point (x_i, y_i) to the L th boundary segment

CHAPTER VIII

EXTENSION TO HETEROGENEOUS POROUS MEDIA

Two kinds of heterogeneities can be considered.

They are:

1. Piece-wise homogeneous systems
2. Homogeneous systems with impermeable inclusions

VIII.1 PIECEWISE HOMOGENEOUS POROUS MEDIA

The domain of interest is made up of a number of homogeneous sub-domains. For the purposes of illustration, only two sub-domains will be treated in this work, but the theory applies to any finite number of sub-domains. Each sub-domain has a different permeability and each sub-domain has three kinds of boundaries enclosing it; namely (Figure 6) for the p th subdomain:

1. S_{ϕ}^p boundary where values of potentials are specified.
2. S_n^p boundary where values of the normal derivatives of potential $\left(\frac{\partial \phi}{\partial n}\right)$ are specified.
3. S_I^p boundary is the interface with the adjacent sub-domain. On S_I^p both ϕ and $\frac{\partial \phi}{\partial n}$ are unknown.

where $p = 1, 2, \dots$, number of sub-domains.

All the equations derived for a single domain (or sub-domain) in the previous sections apply to each one of the p sub-domains in turn.

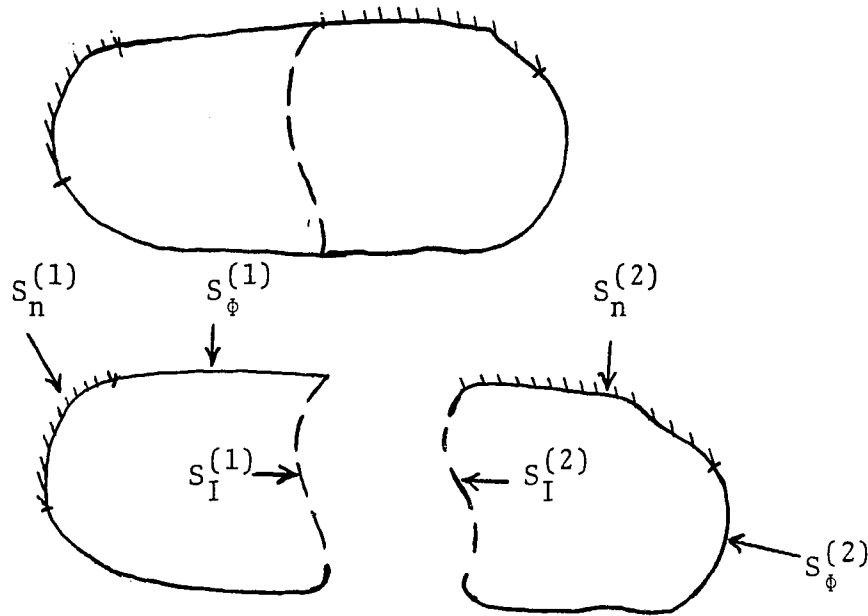


FIGURE 6. Piecewise homogeneous domain showing types of boundary conditions

Considering any arbitrary p 'th sub-domain, Equation (VII.0.1) can be written as:

$$[H]^{(p)} \{\phi\}^{(p)} = [G]^{(p)} \{\phi_n\}^{(p)} + B \quad (\text{VIII.1.1})$$

Equation (VIII.1.1) is applied to each sub-domain as if it exists in isolation, and then all the sub-domain matrices are assembled into one by utilizing continuity of flux conditions and compatibility conditions on the common boundaries. The continuity condition on the common boundary shown in Figure 6 as:

$$-T_1 \left(\frac{\partial \phi_I}{\partial n_I} \right)^{(1)} = -T_2 \left(\frac{\partial \phi_I}{\partial n_I} \right)^{(2)} \quad (\text{VIII.1.2})$$

The compatibility condition on the common boundary is:

$$\phi_I^{(1)} = \phi_I^{(2)} \quad (\text{VIII.1.3})$$

For each sub-domain, the boundary elements are numbered such that the sub-domain is always on the left as we go round the boundary. This ensures that the direction of the normal on the boundary surface of each element is always pointing away from the interior of the sub-domain.

For each sub-domain, the vector $\{\phi\}^{(p)}$ is made up of the unknown ϕ 's from the S_n boundary (where $\frac{\partial \phi}{\partial n}$ is specified) and from the common boundary S_I (where both ϕ and $\frac{\partial \phi}{\partial n}$ are unknown). Therefore:

$$\{\phi\}^{(p)} = \begin{Bmatrix} \phi_{S_n} \\ \phi_{S_I} \end{Bmatrix}^{(p)}$$

Hence the L.H.S. of (VIII.1.1) is expanded as:

$$[H_{S_n} \quad H_{S_\phi} \quad H_{S_I}]^{(p)} \begin{Bmatrix} \phi_{S_n} \\ \bar{\phi}_{S_\phi} \\ \phi_{S_I} \end{Bmatrix}^{(p)}$$

where $\bar{\phi}_{S_\phi}$ indicates known (specified) values of ϕ . Similar treatment of the $\left\{ \frac{\partial \phi}{\partial n} \right\}^{(p)}$ gives the following

$$[G_{S_n} \quad G_{S_\phi} \quad G_{S_I}] \begin{Bmatrix} (\bar{\phi}_n)_{S_n} \\ (\phi_n)_{S_\phi} \\ (\phi_n)_{S_I} \end{Bmatrix}^{(p)}$$

where $\bar{\phi}_n$ denotes known (specified) values of $\frac{\partial \phi}{\partial n}$.

Equation (VIII.1.1) then becomes, for a single sub-region (the p'th subregion),

$$[H_{S_n} \quad H_{S_\phi} \quad H_{S_I}]^{(p)} \begin{Bmatrix} \phi_{S_n} \\ \bar{\phi}_{S_\phi} \\ \phi_{S_I} \end{Bmatrix}^{(p)} = [G_{S_n} \quad G_{S_\phi} \quad G_{S_I}]^{(p)} \begin{Bmatrix} (\bar{\phi}_n)_{S_n} \\ (\phi_n)_{S_\phi} \\ (\phi_n)_{S_I} \end{Bmatrix}^{(p)} + \{B\}$$

(VIII.1.4)

The system of matrices represented by Equation (VIII.1.4) is set up for every one of the sub-domains and assembled using Equations (VIII.1.2) and (VIII.1.3). (consider the case where the number of sub-domains = 2, i.e., $p = 2$). For the first sub-domain, $p = 1$, putting all unknowns in the G-matrix gives:

$$[G_{S_\phi}^{(1)} \quad G_{S_I}^{(1)} - H_{S_I}^{(1)} - H_{S_n}^{(1)}] \begin{Bmatrix} (\phi_n)_{S_\phi}^{(1)} \\ (\phi_n)_{S_I}^{(1)} \\ \phi_{S_I}^{(1)} \\ \phi_{S_n}^{(1)} \end{Bmatrix} = [H_S^{(1)} - G_{S_n}^{(1)}] \begin{Bmatrix} \bar{\phi}_{S_\phi}^{(1)} \\ \bar{\phi}_n^{(1)} \end{Bmatrix} - \{B^{(1)}\} \quad \text{(VIII.1.5)}$$

For the second sub-domain, $p = 2$.

$$[G_{S_\phi}^{(2)} \quad G_{S_I}^{(2)} \quad -H_{S_I}^{(2)} \quad -H_{S_n}^{(2)}] \begin{Bmatrix} \phi_{n_{S_\phi}}^{(2)} \\ \phi_{n_{S_I}}^{(2)} \\ \phi_{S_I}^{(2)} \\ \phi_{S_n}^{(2)} \end{Bmatrix} = [H_{S_\phi}^{(2)} \quad -G_{S_n}^{(2)}] \begin{Bmatrix} \bar{\phi}_{S_\phi}^{(2)} \\ \bar{\phi}_n^{(2)} \end{Bmatrix} - \{B^{(1)}\} \quad (\text{VIII.1.6})$$

The continuity condition on the common boundary is:

$$\phi_{S_I}^{(1)} = \phi_{S_I}^{(2)} \quad (\text{VIII.1.7})$$

The compatibility condition is:

$$\begin{aligned} -T_1 \phi_{n_{S_I}}^{(1)} &= T_2 \phi_{n_{S_I}}^{(2)} \\ \therefore \phi_{n_{S_I}}^{(2)} &= -\frac{T_1}{T_2} \phi_{n_{S_I}}^{(1)} \end{aligned} \quad (\text{VIII.1.8})$$

The assembled system matrix is:

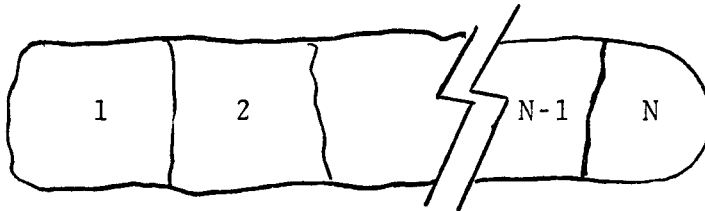
$$\begin{aligned} & \begin{bmatrix} -H_{S_n}^{(1)} & G_{S_\phi}^{(1)} & -H_{S_I}^{(1)} & G_{S_I}^{(1)} & 0 & 0 \\ 0 & 0 & -H_{S_I}^{(2)} - \frac{T_1}{T_2} G_{S_I}^{(2)} & -H_{S_n}^{(2)} & G_{S_\phi}^{(2)} \end{bmatrix} \begin{Bmatrix} \phi_{S_n}^{(1)} \\ \phi_{n_{S_\phi}}^{(1)} \\ \phi_{S_I}^{(1)} \\ \phi_{n_{S_I}}^{(1)} \\ \phi_{S_n}^{(2)} \\ \phi_{n_{S_\phi}}^{(2)} \end{Bmatrix} \\ &= \begin{bmatrix} H_{S_\phi}^{(1)} & -G_{S_n}^{(1)} & 0 & 0 \\ 0 & 0 & H_{S_I}^{(2)} & -G_{S_n}^{(2)} \end{bmatrix} \begin{Bmatrix} \bar{\phi}_{S_\phi}^{(1)} \\ \bar{\phi}_{n_{S_n}}^{(1)} \\ \bar{\phi}_{S_\phi}^{(2)} \\ \bar{\phi}_{n_{S_n}}^{(2)} \end{Bmatrix} - \{B\} \end{aligned} \quad (\text{VIII.1.9})$$

The matrix system given by (VIII.1.9) is solved for the boundary values of ϕ and $\frac{\partial \phi}{\partial n}$. These values are used to calculate the interior values of $\phi(x,y)$, $\frac{\partial \phi}{\partial x}$, and $\frac{\partial \phi}{\partial y}$ within any of the sub-domains of interest by applying Equations (VII.1.3), (VII.2.14) and (VII.2.15).

The matrix system given by (VIII.1.9) can be expressed schematically as:

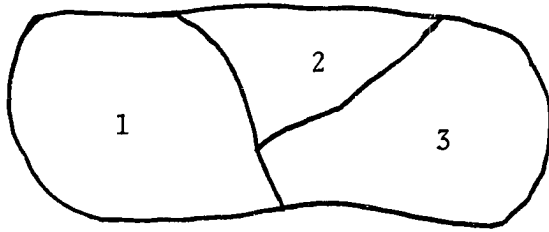
$$\begin{bmatrix} X & X & X & X & 0 & 0 \\ 0 & 0 & X & X & X & X \end{bmatrix} \begin{Bmatrix} \\ \\ \\ \\ \\ \end{Bmatrix} = \begin{Bmatrix} Y \\ Y \end{Bmatrix}$$

This can be generalized for other kinds of systems. For example, for a system having n regions, we have



$$\begin{bmatrix} X & X & X & X & & & & \\ 0 & 0 & X & X & X & X & & \\ 0 & 0 & 0 & 0 & X & X & X & X \\ & & & & \cdot & \cdot & \cdot & \\ & & & & \cdot & \cdot & \cdot & \\ & & & & \cdot & \cdot & \cdot & \\ & & & & & & & X & X & X & X \end{bmatrix} \begin{Bmatrix} \\ \\ \\ \\ \\ \\ \\ \end{Bmatrix} = \begin{Bmatrix} \\ \\ \\ \\ \\ \\ \\ \end{Bmatrix}$$

For a system of the type



the matrix system will be of the form

$$\begin{bmatrix}
 \text{(a)} & \text{(b)} & \text{(c)} & & & & & & & & & \\
 X & X & X & X & X & X & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & X & X & 0 & 0 & X & X & X & X & 0 & 0 \\
 0 & 0 & 0 & 0 & X & X & X & X & 0 & 0 & X & X
 \end{bmatrix}
 \begin{Bmatrix} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{Bmatrix}
 =
 \begin{Bmatrix} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \end{Bmatrix}$$

(a) = interface between regions 1 and 2

(b) = interface between regions 1 and 3

(c) = interface between regions 2 and 3

VIII.2 HOMOGENEOUS SYSTEM WITH IMPERMEABLE INCLUSIONS

This kind of system can easily be handled by the Boundary Element Method. Instead of one, there are now three surfaces that bound the domain of interest (Figure 8). All that is required is to maintain the numbering convention on each of the surfaces so that the domain is on the left hand side.

It is observed in Figure 8 that, to satisfy the numbering convention,

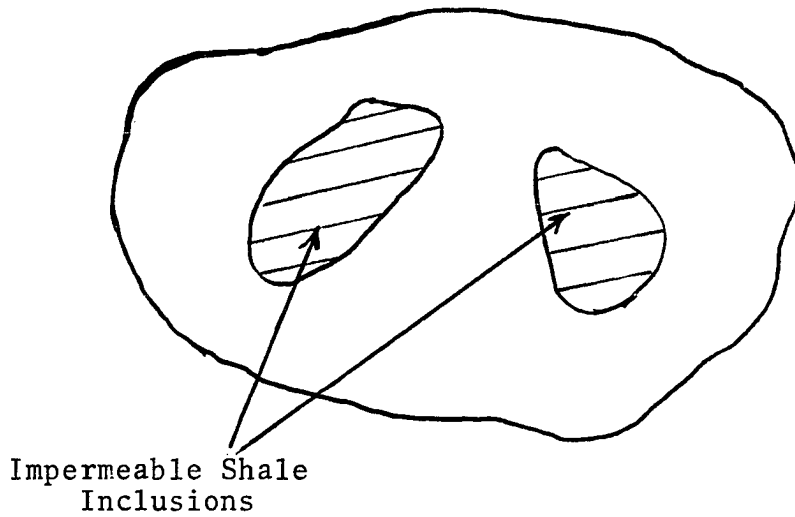


FIGURE 7. Homogeneous porous medium with impermeable shale inclusions.

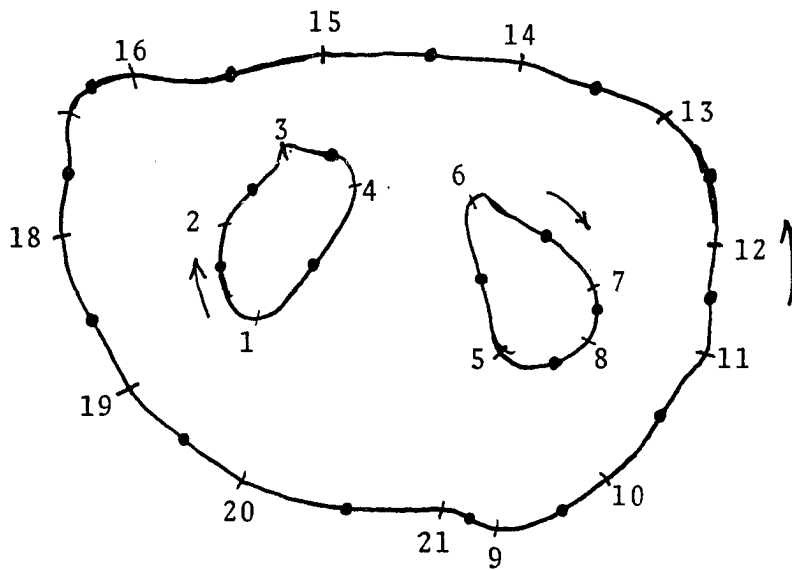


FIGURE 8. Homogeneous porous medium with impermeable inclusions showing discretization and numbering scheme.

- Internal boundaries are numbered in the clockwise direction.
- External boundaries are numbered in the counter-clockwise direction.

VIII.3 STREAMLINE AND STREAMTUBE CALCULATION PROCEDURE

The streamlines are determined by following the path of a hypothetical fluid particle from the injector to the producer. For any injection well, the number of streamlines emanating from that well is independent of its injection rate. For a system with multiple injection wells, a fixed number of streamlines is assigned to each well. Thus, assigning the fixed number of streamlines (N_L) to an injector with rate q_L , the injection rate into each streamtube emanating from it is:

$$q_s = \frac{q_L}{N_L} \quad (\text{VIII.3.1})$$

The injection rate into the streamtubes emanating from any other injector with injection rate q_1 becomes:

$$q_{s_1} = \frac{q_1}{q_L} \times q_s \quad (\text{VIII.3.2})$$

The points of origin of the fluid particles (and therefore the streamlines) are

distributed equidistant along a small arbitrary circle centered at the source. At any current position, the potential is determined using Equation (VII.1.3) and velocity vectors v_x and v_y in the x and y directions are determined using Equations (VII.2.1) to (VII.2.15). In order to determine the next location of the fluid particle, the assumption is made that the velocity remains constant over a finite distance increment Δs . At any point, the resultant particle velocity (v_t) is the vectorial sum of v_x and v_y . Since v_x and v_y and thus v_t are assumed constant over Δs , the time of travel can be calculated as:

$$\Delta t = \frac{\Delta s}{v_t} \quad (\text{VIII.3.3})$$

Hence, from any starting point denoted by (x_i, y_i) , the particle would move to a new point given by:

$$x_{i+1} = x_i + v_{x_i} \Delta t \quad (\text{VIII.3.4})$$

$$y_{i+1} = y_i + v_{y_i} \Delta t \quad (\text{VIII.3.5})$$

The distance step increment Δs is chosen to be constant for convenience. As successive steps are taken, a curve is generated and is commonly called a streamline.

The width of the streamtube at any position of the fluid particle can be determined from the injection rate and the velocity. For each injector, the volume of injected fluid associated with a streamline is:

$$q_s = \frac{q}{N} \quad (\text{VIII.3.6})$$

where N is the number of streamlines emanating from the injector and is given by Equation (VIII.3.2). Since this volume of fluid must be conserved at all positions along the streamline, then for any arbitrary position i of coordinates (x_i, y_i)

$$q_s = v_{t_i} A_i \quad (\text{VIII.3.7})$$

where v_{t_i} is the resultant velocity at position i
 A_i is the vertical cross sectional area at position i , and

ϕ is the porosity, assumed constant

Making the assumption that the vertical thickness of the reservoir (and therefore the streamtube) is a known constant h , the width of the streamtube at any position i (x_i, y_i) can be determined as:

$$W_i = \frac{q_s}{v_{t_i} h} \quad (\text{VIII.3.7})$$

Due to the fact that analytical solutions to the equations that describe the rate of advance of a steam front in porous medium have only been obtained for linear or radial models, the streamtube is modified to be linear as follows:

Using Equation (VIII.3.7), the width of the streamtube is determined at any given streamline location and averaged along the length of the entire streamline to give a single width for every streamtube. The streamtubes are then assumed to have rectangular cross sections. The streamlines and associated streamtubes are terminated by defining an arbitrary capture radius around the producer. When the calculated position of the streamline gets within the capture radius, the streamline is terminated.

VIII.4 CONTINUATION OF STREAMLINES IN ADJACENT REGIONS

For piece-wise homogeneous reservoirs, some of the streamlines may cross into adjacent regions where the permeabilities are different. In such cases, continuation of the streamlines into adjacent regions can be achieved by noting the interface boundary segment where the streamline crossed as well as the coordinates of this crossing point. Since the potential derivative normal to every segment is known as part of the solution, this normal gradient can be used with Darcy's law to determine the velocities and distances needed to continue tracing the streamlines.

In practice, however, the streamlines could not be traced all the way until they met any boundary. The potentials exhibited fluctuations at distances closer than half an element length away from any boundary. The reason for this is the singular nature of the fundamental

solution $\ln \frac{1}{r_{i,j}}$ as i approaches j . Therefore, the method adopted here was to test every new position to see if it is less than half an element length away from the mid point of any boundary. If it is, a perpendicular line was dropped from the previous position to the interface segment. This is consistent with the fact that the streamline must be normal to a flow boundary. The point of intersection of this normal with the interface segment is taken as the starting point for continuation of the streamline into the adjacent region and is given as:

$$x_c = \frac{(c_1 - y_{1L})x_i - c_2x_{1L} + y_{1L} - y_i}{c_1 - c_2}$$

$$y_c = c_1(x_{1L} - x_i) + y_i$$

where:

(x_c, y_c) = coordinates of the point where the streamline crosses the interface boundary

$$c_1 = \frac{x_{1L} - x_{2L}}{y_{2L} - y_{1L}}$$

$$c_2 = \frac{y_{2L} - y_{1L}}{x_{2L} - x_{1L}}$$

(x_i, y_i) = coordinates of last position of streamline

$(x_{1L}, y_{1L}), (x_{2L}, y_{2L})$ = coordinates of the extreme points of the Lth interface element

The normal derivative of the potential at the L th interface is already known as part of the solution. This normal derivative is introduced into Darcy's law to determine the next position away from the boundary into the adjacent region. A check is made to ensure that this position is further from the boundary than half the interface segment length.

A package of computer programs has been written in FORTRAN to generate and plot the streamlines, and to calculate the dimensions of the associated streamtubes. Details of the algorithms and flow charts are listed in Appendix E. The computer program itself is listed in Appendix O.

The final part of this modelling effort is its application to a continuous steam drive process. Since some of the streamtubes in a piecewise homogeneous reservoir will contain two or more permeability regions, the next section develops the appropriate equations for the rate of advance of the steam front in a linear streamtube of constant cross section having two or more permeability regions.

CHAPTER IX

RATE OF ADVANCE OF STEAM FRONT IN A PIECE- WISE HOMOGENEOUS LINEAR POROUS MEDIUM

In order to design a steamflood project or to predict recovery from such projects, a knowledge of the distribution and movement of temperature fronts in the reservoir is required. This information can be obtained by taking a control volume within the medium (reservoir) and applying the general and particular laws of physics to it. A general law is one whose application is independent of the nature of the medium under consideration. A particular law is one whose application is dependent on the nature of the medium. In the formulation of a steam drive model, the general laws employed are:

- a. the law of conservation of energy (heat balance)
- b. the law of conservation of mass

and the particular laws employed are:

- c. Fourier's law of conduction

In modelling steam drive processes, some basic assumptions are made such as:

- 1. Changes in viscosity and density due to changes in temperature have negligible effect on energy transfer. This

assumption enables energy balance equations to be uncoupled from mass balance equations.³⁹

2. The thermal conductivity in the direction perpendicular to the direction of fluid flow within the reservoir is infinite. This assumption means that the temperature in the reservoir at any cross section perpendicular to the direction of fluid flow is uniform.

3. The sand grains and the reservoir fluids maintain instantaneous thermal equilibrium. This means that the sand grains and surrounding fluids were always at the same temperature.

A general heat balance for thermal recovery projects may be expressed as:

$$\left\{ \begin{array}{l} \text{The amount of heat} \\ \text{(enthalpy) contained} \\ \text{within a given vol-} \\ \text{ume at any time} \end{array} \right\} = \left\{ \begin{array}{l} \text{heat} \\ \text{injected} \end{array} \right\} - \left\{ \begin{array}{l} \text{heat} \\ \text{produced} \end{array} \right\}$$

$$+ \left\{ \begin{array}{l} \text{heat} \\ \text{generated} \end{array} \right\} - \left\{ \begin{array}{l} \text{heat} \\ \text{lost} \end{array} \right\}$$

Mathematically, it can be expressed in integral form as⁴²

$$\int_V \rho c T \, dV = \int_0^t \dot{Q}(\tau) \, d\tau - \int_0^t \left[\int_A \left(-k \frac{\partial T}{\partial n} + V_f \rho_f c_f T \right) dA \right] dt$$

In this equation, V is the volume over which the balance is made (the entire steam zone). Since V varies with time for the present problem, it will be more convenient to make an instantaneous heat balance as:

$$\int_V \rho c T \, dV = \dot{Q}(\tau) - \int_A \left(-k \frac{\partial T}{\partial n} + V_f \rho_f c_f T \right) dt$$

This equation and the one preceding it are completely general and include the effects of variations in formation properties with location and time.

Assuming the confining layers are semi-infinite planes with constant, homogeneous thermal properties, assuming further that the temperature difference imposed on the boundaries of the confining layers is directly proportional to the amount of heat stored in the reservoir, the solution of the instantaneous heat balance equation is:⁴²

$$H(t) = \int_0^t \dot{Q}(\tau) e^{\theta(t-\tau)} \operatorname{erfc} \sqrt{\theta(t-\tau)} \, d\tau$$

$\dot{Q}(\tau)$ = rate of heat injection - rate of heat produced

$H(t)$ = heat contained in the reservoir at time t

= (heat content of rock + fluids) vol/sec of steam zone

Since vol/sec = velocity x cross-sectional area, this equation can be used to calculate the velocity of the steam front.

The next section makes such an instantaneous heat balance across the condensation front for a piecewise homogeneous linear porous medium.

IX.1 INSTANTANEOUS HEAT AND MASS BALANCE ACROSS THE MOVING CONDENSATION FRONT

Consider a condensation front located at position $z(t)$ in the N 'th region of a multiregion composite porous medium. Take two arbitrary fixed cross sections at positions z_{N_a} and z_{N_b} on the upstream and downstream sides of the condensation front respectively. (See Figure 9.)

The region behind the condensation front is completely covered by steam. In this work, the terms condensation front and steam front will be used interchangeably.

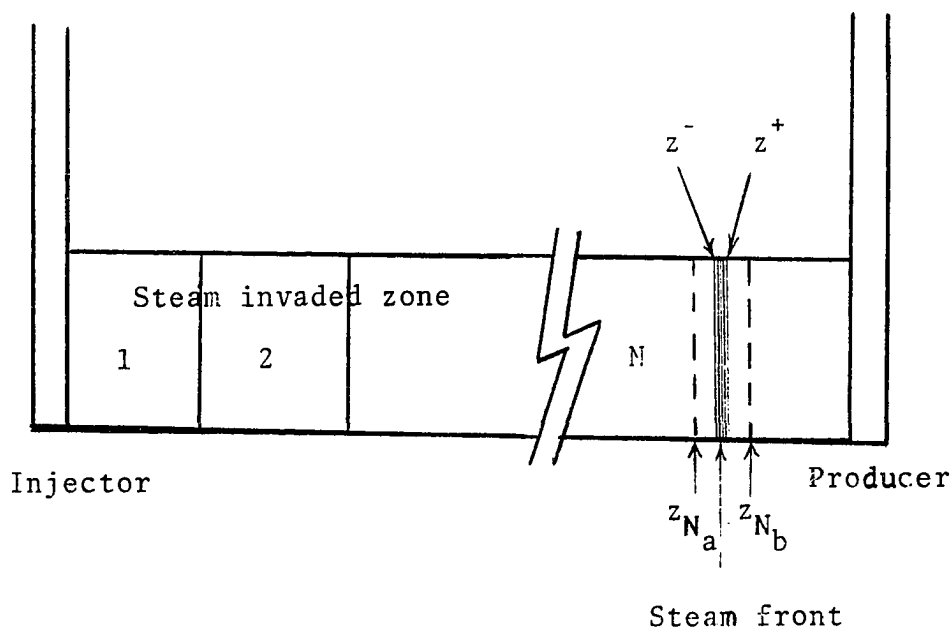


FIGURE 9. Reservoir cross section showing location of steam front.

The objective is to determine the velocity of the steam or condensation front as it moves within any arbitrary region of a multi-region heterogeneous linear porous medium. To achieve this, an elemental volume is taken on the steam front such that one half of it is behind the steam front and the other half is ahead of the steam front. The law of conservation of energy is applied in the form of a heat balance to the elemental volume enclosed by the two fixed cross sections at z_{N_a} and z_{N_b} . The upstream plane of the condensation front is denoted z^- , while the downstream plane is denoted z^+ .

A heat balance in the fixed elemental area between the two fixed cross sections z_{N_a} and z_{N_b} is:

$$\left[\begin{array}{c} \text{Heat flux} \\ \text{in at } z_{N_a} \end{array} \right] - \left[\begin{array}{c} \text{Heat flux} \\ \text{out at } z_{N_b} \end{array} \right] = \left[\begin{array}{c} \text{Rate of} \\ \text{accumulation} \\ \text{of heat} \end{array} \right] + \left[\begin{array}{c} \text{rate of loss} \\ \text{of heat from} \\ \text{the volume} \end{array} \right]$$

Following Mandle and Volek,³² the heat balance can be expressed mathematically as:

$$\begin{aligned} Q_N(z_{N_a}, t) - Q_N(z_{N_b}, t) &= \frac{d}{dt} \int_{z_{N_a}}^{z_N^-(t)} H_N(z_{N_a}, t) dz_N \\ &+ \frac{d}{dt} \int_{z_N^+(t)}^{z_{N_b}} H_N(z_{N_b}, t) dz_N + Q_{L_N}(t) \quad (\text{IX.1.1}) \end{aligned}$$

where, for the steam front in the N'th region,

$H_N(z_{N_a}, t)$ denotes the heat content per unit volume

of fluid/solid system between z_{N_a} and z_N^- .

$H_N(z_{N_b}, t)$ denotes the heat content per unit volume
of fluid/solid system between z^+ and z_{N_b}

$Q_N(z_{N_a}, t)$ = the heat flux in at z_{N_a} per unit cross
sectional area

$Q_N(z_{N_b}, t)$ = the heat flux out at z_{N_b} per unit cross
sectional area

The application of the Leibnitz rule to Equation (IX.1.1)
and considering that z_{N_a} and z_{N_b} are fixed values, gives

$$\begin{aligned} Q_N(z_{N_a}, t) - Q_N(z_{N_b}, t) &= \int_{z_{N_a}}^{z_N^-(t)} \frac{d}{dt} H_N(z_{N_a}, t) dz \\ &+ H_N(z_N^-(t), t) \frac{d}{dt} z_N^-(t) + \int_{z_N^+(t)}^{z_{N_b}} \frac{d}{dt} H_N(z_{N_b}, t) dz \\ &- H_N(z_N^+(t), t) \frac{d}{dt} z_N^+(t) + Q_{L_N}(t) \end{aligned} \quad (\text{IX.1.2})$$

Equation (IX.1.2) is the heat balance across two fixed cross
sections containing the steam front in the N'th region. In
the limit as z_{N_a} tends to $z_N^-(t)$ and z_{N_b} tends to $z_N^+(t)$,
the heat balance equation becomes:

$$\begin{aligned} Q_N(z_N^-(t), t) - Q_N(z_N^+(t), t) &= H_N(z_N^-(t), t) \frac{d}{dt} z_N^-(t) \\ &- H_N(z_N^+(t), t) \frac{d}{dt} z_N^+(t) + Q_{L_N}(t) \end{aligned}$$

Making the assumption that $\frac{d}{dt} z_N^-(t) = \frac{d}{dt} z_N^+(t) = \frac{d}{dt} z_N(t)$
and no heat is lost within the thin condensation front,

and dropping the time function of z^- and z^+ for convenience, the heat balance equation becomes the heat balance across a moving condensation front which is instantaneously frozen in time.

$$Q_N(z_N^-, t) - Q_N(z_N^+, t) = [H_N(z_N^-, t) - H_N(z_N^+, t)]v_N(t) \quad (\text{IX.1.3})$$

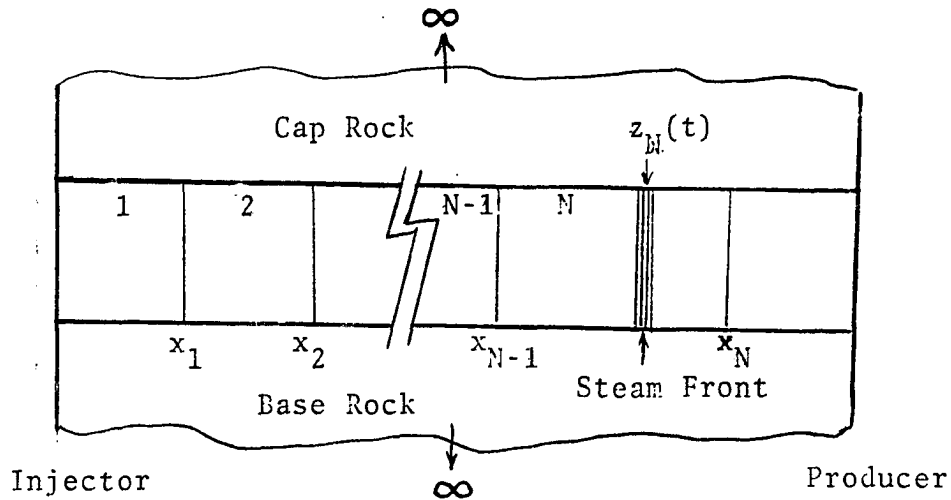


FIGURE 10. Reservoir cross section for heat balance.

That is, the instantaneous heat balance across a moving condensation front is described as:

$$\left[\text{Heat flux} \right]_{\text{in at } z_N^-} - \left[\text{Heat flux} \right]_{\text{out at } z_N^+} = \left[\text{Rate of change of heat content between } z_N^- \text{ and } z_N^+ \right] \quad (\text{IX.1.4})$$

where, since there are no heat sources or sinks in the steam zone,

$$\begin{aligned}
Q_N(z_N^-, t) = & \left[\text{Heat flux} \right] = \left[\text{Rate of heat} \right] - \left[\text{Rate of heat loss} \right] \\
& \left[\text{in at } z_N^- \right] \left[\text{injection} \right] \left[\text{in } N - 1 \text{ regions} \right] \\
& - \left[\text{Rate of heat loss} \right] - \left[\text{Rate of change of heat} \right] \\
& \left[\text{in } N\text{'th region} \right] \left[\text{content in } N - 1 \text{ regions} \right] \\
& - \left[\text{Rate of change of heat} \right] \quad (IX.1.5) \\
& \left[\text{content in } N\text{'th region} \right]
\end{aligned}$$

$$\begin{aligned}
Q_N(z_N^+, t) = & \left[\text{Heat flux} \right] = \left[\text{Heat flux out} \right] + \left[\text{Heat flux out} \right] \\
& \left[\text{out at } z_N^+ \right] \left[\text{by convection} \right] \left[\text{of the fluids} \right] \left[\text{by conduction} \right] \\
& \quad (IX.1.6)
\end{aligned}$$

and

$$\begin{aligned}
\left[\text{Rate of change} \right] & = \left[\text{Heat} \right] - \left[\text{Heat} \right] \left[\text{Velocity of} \right] \\
\left[\text{of heat content} \right] & \left[\text{content} \right] \left[\text{content} \right] \left[\text{the steam} \right] \\
\left[\text{between } z^- \text{ and } z^+ \right] & \left[\text{at } z^- \right] \left[\text{at } z^+ \right] \left[\text{front} \right] \\
& \quad (IX.1.7)
\end{aligned}$$

Equations (IX.1.4) through (IX.1.7) apply to mass balances as well as to heat balances. When making mass balances "heat" is replaced by "mass". Furthermore, for mass balances, the second and third terms on the right hand side of Equation (IX.1.5) disappear because there are no mass losses behind the steam front as is the case with heat. Also, Equation (IX.1.5) shows that the regions behind the steam front are handled in two parts. The first part is made up of all the $N - 1$ regions behind the steam front. The second part is the part of the N 'th region behind the steam front.

IX.2 HEAT BALANCE

For a condensation front at any arbitrary location $z_N(t)$ in the N'th region of a multi-region composite medium, the heat balance for an elemental area taken across the condensation front (Figure 10) has been shown by Equation (IX.1.3) to be

$$Q_N(z_N^-, t) - Q_N(z_N^+, t) = v_N(t) [H_N(z_N^-, t) - H_N(z_N^+, t)] \quad (\text{IX.2.1})$$

for a unit reservoir cross section area. Following Equation (IX.1.5) the rate of heat input $Q_N(z_N^-, t)$ per unit cross section area at z_N^- is:

$$\begin{aligned} Q_N(z_N^-, t) = & Q(0, t) - \sum_{j=1}^{N-1} (Q_L(t))_j - (Q_L(t))_N \\ & - \rho_{st}(L_{st} + c_w \Delta T(z^-)) \left[\sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{stj}(z_j, t)] dz_j \right. \\ & \left. - \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{stj}(z_N, t)] dz_N \right] \\ & - \rho_w(T_{st}) c_w \Delta T(z^-) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{wj}(z_j, t)] dz_j \\ & - \rho_w(T_{st}) c_w \phi_N \Delta T(z^-) \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{wj}(z_N, t)] dz_N \\ & - \rho_o(T_{st}) c_o \Delta T(z^-) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{oj}(z_j, t)] dz_j \\ & - \rho_o(T_{st}) c_o \Delta T(z^-) \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{oj}(z_N, t)] dz_N \quad (\text{IX.2.2}) \end{aligned}$$

where:

$Q_N(z_N^-, t)$ = the rate of flow of heat per unit cross sectional area through the vertical cross sectional area at z_N^- into the elemental volume in the N'th region

$Q_N(z_N^+, t)$ = the rate of flow of heat per unit cross sectional area out of the elemental volume through the vertical cross section at z_N^+ in the N'th region

$v_N(t)$ = the velocity of the steam front in the N'th region

$H_N(z^-, t)$ = heat content per unit volume of reservoir rock/fluid system at the upstream face of the steam front in the N'th region

$H_N(z^+, t)$ = heat content per unit volume of reservoir rock/fluid system at the downstream face of the steam front in the N'th region

$Q(0, t)$ = rate of injection of heat per unit cross sectional area

$(Q_L(t))_j$ = rate of loss of heat per unit cross sectional area to overburden and underburden from the j'th region

$\rho_m(T)$ = density of m at temperature T, where m can be oil, water, or steam

L_{st} = latent heat of steam

c_m = specific heat of m, where m = oil, water, steam

- $\Delta T(z^-)$ = change in temperature at the upstream face of the steam front
 $\Delta T(z^+)$ = change in temperature at the downstream face of the steam front
 ϕ_j = porosity of the rock in the j'th region
 $S_m(z_j, t)$ = saturation distribution of m in the j'th region, where m = oil, water, steam

The rate of heat injection per unit cross sectional area $Q(0, t)$ is obtained as the sum of the rate of heat injection in the form of water and the rate of heat injection in the form of steam.

$$Q(0, t) = M_{st}(0, t) [L_{st} + c_w T_{st}] + M_w(0, t) c_w T_{st} \quad (IX.2.3)$$

The heat flux out of the condensation front

$$Q_N(z_N^+, t) = \text{the heat flux out by horizontal conduction} \\ + \text{convective heat flux associated with mass flux of fluids out of the front}$$

Mathematically, the heat flux out of the condensation front is written as:

$$Q_N(z_N^+, t) = M_w(z_N^+, t) c_w \Delta T(z_N^+) \quad (IX.2.4) \\ + M_o(z_N^+, t) c_o \Delta T(z_N^+) + \left[-\hat{K} \frac{\partial T}{\partial x} \right]_{z_N^+}$$

The mass rates $M_w(z_N^+, t)$ and $M_o(z_N^+, t)$ are obtained from mass balances across the condensation front.

IX.3 MASS BALANCES

A mass balance for any fluid on an elemental volume taken across the condensation front from z_N^- to z_N^+ in the N'th region is given as:

$$\begin{aligned} & \text{mass flux in at } z_N^- - \text{mass flux out at } z_N^+ \\ & = \text{rate of change of mass content in the elemental} \\ & \quad \text{area} \end{aligned}$$

The rate of change of mass content is expressed in terms of the velocity of the condensation front. The mass flux in at z_N^- is the difference between the mass rate of injection, and the rate of accumulation of mass in the swept portion.

IX.3.1 Oil

For oil, the mass rate of injection = 0, and the expression for the mass balance is:

$$\begin{aligned} & - \left\{ \rho_o(T_{st}) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_o(z_j, t)] dz_j \right. \\ & \quad \left. + \rho_o(T_{st}) \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_o(z_N, t)] dz_N \right\} - M_o(z_N^+, t) \\ & = \phi_N v_N(t) [\rho_o(T_{z^-}) S_{oN}(z_N^-, t) - \rho_o(T_{z_N^+}) S_{oN}(z_N^+, t)] \end{aligned} \tag{IX.3.1}$$

from which,

$$\begin{aligned}
M_o(z_N^+, t) = & -\rho_o(T_{st}) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_o(z_j, t)] dz_j \\
& - \rho_o(T_{st}) \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_o(x, t)] dz_N \\
& - \phi_N v_N(t) [\rho_o(T_{z-}) S_{oN}(z_N^-, t) - \rho_o(T_{z+}) S_{oN}(z_N^+, t)]
\end{aligned}
\tag{IX.3.2}$$

IX.3.2 Water

The mass flux of water into the elemental area at z_N^- is from a combination of sources:

(a) The mass fraction of steam that was injected as water.

(b) The mass of water that condensed out of the injected steam.

(c) Displaced connate water.

The mass flux of water that condensed from the injected steam is the difference between the mass rate of injection of steam and the rate of change of steam saturation in the steam zone. Mathematically, the mass balance for water is:

$$\begin{aligned}
q(0, t) [x \rho_{st} + 1 - x \rho_w(T_{st})] - \rho_w(T_{st}) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_w(z_j, t)] dz_j \\
- \rho_w(T_{st}) \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_w(z_N, t)] dz_N - \rho_{st} \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{st}(z_j, t)] dz_j \\
- \rho_{st} \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{st}(z_N, t)] dz_N - M_w(z_N^+, t) = \phi_N v_N(t) [\rho_w(T_{z_N^-}) S_w(z_N^-, t) \\
+ \rho_{st} S_{st}(z_N^-, t) - \rho_w(T_{z_N^+}) S_w(z_N^+, t)]
\end{aligned}
\tag{IX.3.3}$$

from which, the mass flux of water out of the elemental area at z_N^+ can be derived as:

$$\begin{aligned}
 M_w(z_N^+, t) = & Q(0, t) [x_{\rho_{st}} + 1 - x_{\rho_w}] - \rho_w(T_{st}) \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_w(z_j, t)] dz_j \\
 & - \rho_w(T_{st}) \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_w(z_N, t)] dx \\
 & - \rho_{st} \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{st}(z_j, t)] dx \\
 & - \rho_{st} \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{st}(z_N, t)] dz_N \\
 & - \phi_N v_N(t) [\rho_w(T_{z_N^-}) S_w(z_N^-, t) + \rho_{st} S_{st}(z_N^-, t) \\
 & - \rho_w(T_{z_N^+}) S_w(z_N^+, t)] \quad (IX.3.4)
 \end{aligned}$$

All the component terms that make up the heat balance have now been defined. Equations (IX.3.4) and (IX.3.2) can be substituted into Equation (IX.2.4). The resulting expression, as well as Equation (IX.2.2) can be substituted into Equation (IX.2.1) to give a general heat balance equation. But first, the integrals in the heat and mass balance equations are evaluated so as to simplify those equations.

Applying the result of Appendix B, we observe that, for the $N - 1$ regions behind the steam front, their limits of integration are constants. Therefore,

$$\int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_m(z_j, t)] dz_j = (x_{j-1} - x_j) \frac{d}{dt} \bar{S}_{m_j}(t) \quad (IX.3.6)$$

For the N'th region (which contains the steam front), the lower limit of integration, x_{N-1} , is a constant, but the upper limit, $z_N(t)$, varies with time. Therefore,

$$\int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_m(z_N, t)] dz_N = v_N(t) [\bar{S}_{m_N}(t) - S_m(z_N^-, t)] + \frac{d}{dt} \bar{S}_{m_N}(t) [z_N(t) - x_{N-1}] \quad (\text{IX.3.7})$$

Combining (IX.3.6) and (IX.3.7) gives:

$$\begin{aligned} & \sum_{j=1}^{N-1} \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_m(z_j, t)] dz_j + \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_m(z_N, t)] dz_N \\ & + \sum_{j=1}^{N-1} \phi_j [(x_j - x_{j-1}) \frac{d}{dt} \bar{S}_{m_j}(t)] + \phi_N v_N(t) [\bar{S}_{m_N}(t) - S_{m_N}(z_N^-, t)] + \phi_N \frac{d}{dt} \bar{S}_{m_N}(t) [z_N(t) - x_{N-1}] \end{aligned} \quad (\text{IX.3.8})$$

where m = oil, water, or steam. $S_{m_j} = (S_o + S_w + S_t)_j$

Making the assumption that $\frac{d}{dt} \bar{S}_m(z, t) = 0$ for each region

$$\begin{aligned} & \sum_{j=1}^N \phi_j \int_{x_{j-1}}^{x_j} \frac{d}{dt} [S_{m_j}(z_j, t)] dz_j + \phi_N \int_{x_{N-1}}^{z_N(t)} \frac{d}{dt} [S_{m_N}(z_N, t)] dz_N \\ & = \phi_N v_N(t) [\bar{S}_{m_N} - S_m(z_N^-, t)] \end{aligned} \quad (\text{IX.3.9})$$

Introducing Equation (IX.3.9) into Equation (IX.2.2) gives the heat flux per unit cross section at the inlet of the steam front as:

$$\begin{aligned}
Q_N(z_N^-, t) = & Q(0, t) - \sum_{j=1}^N Q_{L_j}(t) - \rho_{st}[L_{st} + c_{st}(T_{st} - T_i)] \\
& \times \phi_N v_N(t) [\bar{S}_{st_N} - S_{st}(z_N^-, t)] - \rho_w(T_{st})c_w(T_{st} - T_i)\phi_N v_N(t) \\
& \times [\bar{S}_{w_N} - S_{w_N}(z_N^-, t)] - \rho_o(T_{st})c_o(T_{st} - T_i)\phi_N v_N(t) \\
& \times [\bar{S}_{o_N} - S_{o_N}(z_N^-, t)] \quad (IX.3.10)
\end{aligned}$$

Note that this equation still contains the unknown heat lost rates per unit cross sectional area $Q_{L_j}(t)$ for $j = 1, 2, \dots$, number of regions.

Noting that $T_{z^-} = T_{st}$, Equation (IX.3.9) is substituted into the equations describing the mass fluxes out of the condensation front as:

$$M_{o_N}(z_N^+, t) = \phi_N v_N(t) [\rho_o(T_{st})\bar{S}_{o_N} - \rho_o(T_{z^+})S_{o_N}(z_N^+, t)] \quad (IX.3.10a)$$

$$\begin{aligned}
M_{w_N}(z_N^+, t) = & Q(0, t)[x\rho_{st} + 1 - x\rho_w(T_{st})] - \phi_N v_N(t) [\rho_w(T_{st})\bar{S}_{w_N} \\
& + \rho_{st}\bar{S}_{st_N} - \rho_w(T_{z^+})S_{w_N}(z_N^+, t)] \quad (IX.3.11)
\end{aligned}$$

Equations (IX.3.10) and (IX.3.11) are substituted into Equation (IX.2.4) to give the heat flux out of the condensation front as:

$$\begin{aligned}
Q_N(z_N^+, t) = & c_o \Delta T(z_N^+) \phi_N v_N(t) [\rho_o(T_{st})\bar{S}_{o_N} - \rho_o(T_{z^+})S_{o_N}(z_N^+, t)] \\
& + \{c_w \Delta T(z_N^+) (1-x)\rho_w + x\rho_{st}L_{st}\} Q(0, t) - c_w \Delta T(z_N^+) \phi_N v_N(t) [\rho_w(T_{st})\bar{S}_{w_N} \\
& + \rho_{st}\bar{S}_{st_N} - \rho_w(T_{st})S_{w_N}(z_N^+, t)] \quad (IX.3.12)
\end{aligned}$$

IX.3.3 Enthalpy

The enthalpy of the rock and associated fluids at the inlet and outlet faces of the elemental area are:

$$\begin{aligned}
 H_N(z_N^-, t) = & \rho_{r_N} c_{r_N} (1 - \phi_N) (T_{st} - T_i) + \rho_{st} \phi_N L_{st} S_{st_N} (z_N^-, t) \\
 & + \rho_{st} c_w (T_{st} - T_i) \phi_N S_{w_N} (z_N^-, t) + \rho_w c_w (T_{st} - T_i) \phi_N S_{w_N} (z_N^-, t) \\
 & + \rho_o c_o (T_{st}) (T_{st} - T_i) \phi_N S_{o_N} (z_N^-, t) \quad (IX.3.13)
 \end{aligned}$$

$$\begin{aligned}
 H(z_N^+, t) = & \rho_r c_{r_N} (1 - \phi_N) \Delta T(z_N^+) + \rho_w (T_{z_N^+}) c_w \Delta T(z_N^+) \phi_N S_{w_N} (z_N^+, t) \\
 & + \rho_o (T_{z_N^+}) c_o \Delta T(z_N^+) \phi_N S_{o_N} (z_N^+, t) \quad (IX.3.14)
 \end{aligned}$$

Recall that the heat balance equation is:

$$Q_N(z_N^-, t) - Q_N(z_N^+, t) = v_N(t) [H_N(z_N^-, t) - H_N(z_N^+, t)] \quad (IX.3.15)$$

where:

$Q_N(z_N^-, t)$ is given by Equation (IX.3.10) but contains the still unknown rates of heat loss $Q_{L_j}(t)$ for $j = 1, \dots, N$

$Q_N(z_N^+, t)$ is given by Equation (IX.3.12)

$H_N(z_N^-, t)$ is given by Equation (IX.3.11)

$H_N(z_N^+, t)$ is given by Equation (IX.3.12)

The assumption is made that there is no convective heat flux across the steam front. This means that all the heat arriving at the steam front is used up in heating the rock and fluids. The condensate leaves at z_N^+ having the original reservoir temperature T_i . Thus $\Delta T(z_N^+) = 0$.

This assumption makes $Q_N(z_N^+, t) = 0$, and $H_N(z_N^+, t) = 0$. Therefore, the heat balance equation becomes:

$$Q_N(z_N^-, t) = v_N(t)H_N(z_N^-, t) \quad (\text{IX.3.16})$$

The unknown heat loss rates contained in the equation for $Q_N(z_N^-, t)$ is calculated next.

IX.4 RATE OF HEAT LOSS TO CAP AND BASE ROCK IN A PIECE-WISE HOMOGENEOUS LINEAR POROUS MEDIUM

The properties of the cap and base rock are assumed to be identical and constant even though those of the porous medium vary in regions (Figure 11). The cap and base rock are assumed to extend to infinity in either direction. The objective here becomes the calculation of the rate of heat loss to semi-infinite slabs on each side that represent the cap and base rocks.

Since the cap and base rock are assumed identical, the heat loss derivations will be made using only the cap rock, but the result is simply doubled to account for the base rock as well.

The steam front is assumed to be located in the N 'th region of a piece-wise homogeneous linear porous medium. Each of the N steam-invaded regions behind the front are divided into equal elemental areas (Figure 11). Let there be P_j such elemental areas in each j 'th region where $j = 1, 2, \dots, N$.

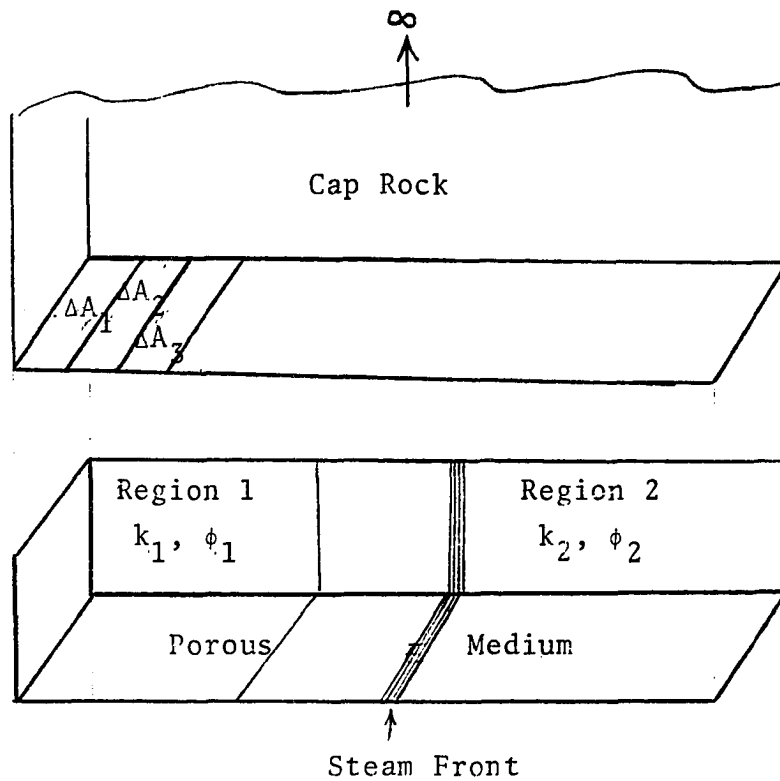


FIGURE 11. Schematic of linear porous medium for heat loss calculations.

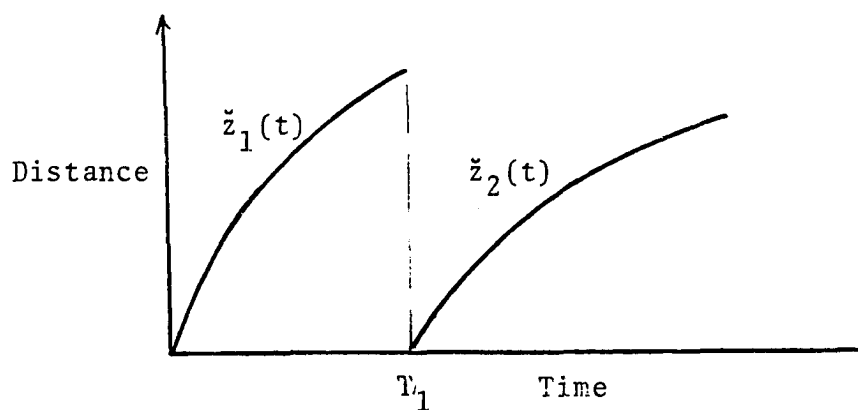


FIGURE 12. Distance as a function of time in each region of the linear piece-wise homogeneous porous medium.

The rate of heat loss per unit cross sectional area to the cap and base rock from an elemental area in an arbitrary j 'th region is:

$$Q_{L_j}(t) \Big|_{\Delta A_j} = -2k\Delta A_j \left. \frac{dT}{dz} \right|_{z=0} \quad (\text{IX.4.1})$$

which, after substituting for $\left. \frac{dT}{dz} \right|_{z=0}$ from Appendix D, becomes:

$$Q_{L_j}(t) \Big|_{\Delta A_j} = - \frac{2k\Delta A_j}{Wh} \left(\frac{T_s - T_i}{\sqrt{\pi\alpha(t - \tau_j)}} \right) \quad (\text{IX.4.1a})$$

Summing for all P_j such elemental areas in the j 'th region gives:

$$Q_{L_j}(t) = - \sum_{m=1}^{P_j} \frac{2k\Delta A_{j,m}}{Wh} \left(\frac{T_s - T_i}{\sqrt{\pi\alpha(t - \tau_{jm})}} \right) \quad (\text{IX.4.2})$$

Since $A_j = W z_j$, in the limit as A_j tends to zero,

$$\lim_{\Delta A \rightarrow 0} Q_{L_j}(t) = - \frac{2k(T_s - T_i)}{L} \int_0^{L_j} \frac{1}{\sqrt{\pi\alpha(t - \tau_j)}} dz_j(\tau_j) \quad (\text{IX.4.3})$$

where L_j is the length of the j 'th region.

Summing for all N regions behind the steam front gives the rate of heat loss per unit cross sectional area to cap and base rock from all N regions behind the steam front as:

$$Q_L(t) = \sum_{j=1}^N Q_{L_j}(t) = - \frac{2k(T_s - T_i)}{L} \sum_{j=1}^N \int_{t_{j-1}}^{t_j} \frac{\frac{dz_j(\tau_j)}{d\tau_j}}{\sqrt{\pi\alpha(t - \tau_j)}} d\tau_j \quad (\text{IX.4.4})$$

where $A_j(\tau) = Wz_j(\tau)$

W = width of the porous medium

$$\alpha = \frac{k}{\rho_{cb} c_{cb}}$$

ρ_{cb} = density of cap and base rock

c_{cb} = heat capacity of cap and base rock

k = thermal conductivity of cap and base rock

h = thickness of porous medium

In Equation (IX.4.4), each $z_j(\tau_j)$ is only defined and continuous in the region $t_{j-1} < \tau_j < t_j$. Since a function can have a differential quotient only at points where the function is continuous, the differential quotient $\frac{dz_j(\tau_j)}{d\tau_j}$ can only be taken at points within the region $t_{j-1} < \tau < t_j$ and not outside the region because $z_j(\tau_j)$ is not defined there. However, if $z_i(\tau_j)$ is expressed in terms of the Heaviside unit step function $H(t)$, then each $z_j(\tau_j)$ can be defined for all $\tau > 0$ with discontinuities at $\tau = t_{j-1}$ and $\tau = t_j$. A "generalized derivative" can be taken by product differentiation where generalization means a definition of the derivative at the points of discontinuity as well. Thus, introducing the Heaviside unit step function $H(t)$ into Equation (IX.4.4) gives:

$$Q_L(t) =$$

$$= \frac{-2k(T_s - T_i)}{L} \sum_{j=1}^N \int_0^t \frac{\frac{d}{d\tau} \{ [H(\tau - t_{j-1}) - H(\tau - t_j)] \xi_j(\tau) \}}{\sqrt{\pi\alpha(t - \tau)}} d\tau$$

(IX.4.5)

The summation can be taken inside the integral and $Q_L(t)$ can be expressed as:

$$Q_L(t) = \frac{-2k(T_s - T_i)}{h} \int_0^t \frac{\frac{d}{d\tau} \{ \sum_{j=1}^N [H(\tau - \tau_{j-1}) - H(\tau - \tau_j)] \xi_j(\tau) \}}{\sqrt{\pi\alpha(t - \tau)}} d\tau \quad (\text{IX.4.6})$$

Let

$$z(\tau) = \sum_{j=1}^N [H(\tau - \tau_{j-1}) - H(\tau - \tau_j)] \xi_j(\tau) \quad (\text{IX.4.7})$$

where each $\xi_j(t)$ is continuous and defined for all $\tau > 0$.

$$Q_L(t) = \frac{-2k(T_s - T_i)}{h} \int_0^t \frac{\frac{d}{d\tau} z(\tau)}{\sqrt{\pi\alpha(t - \tau)}} d\tau \quad (\text{IX.4.8})$$

$\frac{dz(\tau)}{d\tau}$ is obtained by the product differentiation of Equation (IX.4.7) to give

$$\begin{aligned} \frac{dz(\tau)}{d\tau} &= \sum_{j=1}^N \{ [H(\tau - \tau_{j-1}) - H(\tau - \tau_j)] \frac{d\xi_j}{d\tau} \\ &\quad + v(\tau) [\delta(\tau - \tau_{j-1}) - \delta(\tau - \tau_j)] \xi_j(\tau) \} \end{aligned} \quad (\text{IX.4.9})$$

$$\begin{aligned} &= \{ [H(\tau - \tau_{N-1}) - H(\tau - \tau_N)] \frac{d\xi_N}{d\tau} \\ &\quad + [\delta(\tau - \tau_{N-1}) - \delta(\tau - \tau_N)] \xi_N(\tau) \} \end{aligned} \quad (\text{IX.4.10})$$

where $H(\tau)$ is the Heaviside unit step function

$\delta(\tau)$ is the Dirac delta function

Equation (IX.4.8) is substituted into Equation (IX.3.10). The result is $Q_N(z_N^-, t)$ which is substituted into the left side of Equation (IX.3.16). When the expression for $H_N(z_N^-, t)$ is substituted into the right side of Equation (IX.3.16), the result is:

$$\begin{aligned}
 Q(0, t) &+ \frac{2k(T_s - T_i)}{h} \int_0^t \frac{\frac{d}{d\tau} z(\tau)}{\sqrt{\pi\alpha(t - \tau)}} d\tau \\
 &- \rho_{st} [L_{st} + c_w(T_{st} - T_i)] \phi_N v_N(t) [\bar{S}_{st_N} - S_{st}(z_N^-, t)] \\
 &- \rho_w(T_{st}) c_w(T_{st} - T_i) \phi_N v_N(t) [\bar{S}_{w_N} - S_{w_N}(z_N^-, t)] \\
 &- \rho_o(T_{st}) c_o(T_{st}) (T_{st} - T_i) \phi_N v_N(t) [\bar{S}_{o_N} - S_{o_N}(z_N^-, t)] \\
 &= v_N(t) [\rho_{\gamma_N} c_{\gamma_N} (1 - \phi_N) (T_{st} - T_i) + \rho_{st} L_{st} \phi_N S_{st_N}(z_N^-, t) \\
 &\quad + \rho_{st} c_s (T_{st} - T_i) \phi_N S_{w_N}(z_N^-, t) \\
 &\quad + \rho_w c_w (T_{st} - T_i) \phi_N S_{w_N}(z_N^-, t) \\
 &\quad + \rho_o(T_{z_N^+}) c_o(T_{st}) (T_{st} - T_i) \phi_N S_{o_N}(z_N^-, t)]
 \end{aligned}
 \tag{IX.4.10}$$

Considering that

$$Q(0, t) = M_{st}(0, t) [L_{st} + c_w(T_{st} - T_i)] + M_w(0, t) c_w(T_{st} - T_i)
 \tag{IX.4.11}$$

where

$$M_{st}(0, t) = \rho_{st} x_q(0, t)
 \tag{IX.4.12}$$

$$M_w(0,t) = \rho_w(T_{st})(1-x)q(0,t) \quad (IX.4.13)$$

then Equation (IX.4.10) simplifies to:

$$\begin{aligned} q(0,t) [xL_{st}\rho_{st} + \rho_w c_w(T_{st} - T_i)(1-x)] \\ + \frac{2k(T_s - T_i)}{h\sqrt{\pi\alpha}} \int_0^t \frac{\frac{d}{d\tau} z(\tau)}{\sqrt{t-\tau}} d\tau = V_N(t) [\rho_{r_N} c_{r_N}(1 - \phi_N) \\ \times (T_{st} - T_i) + \rho_{st} L_{st} \phi_N \bar{S}_{st_N} + \rho_w(T_{st}) c_w(T_{st} - T_i) \phi_N \bar{S}_{w_N} \\ + \rho_o(T_{st}) c_o(T_{st})(T_{st} - T_i) \phi_N \bar{S}_{o_N}] \quad (IX.4.13) \end{aligned}$$

Let

$$\begin{aligned} \bar{\rho c} = \rho_{r_N} c_{r_N}(1 - \phi_N) + \rho_w(T_{st}) c_w \bar{S}_{w_N} \phi_N \\ + \rho_o(T_{st}) c_o(T_{st}) \bar{S}_{o_N} \phi_N + \rho_{st} c_w \bar{S}_{st_N} \phi_N \quad (IX.4.14) \end{aligned}$$

Also let

$$\frac{dz(\tau)}{d\tau} = v(\tau) \quad (IX.4.15)$$

Then, introducing (IX.4.14) and (IX.4.15) into (IX.4.13) gives

$$\begin{aligned} q(0,t) [x\rho_{st} L_{st} + \rho_w c_s(1-x)(T_{st} - T_i)] + \frac{2k(T_{st} - T_i)}{h\sqrt{\pi\alpha}} \int_0^t \frac{v(\tau)}{\sqrt{t-\tau}} d\tau \\ = V_N(t) \left[\bar{\rho c} + \frac{\rho_{st} L_{st} \bar{S}_{st_N} \phi_N}{T_{st} - T_i} \right] (T_{st} - T_i) \quad (IX.4.16) \end{aligned}$$

where

$$v(\tau) = \sum_{j=1}^N \{ [H(t-t_{j-1}) - H(t-t_j)] \frac{d\xi_j(\tau)}{d\tau} + [\delta(t-t_{j-1}) - \delta(t-t_j)] \xi_j(\tau) \} \quad (\text{IX.4.17})$$

Equation (IX.4.16) is a generalization of the equation of Marx and Langenheim,²⁵ and Mandl and Volek³² derived for a piece-wise homogeneous linear porous medium. If the steam front is within the first region of a piece-wise homogeneous linear porous medium, the $N = 1$, and

$$\left. \begin{aligned} V_N(\tau) &= V_1(\tau) \\ V(\tau) &= V_1(\tau) \end{aligned} \right\} \begin{array}{l} \text{A continuous} \\ \text{function of time} \end{array}$$

Then Equation (IX.4.16) becomes

$$\begin{aligned} q(0,t) [\rho_{st} x L_{st} + c_w (T_{st} - T_i) (1 - x) \rho_w (T_{st})] \\ + \frac{2k(T_{st} - T_i)}{h\sqrt{\pi\alpha}} \int_0^t \frac{V_1(\tau)}{\sqrt{t-\tau}} d\tau = V_1(t) \left[\sum \frac{\overline{\rho c}}{\rho_{st} L_{st} \bar{S}_{st} \phi} \right] (T_{st} - T_i) \end{aligned} \quad (\text{IX.4.18})$$

which is the same as the equation of Mandl and Volek.

However, for a piece-wise homogeneous linear porous medium, $V(\tau)$ is a sectionally continuous function as given by (IX.4.17). After transformation to dimensionless form and simplification, Equation (IX.4.16) can be solved to give

the velocity and distance of the steam front in any region of a piece-wise homogeneous porous medium.

Since $\alpha = \frac{k}{\rho_{cb} c_{cb}}$, Equation (IX.4.16) can be expressed as

$$\begin{aligned}
 q(0,t) & [\rho_{st} x L_{st} + c_w (T_{st} - T_i) (1 - x) \rho_w (T_{st})] \\
 & - \sqrt{\frac{4k \rho_{cb} c_{cb} (T_{st} - T_i)^2}{\pi h^2}} \int_0^t \frac{v(\tau)}{\sqrt{t - \tau}} d\tau \\
 & = v_N(t) (T_{st} - T_i) \left[\bar{\rho} c + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right] \quad (IX.4.19)
 \end{aligned}$$

Define dimensionless variables.

Dimensionless time:

$$t_D = \left(\frac{4k}{\rho_{cb} c_{cb} h^2} \right) t \quad \text{or} \quad \tau_D = \left(\frac{4k}{\rho_{cb} c_{cb} h^2} \right) \tau \quad (IX.4.20)$$

Dimensionless distance:

$$z_D(t) = \frac{z(t)}{L} \quad (IX.4.21)$$

where L is the length of the entire porous medium.

Dimensionless velocity:

$$\begin{aligned}
 v_D(t) & = \frac{dz_D(t)}{dt_D} = \frac{dz_D(t)}{dt} \frac{dt}{dt_D} \\
 & = \frac{1}{L} \frac{dz(t)}{dt} \frac{\rho_{cb} c_{cb} h^2}{4k} \\
 & = \left(\frac{\rho_{cb} c_{cb} h^2}{4kL} \right) v(t) \quad (IX.4.22)
 \end{aligned}$$

Introducing the dimensionless variables into Equation (IX.4.19) gives:

$$\begin{aligned}
 q(0,t) & [\rho_{st} x L_{st} + c_w (T_{st} - T_i) (1 - x) \rho_w (T_{st})] \\
 & - \left(\frac{4k \rho_{cb} c_{cb}}{\pi h^2} \right)^{1/2} \int_0^{t_D} \left(\frac{4kL}{\rho_{cb} c_{cb} h^2} \right) \frac{\frac{dz(\tau_D)}{d\tau_D} \left(\frac{\rho_{cb} c_{cb} h^2}{4k} \right)}{\left(\frac{\rho_{cb} c_{cb} h^2}{4k} \right)^{1/2} \sqrt{t_D - \tau_D}} d\tau_D \\
 & = \left(\frac{4kL}{\rho_{cb} c_{cb} h^2} \right) (T_{st} - T_i) \frac{dz_N(t_D)}{dt_D} \left[\bar{\rho c} + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right]
 \end{aligned}
 \tag{IX.4.23}$$

which is expressed as

$$\begin{aligned}
 & \frac{q(0,t) [\rho_{st} x L_{st} + \rho_w c_w (T_{st} - T_i) (1 - x)] \rho_{cb} c_{cb} h^2}{4kL(T_{st} - T_i) \left[\bar{\rho c} + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right]} \\
 & - \frac{\rho_{cb} c_{cb}}{\sqrt{\pi} \left[\bar{\rho c} + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right]} \int_0^{t_D} \frac{\frac{dz_D(\tau_D)}{d\tau_D}}{\sqrt{t_D - \tau_D}} d\tau_D \\
 & = \frac{d}{dt_D} z_{N_D}(t_D)
 \end{aligned}
 \tag{IX.4.24}$$

Let

$$F_N = \frac{[\rho_{st} x L_{st} + \rho_w c_w (T_{st} - T_i) (1 - x)] \rho_{cb} c_{cb} h^2}{4kL(T_{st} - T_i) \left[\bar{\rho c} + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right]} \tag{IX.4.25}$$

Let

$$\lambda_N = \frac{\rho_{cb} c_{cb}}{\left[\sum \frac{\rho_{st}}{\rho_{cb}} + \frac{\rho_{st} L_{st} \bar{S}_{st} \phi_N}{T_{st} - T_i} \right]} \quad (\text{IX.4.26})$$

Introducing (IX.4.25) and (IX.4.26) into (IX.4.24) gives:

$$F_N q(0, t) - \frac{\lambda_N}{\sqrt{\pi}} \int_0^{t_D} \frac{v_D(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = v_{ND}(t_D) \quad (\text{IX.4.27})$$

Equation (IX.4.27) is the dimensionless form of Equation (IX.4.16) describing the velocity of the steam front in the Nth region of a piece-wise homogeneous linear porous medium. Note that $v_D(t_D) = \sum_{j=1}^m v_{jD}(t_D)$. Therefore, when the steam front is within the first region, $N = 1$, Equation (IX.4.27) becomes:

$$F_1 q_{st}(0, t) - \frac{\lambda_1}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{1D}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = v_{1D}(t_D) \quad (\text{IX.4.28})$$

where v_{1D} is a sectionally continuous function. If the reservoir is homogeneous, then v_{1D} becomes the continuous function $\check{v}_{1D}$. Dropping the subscripts, we have, for a homogeneous reservoir,

$$Fq(0, t) - \frac{\lambda}{\sqrt{\pi}} \int_0^{t_D} \frac{\check{v}_D(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = \check{v}_D(t_D) \quad (\text{IX.4.29})$$

The solution to Equation (IX.4.29) is obtained by Laplace transform method and is given in several references.^{25,32}

For the case of constant injection rate $[q(0,t) = q]$,

$$\xi_D(t_D) = \frac{Fq}{\lambda^2} \left(e^{\lambda^2 t_D} \operatorname{erfc} \sqrt{\lambda^2 t_D} + \frac{2\sqrt{\lambda^2 t_D}}{\sqrt{\pi}} - 1 \right) \quad (\text{IX.4.30})$$

$$v_D(t_D) = Fq \left(e^{\lambda^2 t_D} \operatorname{erfc} \sqrt{\lambda^2 t_D} \right) \quad (\text{IX.4.31})$$

The solution to Equation (IX.4.28) for the case where the steam front is within the first region of a piece-wise homogeneous linear system can also be obtained by the Laplace transform method as well.

Consider the case of a linear piece-wise homogeneous porous medium having only two regions ($N = 2$). When the steam front is in the second region, Equation (IX.4.27) becomes:

$$F_2 q(0,t) - \frac{\lambda_2}{\sqrt{\pi}} \int_0^{t_D} \frac{v_D(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = v_{2D}(t_D) \quad (\text{IX.4.32})$$

where

$$v_{2D}(t_D) = \{ [H(t_D - t_{D1})] \check{v}_{2D}(t_D) + [\delta(t_D - t_{D1}) - \delta(t_D - t_{D2})] \xi_{2D}(t_D) \} \quad (\text{IX.4.33})$$

and

$$\begin{aligned}
v_D(t_D) &= \frac{d\xi_D(t_D)}{dt_D} = \sum_{j=1}^2 \{ [H(t_D - t_{Dj-1}) - H(t_D - t_{Dj})] \\
&\quad \times \check{v}_{jD}(t_D) + [\delta(t_D - t_{Dj-1}) - \delta(t_D - t_{Dj})] \xi_{jD} \} \\
&= v_{1D}(t_D) + v_{2D}(t_D) \quad (IX.4.34)
\end{aligned}$$

The expressions for $\xi_{1D}(t_D)$ and $v_{1D}(t_D)$ are known and given by Equations (IX.4.30) and (IX.4.31), respectively. Substituting for $v_D(t_D)$ from Equation (IX.4.34) into (IX.4.32) gives:

$$\begin{aligned}
F_{2q}(0,t) &= \frac{\lambda_2}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{1D}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D - \frac{\lambda_2}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{2D}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D \\
&= v_{2D}(t_D) \quad (IX.4.35)
\end{aligned}$$

The second term on the left of Equation (IX.4.34) gives the rate of heat loss from the first region while the front is in the second region. In order to obtain an analytical solution to Equation (IX.4.35), the assumption is made that no heat is lost from the first region after the steam front clears it. Equation (IX.4.35) simplifies to:

$$F_{2q_{st}}(0,t) = \frac{\lambda_2}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{2D}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = v_{2D}(t_D) \quad (IX.4.36)$$

In general, the assumption is made that if the steam front is moving within the Nth region, the rate of heat loss from the N-1 regions behind the steam front is negligible. In

other words, only the region containing the steam front is losing heat at any given time. With this assumption, Equation (X.0.27) becomes:

$$F_N q(0, t) - \frac{\lambda_N}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{ND}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D = v_{ND}(t_D) \quad (\text{IX.4.37})$$

IX.5 SOLUTION BY LAPLACE TRANSFORM METHOD

The Laplace transform of each of the terms of Equation (IX.4.37) gives

$$\mathcal{L}\{F_N q_{st}(0, t)\} = F_N q(s) \quad (\text{IX.5.1})$$

$$\int_0^{t_D} \frac{v_{ND}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D$$

is a convolution integral. Therefore,

$$\mathcal{L}\left\{\frac{-\lambda_N}{\sqrt{\pi}} \int_0^{t_D} \frac{v_{ND}(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D\right\} = \frac{-\lambda_N}{\sqrt{\pi}} \left[\mathcal{L}\{v_{ND}(t_D)\} \cdot \mathcal{L}\left\{\frac{1}{\sqrt{t_D}}\right\} \right] \quad (\text{IX.5.2})$$

Since

$$\begin{aligned} v_{ND}(t_D) &= \{[H(t_D - t_{D_{N-1}})]v_{ND}(t_D) \\ &\quad + [\delta(t_D - t_{D_{N-1}}) - \delta(t_D - t_{D_N})]\xi_{ND}(t_D)\} \end{aligned} \quad (\text{IX.5.3})$$

$$\begin{aligned} \mathcal{L}\{v_D(t_D)\} &= \mathcal{L}\{[H(t_D - t_{D_{N-1}})] \frac{d\xi_{ND}(t_D)}{dt_D} \\ &+ [\delta(t_D - t_{D_{N-1}}) - \delta(t_D - t_{D_N})]\xi_{ND}(t_D)\} \quad (\text{IX.5.3a}) \end{aligned}$$

From Appendix D,

$$\mathcal{L}\{v_D(t_D)\} = s\xi_D(s) + \xi_D(0^-) \quad (\text{IX.5.4})$$

where

$$\begin{aligned} \xi_{ND}(0^-) &= 0 \\ \xi_{ND}(s) &= \int_{t_{D_{N-1}}}^{t_{D_N}} \xi_{ND}(\tau_D) e^{-s\tau_D} d\tau_D \end{aligned}$$

From standard tables of Laplace transforms,⁴⁹

$$\mathcal{L}\left\{\frac{1}{\sqrt{t_D}}\right\} = \frac{\sqrt{\pi}}{\sqrt{s}} \quad (\text{IX.5.5})$$

Introducing (IX.5.4) and (IX.5.5) into the left of Equation (IX.4.28) gives:

$$\begin{aligned} \mathcal{L}\left\{\frac{-\lambda_N}{\sqrt{\pi}} \int_0^{t_D} \frac{v_D(\tau_D)}{\sqrt{t_D - \tau_D}} d\tau_D\right\} &= \frac{-\lambda_N}{\sqrt{\pi}} s\xi_D(s) \frac{\sqrt{\pi}}{\sqrt{s}} \\ &= -\lambda_N \sqrt{s} \xi_D(s) \quad (\text{IX.5.6}) \end{aligned}$$

Considering Equations (IX.5.4), (IX.5.5) and (IX.5.6) the Laplace transform of Equation (IX.4.28) becomes:

$$F_N q(s) - \lambda_N \sqrt{s} \xi_{ND}(s) = s\xi_D(s) \quad (\text{IX.5.7})$$

$$\xi_D(s) = \frac{F_N q(s)}{\sqrt{s}(\sqrt{s} + \lambda_N)} \quad (\text{IX.5.8})$$

In order to take the inverse transform of Equation (IX.5.8) it is rewritten as:

$$\xi_{ND}(s) = \frac{F_N}{\lambda_N^2} \frac{\lambda_N^2 q(s)}{\sqrt{s}(\sqrt{s} + \lambda_N)} \quad (\text{IX.5.9})$$

The inverse transform of Equation (IX.5.9) is:

$$z_{ND}(t_D) = \frac{F_N}{\lambda_N^2} \int_0^{t_D} q(0,t) e^{\lambda_N(t_D - \tau_D)} \operatorname{erfc} \sqrt{\lambda_N(t_D - \tau_D)} d\tau_D \quad (\text{IX.5.10})$$

where

$$z_D(t_D) = \sum_{j=1}^N [H(t_D - t_{Dj-1}) - H(t_D - t_{Dj})] \xi_{Dj}(t_D) \quad (\text{IX.5.11})$$

IX.6 CONSTANT INJECTION RATE

In the case that the injection rate is constant,

$$q_{st}(0,t) = \text{constant} = q$$

Therefore, the Laplace transform of Equation (IX.4.28) is:

$$\frac{F_N q}{s} - \lambda_N \sqrt{s} \xi_D(s) = s \xi_D(s) \quad (\text{IX.6.1})$$

from which,

$$\xi_D(s) = \frac{F_N q}{\sqrt{s}(s)(\sqrt{s} + \lambda_N)} \quad (\text{IX.6.3})$$

The inverse transform of which is:

$$z_{ND}(t_D) = \frac{F_{Nq}}{\lambda_N^2} \left[e^{\lambda_N^2 t_D} \operatorname{erfc} \sqrt{\lambda_N^2 t_D} + \frac{2\sqrt{\lambda_N^2 t_D}}{\sqrt{\pi}} - 1 \right] \quad (\text{IX.6.4})$$

For the case of constant injection rates, the velocity is:

$$v_{ND}(t_D) = \frac{dz_{ND}(t_D)}{dt_D} = F_{Nq} [e^{\lambda_N^2 t_D} \operatorname{erfc} \sqrt{\lambda_N^2 t_D}] \quad (\text{IX.6.5})$$

where

$$v_D(t_D) = \sum_{j=1}^N \{ [H(t_D - t_{D_{j-1}}) - H(t_D - t_{D_j})] \frac{d\xi_{Dj}(t_D)}{dt_D} + [\delta(t_D - t_{D_{j-1}}) - \delta(t_D - t_{D_j})] \xi_{Dj}(t_D) \}$$

IX.7 OIL RECOVERY CALCULATION

Oil recovery calculations are made based on the knowledge of the steam invaded volume and the residual oil saturation. This is a very simplified approach which ignores the oil recovery by the other mechanisms such as recovery by the hot and cold water zones ahead of the steam front; recovery by steam distillation; recovery by viscosity reduction; recovery by thermal expansion; and recovery by the gas drive mechanism of the steam. If it is assumed that the oil saturation is reduced from the initial saturation (S_{oi}) to the residual saturation (S_{or}) for each of N permeability zones in a streamtube swept by steam, and the steam front is at position $z_N(t)$ in the N th zone at time

t, then the oil recovered up to time t in a streamtube is:

$$\text{Rec} = \frac{1}{5.61} \left[\sum_{j=1}^{N-1} \{xWh\phi(S_{o_i} - S_{o_r})\}_j + z_N W_N h \phi_N (S_{o_i} - S_{o_r})_N \right] \quad (\text{barrels}) \quad (\text{IX.7.1})$$

If steam has broken through in a streamtube at time t_b , Equation IX.7.1 becomes:

$$\text{Rec} = \frac{1}{5.61} \sum_{j=1}^N \{xWh\phi(S_{o_i} - S_{o_r})\}_j \quad (\text{barrels}) \quad (\text{IX.7.2})$$

After the breakthrough time (t_b) for any streamtube, no more oil is recovered from that streamtube. Thus, the breakthrough times for each of the zones in a streamtube are required for the recovery calculation. Since the length of each permeability zone in a streamtube is known, by starting from time $t = 0$, the time for the steam front to clear each zone in sequence is calculated through the knowledge of the steamfront velocity with time. Equation IX.7.1 is applied to each streamtube in a producer and summed for all streamtubes in the producer. The results are added for all wells to give the recovery from the field at the specified time. For any streamtube where steam has broken through, Equation IX.7.2 is applied.

A computer program was written in FORTRAN to perform the oil recovery calculations with time. Appendix F gives a short description of the program, a listing of the flow chart. A listing of the program itself is given in Appendix O.

CHAPTER X

TEST FOR THE VALIDITY OF BOUNDARY ELEMENT METHOD OF SOLUTION

The validity of the boundary element technique and the correctness of the solution algorithm were verified by applying the method to a simple problem where the analytical solution is known. The example is taken from Street⁵³ and involves the steady two-dimensional irrotational flow of an ideal fluid in a corner (Figure 13).

The mathematical model is:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0, \quad 0 < x < \pi, \quad 0 < y < \pi$$

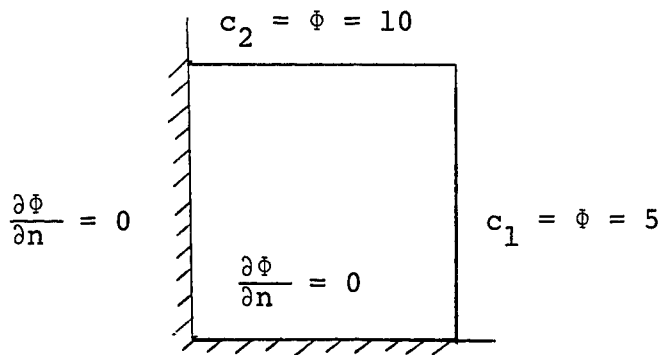


FIGURE 13. Steady flow of an ideal fluid in a corner.

with boundary conditions:

$$\frac{\partial \phi(0,y)}{\partial x} = 0 \quad 0 < y < \pi$$

$$\phi(\pi,y) = 5 \quad 0 < y \leq \pi$$

$$\frac{\partial \phi(x,0)}{\partial y} = 0 \quad 0 < x < \pi$$

$$\phi(x,\pi) = 10 \quad 0 < x < \pi$$

The solution to this problem is given⁵⁸ as:

$$\phi(x,y) = C_1 + \frac{4(C_2 - C_1)}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n-1} \cosh[(n-\frac{1}{2})y] \cos[(n-\frac{1}{2})x]}{(2n-1) \cosh[(n-\frac{1}{2})\pi]}$$

where $C_1 = 5$ and $C_2 = 10$.

The boundary element method of solution is obtained by dividing each side into three equal segments. The nodes are taken at the mid-points of the segments as shown in Figure 14.

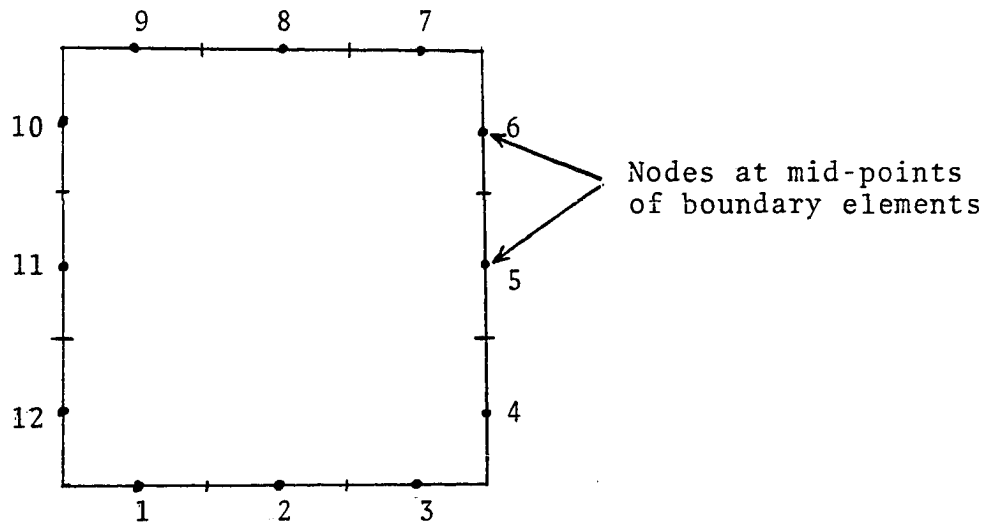


FIGURE 14. Division of square domain into segments and numbering scheme.

The results from the boundary element solution method are shown below. Table 1 lists the input data. Table 2 shows the calculated potentials at the boundary nodes and the potentials calculated at selected interior points. Table 3 compares the solutions obtained by the boundary element method at the interior points to that obtained analytically. The distribution of the absolute values of the errors $|\epsilon|$ at the interior points are calculated and plotted in Figure 15 where $|\epsilon|$ is defined as:

$$|\epsilon| = \left| \frac{\phi(\text{BEM}) - \phi(\text{Analytical})}{\phi(\text{Analytical})} \right|$$

It can be seen from Table 3 that at points farther than half an element length away from any boundary, the maximum error was less than 0.4 percent. Figure 15 shows that the error gets very large on and near the boundaries up to half an element length away from the boundaries.

TABLE 1. Input Data for Test Program

TEST FOR THE VALIDITY OF THE BOUNDARY ELEMENT METHOD (BEM)
 CALCULATION OF POTENTIALS IN A SQUARE USING THE (BEM)
 RESULTS COMPARED WITH ANALYTICAL SOLUTION GIVEN BY STREET

INPUT DATA

NUMBER OF BOUNDARY ELEMENTS=12
 NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED=15
 NUMBER OF SOURCES AND SINKS= 0

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X	Y
1	0.0000E 00	0.0000E 00
2	0.1047E 01	0.0000E 00
3	0.2094E 01	0.0000E 00
4	0.3142E 01	0.0000E 00
5	0.3142E 01	0.1047E 01
6	0.3142E 01	0.2094E 01
7	0.3142E 01	0.3142E 01
8	0.2094E 01	0.3142E 01
9	0.1047E 01	0.3142E 01
10	0.0000E 00	0.3142E 01
11	0.0000E 00	0.2094E 01
12	0.0000E 00	0.1047E 01

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	1	0.0000E 00
2	1	0.0000E 00
3	1	0.0000E 00
4	0	0.5000E 01
5	0	0.5000E 01
6	0	0.5000E 01
7	0	0.1000E 02
8	0	0.1000E 02
9	0	0.1000E 02
10	1	0.0000E 00
11	1	0.0000E 00
12	1	0.0000E 00

TABLE 2 RESULTS AT BOUNDARY NODES

BOUNDARY NODES		POTENTIAL	DERIVATIVE
X	Y		
0.5235E 00	0.0000E 00	0.7431629E 01	0.0000000E 00
0.1570E 01	0.0000E 00	0.6809022E 01	0.0000000E 00
0.2618E 01	0.0000E 00	0.5646283E 01	0.0000000E 00
0.3142E 01	0.5235E 00	0.5000000E 01	-0.1448967E 01
0.3142E 01	0.1570E 01	0.5000000E 01	-0.1798650E 01
0.3142E 01	0.2618E 01	0.5000000E 01	-0.7014348E 01
0.2618E 01	0.3142E 01	0.1000000E 02	0.7014794E 01
0.1570E 01	0.3142E 01	0.1000000E 02	0.1794951E 01
0.5235E 00	0.3142E 01	0.1000000E 02	0.1448620E 01
0.0000E 00	0.2618E 01	0.9354136E 01	0.0000000E 00
0.0000E 00	0.1570E 01	0.8193880E 01	0.0000000E 00
0.0000E 00	0.5235E 00	0.7569907E 01	0.0000000E 00

POTENTIALS AT SELECTED INTERNAL POINTS

X	Y	POTENTIAL
0.1000E 01	0.1000E 01	0.7501E 01
0.1500E 01	0.1500E 01	0.7501E 01
0.2000E 01	0.2000E 01	0.7500E 01
0.2500E 01	0.2500E 01	0.7500E 01
0.3000E 01	0.3000E 01	0.7033E 01
0.1000E 01	0.1571E 01	0.7927E 01
0.1500E 01	0.1571E 01	0.7566E 01
0.2000E 01	0.1571E 01	0.7013E 01
0.2500E 01	0.1571E 01	0.6232E 01
0.3000E 01	0.1571E 01	0.4877E 01
0.1000E 01	0.5000E 00	0.7282E 01
0.1500E 01	0.5000E 00	0.6928E 01
0.2000E 01	0.5000E 00	0.6446E 01
0.2500E 01	0.5000E 00	0.5841E 01
0.3000E 01	0.5000E 00	0.4825E 01

TABLE 3 COMPARISON WITH ANALYTICAL SOLUTION

INTERNAL POINTS		POTENTIAL (BEM)		POTENTIAL (ANALYTICAL)	
X	Y				
0.1000E 01	0.1000E 01	0.7501070E 01	0.7500757E 01		
0.1500E 01	0.1500E 01	0.7500750E 01	0.7500802E 01		
0.2000E 01	0.2000E 01	0.7500497E 01	0.7500597E 01		
0.2500E 01	0.2500E 01	0.7500035E 01	0.7500886E 01		
0.3000E 01	0.3000E 01	0.7033089E 01	0.7504594E 01		
0.1000E 01	0.1571E 01	0.7926839E 01	0.7919108E 01		
0.1500E 01	0.1571E 01	0.7566373E 01	0.7565387E 01		
0.2000E 01	0.1571E 01	0.7012587E 01	0.7021172E 01		
0.2500E 01	0.1571E 01	0.6232046E 01	0.6254806E 01		
0.3000E 01	0.1571E 01	0.4877094E 01	0.5291536E 01		
0.1000E 01	0.5000E 00	0.7282336E 01	0.7292308E 01		
0.1500E 01	0.5000E 00	0.6928425E 01	0.6944761E 01		
0.2000E 01	0.5000E 00	0.6445957E 01	0.6463688E 01		
0.2500E 01	0.5000E 00	0.5841250E 01	0.5867875E 01		
0.3000E 01	0.5000E 00	0.4825499E 01	0.5196271E 01		

DISTRIBUTION OF ABSOLUTE ERROR

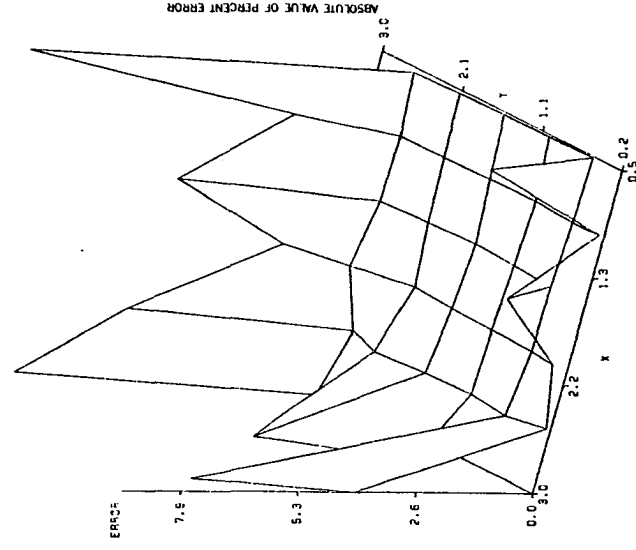


FIGURE 15. Distribution of error.

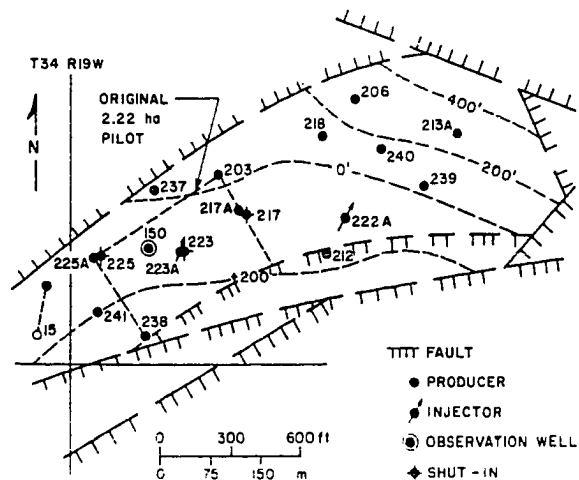
CHAPTER XI

APPLICATIONS AND RESULTS

XI.1 SAMPLE FIELD PROBLEM: SHIELLS CANYON FIELD⁵⁰

A steam-distillation drive pilot project in zone 203 of the Shiells Canyon field, Ventura County, California, has been reported by Konopnicki, et al.⁵⁰ from which this example is taken. Zone 203 is described as part of the Sespe formation of Oligocene age located in the Shiells Canyon field. The average sand thickness is 48.8 m (160 ft). The reservoir is bordered on all sides by faults that are assumed to be fluid sealing at the flood pressure (Fig. 16) The reservoir has a 35° dip, but for purposes of this application, it is assumed to be horizontal.

FIGURE 16. Shiells Canyon Field--Zone 203
(After Konopnicki, et al.⁵⁰)



The primary producing mechanism has been solution-gas drive. The cumulative production before steam injection was $42,500 \text{ m}^3$ (267,000 bbl) which is 9.5% of the estimated $449,175 \text{ m}^3$ (2,825,000 bbl) of oil originally in place. The pilot is an inverted type pattern originally consisting of four producing wells (217A, 203, 237 and 225A) down-dip from the injection well, one producing well (238) up-dip from the injection well, and one thermal observation well (150). Two additional producers (241 and 15) were drilled during the flood (Fig. 16).

XI.1.1 Boundary Conditions

Konopnicki et al. report that the reservoir is bordered on all sides by faults that are assumed to be fluid sealing. However, their figure (Fig. 16) shows the west boundary to be undefined. In the present analysis, two kinds of conditions are assigned at the west boundary for comparative analysis. The two types of boundary conditions are:

- (a) The west boundary is assumed to be an oil-water contact; thus, it is a constant-pressure boundary.
- (b) The west boundary is assumed to be fluid sealing, therefore making all the boundaries fluid sealing.

In what follows, the boundary element method of generating streamtubes and the steamflood calculation method are applied to the Shiells Canyon reservoir by considering it as:

1. A homogeneous reservoir with
 - a. sealed outer boundary
 - b. the west boundary assumed to be an oil/water contact and therefore at original reservoir pressure, the rest of the boundary sealed.
2. A piece-wise homogeneous reservoir consisting of two regions that have unequal permeabilities with
 - a. sealed outer boundary
 - b. the west boundary at constant pressure, the remainder of the boundary sealed.

XI.2 CASE 1:

XI.2.1 Shiells Canyon Field Analysed as a Single Region Homogeneous Reservoir

The discretization of the boundary and the scheme of numbering are shown in Figure 17. The numbering scheme was such that the domain under consideration was always on the left hand side. This scheme ensures that the direction of the normal gradient at any point of the boundary (or boundaries) is always away from the domain. There are a total of sixty elements taken on the boundary, but only a select few of them are indicated in Figure 17. The coordinates of the extreme points of all the boundary elements are shown in Table 4, while Table 5 gives the coordinates and strengths of the injectors and producers.

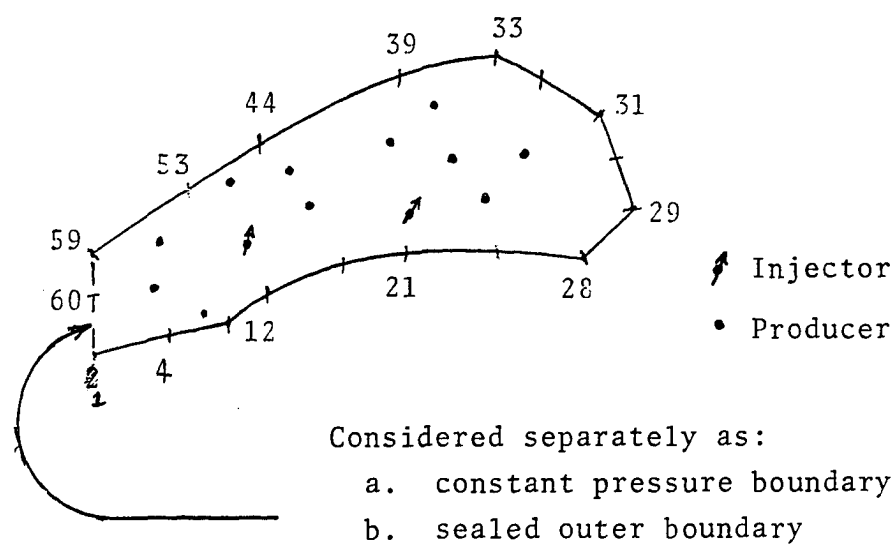


FIGURE 17. Discretization and numbering scheme for Shiells Canyon Field as a homogeneous reservoir.⁵⁰

TABLE 4. The Coordinates of the End Points of the Boundary Elements for Shiells Canyon Reservoir Analysed as a Single Homogeneous Region

POINT	X(INCH)	Y(INCH)
1	0.1000	0.1500
2	0.2500	0.2000
3	0.4000	0.2500
4	0.5500	0.2800
5	0.6500	0.3000
6	0.7000	0.3250
7	0.7500	0.3300
8	0.8000	0.3400
9	0.8500	0.3500
10	0.9000	0.4000
11	1.0000	0.4500
12	1.1000	0.5250
13	1.2500	0.6000
14	1.4000	0.6500
15	1.5500	0.6800
16	1.6500	0.6900
17	1.7000	0.7000
18	1.7500	0.7100
19	1.8000	0.7200
20	1.8500	0.7250
21	1.9000	0.7300
22	1.9500	0.7350
23	2.1000	0.7450
24	2.2500	0.7500
25	2.4000	0.7500
26	2.6000	0.7400
27	2.7500	0.7300
28	2.9000	0.7250
29	3.1600	1.0000
30	3.0600	1.2500
31	2.9500	1.5500
32	2.7500	1.6500
33	2.4000	1.8200
34	2.2500	1.8000
35	2.1000	1.7500
36	2.0500	1.7450
37	2.0000	1.7250
38	1.9500	1.7100
39	1.8500	1.6900
40	1.7500	1.6500
41	1.6500	1.6100
42	1.5000	1.5500
43	1.3500	1.4800
44	1.2500	1.4400
45	1.2000	1.4000
46	1.1500	1.3800
47	1.0500	1.3250
48	0.9500	1.2600
49	0.9000	1.2500
50	0.8500	1.2000
51	0.8000	1.1700
52	0.7500	1.1500
53	0.6500	1.0800
54	0.5500	1.0200
55	0.5000	0.9900
56	0.4500	0.9500
57	0.4000	0.9200
58	0.3000	0.8500
59	0.1000	0.7300
60	0.1000	0.5000

TABLE 5. Coordinates of Sources and Sinks for Shiells Canyon Field Analysed as a Single Homogeneous Region

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
 THE CHARACTERISTIC PRESSURE(Psi) IS= 85.0000

XI.3 RESULTS FOR CASE 1: SHIELLS CANYON FIELD MODELLED AS A HOMOGENEOUS RESERVOIR

XI.3.1 Sealed Outer Boundary

Table 6 lists the prescribed boundary conditions according to a preassigned code. If the potential is prescribed on the boundary segment, the code = 0. If the potential gradient is prescribed, the code = 1.

A plot of the streamlines generated is shown in Figure 18, while Table 7 gives the dimensions of the streamtubes generated which are read as input to the steamflood recovery calculation program. The input data to the steamflood program are given in Table 8. Figure 19 gives a plot of the oil recovery as a function of time. A complete listing of the outputs from the streamline modelling program is presented in Appendix G while the detailed output from the steamflood program is listed in Appendix H.

XI.3.2 Part of the Boundary at Constant Pressure, the Remainder Sealed

The only difference between the data for this case and that of the sealed boundary is the constant potential specified to boundary segments numbered 59 and 60. These two segments are assigned the initial pressure of 85 psi. Figure 20 shows a plot of the streamlines generated. Table 9 gives the dimensions of the streamtubes generated and Figure 21 gives a plot of the oil recovery as a function

TABLE 6. The Prescribed Boundary Conditions for Shiells
Canyon Reservoir Analysed as a Homogeneous Reservoir
with Sealed Boundary

BOUNDARY CONDITIONS		
NODE	CODE	PRESCRIBED VALUE
1	1	0.0
2	1	0.0
3	1	0.0
4	1	0.0
5	1	0.0
6	1	0.0
7	1	0.0
8	1	0.0
9	1	0.0
10	1	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	1	0.0
35	1	0.0
36	1	0.0
37	1	0.0
38	1	0.0
39	1	0.0
40	1	0.0
41	1	0.0
42	1	0.0
43	1	0.0
44	1	0.0
45	1	0.0
46	1	0.0
47	1	0.0
48	1	0.0
49	1	0.0
50	1	0.0
51	1	0.0
52	1	0.0
53	1	0.0
54	1	0.0
55	1	0.0
56	1	0.0
57	1	0.0
58	1	0.0
59	1	0.0
60	1	0.0

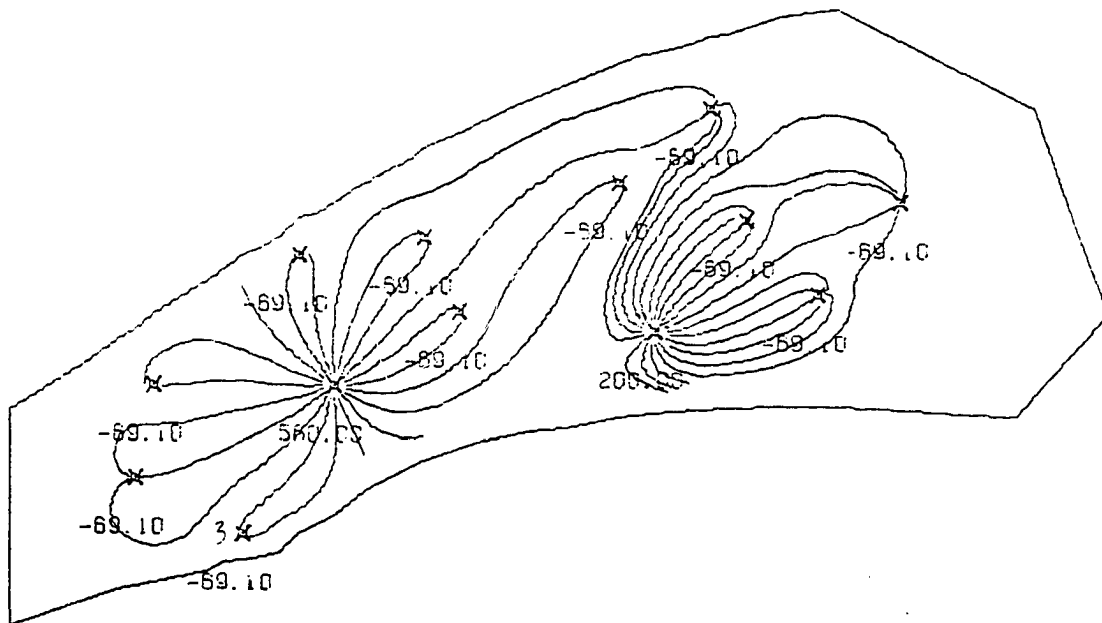


FIGURE 18. Boundary element modelling of Shiells Canyon Field as a single homogeneous reservoir with sealed boundary.

TABLE 7. Calculated Streamtube Dimensions for Shiells Canyon Field Analysed as a Homogeneous Reservoir with Sealed Outer Boundary

WELL NO.	S/L NO.	CODE	LENGTH (ft)	WIDTH (ft)	RATE (bbl/d)
1	1	1	538.10	45.18	28.00
1	2	1	425.00	23.75	28.00
1	3	1	686.10	79.51	28.00
2	1	1	441.40	27.76	28.00
2	2	1	344.70	19.18	28.00
3	1	1	334.80	18.06	28.00
3	2	1	384.80	21.76	28.00
4	1	1	269.70	13.51	28.00
4	2	1	269.70	14.88	28.00
5	1	1	359.80	15.89	28.00
5	2	1	338.10	14.26	28.00
6	1	1	294.70	12.69	28.00
6	2	1	288.10	11.96	28.00
7	1	1	809.80	29.41	28.00
7	2	1	759.80	28.63	28.00
8	1	1	112.00	38.76	28.00
8	2	1	950.00	34.34	28.00
8	3	1	588.10	15.61	10.00
8	4	1	569.70	12.31	10.00
8	5	1	598.00	13.78	10.00
9	1	1	266.40	6.27	10.00
9	2	1	291.40	6.74	10.00
9	3	1	378.30	8.90	10.00
9	4	1	288.10	6.46	10.00
10	1	1	463.10	10.18	10.00
10	2	1	369.70	8.22	10.00
10	3	1	331.60	7.53	10.00
10	4	1	334.80	7.39	10.00
10	5	1	403.30	10.92	10.00
11	1	1	716.40	22.42	10.00
11	2	1	903.30	34.43	10.00
11	3	1	713.10	19.76	10.00
11	4	1	538.10	15.60	10.00
11	5	1	606.60	22.23	10.00

TABLE 8. Input Data to Steamflood Recovery Prediction Program

TOTAL NUMBER OF PRODUCERS=	11
STEAM TEMPERATURE (DEG. F)=	430.0000
STEAM QUALITY=	0.7000
CONVERGENCE LIMIT=	0.0010
RESERVOIR THICKNESS (FT)=	160.0000
STEAM PRESSURE (PSI)=	344.0000
INITIAL RESERVOIR TEMPERATURE(DEG F)=	105.0000

FLUID PROPERTIES

	WATER	OIL	STEAM
DENSITY(LB/CU FT) AT STD. TEMP.	62.4000	55.0000	
DENSITY(LB/CU FT) AT STEAM TEMP.	52.3800	48.7000	0.7431
SPECIFIC HEAT(BTU/LB *F)	1.0000	0.4880	1.0000
LATENT HEAT(BTU/LB)			789.5100

PROPERTIES OF THE CAP AND BASE ROCK

DENSITY(LB/CU. FT)	149.0000
SPECIFIC HEAT(BTU/LB-*F)	0.2130
THERM. COND.(BTU/HR-FT-*F)	1.1000

RESERVOIR ROCK PROPERTIES

	REGION 1	REGION 2
POROSITY	0.2050	0.1800
DEN.(LB/CU FT)	165.0000	165.0000
SPEC. HEAT(BTU/LB*F)	0.2000	0.2000
PERMEABILITY(MD)	140.0000	70.0000

INITIAL FLUID SATURATIONS

	REGION 1	REGION 2
WATER	0.5500	0.5500
OIL	0.4500	0.4500
STEAM	0.0	0.0

AVERAGE RESIDUAL SATURATIONS BEHIND THE FRONT

	REGION 1	REGION 2
WATER	0.5800	0.5800
OIL	0.1800	0.1800
STEAM	0.2400	0.2400

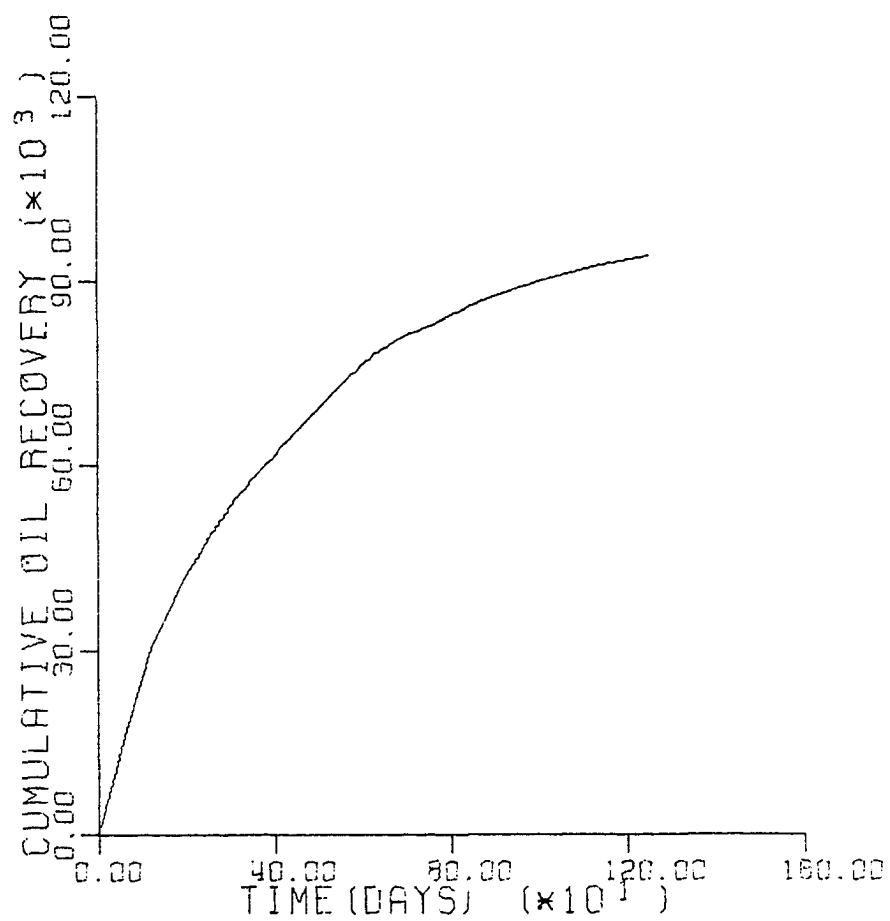


FIGURE 19. Steamflood recovery prediction for Shiells Canyon Field as a homogeneous reservoir with sealed outer boundary.

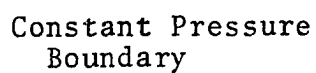


FIGURE 20. Boundary element modelling of Shiells Canyon Field as a homogeneous reservoir. Part of the boundary at constant pressure, part sealed.

TABLE 9. Calculated Streamtube Dimensions for Shiells Canyon Field Analysed as a Homogeneous Reservoir. Part of the Boundary at Constant Pressure, the Remainder Sealed.

WELL NO.	S/L NO.	CODE	LENGTH (ft)	WIDTH (ft)	RATE (bbl/d)
1	1	1	541.40	45.84	28.00
1	2	1	425.00	23.67	28.00
1	3	1	678.30	73.69	28.00
2	1	1	441.40	27.62	28.00
2	2	1	344.70	19.14	28.00
3	1	1	338.10	18.06	28.00
3	2	1	384.80	21.62	28.00
4	1	1	269.70	13.51	28.00
4	2	1	269.70	14.91	28.00
5	1	1	359.80	15.91	28.00
5	2	1	338.10	14.26	28.00
6	1	1	294.70	12.70	28.00
6	2	1	288.10	11.96	28.00
7	1	1	809.80	29.43	28.00
7	2	1	759.80	28.67	28.00
8	1	1	1120.00	38.69	28.00
8	2	1	950.00	34.38	28.00
8	3	1	588.10	15.55	10.00
8	4	1	569.70	12.31	10.00
8	5	1	598.00	13.79	10.00
9	1	1	266.40	6.27	10.00
9	2	1	291.40	6.75	10.00
9	3	1	378.30	8.92	10.00
9	4	1	288.10	6.46	10.00
10	1	1	463.10	10.17	10.00
10	2	1	369.70	8.22	10.00
10	3	1	331.60	7.53	10.00
10	4	1	334.80	7.39	10.00
10	5	1	403.30	10.97	10.00
11	1	1	716.40	22.47	10.00
11	2	1	903.30	34.84	10.00
11	3	1	713.10	19.81	10.00
11	4	1	538.10	15.61	10.00
11	5	1	606.60	22.34	10.00

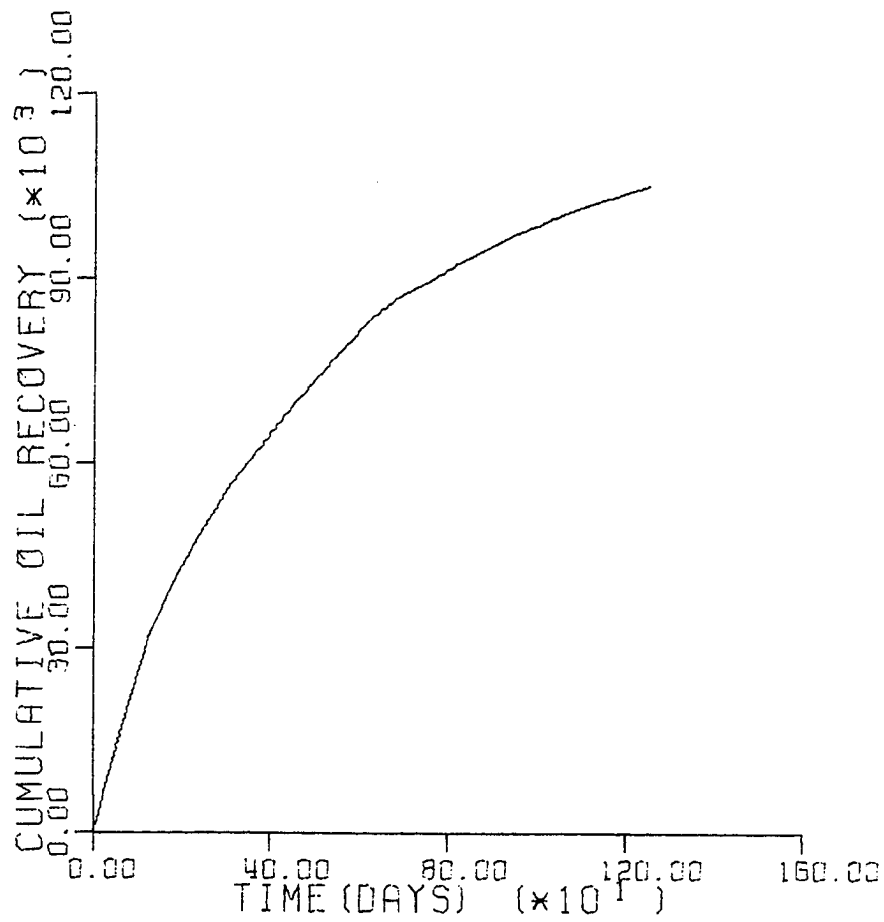


FIGURE 21. Steamflood recovery prediction for Shiells Canyon Field as a homogeneous reservoir with part of the boundary at constant pressure, part sealed.

of time. The detailed output from the streamline modelling program is listed in Appendix I, while Appendix J presents the detailed output of the steamflood recovery program.

XI.4 CASE 2:

XI.4.1 Shiells Canyon Field Analysed as a Piecewise Homogeneous Reservoir

The Shiells Canyon field has been modified to consist of two regions with different permeabilities (Figure 22). Region 1 has an average permeability of $0.138 \mu\text{m}^2$ (140 md) while the second region is assigned a permeability half as much, $0.69 \mu\text{m}^2$ (70 md). Figure 23 shows the numbering scheme employed. The scheme was such that the domain of interest was always on the left as the numbering progressed. It is to be observed that the numbering of the second region starts where that for the first ends on the interface between the two regions. This was done purely for convenience in assembling the two matrices into one. The reservoir is analysed assuming first that the entire reservoir is enclosed by sealing faults, and then by assuming that the west boundary is at constant pressure.

The coordinates of the end points of the boundary elements for both regions are listed in Table 10. The coordinates and strengths of the injectors and producers are listed in Table 11.

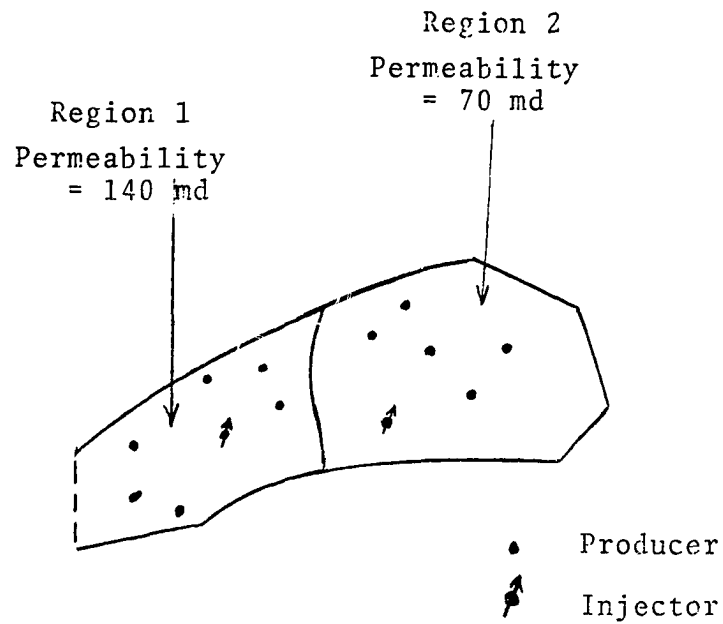
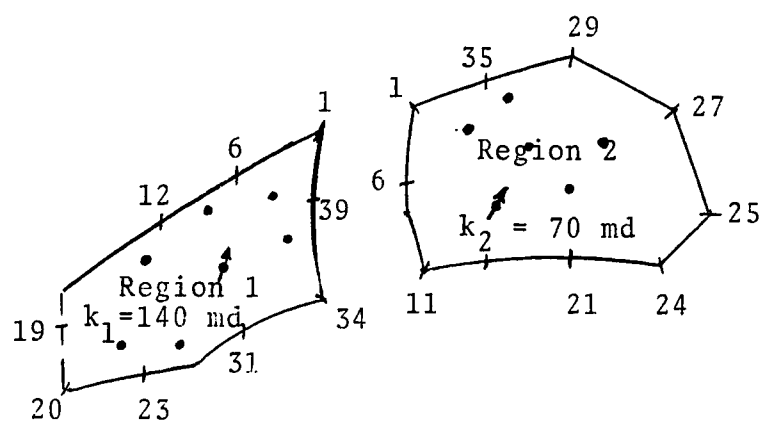


FIGURE 22. Shiells Canyon Field (Zone 203) arbitrarily divided into two regions with different material properties.



Scale: 1" = 738.46 ft

FIGURE 23. Shiells Canyon Field (Zone 203) as a piecewise-homogeneous field showing discretization and numbering scheme.

TABLE 10. Case 2:

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

Region 1			Region 2		
POINT	X(INCH)	Y(INCH)	POINT	X(INCH)	Y(INCH)
1	1.5000	1.5500	1	1.5000	1.5500
2	1.3500	1.4800	2	1.5000	1.5000
3	1.2500	1.4400	3	1.4800	1.4500
4	1.2000	1.4000	4	1.4600	1.3500
5	1.1500	1.3800	5	1.4600	1.2500
6	1.0500	1.3250	6	1.4800	1.1500
7	0.9500	1.2600	7	1.5000	1.0500
8	0.9000	1.2500	8	1.5000	0.9500
9	0.8500	1.2000	9	1.5400	0.8500
10	0.8000	1.1700	10	1.5500	0.7500
11	0.7500	1.1500	11	1.5500	0.6800
12	0.6500	1.0800	12	1.6500	0.6900
13	0.5500	1.0200	13	1.7000	0.7000
14	0.5000	0.9900	14	1.7500	0.7100
15	0.4500	0.9500	15	1.8000	0.7200
16	0.4000	0.9200	16	1.8500	0.7250
17	0.3000	0.8500	17	1.9000	0.7300
18	0.1000	0.7300	18	1.9500	0.7350
19	0.1000	0.5000	19	2.1000	0.7450
20	0.1000	0.1500	20	2.2500	0.7500
21	0.2500	0.2000	21	2.4000	0.7500
22	0.4000	0.2500	22	2.6000	0.7400
23	0.5500	0.2800	23	2.7500	0.7300
24	0.6500	0.3000	24	2.9000	0.7250
25	0.7000	0.3250	25	3.1600	1.0000
26	0.7500	0.3300	26	3.0600	1.2500
27	0.8000	0.3400	27	2.9500	1.5500
28	0.8500	0.3500	28	2.7500	1.6500
29	0.9000	0.4000	29	2.4000	1.8200
30	1.0000	0.4500	30	2.2500	1.8000
31	1.1000	0.5250	31	2.1000	1.7500
32	1.2500	0.6000	32	2.0500	1.7450
33	1.4000	0.6500	33	2.0000	1.7250
34	1.5500	0.6800	34	1.9500	1.7100
35	1.5500	0.7500	35	1.8500	1.6900
36	1.5400	0.8500	36	1.7500	1.6500
37	1.5000	0.9500	37	1.6500	1.6100
38	1.5000	1.0500			
39	1.4800	1.1500			
40	1.4600	1.2500			
41	1.4600	1.3500			
42	1.4800	1.4500			
43	1.5000	1.5000			

TABLE 11. Case 2:

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

Region 1

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
 THE CHARACTERISTIC PRESSURE(Psi) IS= 85.0000

Region 2

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
 THE CHARACTERISTIC PRESSURE(Psi) IS= 85.0000

XI.5 RESULTS FOR CASE 2: SHIELLS CANYON RESERVOIR ANALYSED AS A PIECEWISE HOMOGENEOUS RESERVOIR

XI.5.1 Sealed Boundary

The prescribed boundary conditions are assigned on each boundary segment according to the code described earlier under Case 1. For the interface boundary, a code of 2 is assigned implying that the potential and the normal derivative of the potential are unknown at the interface. A listing of the boundary conditions for both regions is given in Table 12. A plot of the streamlines obtained are shown in Figure 24. Table 13 gives the dimensions of the streamtubes associated with the streamlines. These dimensions are read into the steamflood program to calculate recovery as a function of time which is plotted in Figure 25.

XI.5.2 Part of the Boundary at Constant Pressure, the Remainder Sealed

The coordinates of the end points of the boundary segments are the same as in Table 10. Also, the coordinates and strengths of the injectors and producers are the same as given in Table 11. The left boundary of region 1 is assumed to be at the initial reservoir pressure of 85 psi. Thus, segment numbers 18 and 19 are assigned this constant pressure. The result from the streamline simulation program is plotted in Figure 26. The dimensions of the streamtubes obtained are listed under Table 14. The petrophysical

TABLE 12. Boundary Conditions for Shiells Canyon Reservoir
Analysed as a 2-Region Piecewise Homogeneous
Reservoir

Region 1			Region 2		
BOUNDARY CONDITIONS			BOUNDARY CONDITIONS		
NODE	CODE	PRESCRIBED VALUE	NODE	CODE	PRESCRIBED VALUE
1	1	0.0	1	2	0.0
2	1	0.0	2	2	0.0
3	1	0.0	3	2	0.0
4	1	0.0	4	2	0.0
5	1	0.0	5	2	0.0
6	1	0.0	6	2	0.0
7	1	0.0	7	2	0.0
8	1	0.0	8	2	0.0
9	1	0.0	9	2	0.0
10	1	0.0	10	2	0.0
11	1	0.0	11	1	0.0
12	1	0.0	12	1	0.0
13	1	0.0	13	1	0.0
14	1	0.0	14	1	0.0
15	1	0.0	15	1	0.0
16	1	0.0	16	1	0.0
17	1	0.0	17	1	0.0
18	1	0.0	18	1	0.0
19	1	0.0	19	1	0.0
20	1	0.0	20	1	0.0
21	1	0.0	21	1	0.0
22	1	0.0	22	1	0.0
23	1	0.0	23	1	0.0
24	1	0.0	24	1	0.0
25	1	0.0	25	1	0.0
26	1	0.0	26	1	0.0
27	1	0.0	27	1	0.0
28	1	0.0	28	1	0.0
29	1	0.0	29	1	0.0
30	1	0.0	30	1	0.0
31	1	0.0	31	1	0.0
32	1	0.0	32	1	0.0
33	1	0.0	33	1	0.0
34	2	0.0	34	1	0.0
35	2	0.0	35	1	0.0
36	2	0.0	36	1	0.0
37	2	0.0	37	1	0.0
38	2	0.0			
39	2	0.0			
40	2	0.0			
41	2	0.0			
42	2	0.0			
43	2	0.0			

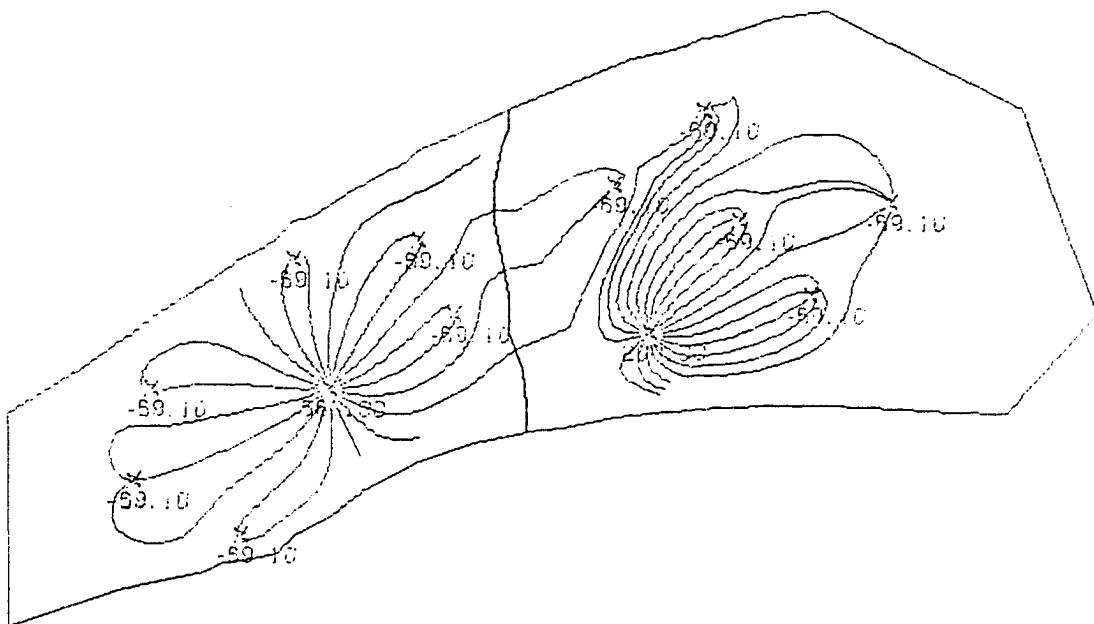


FIGURE 24. Boundary element modelling of Shiells Canyon Field as two-region piecewise homogeneous reservoir. Sealed outer boundary.

TABLE 13. Calculated Streamtube Dimensions for Shiells Canyon Reservoir Analysed as a Piecewise Homogeneous Reservoir with Sealed Boundary

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	534.80	45.44	28.00	0.0	0.0	0.0
1	2	1	425.00	23.89	28.00	0.0	0.0	0.0
1	3	1	700.00	85.96	28.00	0.0	0.0	0.0
2	1	1	441.40	27.95	28.00	0.0	0.0	0.0
2	2	1	344.70	19.24	28.00	0.0	0.0	0.0
3	1	1	334.80	18.07	28.00	0.0	0.0	0.0
3	2	1	391.40	21.94	28.00	0.0	0.0	0.0
4	1	1	269.70	13.51	28.00	0.0	0.0	0.0
4	2	1	269.70	14.84	28.00	0.0	0.0	0.0
5	1	1	359.80	15.86	28.00	0.0	0.0	0.0
5	2	1	338.10	14.26	28.00	0.0	0.0	0.0
6	1	1	294.70	12.69	28.00	0.0	0.0	0.0
6	2	1	288.10	11.95	28.00	0.0	0.0	0.0
7	1	2	376.90	51.38	28.00	403.30	42.31	28.00
7	2	2	411.10	79.31	28.00	313.10	32.17	28.00
8	1	2	602.00	58.92	28.00	319.70	38.38	28.00
8	2	2	427.10	58.81	28.00	484.80	46.20	28.00
8	3	1	581.60	14.81	10.00	0.0	0.0	0.0
8	4	1	566.40	12.23	10.00	0.0	0.0	0.0
8	5	1	594.70	13.71	10.00	0.0	0.0	0.0
9	1	1	266.40	6.20	10.00	0.0	0.0	0.0
9	2	1	291.40	6.74	10.00	0.0	0.0	0.0
9	3	1	378.30	9.01	10.00	0.0	0.0	0.0
9	4	1	288.10	6.44	10.00	0.0	0.0	0.0
10	1	1	463.10	10.12	10.00	0.0	0.0	0.0
10	2	1	369.70	8.20	10.00	0.0	0.0	0.0
10	3	1	331.60	7.52	10.00	0.0	0.0	0.0
10	4	1	334.80	7.39	10.00	0.0	0.0	0.0
10	5	1	403.30	11.19	10.00	0.0	0.0	0.0
11	1	1	719.70	22.68	10.00	0.0	0.0	0.0
11	2	1	913.10	36.20	10.00	0.0	0.0	0.0
11	3	1	713.10	19.91	10.00	0.0	0.0	0.0
11	4	1	538.10	15.62	10.00	0.0	0.0	0.0
11	5	1	609.80	22.74	10.00	0.0	0.0	0.0

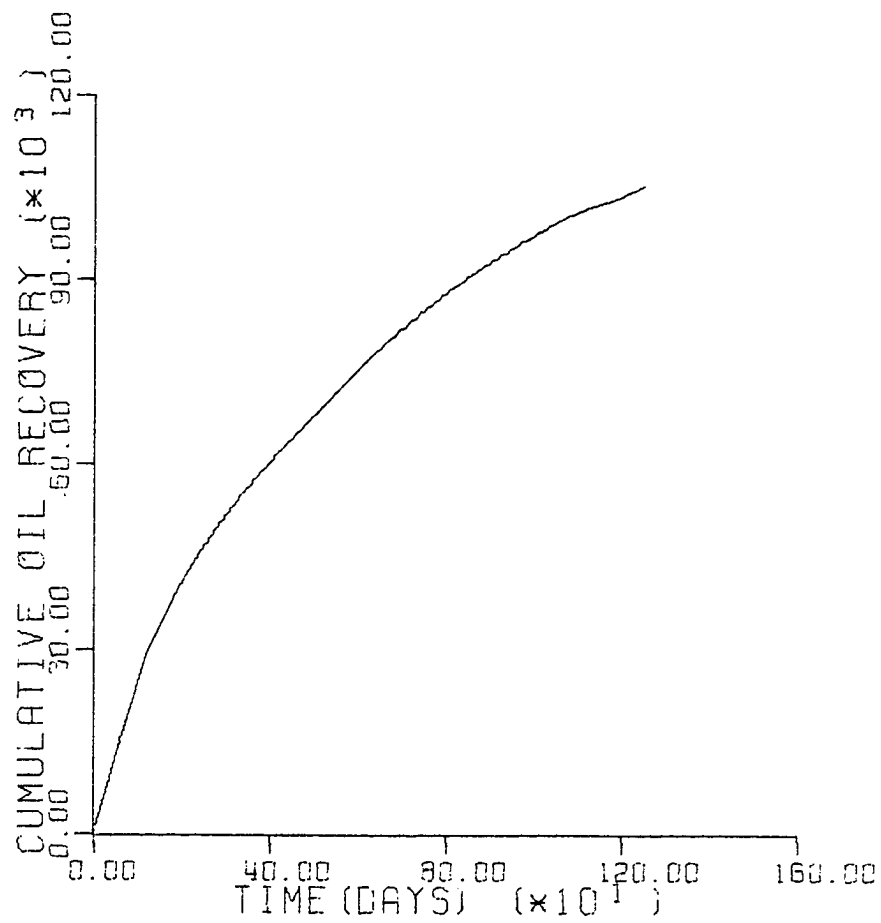


FIGURE 25. Steamflood recovery prediction of Shiells Canyon Field as two-region piecewise homogeneous reservoir. Sealed outer boundary.

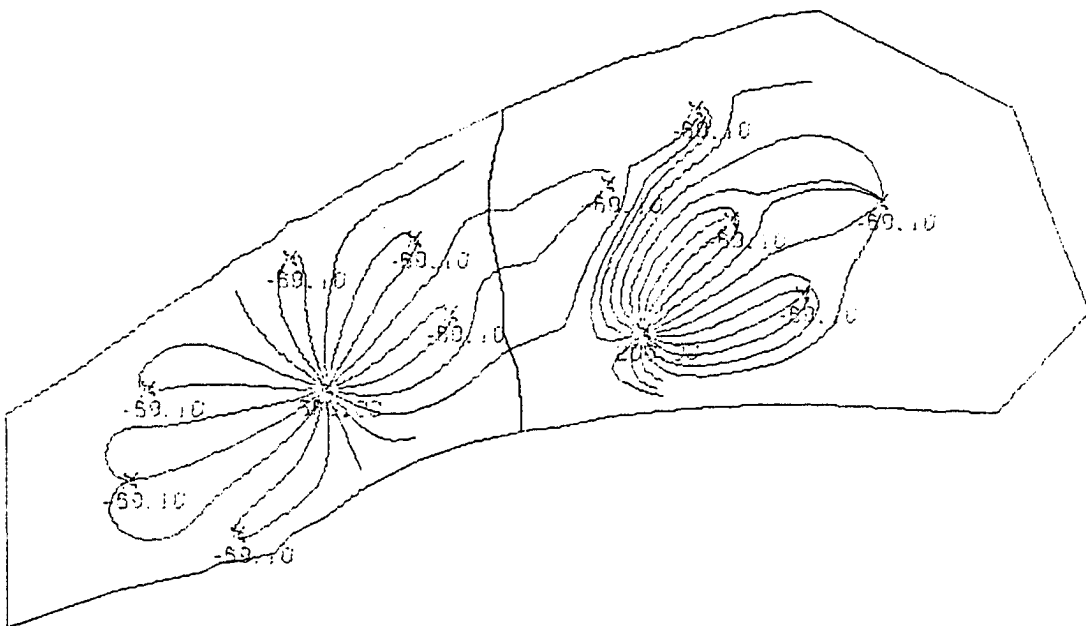


FIGURE 26. Boundary element modelling of Shiells Canyon Field as two-region piecewise homogeneous reservoir. Part of the boundary at constant pressure, part sealed.

TABLE 14. Calculated Streamtube Dimensions for Shiells Canyon Field Analysed as a Piecewise Homogeneous Reservoir. Part of the Boundary at Constant Pressure, the Remainder Sealed

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	553.30	49.91	28.00	0.0	0.0	0.0
1	2	1	425.00	23.40	28.00	0.0	0.0	0.0
1	3	1	656.60	66.11	28.00	0.0	0.0	0.0
2	1	1	444.70	27.95	28.00	0.0	0.0	0.0
2	2	1	344.70	19.06	28.00	0.0	0.0	0.0
3	1	1	338.10	17.95	28.00	0.0	0.0	0.0
3	2	1	381.60	21.11	28.00	0.0	0.0	0.0
4	1	1	269.70	13.31	28.00	0.0	0.0	0.0
4	2	1	269.70	14.75	28.00	0.0	0.0	0.0
5	1	1	366.40	15.86	28.00	0.0	0.0	0.0
5	2	1	331.60	13.79	28.00	0.0	0.0	0.0
6	1	1	291.40	12.62	28.00	0.0	0.0	0.0
6	2	1	294.70	12.12	28.00	0.0	0.0	0.0
7	1	2	451.30	60.06	28.00	216.40	30.08	28.00
7	2	2	478.40	58.33	28.00	248.00	31.38	28.00
8	1	2	376.40	61.82	28.00	541.40	43.88	28.00
8	2	1	591.40	13.37	10.00	0.0	0.0	0.0
8	3	1	619.70	14.81	10.00	0.0	0.0	0.0
9	1	1	266.40	6.40	10.00	0.0	0.0	0.0
9	2	1	288.10	6.87	10.00	0.0	0.0	0.0
9	3	1	363.10	8.27	10.00	0.0	0.0	0.0
9	4	1	288.10	6.64	10.00	0.0	0.0	0.0
10	1	1	478.30	11.11	10.00	0.0	0.0	0.0
10	2	1	384.80	8.51	10.00	0.0	0.0	0.0
10	3	1	334.80	7.71	10.00	0.0	0.0	0.0
10	4	1	331.60	7.47	10.00	0.0	0.0	0.0
10	5	1	384.80	9.71	10.00	0.0	0.0	0.0
11	1	1	703.30	21.65	10.00	0.0	0.0	0.0
11	2	1	834.80	27.73	10.00	0.0	0.0	0.0
11	3	1	719.70	19.85	10.00	0.0	0.0	0.0
11	4	1	538.10	16.01	10.00	0.0	0.0	0.0
11	5	1	591.40	20.86	10.00	0.0	0.0	0.0

data input to the steamflood program are the same as those given under Table 8. The calculated recovery of oil as a function of time is presented in Figure 27.

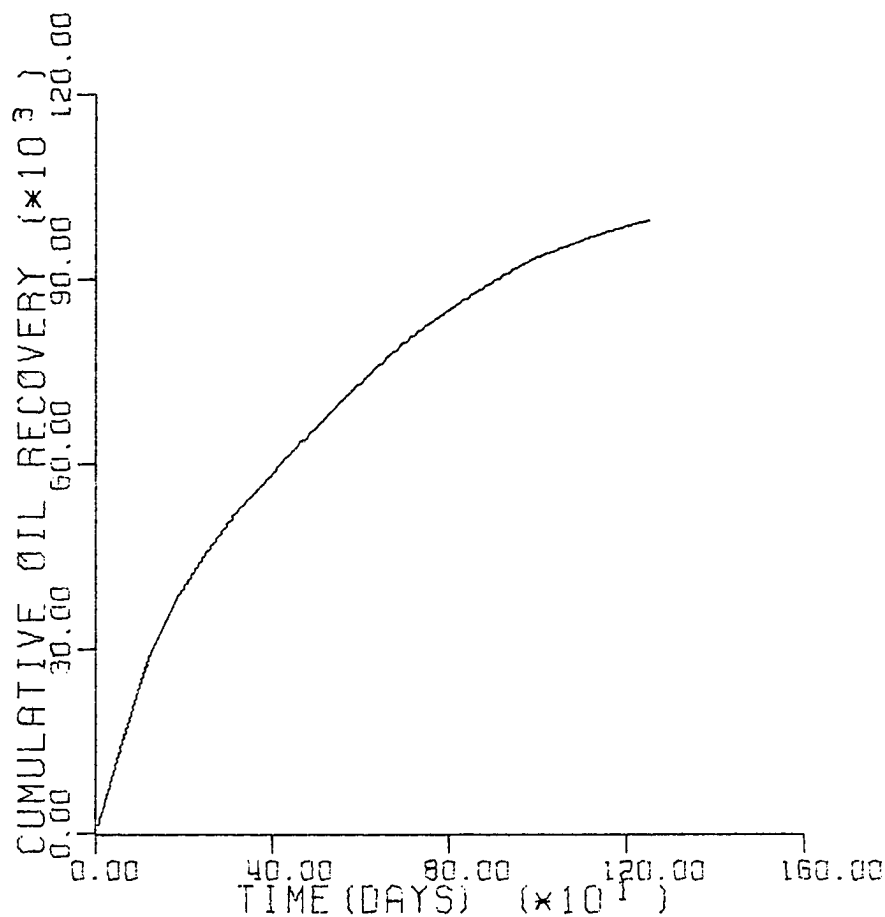


FIGURE 27. Steamflood recovery prediction for Shiells Canyon Field as a two-region piecewise homogeneous reservoir. Part of the boundary at constant pressure, part sealed.

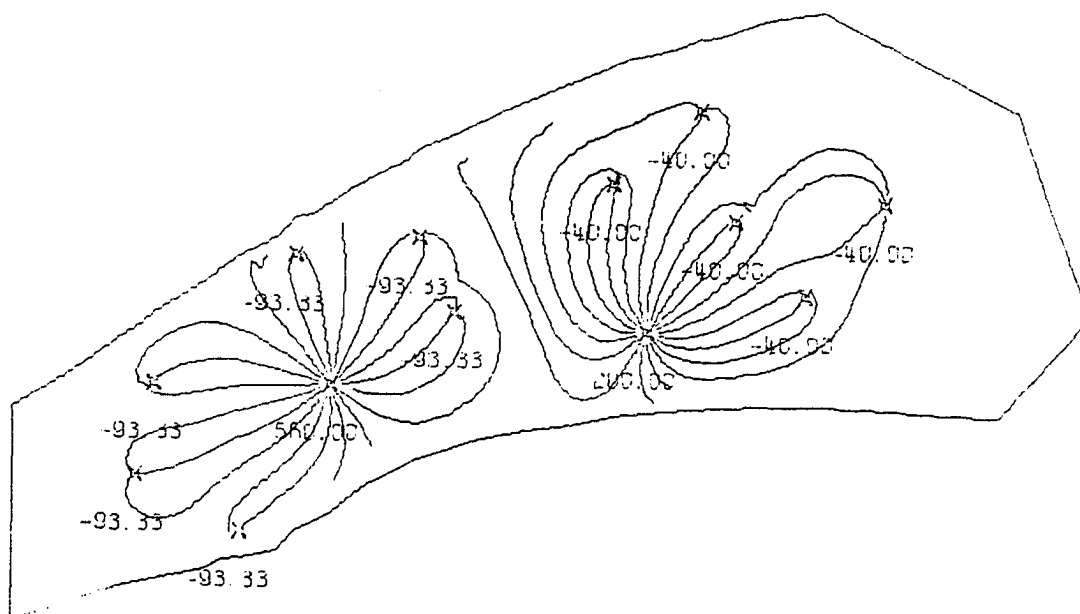


FIGURE 23. Boundary element streamline modelling of Shiells Canyon Field as a homogeneous reservoir with sealed outer boundary. Unequal production rates.

XI.6 DISCUSSION OF RESULTS

The pattern of streamlines obtained by Konopnicki et al, is shown in Figure 28. This pattern differs considerably from those obtained in this study using the Boundary Element Method for a homogeneous reservoir (Figures 18 and 20). However, a direct comparison between the two studies is inappropriate for the following reasons:

a. The production rates assigned to each producer by Konopnicki et al. in their streamline model is unknown. In this study, equal rates, determined as the sum of all the injection rates divided by the number of producers were assigned to each producer. This is important because the rates assigned to the producers significantly affect the pattern of streamlines developed. This effect can be observed in Figure 28 where the producers surrounding each injector have been assigned rates equal to the injection rate divided by the number of producers around the injector.

b. There is one fewer producer in this study than that published by Konopnicki et al.

c. Even though Konopnicki et al. state that the boundary of the reservoir is assumed to be fluid sealing all around, it is evident from Figure 29 that the boundary on

the west was certainly not a no-flow boundary. In this study, the west boundary is first assumed to be fluid sealing and then assumed to be at constant pressure (initial pressure of 85 psi). Considering the level of accuracy obtained for the test case (Section X.1), this author feels that the streamlines obtained in this study are representative of the true streamlines under the assumed boundary conditions.

Comparing the results of the sealed outer boundary and that where the west boundary is kept at constant pressure, it was observed that the values of potentials (ϕ) and the normal gradients of potentials $\left(\frac{\partial \phi}{\partial n}\right)$ obtained on the boundary "nodes" are markedly different for the two cases. However, given such differences, the streamline patterns look identical to each other. Small differences are observed in the widths of the streamtubes generated. To investigate further, more of the boundary segments were changed to the constant pressure conditions. It was observed that the streamlines deviated more from the sealed boundary case as more of the boundary segments were changed to constant boundary segments. Further testing led this author to conclude that changes in the boundary conditions affected the pattern of the streamlines generated. The magnitude of the effect depended on (a) the proportion of the entire boundary that had been changed, (b) the strengths of the sources and sinks present, (c) the proximity to any boundary.

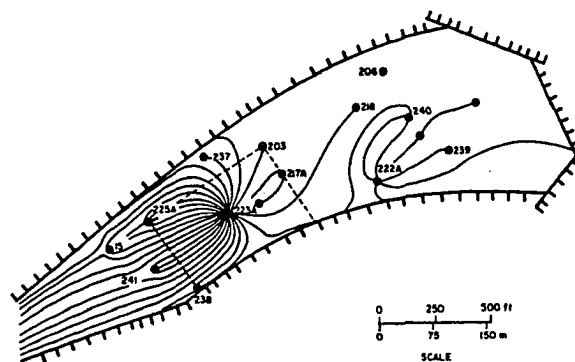


FIGURE 29. Streamlines for Shiells Canyon Steamflood
(After Konopnicki, et al.⁵⁰)

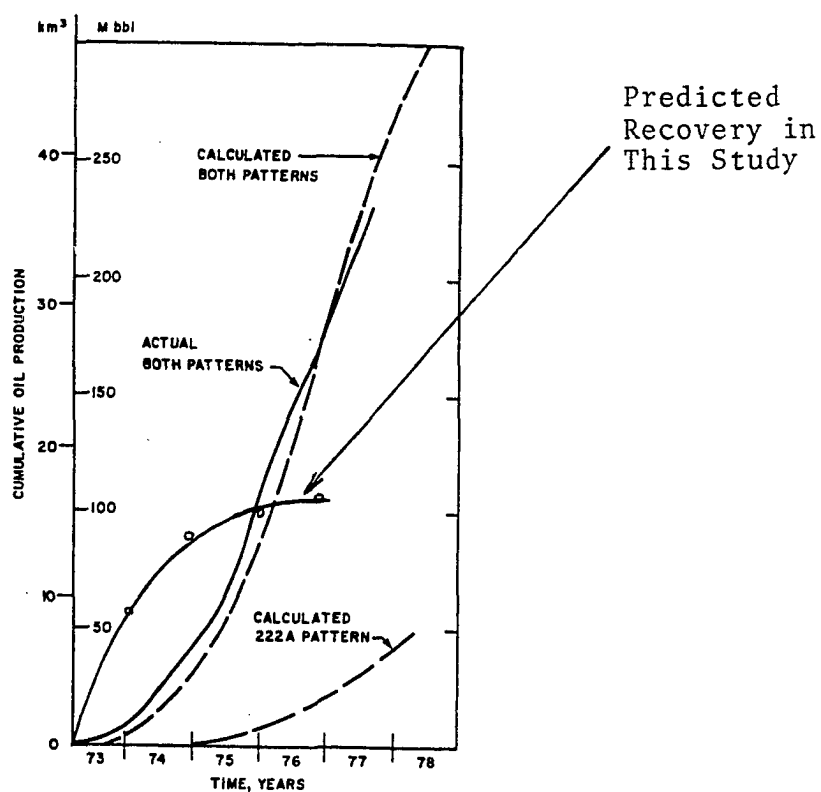


FIGURE 30. Comparison with Field Performance⁽⁵⁰⁾

It was observed that the reason for the streamlines looking nearly identical to each other for the different boundary conditions was because the effects of the sources and sinks was so large in comparison to the boundary effects that the boundary effects were masked. Only near boundaries that are far from sources and sinks were the boundary effects felt.

The results from the analysis of the reservoir as a 2-region piecewise homogeneous reservoir showed the same similarity in streamline patterns for the two kinds of boundary conditions (Figures 24 and 26) as observed earlier for the same reasons. Figure 31 gives a comparison of the oil recovery versus time for the two systems under the two kinds of boundary conditions. It can be observed that the recovery from the single homogeneous reservoir was higher than the recovery from the 2-region piecewise homogeneous reservoir for the two kinds of boundary conditions applied. This agreed with expectations since one of the regions of the 2-region reservoir has a permeability that is half the homogeneous reservoir. For the homogeneous reservoir, the summulative oil recovered was higher when part of the boundary was at constant pressure, the remainder sealed, than that when all the boundary was sealed. The reverse was the case in the 2-region piecewise homogeneous reservoir where the recovery from the sealed boundary condition was higher.

The cause of this is the loss of one streamline (Figure 26) in the case when part of the boundary was at constant pressure. Figure 30 shows a comparison between the predicted performance from this study and the actual field performance. The figure shows a high rate of recovery at early times in this study. This is consistent with the fact that the steam front velocity is relatively high at early times and declines exponentially with time. The recovery prediction by this study does not match the field performance. However, this was expected for several reasons: (a) The actual injection rates were neither steady nor continuous (Figure 32) whereas a constant injection rate was assumed for this study. (b) The prediction model used in this study assumes that the oil saturation reduces to irreducible in the steam swept area. (c) The production rates used to generate the streamlines for this study was assumed. Since the streamlines are very sensitive to the injection and production rates, this could introduce substantial errors.

CUMMULATIVE RECOVERY(BBLS*10E5) VS TIME(DAYS)

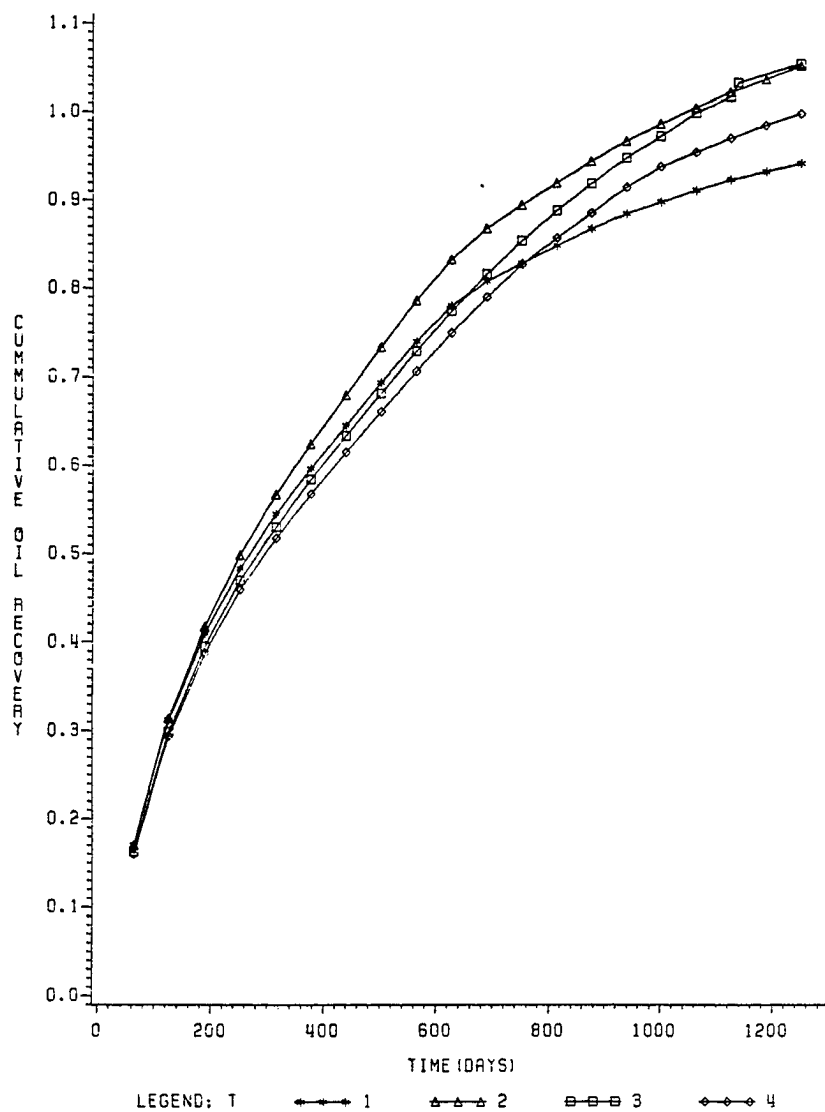


FIGURE 31. Comparison of results

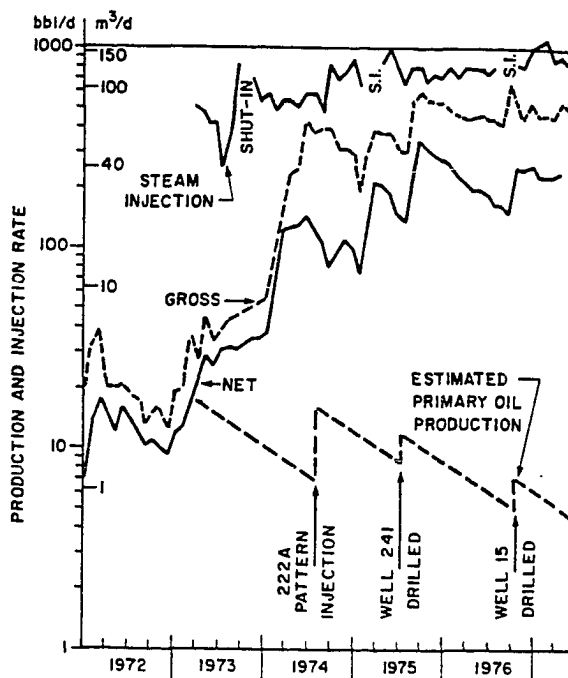


FIGURE 32. Field injection history (50)

XI.7 COMPUTATIONAL EFFICIENCY OF THE BOUNDARY ELEMENT METHOD

In practical terms, the most important question that any numerical method must answer is: "How does it compare with the commonly used methods with regard to accuracy and computational effort?" A comparison with the "method of images" is not appropriate here because the image method is only applicable to homogeneous reservoirs. Furthermore, it is not a numerical method. However, it can be said that the Boundary Element Method is superior

to the image method in terms of accuracy. In fact, it is claimed to be superior to both the finite difference and finite element methods in terms of accuracy and finer resolution in the interior.^{12,54,56} Part of the reason for this is because the error of discretization is confined to points on and near boundaries since only the boundary is discretized as opposed to discretizing the entire domain. Brebbia¹² reports that finite element results are usually accurate for the original variables under consideration (potentials) in this case, but when these variables are differentiated (to obtain fluxes), the results are much less accurate and are usually discontinuous between elements.

Several authors^{21,54,55} have reported on the comparative computational efforts associated with the solution of the coefficient matrix arising in the Boundary Element Method. For a homogeneous medium, the coefficient matrix arising from the Boundary Element Method (Equation VII.0.1) is generally non-symmetric, fully populated, non-singular, well conditioned but not diagonally dominant even though the largest terms are the diagonal terms. Hess²¹ compared the computing effort involved in a direct Gauss elimination procedure to that of the Gauss-Siedel iterative method for both the interior and exterior problem. As reported by Hess, the computing effort involved in a direct elimination is proportional to N^3 , where N is

the number of unknowns. The computing effort for an iterative solution is proportional to IFN^2 , where I is the number of iterations required for convergence, and F is the number of solutions being obtained. Thus, if I was independent of N , the iterative method would require less computing effort than the direct method for sufficiently large N . He further reported that for the exterior flow problem, the Gauss-Siedel procedure converged in a number of iterations that was independent of N but dependent on the body shape. Thus, he found that the iterative procedure required less effort if N was as large as 100-300 depending on F , but had the considerable disadvantage that the required computing time was less predictable than was the case when a direct elimination was used. For the interior flow problem, I increases linearly with N , so that the iterative procedure is usually not competitive with direct elimination.

Bettes⁵⁴ compared the computational work involved in the Boundary Integral Element Method with that of the finite element method. He found that for a square $n \times n$ two-dimensional problem, the finite element method required fewer storage locations for all meshes smaller than $n = 11$, and fewer arithmetic operations for all meshes smaller than $n = 35$. For three dimensional problems (n by n by n cube), the finite element method required fewer storage locations if n is less than 30, and fewer arithmetic operations if n is less than 135. Bettes made the conclusion

that problems have to be very large before the Boundary Integral Element method is computationally cheaper than the finite element method.

Mukherjee and Morjaria⁵⁵ have also compared the accuracy and computational efficiency of the boundary element and finite element methods for problems of time-dependent inelastic torsion of prismatic shafts. After making the comparisons for solid shafts with circular, square, elliptical, and triangular cross sections, they concluded that (a) the CPU times on an IBM 370/168 are of the same order for both BEM and FEM methods but with the BEM program generally running somewhat faster than the FEM program with the same internal mesh; (b) the discretization and input data preparation is much easier for the BEM than for the FEM.

CHAPTER XII

CONCLUSIONS AND RECOMMENDATIONS

XII.1 CONCLUSIONS

1. A streamline simulation model has been developed that is applicable to homogeneous as well as piecewise homogeneous porous media having arbitrarily shaped boundaries and under different kinds of boundary conditions. Such a general capability was not previously possible.

2. The model has successfully been used to predict steamflood recovery in both homogeneous and piecewise homogeneous porous media assuming streamtubes that are thermally isolated from each other.

3. The model can be used to predict oil recovery from any other form of secondary or tertiary recovery.

4. Interpretation of the results from this model should be done with the limitations in mind. It is recommended that the model be used primarily as a diagnostic tool prior to full scale simulation.

XII.2 LIMITATIONS OF THE MODEL

1. Streamlines are invariant with time.

2. Assumption of single phase flow is incorrect.
3. Assumption of zero heat transfer between streamtubes does not represent the physical condition.
4. The boundary element model can only be applied to systems described by linear differential equations.
5. Any region less than half an element size away from the boundary gives wrong results.
6. Cannot be used on domains with rapidly changing heterogeneities.
7. Cannot exploit existing sparse-matrix techniques.

XII.3 SUGGESTIONS FOR FURTHER STUDY

There are several areas in this research work that can be extended, improved upon. Some of the areas where further work is needed are:

1. The application of the Boundary Element Method to other equations such as parabolic equations.
2. Application to anisotropic porous media. This can very easily be done by making the appropriate transformation.
3. The search for a suitable analytical solution to the equation describing the velocity of the steam front in the second region without ignoring the heat losses in the first region.
4. The extension of the Boundary Element Method to handle reservoirs with part of their boundaries at infinity.

NOMENCLATURE

(x,y)	= coordinate axis
(x_i,y_i)	= coordinates of interior points
(x_b,y_b)	= coordinates of boundary points
(x_j,y_j)	= coordinates of sources and sinks
ϕ	= porosity
k	= permeability
K	= thermal conductivity
D	= domain
S	= boundary
Φ	= potential
S_Φ	= boundary where potential is specified
S_n	= boundary where potential gradient $\frac{\partial \Phi}{\partial n}$ is specified
$S_{\Phi,n}$	= boundary where mixed boundary conditions are specified
ρ	= density
u	= volume flux
$\nabla \cdot$	= divergence operator
t	= time
∇	= gradient operator
μ	= viscosity

N	= number of sources and sinks
δ	= Dirac delta function
q_j	= volumetric rate of flow of the j th source or sink
ϕ_n	= potential gradient $\frac{\partial \phi}{\partial n}$
L	= two-dimensional difference operator, $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$
W	= potential distribution due to a unit charge in an infinite system
$\hat{i}$	= unit vector in the x direction
$\hat{j}$	= unit vector in the y direction
$\hat{r}$	= unit vector normal to a surface
r	= radius
r_i	= distance between an observation point (x, y) and an internal charge point (x_i, y_i) $= \sqrt{(x - x_i)^2 + (y - y_i)^2}$
$r_{i,j}$	= distance between a source or sink point (x_j, y_j) and an internal charge point (x_i, y_i) $= \sqrt{(x_j - x_i)^2 + (y_j - y_i)^2}$
$r_{b,j}$	= distance between a source or sink point (x_j, y_j) and a boundary charge point (x_b, y_b) $= \sqrt{(x_j - x_b)^2 + (y_j - y_b)^2}$
$r_{b,L}$	= distance between a boundary node point (x_L, y_L) and a boundary charge point (x_b, y_b) $= \sqrt{(x_L - x_b)^2 + (y_L - y_b)^2}$
ϵ	= radius of a small circle on the surface (Figure 4)

- S_ϵ = the surface of the semicircle of radius ϵ (Figure 4)
- ϕ^* = average value of ϕ on S_ϵ
- (x_{L_k}, y_{L_k}) = coordinates of Gaussian quadrature points
- $\bar{\phi}$ = known value of potential specified as a boundary condition
- $\frac{\partial \bar{\phi}}{\partial n}$ = known value of potential gradient specified as a boundary condition
- M_1 = number of boundary segments in which potentials are specified and potential gradients are unknown
- M_2 = number of boundary segments in which potential derivatives are specified and potentials are unknowns
- $G_{b,L}$ = defined by Equation (VI.0.3)
- $H_{b,L}$ = defined by Equation (VI.0.4)
- $\tilde{G}_{b,L}$ = $G_{b,L}$ when $L \neq b$
- $G_{b,b}$ = $G_{b,L}$ when $L = b$
- $G_{i,L}$ = defined by Equation (VII.1.4)
- $H_{i,L}$ = defined by Equation (VII.1.5)
- v_x = velocity in the x-direction
- v_y = velocity in the y-direction
- S_L = the surface of the Lth line segment
- S_{L_D} = dimensionless form of S_L
- $Q_N(\xi^-, t)$ = the rate of flow of heat per unit cross sectional area through the vertical cross sectional area

at ξ^- into the elemental volume in the Nth region

$Q_N(\xi^+, t)$ = the rate of flow of heat per unit cross sectional area out of the elemental volume through the vertical cross section at ξ^+ in the Nth region

$V_N(t)$ = the velocity of the steam front in the Nth region

$H_N(\xi^-, t)$ = heat content per unit volume of reservoir rock/fluid system at the upstream face of the steam front in the Nth region

$H_N(\xi^+, t)$ = heat content per unit volume of reservoir rock/fluid system at the downstream face of the steam front in the Nth region

$Q(0, t)$ = rate of injection of heat per unit cross sectional area

$[Q_L(t)]_j$ = rate of loss of heat per unit cross sectional area to overburden and underburden from the jth region

$\rho_m(T)$ = density of m at temperature T, where m can be oil, water, or steam

L_{st} = latent heat of steam

c_m = specific heat of m, where m = oil, water, steam

$\Delta T(\xi^-)$ = change in temperature at the upstream face of the steam front

- $\Delta T(\xi^+)$ = change in temperature at the downstream face of the steam front
 ϕ_j = porosity of the rock in the j th region
 $S_m(x,t)$ = saturation distribution of m , where m = oil, water, steam
 $M_{st}(0,t)$ = mass rate of injection of steam per unit cross sectional area
 $M_w(0,t)$ = mass rate of injection of water per unit cross sectional area
 T_{st} = $(T_{st} - T_i)$ where $T_i = 0$ and T_{st} is the temperature of steam
 $\quad = \Delta T(\xi^-)$
 k = thermal conductivity of overburden and underburden
 $\xi_j(t)$ = the location of the steam front in the j th region at time t
 T_{ξ^+} = the temperature at the downstream face of the steam front
 T_{ξ^-} = the temperature at the upstream face of the steam front
 $q_{st}(0,t)$ = rate of injection of steam per unit cross sectional area
 k_j = permeability of the j th region
 h_j = thickness of the j th region
 $\bar{S}_m(t)$ = average saturation of phase m , where m = oil, water, steam

- c_{rN} = specific heat of reservoir rock of the Nth region
 $[Q_L(t)]_{\Delta A}$ = rate of heat loss per unit cross section area
 from an elemental area of size ΔA
 d = thermal diffusivity of the overburden or under-
 burden ($k/\rho_{cb}c_{cb}$)
 P_j = number of elemental areas (or number of time
 steps in the jth region)
 t = time
 τ = variable of time
 Δt = time step
 W_j = width of the jth region of the porous medium
 t_D = dimensionless time
 ξ_D = dimensional distance
 $v_D = \frac{d\xi_D}{dt_D}$ = dimensionless velocity
 L_j = length of the jth region of the porous medium
 x = fraction of injected steam that is water
 λ_j = ratio of heat capacity of cap and base rock to
 that of the jth region of the reservoir
 $\check{z}_j(t)$ = discontinuous distance function defined only
 for the period $t_{j-1} < t < t_j$ and not defined
 outside this interval
 $z_j(t)$ = step-wise continuous distance function defined
 for all time $t > 0$, but equal to zero outside
 the time interval $t_{j-1} < t < t_j$.
 $z_j(t) = \{H(t - t_{j-1}) - H(t - t_j)\}\xi_j(t)$

$\xi_j(t)$ = continuous distance function with time

$z(t)$ = $\sum_{j=1}^N \{H(t - t_{j-1}) - H(t - t_j)\} \xi_j(t)$

Subscripts

+

 = downstream face of the condensation front

-

 = upstream face of the condensation front

w

 = water

st

 = steam

o

 = oil

j

 = denotes the region

ΔA

 = elemental area

cb

 = cap and base rock (overburden and underburden)

D

 = dimensionless

APPENDIX A

DERIVATIONS OF $H_{b,L}$, $G_{b,L}$, $H'_{b,L}$, $G'_{b,L}$

APPENDIX A

DERIVATIONS OF $H_{b,L}$, $G_{b,L}$, $H'_{b,L}$ AND $G'_{b,L}$

$H_{b,L}$ is defined as

$$\frac{1}{2\pi} \int_{S_L} \frac{\partial}{\partial n_L} \left(\ln \frac{1}{r_{b,L}} \right) dS_L$$

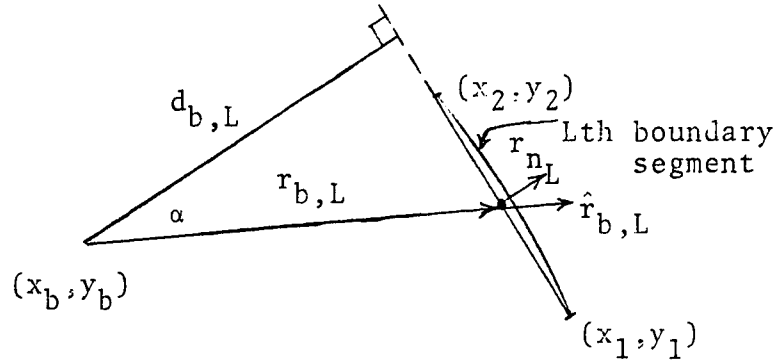
$$\frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) = \frac{\partial}{\partial \hat{r}_{b,L}} \left(\ln \frac{1}{r_{b,L}} \right) \cos \alpha$$

Since

$$\frac{\partial}{\partial n} (\ln u) = \frac{1}{u} \frac{\partial u}{\partial n}$$

then

$$\frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) = - \frac{1}{r_{b,L}} \cos \alpha = \frac{-d_{b,L}}{r_{b,L}^2} \quad (A-1)$$



Therefore,

$$H_{b,L} = - \frac{1}{2\pi} \int_{S_L} \frac{d_{b,L}}{r_{b,L}^2} dS_L \quad (A-2)$$

where $d_{b,L}$ is the perpendicular distance from a point (x_b, y_b) to the line $Ax + By + C = 0$ and is given as:

$$d_{b,L} = \frac{Ax_b + By_b + C}{\pm \sqrt{A^2 + B^2}} \quad (A-3)$$

where the appropriate sign is chosen so as to make $d_{b,L}$ always positive.

The equation of the line joining the points (x_1, y_1) and (x_2, y_2) is:

$$\left(\frac{y_2 - y_1}{x_2 - x_1} \right) x - y + \left[y_1 - x_1 \left(\frac{y_2 - y_1}{x_2 - x_1} \right) \right] = 0$$

Therefore,

$$A = \frac{y_2 - y_1}{x_2 - x_1}$$

$$B = -1$$

$$C = y_1 - x_1 \left(\frac{y_2 - y_1}{x_2 - x_1} \right)$$

$$d_{b,L} = \frac{\left(\frac{y_2 - y_1}{x_2 - x_1} \right) x_b - y_b + y_1 - x_1 \left(\frac{y_2 - y_1}{x_2 - x_1} \right)}{\pm \sqrt{\left(\frac{y_2 - y_1}{x_2 - x_1} \right)^2 + 1}} \quad (A-4)$$

In order that $d_{b,L}$ is always positive, the following convention is used:

If the numerator, $(Ax_b + By_b + C) < 0$, then use $-\sqrt{A^2 + B^2}$ as the denominator.

If the numerator, $(Ax_b + By_b + C) > 0$, then use $+\sqrt{A^2 + B^2}$ as the denominator.

$G_{b,L}$ is defined as

$$\frac{1}{2\pi} \int_{S_L} \ln \frac{1}{r_{b,L}} dS_L \quad (A-5)$$

$\{H'_X\}_{b,L}$ is defined as

$$\int_{S_L} \frac{\partial}{\partial x} \left[\frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) \right] dS_L$$

which is evaluated as follows:

$$\begin{aligned} \frac{\partial}{\partial x} \left[\frac{\partial}{\partial n} \left(\ln \frac{1}{r_{b,L}} \right) \right] &= \frac{\partial}{\partial x} \left(\frac{d_{b,L}}{r_{b,L}^2} \right) = - \left[d_{b,L} \frac{\partial}{\partial x} \left(\frac{1}{r_{b,L}^2} \right) + \frac{1}{r_{b,L}^2} \frac{\partial}{\partial x} (d_{b,L}) \right] \\ &= \left[-d_{b,L} \frac{2(x_b - x_L)}{r_{b,L}^4} + \frac{A}{r_{b,L}^2 \pm \sqrt{A^2 + 1}} \right] \end{aligned}$$

Therefore,

$$\{H'_X\}_{b,L} = \int_{S_L} \left[\frac{-d_{b,L} \cdot 2(x_b - x_L)}{r_{b,L}^4} + \frac{A}{r_{b,L}^2 \pm \sqrt{A^2 + 1}} \right] dS_L \quad (A-6)$$

Similarly,

$$\{H'_y\}_{b,L} = \int_{S_L} \left[\frac{-d_{b,L} \cdot 2(y_b - y_L)}{r_{b,L}^4} + \frac{A}{r_{b,L}^2 \pm \sqrt{A^2 + 1}} \right] dS_L \quad (A-7)$$

$\{G'_x\}_{b,L}$ is defined as

$$\int_{S_L} \frac{\partial}{\partial x} \left(\ln \frac{1}{r_{b,L}} \right) dS_L$$

and is given as follows:

$$\begin{aligned} \frac{\partial}{\partial x} \left(\ln \frac{1}{r_{b,L}} \right) \text{ is given} \\ \frac{-(x_b - x_L)}{r_{b,L}^2}. \\ \therefore \{G'_x\}_{b,L} = - \int_{S_L} \frac{(x_b - x_L)}{r_{b,L}^2} dS_L \end{aligned} \quad (A-8)$$

Similarly,

$$\{G'_y\}_{b,L} = - \int_{S_L} \frac{y_b - y_L}{r_{b,L}^2} dS_L \quad (A-9)$$

The integrals in Equations (A-2) through (A-9) are evaluated by applying the quadrature formula given in Section IX.1 Equation (IX.1.1).

APPENDIX B

EVALUATION OF TERMS UNDER THE INTEGRAL SIGNS OF THE HEAT BALANCE

APPENDIX B

EVALUATION OF THE TERMS UNDER THE
INTEGRAL SIGNS

The average saturation within any arbitrary region from x_a to x_b is:

$$\bar{S}(t) = \frac{1}{x_b - x_a} \int_{x_a}^{x_b} S(x,t) dx \quad (b-1)$$

Since $S(x,t)$ is continuous within each region behind the steam front but discontinuous at the interfaces between the regions, the average saturation in all regions behind

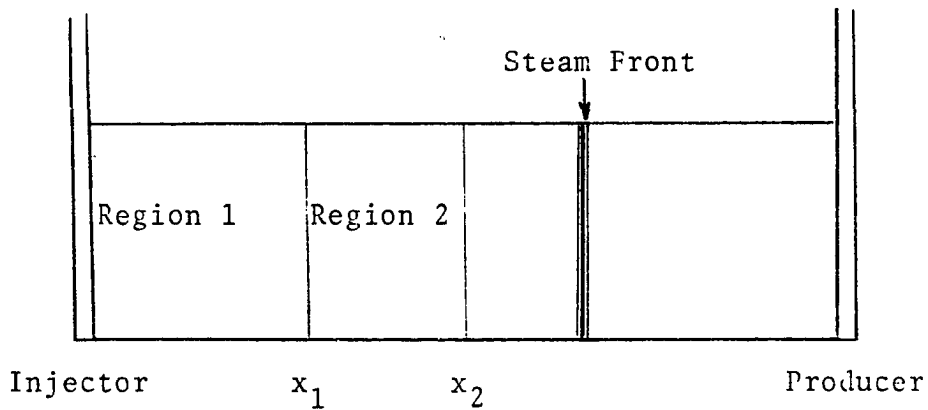


FIGURE B-1

the steam front can be obtained by applying Equation (B-1) to every region behind the front and taking the average of these averages.

Let

$$\bar{S}(t) = \frac{\sum_{k=1}^j S_k(t)}{j} \quad (B-2)$$

Rewrite B-1 as

$$\bar{S}(t)(x_b - x_a) = \int_{x_a}^{x_b} S(x,t) dx$$

and differentiate both sides to obtain

$$\bar{S}(t) \frac{\partial}{\partial t} (x_b - x_a) + (x_b - x_a) \frac{\partial}{\partial t} \bar{S}(t) = \frac{\partial}{\partial t} \int_{x_a}^{x_b} S(x,t) dx \quad (B-3)$$

Apply Leibnitz's theorem to the RHS of Equation (B-3) to obtain

$$\begin{aligned} \frac{\partial}{\partial t} \int_{x_a}^{x_b} S(x,t) dx &= \int_{x_a}^{x_b} \frac{\partial}{\partial t} S(x,t) dx + S(x_b,t) \frac{\partial x_b}{\partial t} \\ &\quad - S(x_a,t) \frac{\partial x_a}{\partial t} \end{aligned} \quad (B-4)$$

Equation B-4 is substituted into B-4 to give

$$\begin{aligned} \bar{S}(t) \frac{\partial}{\partial t} (x_b - x_a) + (x_b - x_a) \frac{\partial}{\partial t} \bar{S}(t) &= \int_{x_a}^{x_b} \frac{\partial}{\partial t} S(x,t) dx \\ &\quad + S(x_b,t) \frac{\partial x_b}{\partial t} - S(x_a,t) \frac{\partial x_a}{\partial t} \end{aligned} \quad (B-5)$$

Now if both x_b and x_a are constants, Equation B-5 reduces to:

$$(x_b - x_a) \frac{\partial}{\partial t} \bar{S}(t) = \int_{x_a}^{x_b} \frac{\partial}{\partial t} S(x,t) dx \quad (B-6)$$

For all regions behind the steam front except the region containing the steam front, the limits x_a and x_b are constants. Therefore, Equation B-6 is applied to as many such regions. For the region containing the condensation front, the limit x_a is a constant, but x_b is the location of the condensation front $\xi(t)$ which is a function of time. Therefore, Equation B-5 becomes:

$$\begin{aligned} \bar{S}(t)V(t) + [\xi(t) - x_a] \frac{\partial}{\partial t} \bar{S}(t) - S(\xi,t)V(t) \\ = \int_{x_a}^{\xi(t)} \frac{\partial}{\partial t} S(x,t) dx \end{aligned} \quad (B-7)$$

APPENDIX C

TEMPERATURE DISTRIBUTION IN THE CAP AND BASE ROCK

APPENDIX C

TEMPERATURE DISTRIBUTION IN THE CAP AND BASE ROCK

A general heat balance on any medium is (Van Poolen, Spillette):

$$\frac{\partial}{\partial t} (MT) + \nabla \cdot (\rho_f c_f T) V = \nabla \cdot \hat{k} \nabla T \quad (C-1)$$

For the cap or base rock, there are no fluids,

$$\therefore M = \rho_{cb} c_{cb}$$

$$\rho_f c_f = \rho_o c_o + \rho_w c_w$$

for an oil water system

$$V = 0 = \text{velocity of the fluids}$$

Thus, Equation (C-1) becomes:

$$\rho_{cb} c_{cb} \frac{\partial T}{\partial t} = \nabla \cdot \hat{k} \nabla T \quad (C-2)$$

For a one-dimensional medium in the z-direction

$$\frac{\partial T}{\partial t} = \frac{\hat{k}}{\rho_{cb} c_{cb}} \frac{\partial^2 T}{\partial z^2} \quad (C-3)$$

Defining $\alpha = \frac{k}{\rho_{cb} c_{cb}}$, then

$$\frac{\partial^2 T}{\partial z^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t} \quad (C-4)$$

The equation describing the temperature distribution in the overburden or underburden is:

$$\frac{\partial^2 T}{\partial z^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t} \quad (C-5)$$

Subject to:

$$T(z,0) = T_i \quad \text{when } t = 0 \quad (C-6)$$

$$T(0,t) = T_{st} \quad \text{at } z = 0, t > 0 \quad (C-7)$$

$$T(\infty,t) = T_i \quad \text{at } z = \infty, t > 0 \quad (C-8)$$

Define $\theta(z,t)$ as

$$\frac{T(z,t) - T_i}{T_{st} - T_i}$$

Thus, the differential equation becomes:

$$\frac{\partial^2 \theta}{\partial z^2} = \frac{1}{\alpha} \frac{\partial \theta}{\partial t} \quad (C-9)$$

subject to:

$$\theta(z,0) = 0 \quad (C-10)$$

$$\theta(0,t) = 1 \quad (C-11)$$

$$\theta(\infty,t) = 0 \quad (C-12)$$

The solution to Equation (C-4) with its boundary conditions (Equations C-5 to C-6) are obtainable from standard texts such as Carslaw and Jaeger.⁵⁸ It is given as:

$$\theta(z,t) = \operatorname{erfc}\left(\frac{z}{\sqrt{4\alpha t}}\right) \quad (\text{C-13})$$

Differentiating (C-13) with respect to z gives

$$\begin{aligned} \frac{d}{dz} \theta(z,t) &= \frac{d}{dz} \left[\operatorname{erfc}\left(\frac{z}{\sqrt{4\alpha t}}\right) \right] \\ &= \frac{d}{dz} \left[-\frac{2}{\sqrt{\pi}} \int_0^{\frac{z}{\sqrt{4\alpha t}}} e^{-t^2} dt \right] = -\frac{2}{\sqrt{\pi}} \left[-e^{-\left(\frac{z}{\sqrt{4\alpha t}}\right)^2} \frac{1}{\sqrt{4\alpha t}} \right] \\ \frac{d}{dz} \theta(z,t) &= -\frac{1}{\sqrt{\pi\alpha t}} e^{-\left(\frac{z}{\sqrt{4\alpha t}}\right)^2} \end{aligned} \quad (\text{C-14})$$

Since

$$\begin{aligned} \frac{dT(z,t)}{dz} &= (T_{st} - T_i) \frac{d\theta(z,t)}{dz} \\ \left. \frac{dT(z,t)}{dz} \right|_{z=0} &= -\frac{T_{st} - T_i}{\sqrt{\pi\alpha t}} \end{aligned} \quad (\text{C-15})$$

For an area that received heat initially at time τ ,

$$\left. \frac{dT(z,t)}{dz} \right|_{z=0} = -\frac{T_{st} - T_i}{\sqrt{\pi\alpha(t - \tau)}} \quad (\text{C-16})$$

APPENDIX D

LAPLACE TRANSFORM OF A SECTIONALLY CONTINUOUS FUNCTION

APPENDIX D

LAPLACE TRANSFORM OF A SECTIONALLY CONTINUOUS FUNCTION

For a function $\xi_D(\tau_D)$ which is sectionally continuous and continuously differentiable in the domain $t > 0$:

$$\xi_D(\tau_D) = \sum_{j=1}^N \xi_{Dj}(\tau_D), \quad t_{Dj-1} < \tau_D < t_{Dj}$$

where the functions $\xi_{Dj}(\tau_D)$ are continuous in the domain $t_{Dj-1} < \tau_D < t_{Dj}$. At the points t_j for $j = 1, \dots, N$, $\xi_D(\tau_D)$ has finite steps. Incorporating the Heaviside unit step function,

$$\xi_D(\tau_D) = \sum_{j=1}^N \{H(t_D - t_{Dj-1}) - H(t_D - t_{Dj})\} \xi_{Dj}(\tau_D)$$

The generalized derivative of this function is:

$$\begin{aligned} V_D(\tau_D) &= \frac{d \xi_D(\tau_D)}{d \tau_D} = \sum_{j=1}^N \{[H(t_D - t_{Dj-1}) - H(t_D - t_{Dj})]\} \\ &\quad \frac{d \xi_{Dj}(\tau_D)}{d \tau_D} + [\delta(t_D - t_{Dj-1}) - \delta(t_D - t_{Dj})] \xi_{Dj}(\tau_D) \\ &= \sum_{j=1}^N \{[H(t_D - t_{Dj-1}) - H(t_D - t_{Dj})] V_{Dj}(\tau_D) \\ &\quad + \xi_{Dj}(\tau_{Dj-1}) - \xi_{Dj}(\tau_{Dj})\} \end{aligned}$$

where H is the Heaviside unit step function

δ is the Dirac delta function.

The Laplace transform of any function $v_D(t_D)$ is defined as

$$\mathcal{L}\{v_D(t_D)\} = \int_0^{\infty} v_D(t_D) e^{-st_D} dt_D$$

that is,

$$\begin{aligned} \mathcal{L}\{V_D(\tau_D)\} = \sum_{j=1}^N & \left\{ \left[\int_{t_{Dj-1}}^{t_{Dj}} V_{Dj}(\tau_D) e^{-s\tau_D} d\tau_D \right] \right. \\ & \left. + \xi_{Dj}(\tau_{Dj-1}) e^{-s\tau_{Dj-1}} - \xi_{Dj}(\tau_{Dj}) e^{-s\tau_{Dj}} \right\} \\ & - \xi_0(0) \end{aligned}$$

Integrate by parts. Let

$$u = e^{-s\tau_D}$$

$$dv = V_{Dj}(\tau_D)$$

$$\begin{aligned} \mathcal{L}\{V_D(\tau_D)\} &= \sum_{j=1}^N \left[e^{-s\tau_D} \xi_{Dj}(\tau_D) \Big|_{t_{Dj-1}}^{t_{Dj}} + s \int_{t_{Dj-1}}^{t_{Dj}} \xi_{Dj}(\tau_D) \right. \\ &\quad \times e^{-s\tau_D} d\tau_D + \xi_{Dj}(\tau_{Dj-1}) e^{-s\tau_{Dj-1}} - \xi_{Dj}(\tau_{Dj}) e^{-s\tau_{Dj}} \Big] - \xi_0(0) \\ &= \sum_{j=1}^N \left[e^{-s\tau_{Dj}} \xi_{Dj}(\tau_{Dj}) - e^{-s\tau_{Dj-1}} \xi_{Dj}(\tau_{Dj-1}) + \xi_{Dj}(\tau_{Dj-1}) \right. \\ &\quad \times e^{-s\tau_{Dj-1}} - \xi_{Dj}(\tau_{Dj}) e^{-s\tau_{Dj}} + sZ(s) \Big] - \xi_0(0) \end{aligned}$$

Since

$$\sum_{j=1}^N \int_{t_{Dj-1}}^{t_{Dj}} \xi_{Dj}(\tau_D) e^{-s\tau_D} d\tau_D = Z_D(s) = \mathcal{L}\{\xi_{Dj}(\tau_D)\}$$

Thus,

$$\mathcal{L}\{V_D(\tau_D)\} = sZ_D(s) - \xi_0(0)$$

where

$$Z_D(s) = \sum_{j=1}^N \int_{t_{Dj-1}}^{t_{Dj}} \xi_{Dj}(\tau_D) e^{-s\tau_D} d\tau_D$$

APPENDIX E

COMPUTER IMPLEMENTATION OF BOUNDARY ELEMENT METHOD OF STREAMLINE SIMULATION

APPENDIX E

COMPUTER IMPLEMENTATION OF BOUNDARY ELEMENT METHOD OF STREAMLINE SIMULATION

A package of computer programs has been written in FORTRAN to perform all the calculations necessary to generate the streamlines and calculate the dimensions of the streamtubes using the Boundary Element Method. The program can handle:

1. Single homogeneous porous medium with or without sources and sinks.
2. A two-region piecewise homogeneous porous medium with or without sources and sinks.

The boundary conditions can be either:

1. sealed
2. constant pressure
3. combination of (1) and (2)

and consists of a main program that calls the following 15 subroutines.

1. INPUT
2. CONVRs
3. MATRIX
4. INTE

5. INLO
6. SOURCE
7. LOWVEL
8. ASEMBL
9. SLNPD
10. SPLIT
11. INTER
12. OUTPUT
13. BDRY
14. STRM
15. COMPAT

Subroutines INPUT, INTE, SLNPD, INTER and INLO are modifications of the subroutines INPUT, INTE, SLNPD, INTER, and INLO published in reference 12.

Subroutine INPUT reads all the input data required, such as the coordinates of the ends of the boundary segments, the coordinates and strengths of the sources and sinks, the boundary conditions, etc.

Subroutine MATRIX performs the integrations to determine the values of $H_{b,L}$ and $G_{b,L}$ which make up the elements of the [H] and [G] matrices. It does this by calling two other subroutines, INTE and INLO. Subroutine INTE calculates the values $\check{H}_{b,L}$ and $\check{G}_{n,L}$ (that is, $H_{b,L}$ and $G_{b,L}$ when $L \neq b$), while subroutine INLO calculates $H_{b,b}$ and $G_{b,b}$. $\check{H}_{b,L}$ and $\check{G}_{b,L}$ form the off-diagonal elements of the [H] and [G] matrices while $H_{b,b}$ and $G_{b,b}$ form the diagonal elements.

Both subroutines are modified from reference 12. Subroutine MATRIX also calls the subroutine SOURCE to calculate the contributions of the sources and sinks, and rearranges the [H] and [G] matrices according to the prescribed boundary conditions such that all the elements of [H] and [G] corresponding to unknown boundary conditions are put in the matrix [G]. Finally, it calculates the $[H'_x]$, $[G'_x]$, $[H'_y]$ and $[G'_y]$ matrices required to calculate the velocities in the x and y directions at interior points.

Subroutine ASEMBL is called for a piecewise homogeneous reservoir having two regions. It arranges the new [G] matrices for both regions (obtained from routine MATRIX) into a single matrix [A]. The arrangement is such that continuity and compatibility conditions at the common interface between the two regions are accounted for.

Subroutine SLNPD solves the resulting [A] matrix in double precision using a Gaussian elimination procedure that allows for row interchange in case of zero pivot elements. After solution, subroutine SPLIT divides the solution vector into the two vectors representing the boundary potentials and gradients for each region.

Subroutine INTER uses these vectors to determine the potential at internal points. It calls both subroutine INTE and subroutine INLO to supply the [H] and [G] matrices needed.

Subroutine OUTPUT is called to give a printed output of the calculated values.

Subroutine BDRY plots the boundaries of the reservoir system.

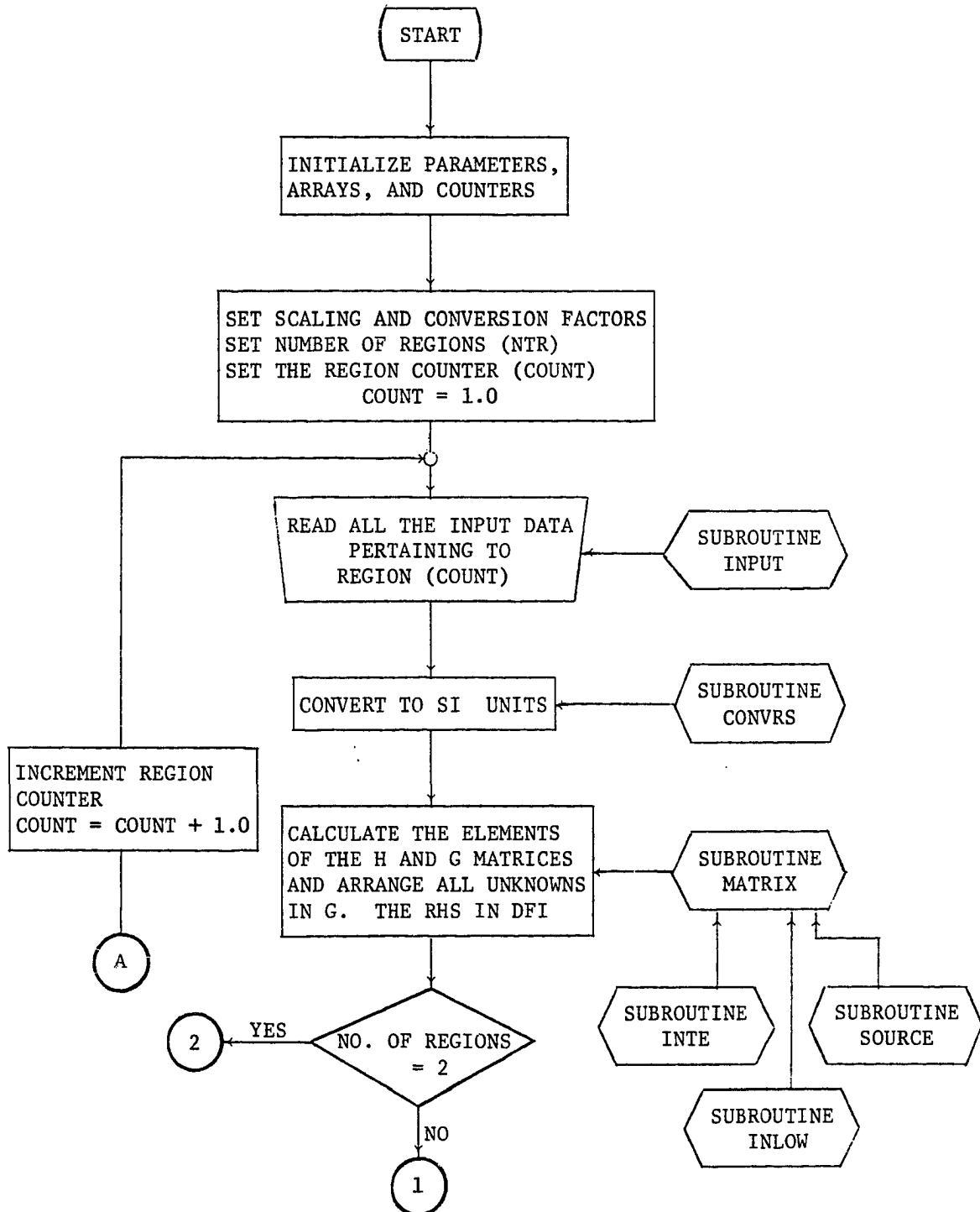
Subroutine STRM generates the streamlines and plots them. It also calculates the lengths and widths of the streamtubes. It calls the following subroutines: (a) subroutine LOWVEL to calculate the velocity near the lowest producer; (b) subroutine INTE and subroutine INLO. The steps of operations used by subroutine STRM can be summarized as follows:

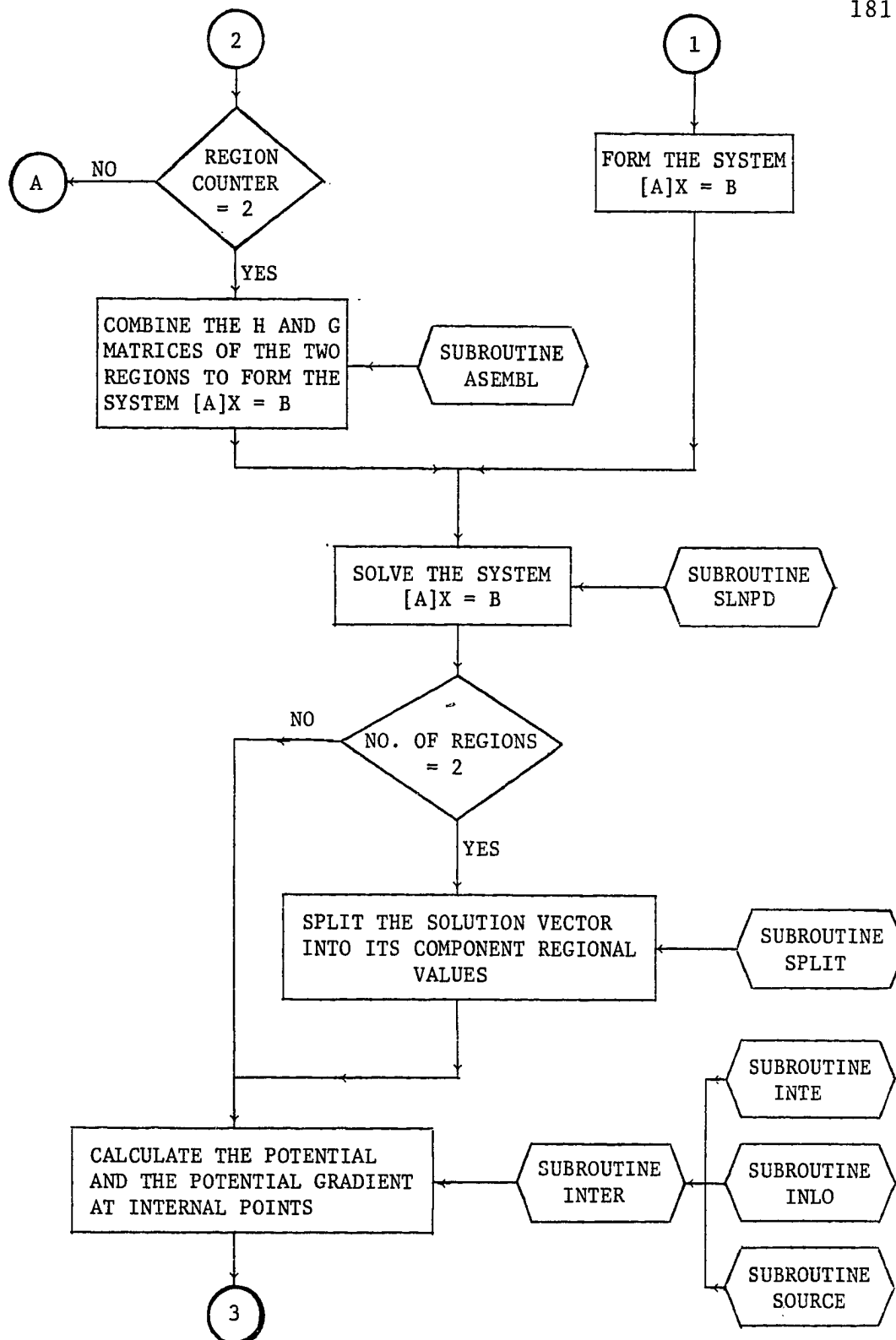
- a. Define a capture radius (say 20 times the wellbore radius) and calculate the velocity at this radius for the lowest producer.
- b. For each injector, calculate the number of streamlines emanating from it. The starting points for each streamline are spaced evenly on the wellbore radius.
- c. Calculate the velocity vectors and the next position. Check to see if the new position is less than a half element length away from any boundary. If the boundary is not an interface, stop the streamline. If it is an interface, calculate the point where the streamline is assumed to cross the interface. This point is obtained as the point of intersection of the normal from the last streamline point to the nearest boundary segment. If the point is not near any boundary, the length and width of the streamline at this point are calculated. The point is plotted and the next position calculated and step (c) repeated for all streamlines and all wells.

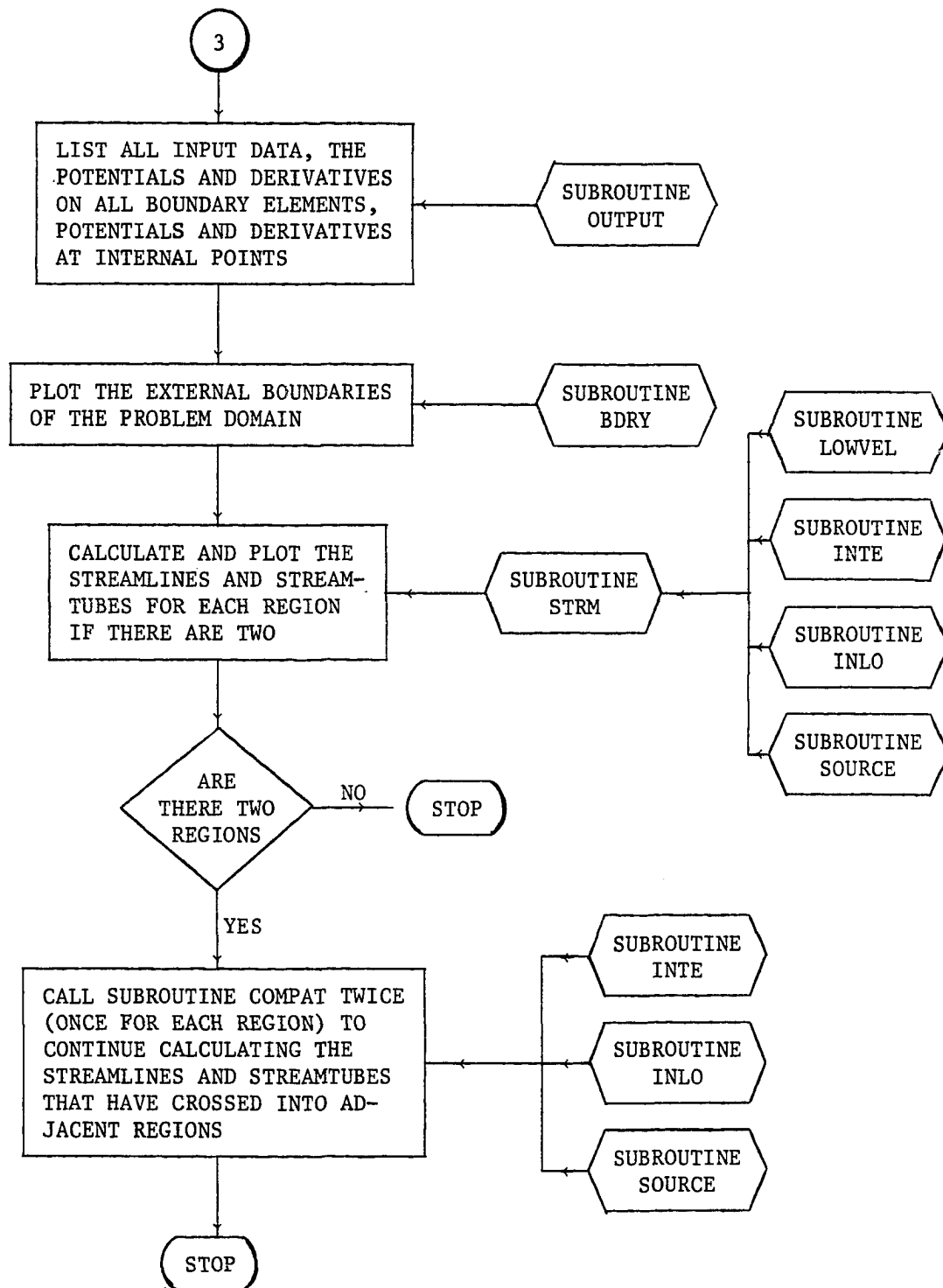
Subroutine COMPAT continues the generation of the streamlines after they cross into adjacent regions. The procedure followed by subroutine COMPAT is similar to that of subroutine STRM except that the starting point is on the interface boundary. The next streamline position away from the interface boundary was calculated using the procedure outlined in Section VIII.4.

In what follows, a listing of the flow charts for all the important subroutines is given.

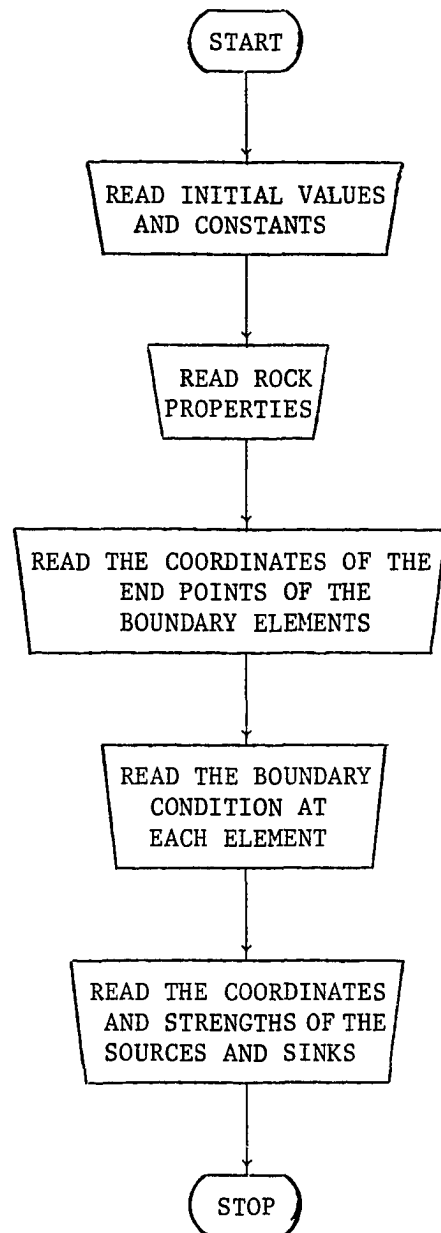
FLOW CHART FOR MAIN PROGRAM OF ELEMENT
METHOD OF STREAMLINE SIMULATION



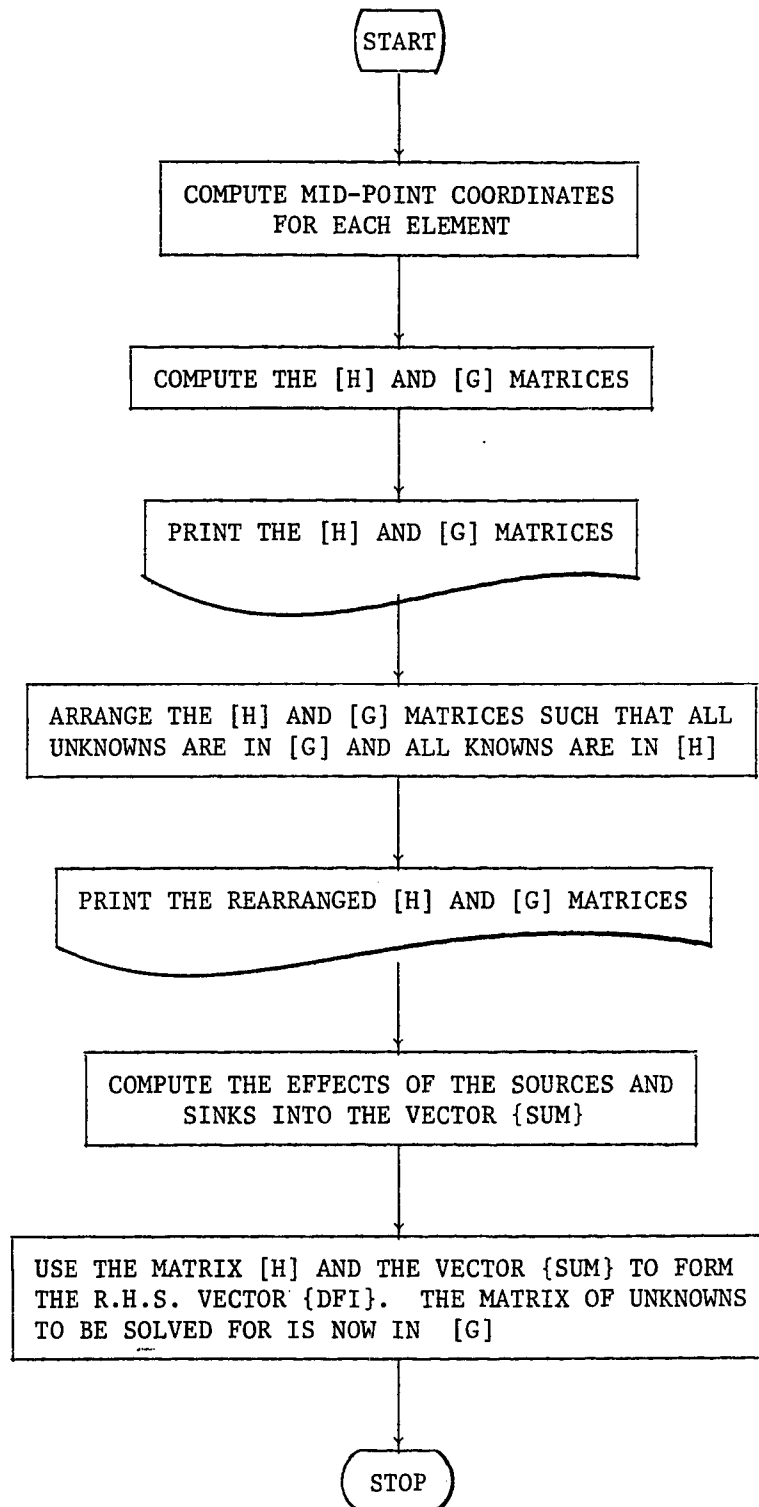




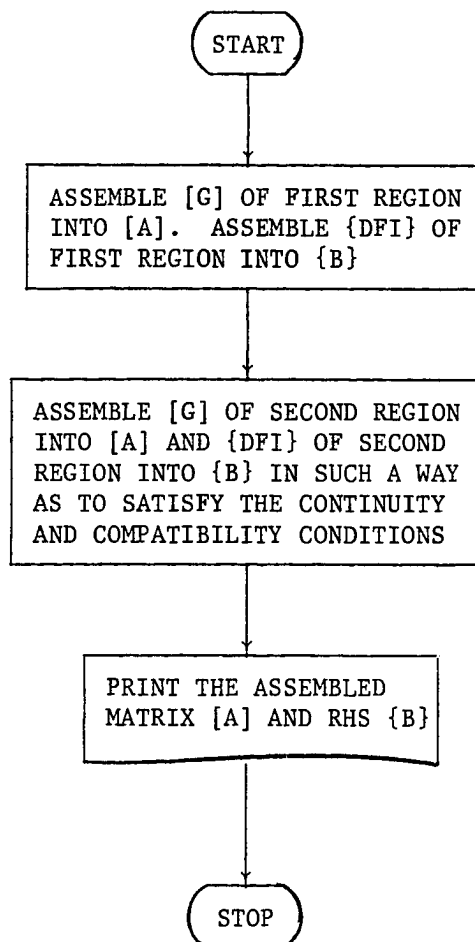
FLOW CHART FOR SUBROUTINE INPUT



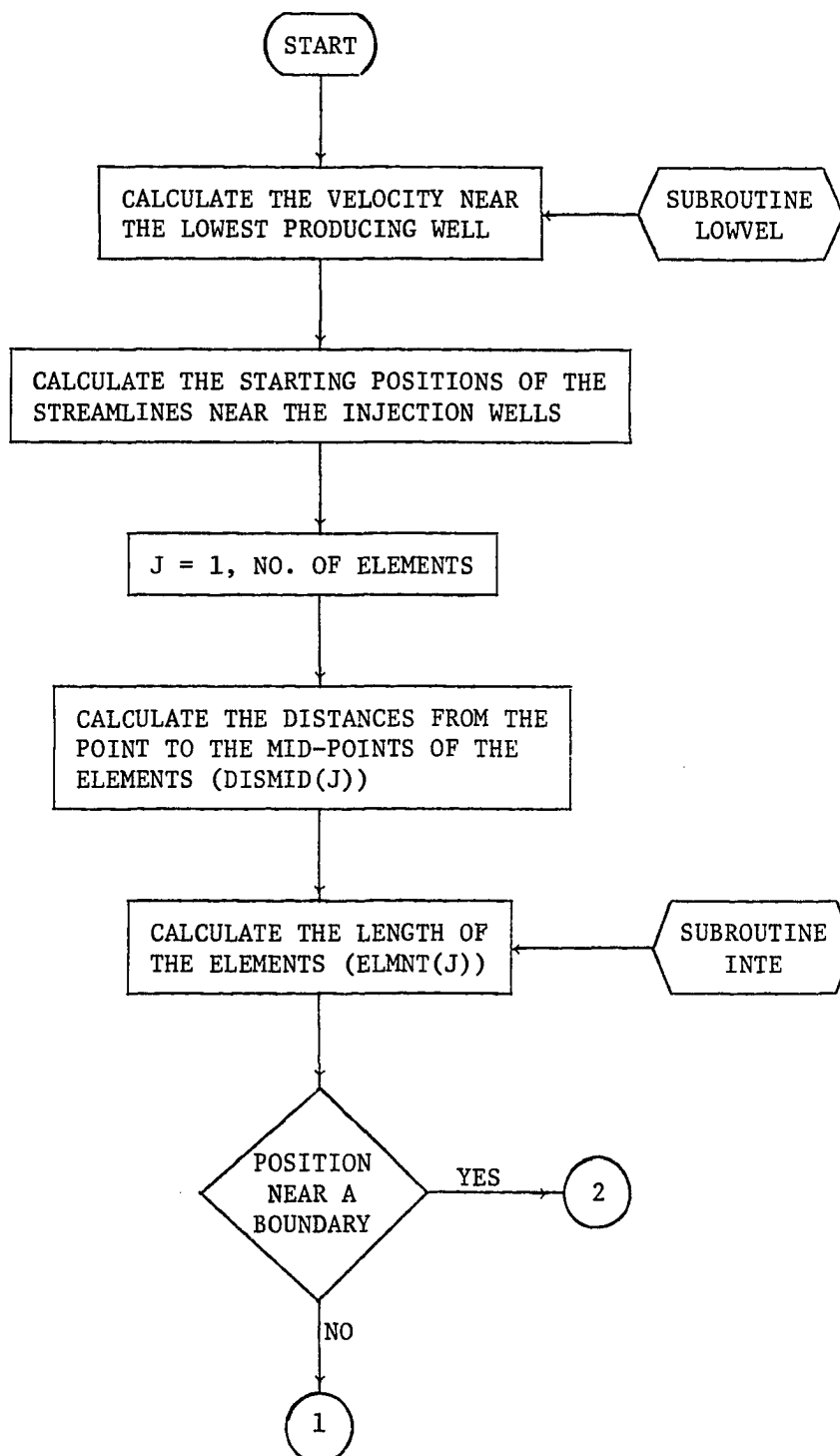
FLOW CHART FOR SUBROUTINE MATRIX

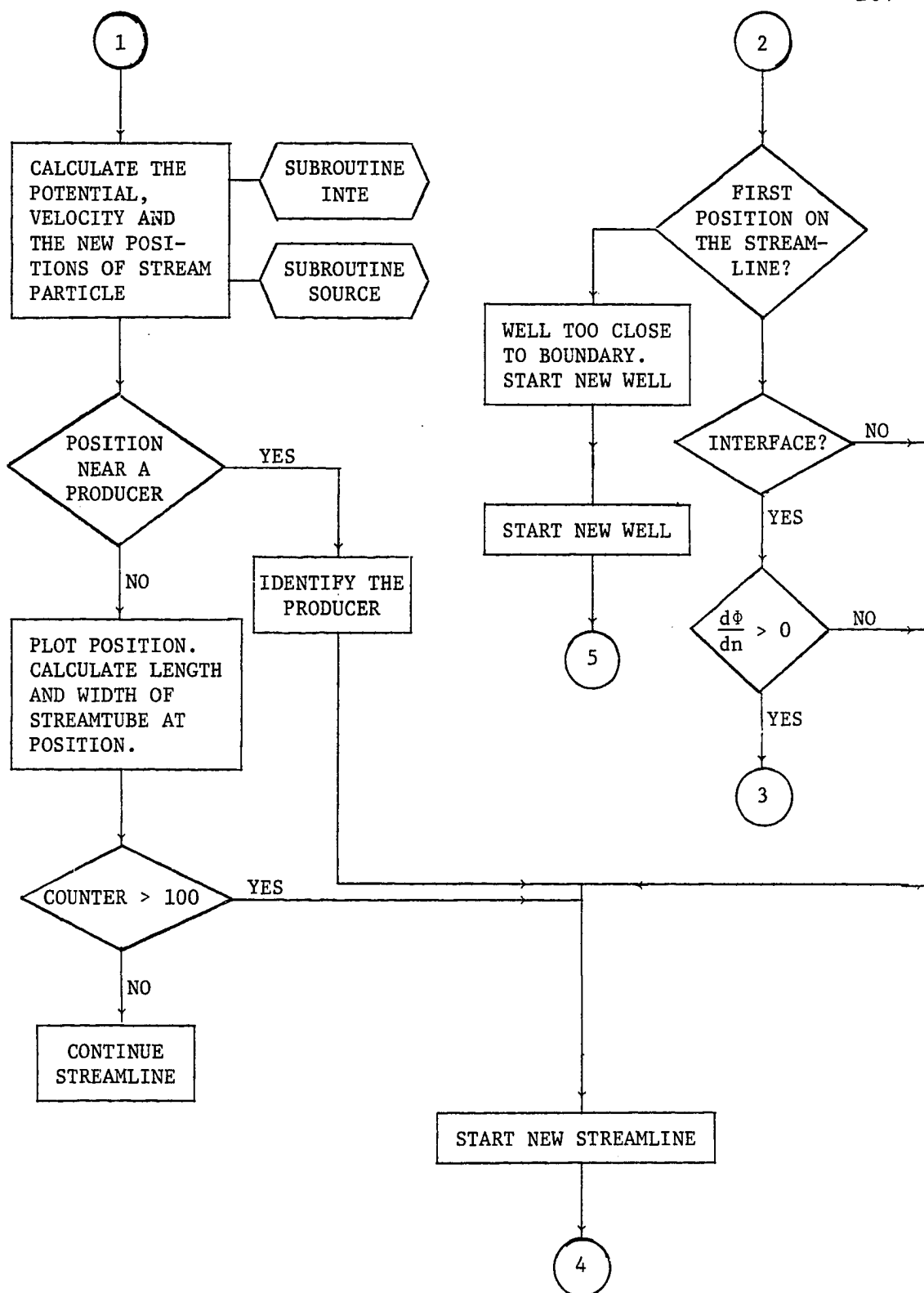


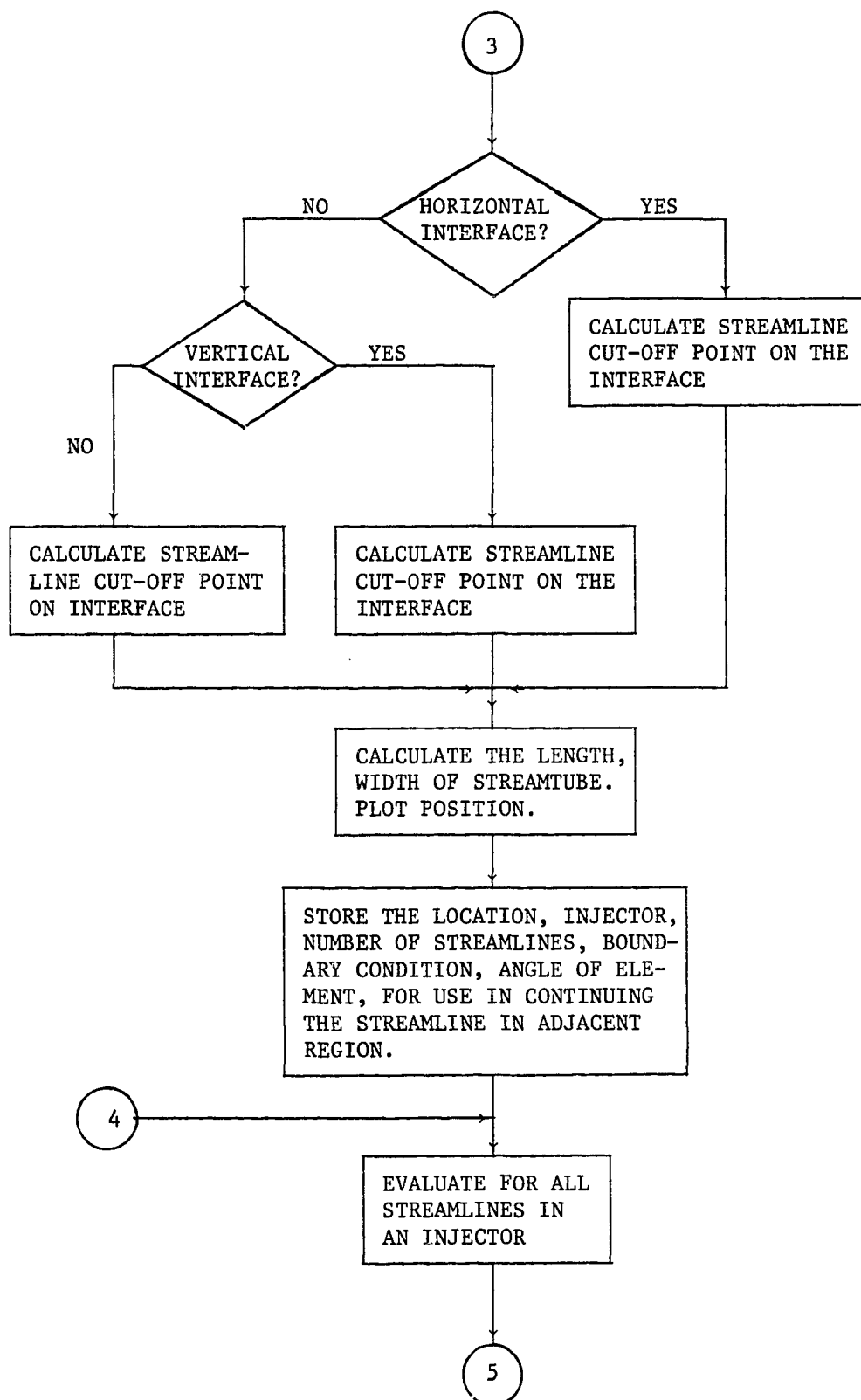
FLOW CHART FOR SUBROUTINE ASEMBL

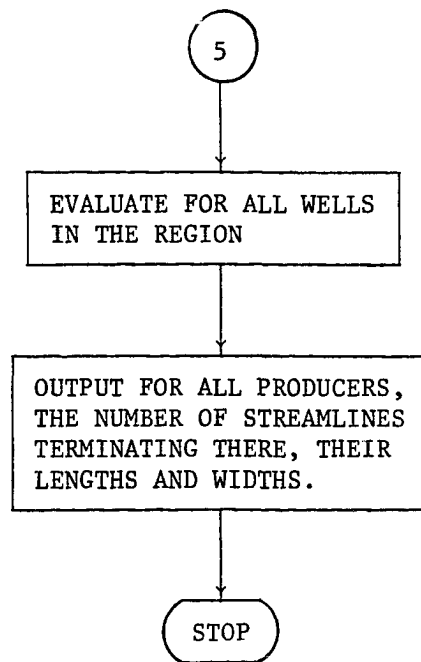


FLOW CHART FOR SUBROUTINE STRM

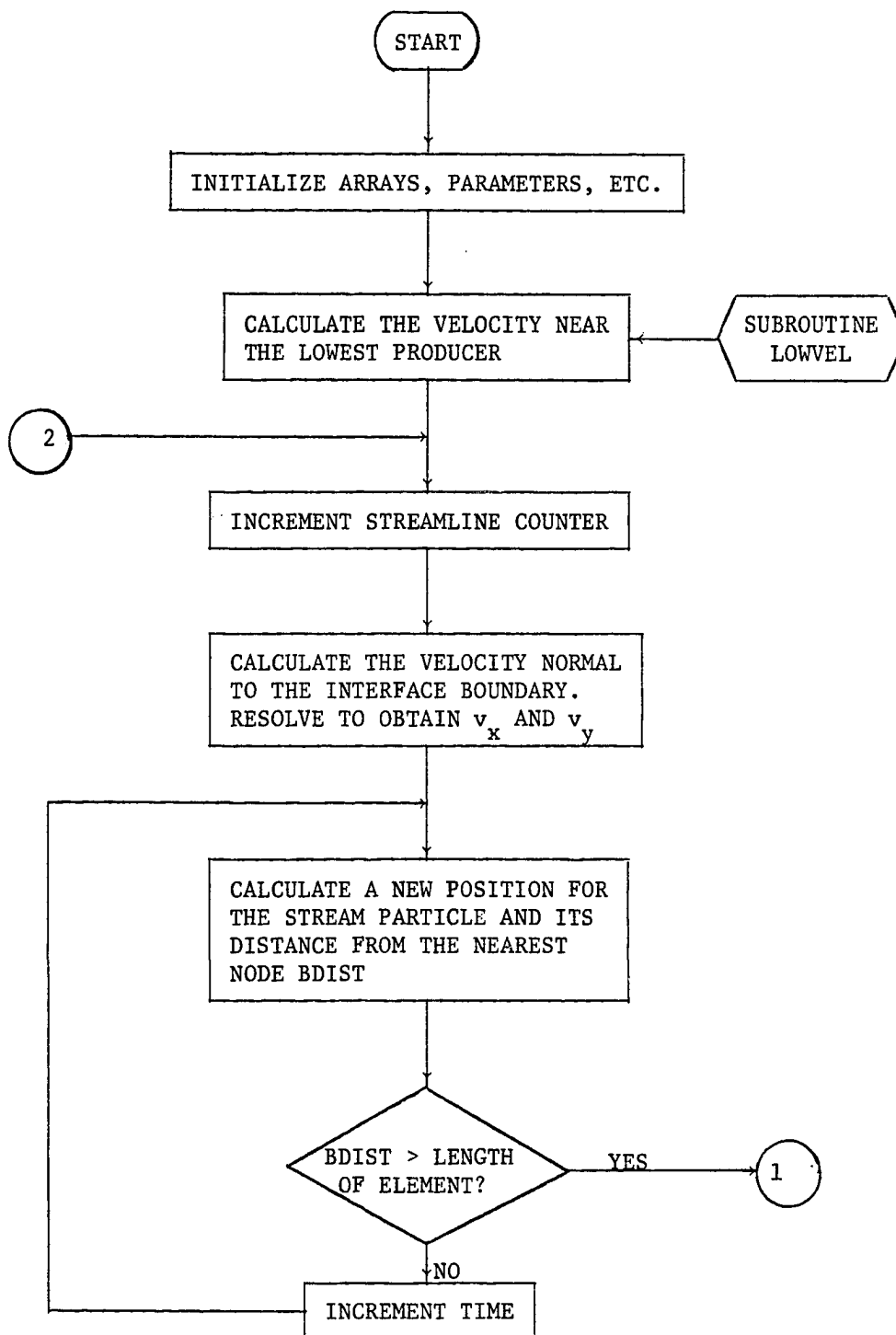


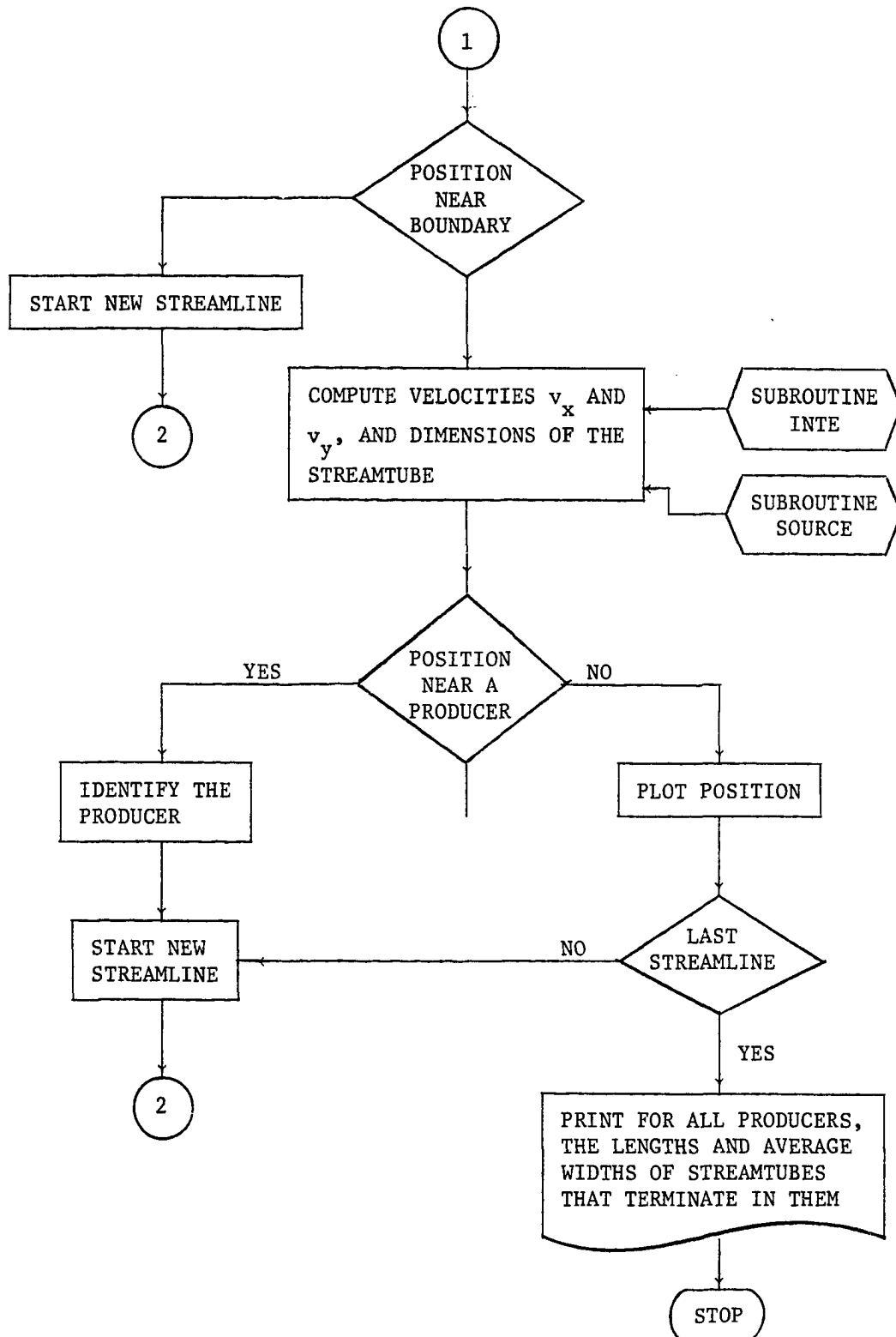






FLOW CHART FOR SUBROUTINE COMPAT





APPENDIX F

COMPUTER IMPLEMENTATION OF STEAMFLOOD RECOVERY CALCULATION

APPENDIX F

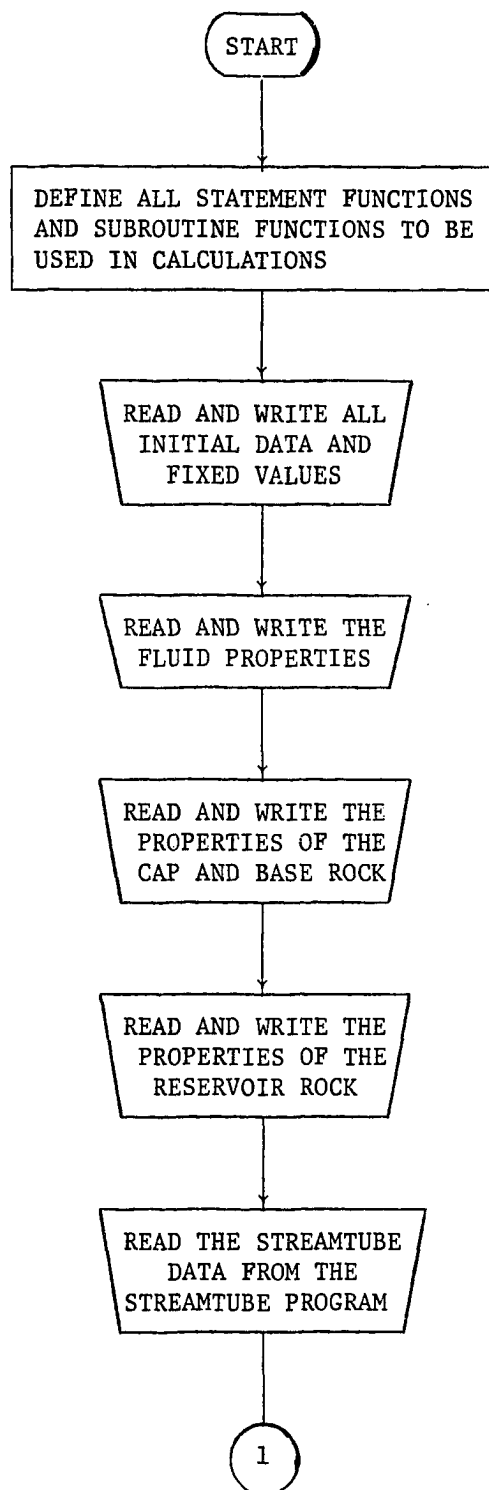
COMPUTER PROGRAM FOR STEAMFLOOD RECOVERY CALCULATION

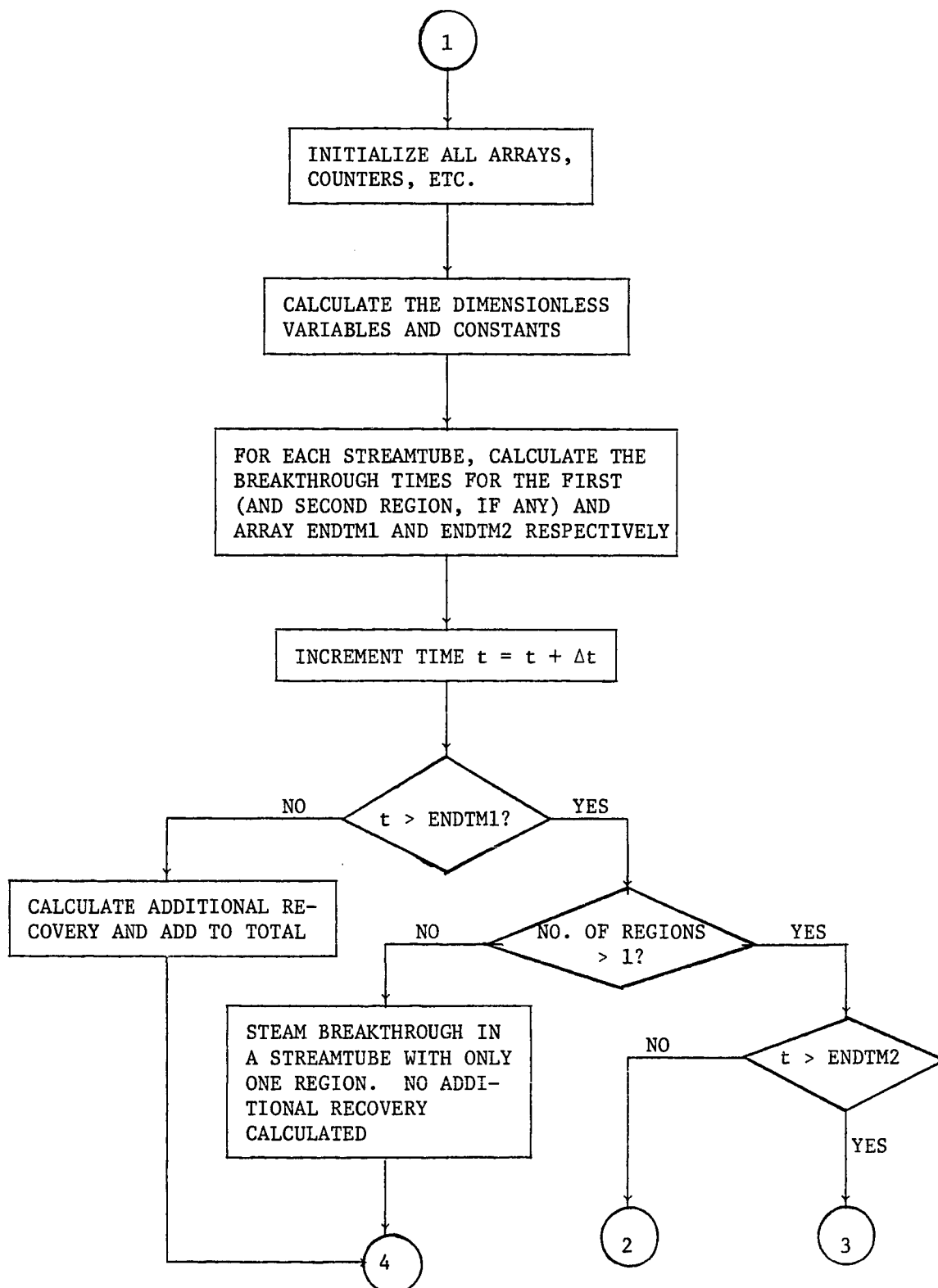
The steamflood recovery program requires as part of its input data, the dimensions of the streamtubes which are obtained from the output of the streamline/streamtube program. Briefly, it performs the following operations:

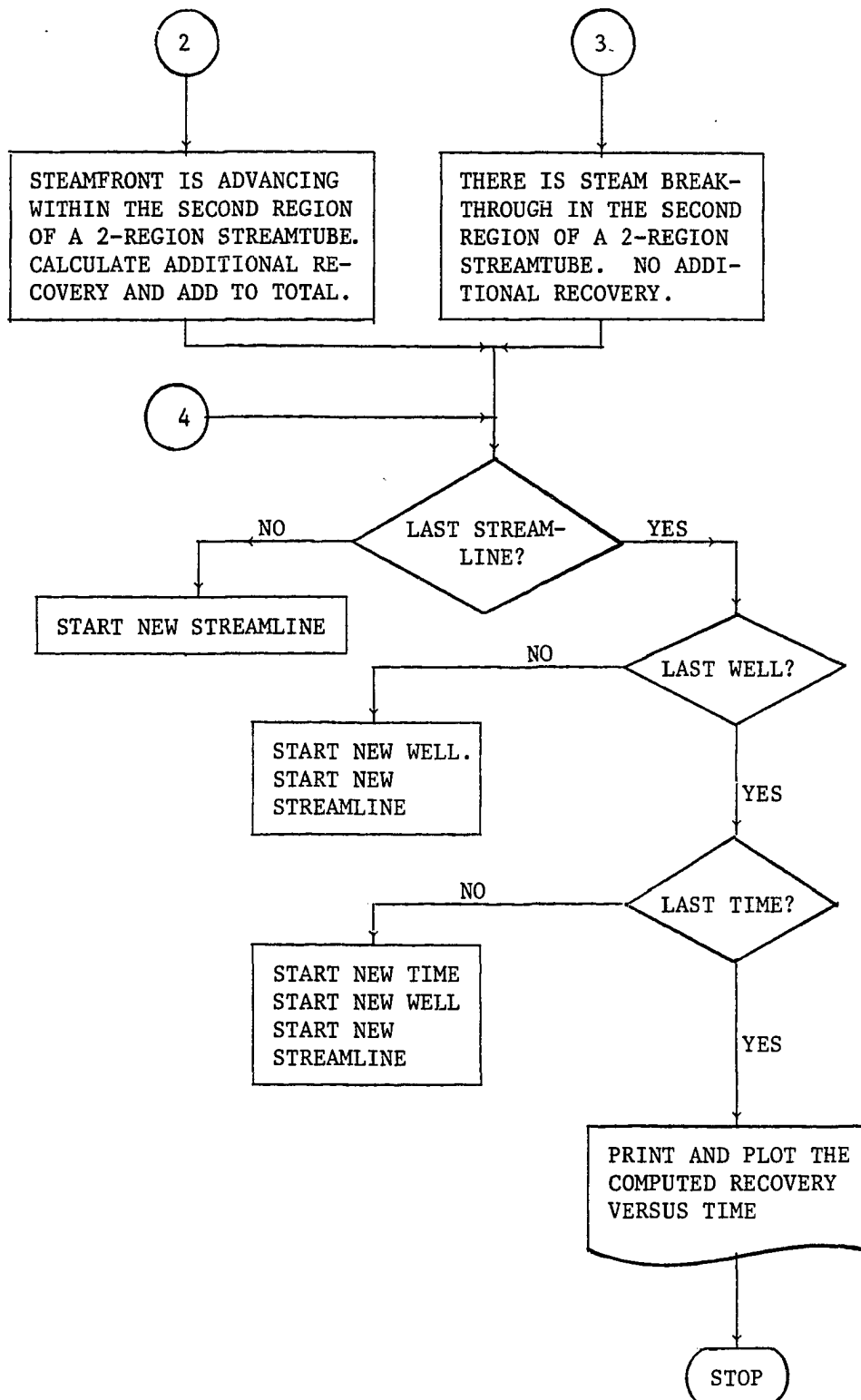
1. Reads the input data under the following format:
 - a. initial and fixed values
 - b. fluid properties
 - c. properties of the cap and base rock
 - d. reservoir rock properties
 - e. initial and residual fluid saturations
 - f. streamtube dimensions from streamline program
2. Calculates the sweep out time for each of the permeability zones in the streamtube.
3. Calculates the recovery from each streamtube, well, and field for any time period. It prints out the results and plots them.

A listing of the flow chart follows.

FLOW CHART FOR STEAMFLOOD RECOVERY CALCULATION







APPENDIX G

RESULTS OF STREAMLINE MODELLING OF SHIELLS CANYON
FIELD AS A SINGLE HOMOGENEOUS RESERVOIR

BOUNDARY ELEMENT MODELLING OF SHIELLS CANYON(203) FIELD
ANALYSIS AS A SINGLE HOMOGENECUS RESERVOIR
SEALED OUTER BOUNDARY
BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, JAN., 1982

REGION 1.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 60
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 13
NUMBER OF INTERFACE NODES= 0

PERMEABILITY(MD)= 140.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X(INCH)	Y(INCH)
1	0.1000	0.1500
2	0.2500	0.2000
3	0.4000	0.2500
4	0.5500	0.2800
5	0.6500	0.3000
6	0.7000	0.3250
7	0.7500	0.3300
8	0.8000	0.3400
9	0.8500	0.3500
10	0.9000	0.4000
11	1.0000	0.4500
12	1.1000	0.5250
13	1.2500	0.6000
14	1.4000	0.6500
15	1.5500	0.6800
16	1.6500	0.6900
17	1.7000	0.7000
18	1.7500	0.7100
19	1.8000	0.7200
20	1.8500	0.7250
21	1.9000	0.7300
22	1.9500	0.7350
23	2.1000	0.7450
24	2.2500	0.7500
25	2.4000	0.7500
26	2.6000	0.7400
27	2.7500	0.7300
28	2.9000	0.7250
29	3.1600	1.0000
30	3.0600	1.2500
31	2.9500	1.5500
32	2.7500	1.6500
33	2.4000	1.8200
34	2.2500	1.8000
35	2.1000	1.7500
36	2.0500	1.7450
37	2.0000	1.7250
38	1.9500	1.7100
39	1.8500	1.6900
40	1.7500	1.6500
41	1.6500	1.6100
42	1.5000	1.5500
43	1.3500	1.4800

44	1.2500	1.4400
45	1.2000	1.4000
46	1.1500	1.3800
47	1.0500	1.3250
48	0.9500	1.2600
49	0.9000	1.2500
50	0.8500	1.2000
51	0.8000	1.1700
52	0.7500	1.1500
53	0.6500	1.0800
54	0.5500	1.0200
55	0.5000	0.9900
56	0.4500	0.9500
57	0.4000	0.9200
58	0.3000	0.8500
59	0.1000	0.7300
60	0.1000	0.5000

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	1	0.0
2	1	0.0
3	1	0.0
4	1	0.0
5	1	0.0
6	1	0.0
7	1	0.0
8	1	0.0
9	1	0.0
10	1	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	1	0.0
35	1	0.0
36	1	0.0
37	1	0.0
38	1	0.0
39	1	0.0
40	1	0.0
41	1	0.0
42	1	0.0
43	1	0.0

44	1	0.0
45	1	0.0
46	1	0.0
47	1	0.0
48	1	0.0
49	1	0.0
50	1	0.0
51	1	0.0
52	1	0.0
53	1	0.0
54	1	0.0
55	1	0.0
56	1	0.0
57	1	0.0
58	1	0.0
59	1	0.0
60	1	0.0

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
 THE CHARACTERISTIC PRESSURE(Psi) IS= 85.0000

RESULTS

BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE(P.S.I)	NORMAL GRADIENT
0.1750E+00	0.1750E+00	0.8469E+02	0.0
0.3250E+00	0.2250E+00	0.8510E+02	0.0
0.4750E+00	0.2650E+00	0.8524E+02	0.0
0.6000E+00	0.2900E+00	0.8547E+02	0.0
0.6750E+00	0.3125E+00	0.8556E+02	0.0
0.7250E+00	0.3275E+00	0.8565E+02	0.0
0.7750E+00	0.3350E+00	0.8568E+02	0.0
0.8250E+00	0.3450E+00	0.8643E+02	0.0
0.8750E+00	0.3750E+00	0.8657E+02	0.0
0.9500E+00	0.4250E+00	0.8830E+02	0.0
0.1050E+01	0.4875E+00	0.8940E+02	0.0
0.1175E+01	0.5625E+00	0.8961E+02	0.0
0.1325E+01	0.6250E+00	0.8772E+02	0.0
0.1475E+01	0.6650E+00	0.8565E+02	0.0
0.1600E+01	0.6850E+00	0.8446E+02	0.0
0.1675E+01	0.6950E+00	0.8398E+02	0.0
0.1725E+01	0.7050E+00	0.8370E+02	0.0
0.1775E+01	0.7150E+00	0.8344E+02	0.0
0.1825E+01	0.7225E+00	0.8315E+02	0.0
0.1875E+01	0.7275E+00	0.8283E+02	0.0
0.1925E+01	0.7325E+00	0.8243E+02	0.0
0.2025E+01	0.7400E+00	0.8129E+02	0.0
0.2175E+01	0.7475E+00	0.7540E+02	0.0
0.2325E+01	0.7500E+00	0.7791E+02	0.0
0.2500E+01	0.7450E+00	0.7689E+02	0.0
0.2675E+01	0.7350E+00	0.7642E+02	0.0
0.2825E+01	0.7275E+00	0.7623E+02	0.0
0.3030E+01	0.8625E+00	0.7610E+02	0.0
0.3110E+01	0.1125E+01	0.7601E+02	0.0
0.3005E+01	0.1400E+01	0.7582E+02	0.0
0.2850E+01	0.1600E+01	0.7579E+02	0.0
0.2575E+01	0.1735E+01	0.7598E+02	0.0
0.2325E+01	0.1810E+01	0.7614E+02	0.0
0.2175E+01	0.1775E+01	0.7637E+02	0.0
0.2075E+01	0.1747E+01	0.7668E+02	0.0
0.2025E+01	0.1735E+01	0.7673E+02	0.0
0.1975E+01	0.1717E+01	0.7699E+02	0.0
0.1900E+01	0.1700E+01	0.7748E+02	0.0
0.1800E+01	0.1670E+01	0.7816E+02	0.0
0.1700E+01	0.1630E+01	0.7900E+02	0.0
0.1575E+01	0.1580E+01	0.8013E+02	0.0
0.1425E+01	0.1515E+01	0.8143E+02	0.0
0.1300E+01	0.1460E+01	0.8246E+02	0.0
0.1225E+01	0.1420E+01	0.8296E+02	0.0
0.1175E+01	0.1390E+01	0.8350E+02	0.0
0.1100E+01	0.1352E+01	0.8418E+02	0.0
0.1000E+01	0.1292E+01	0.8510E+02	0.0
0.9250E+00	0.1255E+01	0.8565E+02	0.0
0.8750E+00	0.1225E+01	0.8572E+02	0.0
0.8250E+00	0.1185E+01	0.8626E+02	0.0
0.7750E+00	0.1160E+01	0.8677E+02	0.0
0.7000E+00	0.1115E+01	0.8714E+02	0.0
0.6000E+00	0.1050E+01	0.8695E+02	0.0
0.5250E+00	0.1005E+01	0.8648E+02	0.0
0.4750E+00	0.9700E+00	0.8610E+02	0.0
0.4250E+00	0.9350E+00	0.8569E+02	0.0

```
0.3500E+00      0.8850E+00      0.4501E+02      0.0
0.2000E+00      0.7900E+00      0.4476E+02      0.0
0.1000E+00      0.6150E+00      0.4469E+02      0.0
0.1000E+00      0.3250E+00      0.4464E+02      0.0
*****
```

PRODUCER NUMBER 1

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.5381E+03	0.4518E+02	0.2800E+02
2	0.4250E+03	0.2375E+02	0.2800E+02
3	0.6881E+03	0.7951E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4414E+03	0.2776E+02	0.2800E+02
2	0.3447E+03	0.1918E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3348E+03	0.1806E+02	0.2800E+02
2	0.3848E+03	0.2176E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2697E+03	0.1351E+02	0.2800E+02
2	0.2697E+03	0.1488E+02	0.2800E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3598E+03	0.1589E+02	0.2800E+02
2	0.3381E+03	0.1426E+02	0.2800E+02

PRODUCER NUMBER 7

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2947E+03	0.1269E+02	0.2800E+02
2	0.2881E+03	0.1196E+02	0.2800E+02

PRODUCER NUMBER 9

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.8098E+03	0.2941E+02	0.2800E+02
2	0.7598E+03	0.2863E+02	0.2800E+02

PRODUCER NUMBER 10

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.1120E+04	0.3876E+02	0.2800E+02
2	0.9500E+03	0.3434E+02	0.2800E+02
3	0.5881E+03	0.1561E+02	0.1000E+02
4	0.5697E+03	0.1231E+02	0.1000E+02
5	0.5980E+03	0.1378E+02	0.1000E+02

PRODUCER NUMBER 11

NUMBER OF STREAMLINES= 4

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2664E+03	0.6268E+01	0.1000E+02
2	0.2914E+03	0.6742E+01	0.1000E+02
3	0.3783E+03	0.8902E+01	0.1000E+02
4	0.2881E+03	0.6459E+01	0.1000E+02

PRODUCER NUMBER 12

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4631E+03	0.1018E+02	0.1000E+02
2	0.3697E+03	0.8221E+01	0.1000E+02
3	0.3316E+03	0.7526E+01	0.1000E+02
4	0.3348E+03	0.7388E+01	0.1000E+02
5	0.4033E+03	0.1092E+02	0.1000E+02

PRODUCER NUMBER 13

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.7164E+03	0.2242E+02	0.1000E+02
2	0.9033E+03	0.3443E+02	0.1000E+02
3	0.7131E+03	0.1976E+02	0.1000E+02
4	0.5381E+03	0.1560E+02	0.1000E+02
5	0.6066E+03	0.2223E+02	0.1000E+02

APPENDIX H

RESULTS OF STEAMFLOOD PREDICTION OF SHIELLS
CANYON FIELD AS A HOMOGENEOUS RESERVOIR
WITH SEALED OUTER BOUNDARY

APPENDIX H

STEAMFLOOD PREDICTION OF SHIELLS CANYON(203) FIELD

METHOD: STREAMLINE/STREAMTUBE METHOD USING (BEM)TECHNIQUE

ANALYSIS AS A SINGLE HOMOGENOUS RESERVOIR

BOUNDARY CONDITION: SEALED OUTER BOUNDARY

PRODUCTION RATES: EQUAL RATES ASSIGNED TO ALL

PRODUCERS IN ENTIRE REGION

BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, 1982

DATA

TOTAL NUMBER OF PRODUCERS=	11
STEAM TEMPERATURE (DEG. F)=	430.0000
STEAM QUALITY=	0.7000
CONVERGENCE LIMIT=	0.0010
RESERVOIR THICKNESS (FT)=	160.0000
STEAM PRESSURE (PSI)=	344.0000
INITIAL RESERVOIR TEMPERATURE(DEG F)=	105.0000

FLUID PROPERTIES

	WATER	OIL	STEAM
DENSITY(LB/CU FT) AT STD. TEMP.	62.4000	55.0000	
DENSITY(LB/CU FT) AT STEAM TEMP.	52.3800	48.7000	0.7431
SPECIFIC HEAT(BTU/LB *F)	1.0000	0.4880	1.0000
LATENT HEAT(BTU/LB)			789.5100

PROPERTIES OF THE CAP AND BASE ROCK

DENSITY(LB/CU. FT)	149.0000
SPECIFIC HEAT(BTU/LB-*F)	0.2130
THERM. COND.(BTU/HR-FT-*F)	1.1000

RESERVOIR ROCK PROPERTIES

	REGION 1	REGION 2
	-----	-----
POROSITY	0.2050	0.1800
DEN.(LB/CU FT)	165.0000	165.0000
SPEC. HEAT(BTU/LB*F)	0.2000	0.2000
PERMEABILITY(MD)	140.0000	70.0000

INITIAL FLUID SATURATIONS

	REGION 1	REGION 2
	-----	-----
WATER	0.5500	0.5500
OIL	0.4500	0.4500
STEAM	0.0	0.0

AVERAGE RESIDUAL SATURATIONS BEHIND THE FRONT

	REGION 1	REGION 2
	-----	-----
WATER	0.5800	0.5800
OIL	0.1800	0.1800
STEAM	0.2400	0.2400

INPUT DATA FROM STREAMLINE PROGRAM

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	538.10	45.18	28.00	0.0	0.0	0.0
1	2	1	425.00	23.75	28.00	0.0	0.0	0.0
1	3	1	688.10	79.51	28.00	0.0	0.0	0.0
2	1	1	441.40	27.76	28.00	0.0	0.0	0.0
2	2	1	344.70	19.18	28.00	0.0	0.0	0.0
3	1	1	334.80	18.06	28.00	0.0	0.0	0.0
3	2	1	384.80	21.76	28.00	0.0	0.0	0.0
4	1	1	269.70	13.51	28.00	0.0	0.0	0.0
4	2	1	269.70	14.88	28.00	0.0	0.0	0.0
5	1	1	359.80	15.89	28.00	0.0	0.0	0.0
5	2	1	338.10	14.26	28.00	0.0	0.0	0.0
6	1	1	294.70	12.69	28.00	0.0	0.0	0.0
6	2	1	288.10	11.96	28.00	0.0	0.0	0.0
7	1	1	809.80	29.41	28.00	0.0	0.0	0.0
7	2	1	759.80	28.63	28.00	0.0	0.0	0.0
8	1	1	112.00	38.76	28.00	0.0	0.0	0.0
8	2	1	950.00	34.34	28.00	0.0	0.0	0.0
8	3	1	588.10	15.61	10.00	0.0	0.0	0.0
8	4	1	569.70	12.31	10.00	0.0	0.0	0.0
8	5	1	598.00	13.78	10.00	0.0	0.0	0.0
9	1	1	266.40	6.27	10.00	0.0	0.0	0.0
9	2	1	291.40	6.74	10.00	0.0	0.0	0.0
9	3	1	378.30	8.90	10.00	0.0	0.0	0.0
9	4	1	288.10	6.46	10.00	0.0	0.0	0.0
10	1	1	463.10	10.18	10.00	0.0	0.0	0.0
10	2	1	369.70	8.22	10.00	0.0	0.0	0.0
10	3	1	331.60	7.53	10.00	0.0	0.0	0.0
10	4	1	334.80	7.39	10.00	0.0	0.0	0.0
10	5	1	403.30	10.92	10.00	0.0	0.0	0.0
11	1	1	716.40	22.42	10.00	0.0	0.0	0.0
11	2	1	903.30	34.43	10.00	0.0	0.0	0.0
11	3	1	713.10	19.76	10.00	0.0	0.0	0.0
11	4	1	538.10	15.60	10.00	0.0	0.0	0.0
11	5	1	606.60	22.23	10.00	0.0	0.0	0.0

CALCULATED BREAKTHROUGH TIMES(HOURS) FOR EACH STREAMTUBE

PROD. NO.	STREAMLINE NO.	CODE	ENDTIME(1)	ENDTIME(2)
1	1	1	15868.5625	0.0
1	2	1	6144.6992	0.0
1	3	1	39442.2148	0.0
2	1	1	7555.3711	0.0
2	2	1	3929.7900	0.0
3	1	1	3578.2664	0.0
3	2	1	5040.6680	0.0
4	1	1	2110.3652	0.0
4	2	1	2332.9001	0.0
5	1	1	3374.3901	0.0
5	2	1	2823.9495	0.0
6	1	1	2163.1919	0.0
6	2	1	1991.6033	0.0
7	1	1	15514.1211	0.0
7	2	1	14048.2070	0.0
8	1	1	2531.5093	0.0
8	2	1	21971.1758	0.0
8	3	1	16871.8437	0.0
8	4	1	12564.0469	0.0
8	5	1	14984.0547	0.0
9	1	1	2734.8716	0.0
9	2	1	3240.8657	0.0
9	3	1	5716.1836	0.0
9	4	1	3062.2073	0.0
10	1	1	8182.2500	0.0
10	2	1	5127.8672	0.0
10	3	1	4165.2852	0.0
10	4	1	4126.5156	0.0
10	5	1	7606.9062	0.0
11	1	1	31527.6133	0.0
11	2	1	67928.6875	0.0
11	3	1	27188.5703	0.0
11	4	1	15291.1250	0.0
11	5	1	25876.1133	0.0

REAL TIME(HCURS)= 6000.0000

DIMENSIONLESS TIME= 0.0325

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TCTAL
1	1	1	0.1557E+05		
1	2	1	0.1557E+05		
1	3	1	0.1557E+05		
				0.4670E+05	
2	1	1	0.1557E+05		
2	2	1	0.1043E+05		
				0.2599E+05	
3	1	1	0.9537E+04		
3	2	1	0.1321E+05		
				0.2274E+05	
4	1	1	0.5747E+04		
4	2	1	0.6330E+04		
				0.1208E+05	
5	1	1	0.9017E+04		
5	2	1	0.7604E+04		
				0.1662E+05	
6	1	1	0.5898E+04		
6	2	1	0.5435E+04		
				0.1133E+05	
7	1	1	0.1557E+05		
7	2	1	0.1557E+05		
				0.3113E+05	
8	1	1	0.6847E+04		
8	2	1	0.1557E+05		
8	3	1	0.5560E+04		
8	4	1	0.5560E+04		
8	5	1	0.5560E+04		
				0.3909E+05	
9	1	1	0.2634E+04		
9	2	1	0.3099E+04		
9	3	1	0.5311E+04		
9	4	1	0.2935E+04		
				0.1398E+05	
10	1	1	0.5560E+04		
10	2	1	0.4794E+04		
10	3	1	0.3936E+04		
10	4	1	0.3901E+04		
10	5	1	0.5560E+04		
				0.2375E+05	
11	1	1	0.5560E+04		
11	2	1	0.5560E+04		
11	3	1	0.5560E+04		
11	4	1	0.5560E+04		
11	5	1	0.5560E+04		
				0.2780E+05	
					0.2712E+06

REAL TIME(HOURS)= 12000.0000

DIMENSIONLESS TIME= 0.0650

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.2969E+05		
1	2	1	0.1592E+05		
1	3	1	0.2969E+05		
				0.7530E+05	
2	1	1	0.1933E+05		
2	2	1	0.1043E+05		
				0.2975E+05	
3	1	1	0.9537E+04		
3	2	1	0.1321E+05		
				0.2274E+05	
4	1	1	0.5747E+04		
4	2	1	0.6330E+04		
				0.1208E+05	
5	1	1	0.9017E+04		
5	2	1	0.7604E+04		
				0.1662E+05	
6	1	1	0.5898E+04		
6	2	1	0.5435E+04		
				0.1133E+05	
7	1	1	0.2969E+05		
7	2	1	0.2969E+05		
				0.5938E+05	
8	1	1	0.6847E+04		
8	2	1	0.2969E+05		
8	3	1	0.1060E+05		
8	4	1	0.1060E+05		
8	5	1	0.1060E+05		
				0.6835E+05	
9	1	1	0.2634E+04		
9	2	1	0.3099E+04		
9	3	1	0.5311E+04		
9	4	1	0.2935E+04		
				0.1398E+05	
10	1	1	0.7436E+04		
10	2	1	0.4794E+04		
10	3	1	0.3936E+04		
10	4	1	0.3901E+04		
10	5	1	0.6946E+04		
				0.2701E+05	
11	1	1	0.1060E+05		
11	2	1	0.1060E+05		
11	3	1	0.1060E+05		
11	4	1	0.1060E+05		
11	5	1	0.1060E+05		
				0.5302E+05	
					0.3896E+06

REAL TIME(HCURS)= 18000.0000

DIMENSIONLESS TIME= 0.0975

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3834E+05		
1	2	1	0.1592E+05		
1	3	1	0.4299E+05		
				0.9726E+05	
2	1	1	0.1933E+05		
2	2	1	0.1043E+05		
				0.2975E+05	
3	1	1	0.9537E+04		
3	2	1	0.1321E+05		
				0.2274E+05	
4	1	1	0.5747E+04		
4	2	1	0.6330E+04		
				0.1208E+05	
5	1	1	0.9017E+04		
5	2	1	0.7604E+04		
				0.1662E+05	
6	1	1	0.5898E+04		
6	2	1	0.5435E+04		
				0.1133E+05	
7	1	1	0.3756E+05		
7	2	1	0.3431E+05		
				0.7187E+05	
8	1	1	0.6847E+04		
8	2	1	0.4299E+05		
8	3	1	0.1448E+05		
8	4	1	0.1106E+05		
8	5	1	0.1300E+05		
				0.8838E+05	
9	1	1	0.2634E+04		
9	2	1	0.3099E+04		
9	3	1	0.5311E+04		
9	4	1	0.2935E+04		
				0.1398E+05	
10	1	1	0.7436E+04		
10	2	1	0.4794E+04		
10	3	1	0.3936E+04		
10	4	1	0.3901E+04		
10	5	1	0.6946E+04		
				0.2701E+05	
11	1	1	0.1535E+05		
11	2	1	0.1535E+05		
11	3	1	0.1535E+05		
11	4	1	0.1324E+05		
11	5	1	0.1535E+05		
				0.7466E+05	
					0.4657E+06

REAL TIME(HOURS)= 24000.0000

DIMENSIONLESS TIME= 0.1300

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3834E+05		
1	2	1	0.1592E+05		
1	3	1	0.5569E+05		
				0.1099E+06	
2	1	1	0.1933E+05		
2	2	1	0.1043E+05		
				0.2975E+05	
3	1	1	0.9537E+04		
3	2	1	0.1321E+05		
				0.2274E+05	
4	1	1	0.5747E+04		
4	2	1	0.6330E+04		
				0.1208E+05	
5	1	1	0.9017E+04		
5	2	1	0.7604E+04		
				0.1662E+05	
6	1	1	0.5898E+04		
6	2	1	0.5435E+04		
				0.1133E+05	
7	1	1	0.3756E+05		
7	2	1	0.3431E+05		
				0.7187E+05	
8	1	1	0.6847E+04		
8	2	1	0.5145E+05		
8	3	1	0.1448E+05		
8	4	1	0.1106E+05		
8	5	1	0.1300E+05		
				0.9684E+05	
9	1	1	0.2634E+04		
9	2	1	0.3099E+04		
9	3	1	0.5311E+04		
9	4	1	0.2935E+04		
				0.1398E+05	
10	1	1	0.7436E+04		
10	2	1	0.4794E+04		
10	3	1	0.3936E+04		
10	4	1	0.3901E+04		
10	5	1	0.6946E+04		
				0.2701E+05	
11	1	1	0.1989E+05		
11	2	1	0.1989E+05		
11	3	1	0.1989E+05		
11	4	1	0.1324E+05		
11	5	1	0.1989E+05		
				0.9279E+05	
					0.5050E+06

REAL TIME(HOURS)= 30000.0000

DIMENSIONLESS TIME= 0.1625

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3834E+05		
1	2	1	0.1592E+05		
1	3	1	0.6789E+05		
				0.1222E+06	
2	1	1	0.1933E+05		
2	2	1	0.1043E+05		
				0.2975E+05	
3	1	1	0.9537E+04		
3	2	1	0.1321E+05		
				0.2274E+05	
4	1	1	0.5747E+04		
4	2	1	0.6330E+04		
				0.1208E+05	
5	1	1	0.9017E+04		
5	2	1	0.7604E+04		
				0.1662E+05	
6	1	1	0.5898E+04		
6	2	1	0.5435E+04		
				0.1133E+05	
7	1	1	0.3756E+05		
7	2	1	0.3431E+05		
				0.7187E+05	
8	1	1	0.6847E+04		
8	2	1	0.5145E+05		
8	3	1	0.1448E+05		
8	4	1	0.1106E+05		
8	5	1	0.1300E+05		
				0.9684E+05	
9	1	1	0.2634E+04		
9	2	1	0.3099E+04		
9	3	1	0.5311E+04		
9	4	1	0.2935E+04		
				0.1398E+05	
10	1	1	0.7436E+04		
10	2	1	0.4794E+04		
10	3	1	0.3936E+04		
10	4	1	0.3901E+04		
10	5	1	0.6946E+04		
				0.2701E+05	
11	1	1	0.2425E+05		
11	2	1	0.2425E+05		
11	3	1	0.2222E+05		
11	4	1	0.1324E+05		
11	5	1	0.2127E+05		
				0.1052E+06	
					0.5296E+06

PREDICTED RECOVERY

TIME(DAYS)	RECOVERY(BBLS)
62.5000	0.1697E+05
125.0000	0.3120E+05
187.5000	0.4079E+05
250.0000	0.4830E+05
312.5000	0.5455E+05
375.0000	0.5967E+05
437.5000	0.6456E+05
500.0000	0.6938E+05
562.5000	0.7399E+05
625.0000	0.7808E+05
687.5000	0.8089E+05
750.0000	0.8293E+05
812.5000	0.8450E+05
875.0000	0.8685E+05
937.5000	0.8858E+05
1000.0000	0.8993E+05
1062.5000	0.9127E+05
1125.0000	0.9245E+05
1187.5000	0.9340E+05
1250.0000	0.9432E+05

APPENDIX I

RESULTS OF STREAMLINE MODELLING OF SHIELLS
CANYON FIELD AS A HOMOGENEOUS RESERVOIR
HAVING A PART OF THE BOUNDARY AT
CONSTANT PRESSURE WHILE THE
REMAINDER IS SEALED

BOUNDARY ELEMENT MODELLING OF SHIELLS CANYON(203) FIELD
ANALYSIS AS A SINGLE HOMOGENEOUS RESERVOIR
PART OF THE BOUNDARY AT CONSTANT PRESSURE. THE REMAINDER SEALED
BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, JAN., 1982

REGION 1.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 60
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 13
NUMBER OF INTERFACE NODES= 0

PERMEABILITY(MD)= 140.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X (INCH)	Y (INCH)
1	0.1000	0.1500
2	0.2500	0.2000
3	0.4000	0.2500
4	0.5500	0.2800
5	0.6500	0.3000
6	0.7000	0.3250
7	0.7500	0.3300
8	0.8000	0.3400
9	0.8500	0.3500
10	0.9000	0.4000
11	1.0000	0.4500
12	1.1000	0.5250
13	1.2500	0.6000
14	1.4000	0.6500
15	1.5500	0.6800
16	1.6500	0.6900
17	1.7000	0.7000
18	1.7500	0.7100
19	1.8000	0.7200
20	1.8500	0.7250
21	1.9000	0.7300
22	1.9500	0.7350
23	2.1000	0.7450
24	2.2500	0.7500
25	2.4000	0.7500
26	2.6000	0.7400
27	2.7500	0.7300
28	2.9000	0.7250
29	3.1600	1.0000
30	3.0600	1.2500
31	2.9500	1.5500
32	2.7500	1.6500
33	2.4000	1.8200
34	2.2500	1.8000
35	2.1000	1.7500
36	2.0500	1.7450
37	2.0000	1.7250
38	1.9500	1.7100
39	1.8500	1.6900
40	1.7500	1.6500
41	1.6500	1.6100
42	1.5000	1.5500
43	1.3500	1.4800

44	1.2500	1.4400
45	1.2000	1.4000
46	1.1500	1.3800
47	1.0500	1.3250
48	0.9500	1.2600
49	0.9000	1.2500
50	0.8500	1.2000
51	0.8000	1.1700
52	0.7500	1.1500
53	0.6500	1.0800
54	0.5500	1.0200
55	0.5000	0.9900
56	0.4500	0.9500
57	0.4000	0.9200
58	0.3000	0.8500
59	0.1000	0.7300
60	0.1000	0.5000

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	1	0.0
2	1	0.0
3	1	0.0
4	1	0.0
5	1	0.0
6	1	0.0
7	1	0.0
8	1	0.0
9	1	0.0
10	1	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	1	0.0
35	1	0.0
36	1	0.0
37	1	0.0
38	1	0.0
39	1	0.0
40	1	0.0
41	1	0.0
42	1	0.0
43	1	0.0

44	1	0.0
45	1	0.0
46	1	0.0
47	1	0.0
48	1	0.0
49	1	0.0
50	1	0.0
51	1	0.0
52	1	0.0
53	1	0.0
54	1	0.0
55	1	0.0
56	1	0.0
57	1	0.0
58	1	0.0
59	0	0.8500000E+02
60	0	0.8500000E+02

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
 THE CHARACTERISTIC PRESSURE(Psi) IS= 85.0000

RESULTS

BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE (PSI)	NORMAL GRADIENT
0.1750E+00	0.1750E+00	0.4448E+02	0.0
0.3250E+00	0.2250E+00	0.4472E+02	0.0
0.4750E+00	0.2650E+00	0.4487E+02	0.0
0.6000E+00	0.2900E+00	0.4510E+02	0.0
0.6750E+00	0.3125E+00	0.4520E+02	0.0
0.7250E+00	0.3275E+00	0.4527E+02	0.0
0.7750E+00	0.3350E+00	0.4551E+02	0.0
0.8250E+00	0.3450E+00	0.4605E+02	0.0
0.8750E+00	0.3750E+00	0.4660E+02	0.0
0.9500E+00	0.4250E+00	0.4793E+02	0.0
0.1050E+01	0.4875E+00	0.4903E+02	0.0
0.1175E+01	0.5625E+00	0.4923E+02	0.0
0.1325E+01	0.6250E+00	0.4734E+02	0.0
0.1475E+01	0.6650E+00	0.4527E+02	0.0
0.1600E+01	0.6850E+00	0.4408E+02	0.0
0.1675E+01	0.6950E+00	0.4360E+02	0.0
0.1725E+01	0.7050E+00	0.4332E+02	0.0
0.1775E+01	0.7150E+00	0.4305E+02	0.0
0.1825E+01	0.7225E+00	0.4277E+02	0.0
0.1875E+01	0.7275E+00	0.4244E+02	0.0
0.1925E+01	0.7325E+00	0.4204E+02	0.0
0.2025E+01	0.7400E+00	0.4090E+02	0.0
0.2175E+01	0.7475E+00	0.3901E+02	0.0
0.2325E+01	0.7500E+00	0.3752E+02	0.0
0.2500E+01	0.7450E+00	0.3650E+02	0.0
0.2675E+01	0.7350E+00	0.3603E+02	0.0
0.2825E+01	0.7275E+00	0.3584E+02	0.0
0.3030E+01	0.8625E+00	0.3571E+02	0.0
0.3110E+01	0.1125E+01	0.3561E+02	0.0
0.3005E+01	0.1400E+01	0.3543E+02	0.0
0.2850E+01	0.1600E+01	0.3540E+02	0.0
0.2755E+01	0.1735E+01	0.3560E+02	0.0
0.2325E+01	0.1810E+01	0.3579E+02	0.0
0.2175E+01	0.1775E+01	0.3599E+02	0.0
0.2075E+01	0.1747E+01	0.3625E+02	0.0
0.2025E+01	0.1735E+01	0.3634E+02	0.0
0.1975E+01	0.1717E+01	0.3659E+02	0.0
0.1900E+01	0.1700E+01	0.3709E+02	0.0
0.1800E+01	0.1670E+01	0.3778E+02	0.0
0.1700E+01	0.1630E+01	0.3862E+02	0.0
0.1575E+01	0.1580E+01	0.3974E+02	0.0
0.1425E+01	0.1515E+01	0.4104E+02	0.0
0.1300E+01	0.1460E+01	0.4208E+02	0.0
0.1225E+01	0.1420E+01	0.4259E+02	0.0
0.1175E+01	0.1390E+01	0.4313E+02	0.0
0.1100E+01	0.1352E+01	0.4381E+02	0.0
0.1000E+01	0.1292E+01	0.4473E+02	0.0
0.9250E+00	0.1255E+01	0.4525E+02	0.0
0.8750E+00	0.1225E+01	0.4535E+02	0.0
0.8250E+00	0.1185E+01	0.4588E+02	0.0
0.7750E+00	0.1160E+01	0.4641E+02	0.0
0.7000E+00	0.1115E+01	0.4677E+02	0.0
0.6000E+00	0.1050E+01	0.4659E+02	0.0
0.5250E+00	0.1005E+01	0.4611E+02	0.0
0.4750E+00	0.9700E+00	0.4574E+02	0.0
0.4250E+00	0.9350E+00	0.4534E+02	0.0

0.3500E+00	0.8850E+00	0.8537E+02	0.0
0.2000E+00	0.7900E+00	0.8509E+02	0.0
0.1000E+00	0.6150E+00	0.8500E+02	-0.4122E-03
0.1000E+00	0.3250E+00	0.8500E+02	0.2567E-03

PRODUCER NUMBER 1

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.5414E+03	0.4584E+02	0.2800E+02
2	0.4250E+03	0.2367E+02	0.2800E+02
3	0.6783E+03	0.7369E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4414E+03	0.2762E+02	0.2800E+02
2	0.3447E+03	0.1914E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3381E+03	0.1806E+02	0.2800E+02
2	0.3848E+03	0.2162E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2697E+03	0.1351E+02	0.2800E+02
2	0.2697E+03	0.1491E+02	0.2800E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3598E+03	0.1591E+02	0.2800E+02
2	0.3381E+03	0.1426E+02	0.2800E+02

PRODUCER NUMBER 7

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2947E+03	0.1270E+02	0.2800E+02
2	0.2881E+03	0.1196E+02	0.2800E+02

PRODUCER NUMBER 9

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.8098E+03	0.2943E+02	0.2800E+02
2	0.7598E+03	0.2867E+02	0.2800E+02

PRODUCER NUMBER 10

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.1120E+04	0.3869E+02	0.2800E+02
2	0.9500E+03	0.3438E+02	0.2800E+02
3	0.5881E+03	0.1555E+02	0.1000E+02
4	0.5697E+03	0.1231E+02	0.1000E+02
5	0.5980E+03	0.1379E+02	0.1000E+02

PRODUCER NUMBER 11

NUMBER OF STREAMLINES= 4

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2664E+03	0.6270E+01	0.1000E+02
2	0.2914E+03	0.6746E+01	0.1000E+02
3	0.3783E+03	0.8922E+01	0.1000E+02
4	0.2881E+03	0.6459E+01	0.1000E+02

PRODUCER NUMBER 12

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4631E+03	0.1017E+02	0.1000E+02
2	0.3697E+03	0.8223E+01	0.1000E+02
3	0.3316E+03	0.7528E+01	0.1000E+02
4	0.3348E+03	0.7392E+01	0.1000E+02
5	0.4033E+03	0.1097E+02	0.1000E+02

PRODUCER NUMBER 13

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.7164E+03	0.2247E+02	0.1000E+02
2	0.9033E+03	0.3484E+02	0.1000E+02
3	0.7131E+03	0.1981E+02	0.1000E+02
4	0.5381E+03	0.1561E+02	0.1000E+02
5	0.6066E+03	0.2234E+02	0.1000E+02

APPENDIX J

RESULTS OF STEAMFLOOD PREDICTION OF SHIELLS
CANYON FIELD AS A HOMOGENEOUS RESERVOIR.

PART OF ITS BOUNDARY AT CONSTANT
PRESSURE, THE REMAINDER SEALED

APPENDIX J

STEAMFLOOD PREDICTION OF SHIELLS CANYON(203) FIELD

METHOD: STREAMLINE/STREAMTUBE METHOD USING (BEM) TECHNIQUE

ANALYSIS AS A SINGLE HOMOGENEOUS RESERVOIR

BOUNDARY CONDITION: PART OF THE BOUNDARY AT CONSTANT
PRESSURE, THE REMAINDER IS SEALED.

PRODUCTION RATES: ASSIGNED EQUAL IN ENTIRE REGION

BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, 1982

DATA

TOTAL NUMBER OF PRODUCERS=	11
STEAM TEMPERATURE (DEG. F)=	430.0000
STEAM QUALITY=	0.7000
CONVERGENCE LIMIT=	0.0010
RESERVOIR THICKNESS (FT)=	160.0000
STEAM PRESSURE (PSI)=	344.0000
INITIAL RESERVOIR TEMPERATURE(DEG F)=	105.0000

FLUID PROPERTIES

	WATER -----	OIL -----	STEAM -----
DENSITY(LB/CU FT) AT STD. TEMP.	62.4000	55.0000	
DENSITY(LB/CU FT) AT STEAM TEMP.	52.3800	48.7000	0.7431
SPECIFIC HEAT(BTU/LB *F)	1.0000	0.4880	1.0000
LATENT HEAT(BTU/LB)			789.5100

PROPERTIES OF THE CAP AND BASE ROCK

DENSITY(LB/CU. FT)	145.0000
SPECIFIC HEAT(BTU/LB-*F)	0.2130
THERM. COND.(BTU/HR-FT-*F)	1.1000

RESERVOIR ROCK PROPERTIES

	REGION 1	REGION 2
POROSITY	0.2050	0.1800
DEN.(LB/CU FT)	165.0000	165.0000
SPEC. HEAT (BTU/LB*F)	0.2000	0.2000
PERMEABILITY (MD)	140.0000	70.0000

INITIAL FLUID SATURATIONS

	REGION 1	REGION 2
WATER	0.5500	0.5500
OIL	0.4500	0.4500
STEAM	0.0	0.0

AVERAGE RESIDUAL SATURATIONS BEHIND THE FRONT

	REGION 1	REGION 2
WATER	0.5800	0.5800
OIL	0.1800	0.1800
STEAM	0.2400	0.2400

INPUT DATA FROM STREAMLINE PROGRAM

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	541.40	45.84	28.00	0.0	0.0	0.0
1	2	1	425.00	23.67	28.00	0.0	0.0	0.0
1	3	1	678.30	73.69	28.00	0.0	0.0	0.0
2	1	1	441.40	27.62	28.00	0.0	0.0	0.0
2	2	1	344.70	19.14	28.00	0.0	0.0	0.0
3	1	1	338.10	18.06	28.00	0.0	0.0	0.0
3	2	1	384.80	21.62	28.00	0.0	0.0	0.0
4	1	1	269.70	13.51	28.00	0.0	0.0	0.0
4	2	1	269.70	14.91	28.00	0.0	0.0	0.0
5	1	1	359.80	15.91	28.00	0.0	0.0	0.0
5	2	1	338.10	14.26	28.00	0.0	0.0	0.0
6	1	1	294.70	12.70	28.00	0.0	0.0	0.0
6	2	1	288.10	11.96	28.00	0.0	0.0	0.0
7	1	1	809.80	29.43	28.00	0.0	0.0	0.0
7	2	1	759.80	28.67	28.00	0.0	0.0	0.0
8	1	1	1120.00	38.69	28.00	0.0	0.0	0.0
8	2	1	950.00	34.38	28.00	0.0	0.0	0.0
8	3	1	588.10	15.55	10.00	0.0	0.0	0.0
8	4	1	569.70	12.31	10.00	0.0	0.0	0.0
8	5	1	598.00	13.79	10.00	0.0	0.0	0.0
9	1	1	266.40	6.27	10.00	0.0	0.0	0.0
9	2	1	291.40	6.75	10.00	0.0	0.0	0.0
9	3	1	378.30	8.92	10.00	0.0	0.0	0.0
9	4	1	288.10	6.46	10.00	0.0	0.0	0.0
10	1	1	463.10	10.17	10.00	0.0	0.0	0.0
10	2	1	369.70	8.22	10.00	0.0	0.0	0.0
10	3	1	331.60	7.53	10.00	0.0	0.0	0.0
10	4	1	334.80	7.39	10.00	0.0	0.0	0.0
10	5	1	403.30	10.97	10.00	0.0	0.0	0.0
11	1	1	716.40	22.47	10.00	0.0	0.0	0.0
11	2	1	903.30	34.84	10.00	0.0	0.0	0.0
11	3	1	713.10	19.81	10.00	0.0	0.0	0.0
11	4	1	538.10	15.61	10.00	0.0	0.0	0.0
11	5	1	606.60	22.34	10.00	0.0	0.0	0.0

CALCULATED BREAKTHROUGH TIMES(HOURS) FOR EACH STREAMTUBE

PROD. NO.	STREAMLINE NO.	CODE	ENDTIME(1)	ENDTIME(2)
1	1	1	16232.3555	0.0
1	2	1	6122.6797	0.0
1	3	1	35560.2891	0.0
2	1	1	7514.6562	0.0
2	2	1	3921.2864	0.0
3	1	1	3615.2166	0.0
3	2	1	5006.4922	0.0
4	1	1	2110.3652	0.0
4	2	1	2337.7622	0.0
5	1	1	3378.9150	0.0
5	2	1	2823.9495	0.0
6	1	1	2169.8333	0.0
6	2	1	1991.6033	0.0
7	1	1	15525.5352	0.0
7	2	1	14069.7266	0.0
8	1	1	30226.8555	0.0
8	2	1	21999.7695	0.0
8	3	1	16800.4453	0.0
8	4	1	12564.0469	0.0
8	5	1	14995.8437	0.0
9	1	1	2735.7625	0.0
9	2	1	3243.0010	0.0
9	3	1	5729.8828	0.0
9	4	1	3062.2073	0.0
10	1	1	8173.4727	0.0
10	2	1	5129.1250	0.0
10	3	1	4166.4844	0.0
10	4	1	4128.8320	0.0
10	5	1	7644.0664	0.0
11	1	1	31607.5586	0.0
11	2	1	68893.5000	0.0
11	3	1	27266.1172	0.0
11	4	1	15301.9453	0.0
11	5	1	26020.0859	0.0

REAL TIME(HCURS)= 6000.0000

DIMENSIONLESS TIME= 0.0325

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.1557E+05		
1	2	1	0.1557E+05		
1	3	1	0.1557E+05		
				0.4670E+05	
2	1	1	0.1557E+05		
2	2	1	0.1041E+05		
				0.2597E+05	
3	1	1	0.9631E+04		
3	2	1	0.1312E+05		
				0.2275E+05	
4	1	1	0.5747E+04		
4	2	1	0.6342E+04		
				0.1209E+05	
5	1	1	0.9029E+04		
5	2	1	0.7604E+04		
				0.1663E+05	
6	1	1	0.5903E+04		
6	2	1	0.5435E+04		
				0.1134E+05	
7	1	1	0.1557E+05		
7	2	1	0.1557E+05		
				0.3113E+05	
8	1	1	0.1557E+05		
8	2	1	0.1557E+05		
8	3	1	0.5560E+04		
8	4	1	0.5560E+04		
8	5	1	0.5560E+04		
				0.4781E+05	
9	1	1	0.2634E+04		
9	2	1	0.3100E+04		
9	3	1	0.5323E+04		
9	4	1	0.2935E+04		
				0.1399E+05	
10	1	1	0.5560E+04		
10	2	1	0.4795E+04		
10	3	1	0.3937E+04		
10	4	1	0.3903E+04		
10	5	1	0.5560E+04		
				0.2375E+05	
11	1	1	0.5560E+04		
11	2	1	0.5560E+04		
11	3	1	0.5560E+04		
11	4	1	0.5560E+04		
11	5	1	0.5560E+04		
				0.2780E+05	
					0.2800E+06

REAL TIME(HCURS)= 12000.0000

DIMENSIONLESS TIME= 0.0650

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.2969E+05		
1	2	1	0.1587E+05		
1	3	1	0.2969E+05		
				0.7525E+05	
2	1	1	0.1923E+05		
2	2	1	0.1041E+05		
				0.2963E+05	
3	1	1	0.9631E+04		
3	2	1	0.1312E+05		
				0.2275E+05	
4	1	1	0.5747E+04		
4	2	1	0.6342E+04		
				0.1209E+05	
5	1	1	0.9029E+04		
5	2	1	0.7604E+04		
				0.1663E+05	
6	1	1	0.5903E+04		
6	2	1	0.5435E+04		
				0.1134E+05	
7	1	1	0.2969E+05		
7	2	1	0.2969E+05		
				0.5938E+05	
8	1	1	0.2969E+05		
8	2	1	0.2969E+05		
8	3	1	0.1060E+05		
8	4	1	0.1060E+05		
8	5	1	0.1060E+05		
				0.9119E+05	
9	1	1	0.2634E+04		
9	2	1	0.3100E+04		
9	3	1	0.5323E+04		
9	4	1	0.2935E+04		
				0.1399E+05	
10	1	1	0.7428E+04		
10	2	1	0.4795E+04		
10	3	1	0.3937E+04		
10	4	1	0.3903E+04		
10	5	1	0.6978E+04		
				0.2704E+05	
11	1	1	0.1060E+05		
11	2	1	0.1060E+05		
11	3	1	0.1060E+05		
11	4	1	0.1060E+05		
11	5	1	0.1060E+05		
				0.5302E+05	
					0.4123E+06

REAL TIME(HOURS)= 18000.0000

DIMENSIONLESS TIME= 0.0975

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3914E+05		
1	2	1	0.1587E+05		
1	3	1	0.4299E+05		
				0.9800E+05	
2	1	1	0.1923E+05		
2	2	1	0.1041E+05		
				0.2963E+05	
3	1	1	0.9631E+04		
3	2	1	0.1312E+05		
				0.2275E+05	
4	1	1	0.5747E+04		
4	2	1	0.6342E+04		
				0.1209E+05	
5	1	1	0.9029E+04		
5	2	1	0.7604E+04		
				0.1663E+05	
6	1	1	0.5903E+04		
6	2	1	0.5435E+04		
				0.1134E+05	
7	1	1	0.3759E+05		
7	2	1	0.3436E+05		
				0.7195E+05	
8	1	1	0.4299E+05		
8	2	1	0.4299E+05		
8	3	1	0.1442E+05		
8	4	1	0.1106E+05		
8	5	1	0.1301E+05		
				0.1245E+06	
9	1	1	0.2634E+04		
9	2	1	0.3100E+04		
9	3	1	0.5323E+04		
9	4	1	0.2935E+04		
				0.1399E+05	
10	1	1	0.7428E+04		
10	2	1	0.4795E+04		
10	3	1	0.3937E+04		
10	4	1	0.3903E+04		
10	5	1	0.6978E+04		
				0.2704E+05	
11	1	1	0.1535E+05		
11	2	1	0.1535E+05		
11	3	1	0.1535E+05		
11	4	1	0.1325E+05		
11	5	1	0.1535E+05		
				0.7466E+05	
					0.5025E+06

REAL TIME(HOURS)= 24000.0000

DIMENSIONLESS TIME= 0.1300

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3914E+05		
1	2	1	0.1587E+05		
1	3	1	0.5569E+05		
				0.1107E+06	
2	1	1	0.1923E+05		
2	2	1	0.1041E+05		
				0.2963E+05	
3	1	1	0.9631E+04		
3	2	1	0.1312E+05		
				0.2275E+05	
4	1	1	0.5747E+04		
4	2	1	0.6342E+04		
				0.1209E+05	
5	1	1	0.9029E+04		
5	2	1	0.7604E+04		
				0.1663E+05	
6	1	1	0.5903E+04		
6	2	1	0.5435E+04		
				0.1134E+05	
7	1	1	0.3759E+05		
7	2	1	0.3436E+05		
				0.7195E+05	
8	1	1	0.5569E+05		
8	2	1	0.5151E+05		
8	3	1	0.1442E+05		
8	4	1	0.1106E+05		
8	5	1	0.1301E+05		
				0.1457E+06	
9	1	1	0.2634E+04		
9	2	1	0.3100E+04		
9	3	1	0.5323E+04		
9	4	1	0.2935E+04		
				0.1399E+05	
10	1	1	0.7428E+04		
10	2	1	0.4795E+04		
10	3	1	0.3937E+04		
10	4	1	0.3903E+04		
10	5	1	0.6978E+04		
				0.2704E+05	
11	1	1	0.1989E+05		
11	2	1	0.1989E+05		
11	3	1	0.1989E+05		
11	4	1	0.1325E+05		
11	5	1	0.1989E+05		
				0.9280E+05	
					0.5546E+06

REAL TIME(HOURS)= 30000.0000

DIMENSIONLESS TIME= 0.1625

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.3914E+05		
1	2	1	0.1587E+05		
1	3	1	0.6789E+05		
				0.1229E+06	
2	1	1	0.1923E+05		
2	2	1	0.1041E+05		
				0.2963E+05	
3	1	1	0.9631E+04		
3	2	1	0.1312E+05		
				0.2275E+05	
4	1	1	0.5747E+04		
4	2	1	0.6342E+04		
				0.1209E+05	
5	1	1	0.9029E+04		
5	2	1	0.7604E+04		
				0.1663E+05	
6	1	1	0.5903E+04		
6	2	1	0.5435E+04		
				0.1134E+05	
7	1	1	0.3759E+05		
7	2	1	0.3436E+05		
				0.7195E+05	
8	1	1	0.6789E+05		
8	2	1	0.5151E+05		
8	3	1	0.1442E+05		
8	4	1	0.1106E+05		
8	5	1	0.1301E+05		
				0.1579E+06	
9	1	1	0.2634E+04		
9	2	1	0.3100E+04		
9	3	1	0.5323E+04		
9	4	1	0.2935E+04		
				0.1399E+05	
10	1	1	0.7428E+04		
10	2	1	0.4795E+04		
10	3	1	0.3937E+04		
10	4	1	0.3903E+04		
10	5	1	0.6978E+04		
				0.2704E+05	
11	1	1	0.2425E+05		
11	2	1	0.2425E+05		
11	3	1	0.2228E+05		
11	4	1	0.1325E+05		
11	5	1	0.2137E+05		
				0.1054E+06	
					0.5916E+06

PREDICTED RECOVERY

TIME(DAYS)	RECOVERY(BBLS)
62.5000	0.1697E+05
125.0000	0.3142E+05
187.5000	0.4171E+05
250.0000	0.4986E+05
312.5000	0.5674E+05
375.0000	0.6249E+05
437.5000	0.6800E+05
500.0000	0.7343E+05
562.5000	0.7865E+05
625.0000	0.8334E+05
687.5000	0.8689E+05
750.0000	0.8950E+05
812.5000	0.9205E+05
875.0000	0.9456E+05
937.5000	0.9687E+05
1000.0000	0.9877E+05
1062.5000	0.1007E+06
1125.0000	0.1024E+06
1187.5000	0.1039E+06
1250.0000	0.1054E+06

APPENDIX K

RESULTS OF STREAMLINE MODELLING OF SHIELLS
CANYON FIELD AS A PIECEWISE HOMOGENEOUS
RESERVOIR WITH SEALED OUTER BOUNDARY

BOUNDARY ELEMENT MODELLING OF SHIELLS CANYON(203) FIELD
ANALYSIS AS A PIECEWISE-HOMOGENEOUS RESERVIOR
SEALED OUTER BOUNDARY
BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, 1981

REGION 1.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 43
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 7
NUMBER OF INTERFACE NODES= 10

PERMEABILITY(MD)= 140.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X(INCH)	Y(INCH)
1	1.5000	1.5500
2	1.3500	1.4800
3	1.2500	1.4400
4	1.2000	1.4000
5	1.1500	1.3800
6	1.0500	1.3250
7	0.9500	1.2600
8	0.9000	1.2500
9	0.8500	1.2000
10	0.8000	1.1700
11	0.7500	1.1500
12	0.6500	1.0800
13	0.5500	1.0200
14	0.5000	0.9900
15	0.4500	0.9500
16	0.4000	0.9200
17	0.3000	0.8500
18	0.1000	0.7300
19	0.1000	0.5000
20	0.1000	0.1500
21	0.2500	0.2000
22	0.4000	0.2500
23	0.5500	0.2800
24	0.6500	0.3000
25	0.7000	0.3250
26	0.7500	0.3300
27	0.8000	0.3400
28	0.8500	0.3500
29	0.9000	0.4000
30	1.0000	0.4500
31	1.1000	0.5250
32	1.2500	0.6000
33	1.4000	0.6500
34	1.5500	0.6800
35	1.5500	0.7500
36	1.5400	0.8500
37	1.5000	0.9500
38	1.5000	1.0500
39	1.4800	1.1500
40	1.4600	1.2500
41	1.4600	1.3500
42	1.4800	1.4500
43	1.5000	1.5000

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	1	0.0
2	1	0.0
3	1	0.0
4	1	0.0
5	1	0.0
6	1	0.0
7	1	0.0
8	1	0.0
9	1	0.0
10	1	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	2	0.0
35	2	0.0
36	2	0.0
37	2	0.0
38	2	0.0
39	2	0.0
40	2	0.0
41	2	0.0
42	2	0.0
43	2	0.0

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
THE CHARACTERISTIC PRESSURE(PSI) IS= 85.0000

REGION 2.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 37
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 6
NUMBER OF INTERFACE NODES= 10

PERMEABILITY(MD)= 70.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X(INCH)	Y(INCH)
1	1.5000	1.5500
2	1.5000	1.5000
3	1.4800	1.4500
4	1.4600	1.3500
5	1.4600	1.2500
6	1.4800	1.1500
7	1.5000	1.0500
8	1.5000	0.9500
9	1.5400	0.8500
10	1.5500	0.7500
11	1.5500	0.6800
12	1.6500	0.6900
13	1.7000	0.7000
14	1.7500	0.7100
15	1.8000	0.7200
16	1.8500	0.7250
17	1.9000	0.7300
18	1.9500	0.7350
19	2.1000	0.7450
20	2.2500	0.7500
21	2.4000	0.7500
22	2.6000	0.7400
23	2.7500	0.7300
24	2.9000	0.7250
25	3.1600	1.0000
26	3.0600	1.2500
27	2.9500	1.5500
28	2.7500	1.6500
29	2.4000	1.8200
30	2.2500	1.8000
31	2.1000	1.7500
32	2.0500	1.7450
33	2.0000	1.7250
34	1.9500	1.7100
35	1.8500	1.6900
36	1.7500	1.6500
37	1.6500	1.6100

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	2	0.0
2	2	0.0
3	2	0.0
4	2	0.0
5	2	0.0
6	2	0.0
7	2	0.0
8	2	0.0
9	2	0.0
10	2	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	1	0.0
35	1	0.0
36	1	0.0
37	1	0.0

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
THE CHARACTERISTIC PRESSURE(PSI) IS= 85.0000

RESULTS BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE(PSI)	NORMAL GRADIENT
0.1425E+01	0.1515E+01	-0.5609E+02	0.0
0.1300E+01	0.1460E+01	-0.5464E+02	0.0
0.1225E+01	0.1420E+01	-0.5398E+02	0.0
0.1175E+01	0.1390E+01	-0.5335E+02	0.0
0.1100E+01	0.1352E+01	-0.5260E+02	0.0
0.1000E+01	0.1292E+01	-0.5159E+02	0.0
0.9250E+00	0.1255E+01	-0.5107E+02	0.0
0.8750E+00	0.1225E+01	-0.5091E+02	0.0
0.8250E+00	0.1185E+01	-0.5035E+02	0.0
0.7750E+00	0.1160E+01	-0.4976E+02	0.0
0.7000E+00	0.1115E+01	-0.4942E+02	0.0
0.6000E+00	0.1050E+01	-0.4958E+02	0.0
0.5250E+00	0.1005E+01	-0.5004E+02	0.0
0.4750E+00	0.9700E+00	-0.5041E+02	0.0
0.4250E+00	0.9350E+00	-0.5080E+02	0.0
0.3500E+00	0.8850E+00	-0.5113E+02	0.0
0.2000E+00	0.7900E+00	-0.5137E+02	0.0
0.1000E+00	0.6150E+00	-0.5143E+02	0.0
0.1000E+00	0.3250E+00	-0.5139E+02	0.0
0.1750E+00	0.1750E+00	-0.5119E+02	0.0
0.3250E+00	0.2250E+00	-0.5141E+02	0.0
0.4750E+00	0.2650E+00	-0.5127E+02	0.0
0.6000E+00	0.2900E+00	-0.5104E+02	0.0
0.6750E+00	0.3125E+00	-0.5092E+02	0.0
0.7250E+00	0.3275E+00	-0.5087E+02	0.0
0.7750E+00	0.3350E+00	-0.5064E+02	0.0
0.8250E+00	0.3450E+00	-0.5009E+02	0.0
0.8750E+00	0.3750E+00	-0.4956E+02	0.0
0.9500E+00	0.4250E+00	-0.4820E+02	0.0
0.1050E+01	0.4875E+00	-0.4710E+02	0.0
0.1175E+01	0.5625E+00	-0.4686E+02	0.0
0.1325E+01	0.6250E+00	-0.4866E+02	0.0
0.1475E+01	0.6650E+00	-0.5054E+02	0.0
0.1550E+01	0.7150E+00	-0.5120E+02	-0.8409E-02
0.1545E+01	0.8000E+00	-0.5147E+02	-0.7983E-02
0.1520E+01	0.9000E+00	-0.5197E+02	-0.8892E-02
0.1500E+01	0.1000E+01	-0.5283E+02	-0.4986E-02
0.1490E+01	0.1100E+01	-0.5365E+02	-0.8297E-02
0.1470E+01	0.1200E+01	-0.5434E+02	-0.9221E-02
0.1460E+01	0.1300E+01	-0.5503E+02	-0.9038E-02
0.1470E+01	0.1400E+01	-0.5580E+02	-0.9774E-02
0.1490E+01	0.1475E+01	-0.5652E+02	-0.1061E-01
0.1500E+01	0.1525E+01	-0.5741E+02	-0.1380E-01

RESULTS

BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE(PSI)	NORMAL GRADIENT
0.1500E+01	0.1525E+01	-0.5741E+02	0.2760E-01
0.1490E+01	0.1475E+01	-0.5652E+02	0.2121E-01
0.1470E+01	0.1400E+01	-0.5580E+02	0.1955E-01
0.1460E+01	0.1300E+01	-0.5503E+02	0.1808E-01
0.1470E+01	0.1200E+01	-0.5434E+02	0.1844E-01
0.1490E+01	0.1100E+01	-0.5365E+02	0.1659E-01
0.1500E+01	0.1000E+01	-0.5283E+02	0.9972E-02
0.1520E+01	0.9000E+00	-0.5197E+02	0.1778E-01
0.1545E+01	0.8000E+00	-0.5147E+02	0.1597E-01
0.1550E+01	0.7150E+00	-0.5120E+02	0.1682E-01
0.1600E+01	0.6850E+00	-0.5169E+02	0.0
0.1675E+01	0.6950E+00	-0.5239E+02	0.0
0.1725E+01	0.7050E+00	-0.5283E+02	0.0
0.1775E+01	0.7150E+00	-0.5326E+02	0.0
0.1825E+01	0.7225E+00	-0.5374E+02	0.0
0.1875E+01	0.7275E+00	-0.5431E+02	0.0
0.1925E+01	0.7325E+00	-0.5507E+02	0.0
0.2025E+01	0.7400E+00	-0.5723E+02	0.0
0.2175E+01	0.7475E+00	-0.6092E+02	0.0
0.2325E+01	0.7500E+00	-0.6386E+02	0.0
0.2500E+01	0.7450E+00	-0.6587E+02	0.0
0.2675E+01	0.7350E+00	-0.6679E+02	0.0
0.2825E+01	0.7275E+00	-0.6714E+02	0.0
0.3030E+01	0.8625E+00	-0.6741E+02	0.0
0.3110E+01	0.1125E+01	-0.6763E+02	0.0
0.3005E+01	0.1400E+01	-0.6795E+02	0.0
0.2850E+01	0.1600E+01	-0.6800E+02	0.0
0.2575E+01	0.1735E+01	-0.6759E+02	0.0
0.2325E+01	0.1810E+01	-0.6705E+02	0.0
0.2175E+01	0.1775E+01	-0.6675E+02	0.0
0.2075E+01	0.1747E+01	-0.6637E+02	0.0
0.2025E+01	0.1735E+01	-0.6602E+02	0.0
0.1975E+01	0.1717E+01	-0.6551E+02	0.0
0.1900E+01	0.1700E+01	-0.6445E+02	0.0
0.1800E+01	0.1670E+01	-0.6301E+02	0.0
0.1700E+01	0.1630E+01	-0.6123E+02	0.0
0.1575E+01	0.1580E+01	-0.5877E+02	0.0

PRODUCER NUMBER 1

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.5533E+03	0.5142E+02	0.2800E+02
2	0.4250E+03	0.2374E+02	0.2800E+02
3	0.6697E+03	0.7554E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4447E+03	0.2874E+02	0.2800E+02
2	0.3447E+03	0.1926E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3381E+03	0.1804E+02	0.2800E+02
2	0.3881E+03	0.2164E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2730E+03	0.1335E+02	0.2800E+02
2	0.2664E+03	0.1458E+02	0.2800E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3664E+03	0.1561E+02	0.2800E+02
2	0.3348E+03	0.1371E+02	0.2800E+02

PRODUCER NUMBER 7

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2881E+03	0.1253E+02	0.2800E+02
2	0.2947E+03	0.1203E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4523E+03	0.5838E+02	0.2800E+02
1	0.2164E+03	0.2901E+02	0.2800E+02
2	0.4779E+03	0.5484E+02	0.2800E+02
2	0.2816E+03	0.4460E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.6447E+03	0.2909E+02	0.1000E+02
2	0.5881E+03	0.1317E+02	0.1000E+02
3	0.6164E+03	0.1461E+02	0.1000E+02
1	0.3765E+03	0.5957E+02	0.2800E+02
1	0.5348E+03	0.4241E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 4

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2664E+03	0.6362E+01	0.1000E+02
2	0.2881E+03	0.6823E+01	0.1000E+02
3	0.3664E+03	0.8241E+01	0.1000E+02
4	0.2881E+03	0.6602E+01	0.1000E+02

PRODUCER NUMBER 5

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4750E+03	0.1096E+02	0.1000E+02
2	0.3848E+03	0.8447E+01	0.1000E+02
3	0.3348E+03	0.7668E+01	0.1000E+02
4	0.3316E+03	0.7443E+01	0.1000E+02
5	0.3881E+03	0.9796E+01	0.1000E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.7033E+03	0.2170E+02	0.1000E+02
2	0.8447E+03	0.2877E+02	0.1000E+02
3	0.7197E+03	0.1975E+02	0.1000E+02
4	0.5381E+03	0.1590E+02	0.1000E+02
5	0.5914E+03	0.2104E+02	0.1000E+02

NO STREAMLINE FROM REGION 2 CROSSES INTO REGION 1

APPENDIX L

STEAMFLOOD PREDICTION OF SHIELLS CANYON FIELD
ANALYSED AS A PIECEWISE HOMOGENEOUS
RESERVOIR HAVING SEALED BOUNDARY

APPENDIX L

STEAMFLOOD PREDICTION OF SHIELLS CANYON(203) FIELD

METHOD: STREAMLINE/STREAMTUBE METHOD USING (BEM)TECHNIQUE

ANALYSIS AS A PIECE-WISE HOMOGENEOUS RESERVOIR

HAVING TWO REGIONS OF UNEQUAL PERMEABILITY

BOUNDARY CONDITION: SEALED OUTER BOUNDARY

PRODUCTION RATES: ASSIGNED EQUAL IN ENTIRE REGION

BY: D. T. NUMBERS, MARCH 1982

DATA

TOTAL NUMBER OF PRODUCERS=	11
STEAM TEMPERATURE (DEG. F)=	430.0000
STEAM QUALITY=	0.7000
CONVERGENCE LIMIT=	0.0010
RESERVOIR THICKNESS (FT)=	160.0000
STEAM PRESSURE (PSI)=	344.0000
INITIAL RESERVOIR TEMPERATURE(DEG F)=	105.0000

FLUID PROPERTIES

	WATER	OIL	STEAM
	-----	-----	-----
DENSITY(LB/CU FT) AT STD. TEMP.	62.4000	55.0000	
DENSITY(LB/CU FT) AT STEAM TEMP.	52.3800	48.7000	0.7431
SPECIFIC HEAT(BTU/LB *F)	1.0000	0.4860	1.0000
LATENT HEAT(BTU/LB)			789.5100

PROPERTIES OF THE CAP AND BASE ROCK

DENSITY(LB/CU. FT)	149.0000
SPECIFIC HEAT(BTU/LB-*F)	0.2130
THERM. COND.(BTU/HR-FT-*F)	1.1000

RESERVOIR ROCK PROPERTIES

	REGION 1	REGION 2
	-----	-----
POROSITY	0.2050	0.1800
DEN.(LB/CU FT)	165.0000	165.0000
SPEC. HEAT(BTU/LB*F)	0.2000	0.2000
PERMEABILITY(MD)	140.0000	70.0000

INITIAL FLUID SATURATIONS

	REGION 1	REGION 2
	-----	-----
WATER	0.5500	0.5500
OIL	0.4500	0.4500
STEAM	0.0	0.0

AVERAGE RESIDUAL SATURATIONS BEHIND THE FRONT

	REGION 1	REGION 2
	-----	-----
WATER	0.5800	0.5800
OIL	0.1800	0.1800
STEAM	0.2400	0.2400

INPUT DATA FROM STREAMLINE PROGRAM

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	553.30	51.42	28.00	0.0	0.0	0.0
1	2	1	425.00	23.74	28.00	0.0	0.0	0.0
1	3	1	669.70	75.54	28.00	0.0	0.0	0.0
2	1	1	444.70	28.74	28.00	0.0	0.0	0.0
2	2	1	344.70	19.26	28.00	0.0	0.0	0.0
3	1	1	338.10	12.04	28.00	0.0	0.0	0.0
3	2	1	388.10	21.64	28.00	0.0	0.0	0.0
4	1	1	273.00	13.35	28.00	0.0	0.0	0.0
4	2	1	266.40	14.58	28.00	0.0	0.0	0.0
5	1	1	366.40	15.51	28.00	0.0	0.0	0.0
5	2	1	334.80	13.71	28.00	0.0	0.0	0.0
6	1	1	288.10	12.53	28.00	0.0	0.0	0.0
6	2	1	294.70	12.03	28.00	0.0	0.0	0.0
7	1	2	452.30	58.38	28.00	216.40	29.01	28.00
7	2	2	477.90	54.84	28.00	281.60	44.60	28.00
8	1	2	376.50	59.57	28.00	534.80	42.41	28.00
8	2	1	644.70	29.09	10.00	0.0	0.0	0.0
8	3	1	588.10	13.17	10.00	0.0	0.0	0.0
8	4	1	616.40	14.61	10.00	0.0	0.0	0.0
9	1	1	266.40	6.36	10.00	0.0	0.0	0.0
9	2	1	288.10	6.32	10.00	0.0	0.0	0.0
9	3	1	366.40	8.24	10.00	0.0	0.0	0.0
9	4	1	288.10	6.60	10.00	0.0	0.0	0.0
10	1	1	475.00	10.96	10.00	0.0	0.0	0.0
10	2	1	384.80	8.45	10.00	0.0	0.0	0.0
10	3	1	334.80	7.67	10.00	0.0	0.0	0.0
10	4	1	331.60	7.44	10.00	0.0	0.0	0.0
10	5	1	388.10	9.80	10.00	0.0	0.0	0.0
11	1	1	703.30	21.70	10.00	0.0	0.0	0.0
11	2	1	844.70	28.77	10.00	0.0	0.0	0.0
11	3	1	719.70	19.75	10.00	0.0	0.0	0.0
11	4	1	538.10	15.90	10.00	0.0	0.0	0.0
11	5	1	591.40	21.04	10.00	0.0	0.0	0.0

CALCULATED BREAKTHROUGH TIMES(HOURS) FOR EACH STREAMTUBE

PROD. NO.	STREAMLINE NO.	CODE	ENDTIME(1)	ENDTIME(2)
1	1	1	18872.6953	0.0
1	2	1	6142.0391	0.0
1	3	1	36053.1523	0.0
2	1	1	7903.7891	0.0
2	2	1	3947.0635	0.0
3	1	1	3611.0315	0.0
3	2	1	5056.8008	0.0
4	1	1	2110.9402	0.0
4	2	1	2255.1360	0.0
5	1	1	3375.7742	0.0
5	2	1	2682.9524	0.0
6	1	1	2090.1030	0.0
6	2	1	2051.3933	0.0
7	1	2	17379.4336	21095.7500
7	2	2	17236.3633	24932.4414
8	1	2	14525.8516	29210.0820
8	2	1	37627.2305	0.0
8	3	1	14001.4336	0.0
8	4	1	16518.7578	0.0
9	1	1	2777.7727	0.0
9	2	1	3242.8152	0.0
9	3	1	5092.4102	0.0
9	4	1	3133.0645	0.0
10	1	1	9101.8555	0.0
10	2	1	5505.5273	0.0
10	3	1	4291.3516	0.0
10	4	1	4117.0391	0.0
10	5	1	6502.2969	0.0
11	1	1	29753.0977	0.0
11	2	1	50784.7969	0.0
11	3	1	27456.6445	0.0
11	4	1	15614.1445	0.0
11	5	1	23645.6953	0.0

REAL TIME(HOURS)= 6000.0000

DIMENSIONLESS TIME= 0.0325

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.1557E+05		
1	2	1	0.1557E+05		
1	3	1	0.1557E+05		
				0.4670E+05	
2	1	1	0.1557E+05		
2	2	1	0.1047E+05		
				0.2604E+05	
3	1	1	0.9620E+04		
3	2	1	0.1325E+05		
				0.2287E+05	
4	1	1	0.5748E+04		
4	2	1	0.6126E+04		
				0.1187E+05	
5	1	1	0.9021E+04		
5	2	1	0.7240E+04		
				0.1626E+05	
6	1	1	0.5694E+04		
6	2	1	0.5592E+04		
				0.1129E+05	
7	1	2	0.1557E+05		
7	2	2	0.1557E+05		
				0.3113E+05	
8	1	2	0.1557E+05		
8	2	1	0.5560E+04		
8	3	1	0.5560E+04		
8	4	1	0.5560E+04		
				0.3225E+05	
9	1	1	0.2673E+04		
9	2	1	0.3100E+04		
9	3	1	0.4762E+04		
9	4	1	0.3000E+04		
				0.1354E+05	
10	1	1	0.5560E+04		
10	2	1	0.5127E+04		
10	3	1	0.4049E+04		
10	4	1	0.3893E+04		
10	5	1	0.5560E+04		
				0.2419E+05	
11	1	1	0.5560E+04		
11	2	1	0.5560E+04		
11	3	1	0.5560E+04		
11	4	1	0.5560E+04		
11	5	1	0.5560E+04		
				0.2780E+05	
					0.2639E+06

REAL TIME(HOURS)= 12000.0000

DIMENSIONLESS TIME= 0.0650

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.2969E+05		
1	2	1	0.1591E+05		
1	3	1	0.2969E+05		
				0.7529E+05	
2	1	1	0.2016E+05		
2	2	1	0.1047E+05		
				0.3063E+05	
3	1	1	0.9620E+04		
3	2	1	0.1325E+05		
				0.2287E+05	
4	1	1	0.5743E+04		
4	2	1	0.6126E+04		
				0.1187E+05	
5	1	1	0.9021E+04		
5	2	1	0.7240E+04		
				0.1626E+05	
6	1	1	0.5694E+04		
6	2	1	0.5592E+04		
				0.1129E+05	
7	1	2	0.2969E+05		
7	2	2	0.2969E+05		
				0.5938E+05	
8	1	2	0.2969E+05		
8	2	1	0.1060E+05		
8	3	1	0.1060E+05		
8	4	1	0.1060E+05		
				0.6150E+05	
9	1	1	0.2073E+04		
9	2	1	0.3100E+04		
9	3	1	0.4762E+04		
9	4	1	0.3000E+04		
				0.1354E+05	
10	1	1	0.8211E+04		
10	2	1	0.5127E+04		
10	3	1	0.4049E+04		
10	4	1	0.3893E+04		
10	5	1	0.5996E+04		
				0.2728E+05	
11	1	1	0.1060E+05		
11	2	1	0.1060E+05		
11	3	1	0.1060E+05		
11	4	1	0.1060E+05		
11	5	1	0.1060E+05		
				0.5302E+05	
					0.3829E+06

PEAL TIME(HOURS)= 18000.0000

DIMENSIONLESS TIME= 0.0975

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4299E+05		
1	2	1	0.1591E+05		
1	3	1	0.4299E+05		
				0.1019E+06	
2	1	1	0.2016E+05		
2	2	1	0.1047E+05		
				0.3063E+05	
3	1	1	0.9620E+04		
3	2	1	0.1325E+05		
				0.2287E+05	
4	1	1	0.5748E+04		
4	2	1	0.6126E+04		
				0.1187E+05	
5	1	1	0.9021E+04		
5	2	1	0.7240E+04		
				0.1626E+05	
6	1	1	0.5694E+04		
6	2	1	0.5592E+04		
				0.1129E+05	
7	1	2	0.4283E+05		
7	2	2	0.4279E+05		
				0.8562E+05	
8	1	2	0.4207E+05		
8	2	1	0.1535E+05		
8	3	1	0.1222E+05		
8	4	1	0.1420E+05		
				0.8384E+05	
9	1	1	0.2673E+04		
9	2	1	0.3100E+04		
9	3	1	0.4762E+04		
9	4	1	0.3000E+04		
				0.1354E+05	
10	1	1	0.8211E+04		
10	2	1	0.5127E+04		
10	3	1	0.4049E+04		
10	4	1	0.3893E+04		
10	5	1	0.5996E+04		
				0.2728E+05	
11	1	1	0.1535E+05		
11	2	1	0.1535E+05		
11	3	1	0.1535E+05		
11	4	1	0.1349E+05		
11	5	1	0.1535E+05		
				0.7491E+05	
					0.4800E+06

REAL TIME(HOURS)= 24000.0000

DIMENSIONLESS TIME= 0.1300

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4487E+05		
1	2	1	0.1591E+05		
1	3	1	0.5569E+05		
				0.1165E+06	
2	1	1	0.2016E+05		
2	2	1	0.1047E+05		
				0.3063E+05	
3	1	1	0.9620E+04		
3	2	1	0.1325E+05		
				0.2287E+05	
4	1	1	0.5748E+04		
4	2	1	0.6126E+04		
				0.1187E+05	
5	1	1	0.9021E+04		
5	2	1	0.7240E+04		
				0.1626E+05	
6	1	1	0.5694E+04		
6	2	1	0.5592E+04		
				0.1129E+05	
7	1	2	0.5034E+05		
7	2	2	0.5395E+05		
				0.1043E+06	
8	1	2	0.5323E+05		
8	2	1	0.1939E+05		
8	3	1	0.1222E+05		
8	4	1	0.1420E+05		
				0.9954E+05	
9	1	1	0.2673E+04		
9	2	1	0.3100E+04		
9	3	1	0.4762E+04		
9	4	1	0.3000E+04		
				0.1354E+05	
10	1	1	0.8211E+04		
10	2	1	0.5127E+04		
10	3	1	0.4049E+04		
10	4	1	0.3893E+04		
10	5	1	0.5996E+04		
				0.2728E+05	
11	1	1	0.1989E+05		
11	2	1	0.1989E+05		
11	3	1	0.1989E+05		
11	4	1	0.1349E+05		
11	5	1	0.1963E+05		
				0.9278E+05	
					0.5468E+06

REAL TIME(HOURS)= 30000.0000

DIMENSIONLESS TIME= 0.1625

WELL	STP/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4487E+05		
1	2	1	0.1591E+05		
1	3	1	0.6789E+05		
				0.1287E+06	
2	1	1	0.2016E+05		
2	2	1	0.1047E+05		
				0.3063E+05	
3	1	1	0.9620E+04		
3	2	1	0.1325E+05		
				0.2287E+05	
4	1	1	0.5748E+04		
4	2	1	0.6126E+04		
				0.1187E+05	
5	1	1	0.9021E+04		
5	2	1	0.7240E+04		
				0.1626E+05	
6	1	1	0.5694E+04		
6	2	1	0.3392E+04		
				0.1129E+05	
7	1	2	0.5034E+05		
7	2	2	0.5873E+05		
				0.1091E+06	
8	1	2	0.6678E+05		
8	2	1	0.2425E+05		
8	3	1	0.1222E+05		
8	4	1	0.1420E+05		
				0.1174E+06	
9	1	1	0.2673E+04		
9	2	1	0.3100E+04		
9	3	1	0.4762E+04		
9	4	1	0.3000E+04		
				0.1354E+05	
10	1	1	0.8211E+04		
10	2	1	0.5127E+04		
10	3	1	0.4049E+04		
10	4	1	0.3893E+04		
10	5	1	0.5996E+04		
				0.2728E+05	
11	1	1	0.2407E+05		
11	2	1	0.2425E+05		
11	3	1	0.2242E+05		
11	4	1	0.1349E+05		
11	5	1	0.1963E+05		
				0.1039E+06	
					0.5928E+06

PREDICTED RECOVERY

TIME(DAYS)	RECOVERY(BBLS)
62.5000	0.1624E+05
125.0000	0.2988E+05
187.5000	0.3953E+05
250.0000	0.4700E+05
312.5000	0.5309E+05
375.0000	0.5847E+05
437.5000	0.6338E+05
500.0000	0.6819E+05
562.5000	0.7294E+05
625.0000	0.7746E+05
687.5000	0.8167E+05
750.0000	0.8548E+05
812.5000	0.8893E+05
875.0000	0.9201E+05
937.5000	0.9490E+05
1000.0000	0.9732E+05
1062.5000	0.1001E+06
1125.0000	0.1019E+06
1187.5000	0.1035E+06
1250.0000	0.1056E+06

APPENDIX M

RESULTS OF STREAMLINE MODELLING OF SHIELLS
CANYON FIELD AS A PIECEWISE HOMOGENEOUS
RESERVOIR. PART OF ITS BOUNDARY AT
CONSTANT PRESSURE, THE REMAINDER
SEALED

APPENDIX M

BOUNDARY ELEMENT MODELLING OF SHIELLS CANYON(203) FIELD
ANALYSIS AS A PIECEWISE-HOMOGENEOUS RESERVIOR
PART OF THE BOUNDARY AT CONSTANT PRESSURE, THE REMAINDER SEALED
BY: D. T. NUMBERE, UNIVERSITY OF OKLAHOMA, 1982

REGION 1.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 43
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 7
NUMBER OF INTERFACE NODES= 10

PERMEABILITY(MD)= 140.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X (INCH)	Y (INCH)
1	1.5000	1.5500
2	1.3500	1.4800
3	1.2500	1.4400
4	1.2000	1.4000
5	1.1500	1.3800
6	1.0500	1.3250
7	0.9500	1.2600
8	0.9000	1.2500
9	0.8500	1.2000
10	0.8000	1.1700
11	0.7500	1.1500
12	0.6500	1.0800
13	0.5500	1.0200
14	0.5000	0.9900
15	0.4500	0.9500
16	0.4000	0.9200
17	0.3000	0.8500
18	0.1000	0.7300
19	0.1000	0.5000
20	0.1000	0.1500
21	0.2500	0.2000
22	0.4000	0.2500
23	0.5500	0.2800
24	0.6500	0.3000
25	0.7000	0.3250
26	0.7500	0.3300
27	0.8000	0.3400
28	0.8500	0.3500
29	0.9000	0.4000
30	1.0000	0.4500
31	1.1000	0.5250
32	1.2500	0.6000
33	1.4000	0.6500
34	1.5500	0.6800
35	1.5500	0.7500
36	1.5400	0.8500
37	1.5000	0.9500
38	1.5000	1.0500
39	1.4800	1.1500
40	1.4600	1.2500
41	1.4600	1.3500
42	1.4800	1.4500
43	1.5000	1.5000

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	1	0.0
2	1	0.0
3	1	0.0
4	1	0.0
5	1	0.0
6	1	0.0
7	1	0.0
8	1	0.0
9	1	0.0
10	1	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	0	0.8500000E+02
19	0	0.8500000E+02
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	2	0.0
35	2	0.0
36	2	0.0
37	2	0.0
38	2	0.0
39	2	0.0
40	2	0.0
41	2	0.0
42	2	0.0
43	2	0.0

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
0.4500	0.5500	-69.1000
0.5000	0.8000	-69.1000
0.7500	0.4000	-69.1000
0.9000	1.1500	-69.1000
1.0000	0.8000	560.0000
1.2500	1.2000	-69.1000
1.3500	1.0000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
THE CHARACTERISTIC PRESSURE(PSI) IS= 85.0000

REGION 2.0

DATA

NUMBER OF BOUNDARY ELEMENTS= 37
NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS CALCULATED= 0
NUMBER OF SOURCES AND SINKS= 6
NUMBER OF INTERFACE NODES= 10

PERMEABILITY(MD)= 70.0000

THICKNESS(FT)= 160.0000

POROSITY= 0.2050

SCALE: 1 INCH = 738.46FEET

THE COORDINATES OF THE EXTREME POINTS OF THE BOUNDARY ELEMENTS

POINT	X (INCH)	Y (INCH)
1	1.5000	1.5500
2	1.5000	1.5000
3	1.4800	1.4500
4	1.4600	1.3500
5	1.4600	1.2500
6	1.4800	1.1500
7	1.5000	1.0500
8	1.5000	0.9500
9	1.5400	0.8500
10	1.5500	0.7500
11	1.5500	0.6800
12	1.6500	0.6900
13	1.7000	0.7000
14	1.7500	0.7100
15	1.8000	0.7200
16	1.8500	0.7250
17	1.9000	0.7300
18	1.9500	0.7350
19	2.1000	0.7450
20	2.2500	0.7500
21	2.4000	0.7500
22	2.6000	0.7400
23	2.7500	0.7300
24	2.9000	0.7250
25	3.1600	1.0000
26	3.0600	1.2500
27	2.9500	1.5500
28	2.7500	1.6500
29	2.4000	1.8200
30	2.2500	1.8000
31	2.1000	1.7500
32	2.0500	1.7450
33	2.0000	1.7250
34	1.9500	1.7100
35	1.8500	1.6900
36	1.7500	1.6500
37	1.6500	1.6100

BOUNDARY CONDITIONS

NODE	CODE	PRESCRIBED VALUE
1	2	0.0
2	2	0.0
3	2	0.0
4	2	0.0
5	2	0.0
6	2	0.0
7	2	0.0
8	2	0.0
9	2	0.0
10	2	0.0
11	1	0.0
12	1	0.0
13	1	0.0
14	1	0.0
15	1	0.0
16	1	0.0
17	1	0.0
18	1	0.0
19	1	0.0
20	1	0.0
21	1	0.0
22	1	0.0
23	1	0.0
24	1	0.0
25	1	0.0
26	1	0.0
27	1	0.0
28	1	0.0
29	1	0.0
30	1	0.0
31	1	0.0
32	1	0.0
33	1	0.0
34	1	0.0
35	1	0.0
36	1	0.0
37	1	0.0

COORDINATES OF SOURCES AND SINKS AND THEIR STRENGTHS

X(INCH)	Y(INCH)	RATE(BBL/D)
1.9000	0.9500	200.0000
1.8000	1.3500	-69.1000
2.0500	1.5500	-69.1000
2.1500	1.2500	-69.1000
2.3500	1.0500	-69.1000
2.5750	1.3000	-69.1000

THE AVERAGE RADIUS OF THE WELLS(INCH) IS= 0.500
THE CHARACTERISTIC PRESSURE(PSI) IS= 85.0000

RESULTS

BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE(PSI)	NORMAL GRADIENT
0.1425E+01	0.1515E+01	0.8076E+02	0.0
0.1300E+01	0.1460E+01	0.8205E+02	0.0
0.1225E+01	0.1420E+01	0.8266E+02	0.0
0.1175E+01	0.1390E+01	0.8327E+02	0.0
0.1100E+01	0.1352E+01	0.8402E+02	0.0
0.1000E+01	0.1292E+01	0.8501E+02	0.0
0.9250E+00	0.1255E+01	0.8560E+02	0.0
0.8750E+00	0.1225E+01	0.8568E+02	0.0
0.8250E+00	0.1185E+01	0.8623E+02	0.0
0.7750E+00	0.1160E+01	0.8675E+02	0.0
0.7000E+00	0.1115E+01	0.8713E+02	0.0
0.6000E+00	0.1050E+01	0.8695E+02	0.0
0.5250E+00	0.1005E+01	0.8648E+02	0.0
0.4750E+00	0.9700E+00	0.8610E+02	0.0
0.4250E+00	0.9350E+00	0.8569E+02	0.0
0.3500E+00	0.8850E+00	0.8537E+02	0.0
0.2000E+00	0.7900E+00	0.8509E+02	0.0
0.1000E+00	0.6150E+00	0.8500E+02	-0.3533E-03
0.1000E+00	0.3250E+00	0.8500E+02	0.3379E-03
0.1750E+00	0.1750E+00	0.8468E+02	0.0
0.3250E+00	0.2250E+00	0.8510E+02	0.0
0.4750E+00	0.2650E+00	0.8524E+02	0.0
0.6000E+00	0.2900E+00	0.8548E+02	0.0
0.6750E+00	0.3125E+00	0.8557E+02	0.0
0.7250E+00	0.3275E+00	0.8566E+02	0.0
0.7750E+00	0.3350E+00	0.8590E+02	0.0
0.8250E+00	0.3450E+00	0.8645E+02	0.0
0.8750E+00	0.3750E+00	0.8700E+02	0.0
0.9500E+00	0.4250E+00	0.8834E+02	0.0
0.1050E+01	0.4875E+00	0.8946E+02	0.0
0.1175E+01	0.5625E+00	0.8972E+02	0.0
0.1325E+01	0.6250E+00	0.8794E+02	0.0
0.1475E+01	0.6650E+00	0.8608E+02	0.0
0.1550E+01	0.7150E+00	0.8560E+02	-0.1048E-01
0.1545E+01	0.8000E+00	0.8519E+02	-0.6458E-02
0.1520E+01	0.9000E+00	0.8469E+02	-0.8446E-02
0.1500E+01	0.1000E+01	0.8383E+02	-0.4569E-02
0.1490E+01	0.1100E+01	0.8302E+02	-0.7717E-02
0.1470E+01	0.1200E+01	0.8234E+02	-0.8584E-02
0.1460E+01	0.1300E+01	0.8167E+02	-0.8253E-02
0.1470E+01	0.1400E+01	0.8094E+02	-0.7529E-02
0.1490E+01	0.1475E+01	0.8047E+02	-0.6497E-02
0.1500E+01	0.1525E+01	0.8102E+02	-0.2877E-01

RESULTS

BOUNDARY NODES

X(INCH)	Y(INCH)	PRESSURE(PSI)	NORMAL GRADIENT
0.1500E+01	0.1525E+01	0.8102E+02	0.5754E-01
0.1490E+01	0.1475E+01	0.8047E+02	0.1299E-01
0.1470E+01	0.1400E+01	0.8094E+02	0.1506E-01
0.1460E+01	0.1300E+01	0.8167E+02	0.1651E-01
0.1470E+01	0.1200E+01	0.8234E+02	0.1717E-01
0.1490E+01	0.1100E+01	0.8302E+02	0.1543E-01
0.1500E+01	0.1000E+01	0.8383E+02	0.9137E-02
0.1520E+01	0.9000E+00	0.8469E+02	0.1689E-01
0.1545E+01	0.8000E+00	0.8519E+02	0.1292E-01
0.1550E+01	0.7150E+00	0.8560E+02	0.2096E-01
0.1600E+01	0.6850E+00	0.8512E+02	0.0
0.1675E+01	0.6950E+00	0.8438E+02	0.0
0.1725E+01	0.7050E+00	0.8395E+02	0.0
0.1775E+01	0.7150E+00	0.8353E+02	0.0
0.1825E+01	0.7225E+00	0.8306E+02	0.0
0.1875E+01	0.7275E+00	0.8251E+02	0.0
0.1925E+01	0.7325E+00	0.8177E+02	0.0
0.2025E+01	0.7400E+00	0.7962E+02	0.0
0.2175E+01	0.7475E+00	0.7595E+02	0.0
0.2325E+01	0.7500E+00	0.7303E+02	0.0
0.2500E+01	0.7450E+00	0.7103E+02	0.0
0.2675E+01	0.7350E+00	0.7012E+02	0.0
0.2825E+01	0.7275E+00	0.6975E+02	0.0
0.3030E+01	0.8625E+00	0.6950E+02	0.0
0.3110E+01	0.1125E+01	0.6931E+02	0.0
0.3005E+01	0.1400E+01	0.6897E+02	0.0
0.2850E+01	0.1600E+01	0.6892E+02	0.0
0.2575E+01	0.1735E+01	0.6933E+02	0.0
0.2325E+01	0.1810E+01	0.6975E+02	0.0
0.2175E+01	0.1775E+01	0.7019E+02	0.0
0.2075E+01	0.1747E+01	0.7075E+02	0.0
0.2025E+01	0.1735E+01	0.7096E+02	0.0
0.1975E+01	0.1717E+01	0.7149E+02	0.0
0.1900E+01	0.1700E+01	0.7255E+02	0.0
0.1800E+01	0.1670E+01	0.7403E+02	0.0
0.1700E+01	0.1630E+01	0.7589E+02	0.0
0.1575E+01	0.1580E+01	0.7864E+02	0.0

PRODUCER NUMBER 1

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1			0.2800E+02
2			0.2800E+02
3			0.2800E+02

PRODUCER NUMBER 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1			0.2800E+02
2			0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3381E+03	0.1795E+02	0.2800E+02
2	0.3816E+03	0.2111E+02	0.2800E+02

PRODUCER NUMBER 1

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.5533E+03	0.4991E+02	0.2800E+02
2	0.4250E+03	0.2343E+02	0.2800E+02
3	0.6566E+03	0.6611E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4447E+03	0.2795E+02	0.2800E+02
2	0.3447E+03	0.1906E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3381E+03	0.1795E+02	0.2800E+02
2	0.3816E+03	0.2111E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2697E+03	0.1331E+02	0.2800E+02
2	0.2697E+03	0.1475E+02	0.2800E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.3664E+03	0.1586E+02	0.2800E+02
2	0.3316E+03	0.1379E+02	0.2800E+02

PRODUCER NUMBER 7

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2914E+03	0.1262E+02	0.2800E+02
2	0.2947E+03	0.1212E+02	0.2800E+02

PRODUCER NUMBER 2

NUMBER OF STREAMLINES= 2

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4513E+03	0.6006E+02	0.2800E+02
1	0.2164E+03	0.3008E+02	0.2800E+02
2	0.4784E+03	0.5833E+02	0.2800E+02
2	0.2480E+03	0.3138E+02	0.2800E+02

PRODUCER NUMBER 3

NUMBER OF STREAMLINES= 3

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.5914E+03	0.1337E+02	0.1000E+02
2	0.6197E+03	0.1481E+02	0.1000E+02
1	0.3764E+03	0.6182E+02	0.2800E+02
1	0.5414E+03	0.4388E+02	0.2800E+02

PRODUCER NUMBER 4

NUMBER OF STREAMLINES= 4

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.2664E+03	0.6399E+01	0.1000E+02
2	0.2881E+03	0.6867E+01	0.1000E+02
3	0.3631E+03	0.8268E+01	0.1000E+02
4	0.2881E+03	0.6640E+01	0.1000E+02

PRODUCER NUMBER 5

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.4783E+03	0.1111E+02	0.1000E+02
2	0.3848E+03	0.8507E+01	0.1000E+02
3	0.3348E+03	0.7707E+01	0.1000E+02
4	0.3316E+03	0.7468E+01	0.1000E+02
5	0.3848E+03	0.9714E+01	0.1000E+02

PRODUCER NUMBER 6

NUMBER OF STREAMLINES= 5

STREAMLINE NUMBER	TOTAL LENGTH	AVERAGE WIDTH	INJ. RATE
1	0.7033E+03	0.2165E+02	0.1000E+02
2	0.8348E+03	0.2773E+02	0.1000E+02
3	0.7197E+03	0.1985E+02	0.1000E+02
4	0.5381E+03	0.1601E+02	0.1000E+02
5	0.5914E+03	0.2086E+02	0.1000E+02

NO STREAMLINE FROM REGION 2 CROSSES INTO REGION 1

APPENDIX N

RESULTS OF STEAMFLOOD PREDICTION FOR SHIELLS CANYON
FIELD AS A PIECEWISE HOMOGENEOUS RESERVOIR.
PART OF ITS BOUNDARY AT CONSTANT
PRESSURE, THE REMAINDER SEALED

APPENDIX N

STEAMFLOOD PREDICTION OF SHIELLS CANYON(203) FIELD

METHOD: STREAMLINE/STREAMTUBE METHOD USING (EEM)TECHNIQUE

ANALYSIS AS A PIECE-WISE HOMOGENEOUS RESERVOIR

HAVING TWO REGIONS OF UNEQUAL PERMEABILITY

BOUNDARY CONDITION: PART OF THE BOUNDARY AT CONSTANT

PRESSURE, THE REMAINDER IS SEALED.

PRODUCTION RATES: ASSIGNED EQUAL IN ENTIRE REGION

DATA

TOTAL NUMBER OF PRODUCERS=	11
STEAM TEMPERATURE (DEG. F)=	430.0000
STEAM QUALITY=	0.7000
CONVERGENCE LIMIT=	0.0010
RESERVOIR THICKNESS (FT)=	160.0000
STEAM PRESSURE (PSI)=	344.0000
INITIAL RESERVOIR TEMPERATURE(DEG F)=	105.0000

FLUID PROPERTIES

	WATER	OIL	STEAM
	-----	-----	-----
DENSITY(LB/CU FT) AT STD. TEMP.	62.4000	55.0000	
DENSITY(LB/CU FT) AT STEAM TEMP.	52.3800	48.7000	0.7431
SPECIFIC HEAT(BTU/LB *F)	1.0000	0.4880	1.0000
LATENT HEAT(BTU/LB)			789.5100

PROPERTIES OF THE CAP AND BASE ROCK

DENSITY(LB/CU. FT)	149.0000
SPECIFIC HEAT(BTU/LB-*F)	0.2130
THERM. COND.(BTU/HR-FT-*F)	1.1000

RESERVOIR ROCK PROPERTIES

	REGION 1	REGION 2
	-----	-----
POROSITY	0.2050	0.1800
DEN.(LB/CU FT)	165.0000	165.0000
SPEC. HEAT(BTU/LB*F)	0.2000	0.2000
PERMEABILITY(MD)	140.0000	70.0000

INITIAL FLUID SATURATIONS

	REGION 1	REGION 2
	-----	-----
WATER	0.5500	0.5500
OIL	0.4500	0.4500
STEAM	0.0	0.0

AVERAGE RESIDUAL SATURATIONS BEHIND THE FRONT

	REGION 1	REGION 2
	-----	-----
WATER	0.5800	0.5800
OIL	0.1800	0.1800
STEAM	0.2400	0.2400

INPUT DATA FROM STREAMLINE PROGRAM

WELL NO.	S/L NO.	CODE	REGION 1			REGION 2		
			-----	-----	-----	-----	-----	-----
			LENGTH	WIDTH	RATE	LENGTH	WIDTH	RATE
1	1	1	553.30	49.91	28.00	0.0	0.0	0.0
1	2	1	425.00	23.40	28.00	0.0	0.0	0.0
1	3	1	656.60	66.11	28.00	0.0	0.0	0.0
2	1	1	444.70	27.95	28.00	0.0	0.0	0.0
2	2	1	344.70	19.06	28.00	0.0	0.0	0.0
3	1	1	338.10	17.95	28.00	0.0	0.0	0.0
3	2	1	381.60	21.11	28.00	0.0	0.0	0.0
4	1	1	269.70	13.31	28.00	0.0	0.0	0.0
4	2	1	269.70	14.75	28.00	0.0	0.0	0.0
5	1	1	366.40	15.86	28.00	0.0	0.0	0.0
5	2	1	331.60	13.79	28.00	0.0	0.0	0.0
6	1	1	291.40	12.62	28.00	0.0	0.0	0.0
6	2	1	294.70	12.12	28.00	0.0	0.0	0.0
7	1	2	451.30	60.06	28.00	216.40	30.08	28.00
7	2	2	478.40	58.33	28.00	248.00	31.38	28.00
8	1	2	376.40	61.82	28.00	541.40	43.88	28.00
8	2	1	591.40	13.37	10.00	0.0	0.0	0.0
8	3	1	619.70	14.81	10.00	0.0	0.0	0.0
9	1	1	266.40	6.40	10.00	0.0	0.0	0.0
9	2	1	288.10	6.87	10.00	0.0	0.0	0.0
9	3	1	363.10	8.27	10.00	0.0	0.0	0.0
9	4	1	288.10	6.64	10.00	0.0	0.0	0.0
10	1	1	478.30	11.11	10.00	0.0	0.0	0.0
10	2	1	384.80	8.51	10.00	0.0	0.0	0.0
10	3	1	334.80	7.71	10.00	0.0	0.0	0.0
10	4	1	331.60	7.47	10.00	0.0	0.0	0.0
10	5	1	384.80	9.71	10.00	0.0	0.0	0.0
11	1	1	703.30	21.65	10.00	0.0	0.0	0.0
11	2	1	834.80	27.73	10.00	0.0	0.0	0.0
11	3	1	719.70	19.85	10.00	0.0	0.0	0.0
11	4	1	538.10	16.01	10.00	0.0	0.0	0.0
11	5	1	591.40	20.86	10.00	0.0	0.0	0.0

CALCULATED BREAKTHROUGH TIMES(HOURS) FOR EACH STREAMTUBE

PRUD. NO.	STREAMLINE NO.	CCODE	ENDTIME(1)	ENDTIME(2)
1	1	1	18260.6484	0.0
1	2	1	6048.5781	0.0
1	3	1	30286.2461	0.0
2	1	1	7671.5547	0.0
2	2	1	3904.0466	0.0
3	1	1	3592.2112	0.0
3	2	1	4838.9414	0.0
4	1	1	2078.0183	0.0
4	2	1	2311.8730	0.0
5	1	1	3432.4089	0.0
5	2	1	2672.3630	0.0
6	1	1	2130.7292	0.0
6	2	1	2067.1299	0.0
7	1	2	17888.3867	21748.8750
7	2	2	18472.7227	23131.5273
8	1	2	15123.5234	30573.1289
8	2	1	14321.5000	0.0
8	3	1	16866.7812	0.0
9	1	1	2794.5066	0.0
9	2	1	3264.7048	0.0
9	3	1	5061.4687	0.0
9	4	1	3151.8491	0.0
10	1	1	9305.1289	0.0
10	2	1	5547.0469	0.0
10	3	1	4314.2812	0.0
10	4	1	4131.6172	0.0
10	5	1	6386.0742	0.0
11	1	1	29675.5391	0.0
11	2	1	47986.7930	0.0
11	3	1	27613.4141	0.0
11	4	1	15732.7773	0.0
11	5	1	23419.4766	0.0

REAL TIME(HOURS)= 6000.0000

DIMENSIONLESS TIME= 0.0325

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.1557E+05		
1	2	1	0.1557E+05		
1	3	1	0.1557E+05		
				0.4670E+05	
2	1	1	0.1557E+05		
2	2	1	0.1036E+05		
				0.2593E+05	
3	1	1	0.9572E+04		
3	2	1	0.1271E+05		
				0.2228E+05	
4	1	1	0.5662E+04		
4	2	1	0.6274E+04		
				0.1194E+05	
5	1	1	0.9165E+04		
5	2	1	0.7212E+04		
				0.1638E+05	
6	1	1	0.5800E+04		
6	2	1	0.5633E+04		
				0.1143E+05	
7	1	2	0.1557E+05		
7	2	2	0.1557E+05		
				0.3113E+05	
8	1	2	0.1557E+05		
8	2	1	0.5560E+04		
8	3	1	0.5560E+04		
				0.2669E+05	
9	1	1	0.2689E+04		
9	2	1	0.3120E+04		
9	3	1	0.4735E+04		
9	4	1	0.3017E+04		
				0.1356E+05	
10	1	1	0.5560E+04		
10	2	1	0.5163E+04		
10	3	1	0.4070E+04		
10	4	1	0.3906E+04		
10	5	1	0.5560E+04		
				0.2426E+05	
11	1	1	0.5560E+04		
11	2	1	0.5560E+04		
11	3	1	0.5560E+04		
11	4	1	0.5560E+04		
11	5	1	0.5560E+04		
				0.2780E+05	
					0.2581E+06

REAL TIME(HOURS)= 12000.0000

DIMENSIONLESS TIME= 0.0650

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.2969E+05		
1	2	1	0.1569E+05		
1	3	1	0.2969E+05		
				0.7506E+05	
2	1	1	0.1960E+05		
2	2	1	0.1036E+05		
				0.2997E+05	
3	1	1	0.9572E+04		
3	2	1	0.1271E+05		
				0.2228E+05	
4	1	1	0.5662E+04		
4	2	1	0.6274E+04		
				0.1194E+05	
5	1	1	0.9165E+04		
5	2	1	0.7212E+04		
				0.1638E+05	
6	1	1	0.5800E+04		
6	2	1	0.5633E+04		
				0.1143E+05	
7	1	2	0.2969E+05		
7	2	2	0.2969E+05		
				0.5938E+05	
8	1	2	0.2969E+05		
8	2	1	0.1060E+05		
8	3	1	0.1060E+05		
				0.5090E+05	
9	1	1	0.2689E+04		
9	2	1	0.3120E+04		
9	3	1	0.4735E+04		
9	4	1	0.3017E+04		
				0.1356E+05	
10	1	1	0.8381E+04		
10	2	1	0.5163E+04		
10	3	1	0.4070E+04		
10	4	1	0.3906E+04		
10	5	1	0.5896E+04		
				0.2742E+05	
11	1	1	0.1060E+05		
11	2	1	0.1060E+05		
11	3	1	0.1060E+05		
11	4	1	0.1060E+05		
11	5	1	0.1060E+05		
				0.5302E+05	
					0.3713E+06

REAL TIME(HOURS)= 18000.0000

DIMENSIONLESS TIME= 0.0975

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4299E+05		
1	2	1	0.1569E+05		
1	3	1	0.4299E+05		
				0.1017E+06	
2	1	1	0.1960E+05		
2	2	1	0.1036E+05		
				0.2997E+05	
3	1	1	0.9572E+04		
3	2	1	0.1271E+05		
				0.2228E+05	
4	1	1	0.5662E+04		
4	2	1	0.6274E+04		
				0.1194E+05	
5	1	1	0.9165E+04		
5	2	1	0.7212E+04		
				0.1638E+05	
6	1	1	0.5800E+04		
6	2	1	0.5633E+04		
				0.1143E+05	
7	1	2	0.4296E+05		
7	2	2	0.4299E+05		
				0.8595E+05	
8	1	2	0.4223E+05		
8	2	1	0.1247E+05		
8	3	1	0.1448E+05		
				0.6918E+05	
9	1	1	0.2689E+04		
9	2	1	0.3120E+04		
9	3	1	0.4735E+04		
9	4	1	0.3017E+04		
				0.1356E+05	
10	1	1	0.8381E+04		
10	2	1	0.5163E+04		
10	3	1	0.4070E+04		
10	4	1	0.3906E+04		
10	5	1	0.5896E+04		
				0.2742E+05	
11	1	1	0.1535E+05		
11	2	1	0.1535E+05		
11	3	1	0.1535E+05		
11	4	1	0.1359E+05		
11	5	1	0.1535E+05		
				0.7500E+05	
					0.4648E+06

REAL TIME(HOURS)= 24000.0000

DIMENSIONLESS TIME= 0.1300

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4355E+05		
1	2	1	0.1569E+05		
1	3	1	0.5569E+05		
				0.1149E+06	
2	1	1	0.1960E+05		
2	2	1	0.1036E+05		
				0.2997E+05	
3	1	1	0.9572E+04		
3	2	1	0.1271E+05		
				0.2228E+05	
4	1	1	0.5662E+04		
4	2	1	0.6274E+04		
				0.1194E+05	
5	1	1	0.9165E+04		
5	2	1	0.7212E+04		
				0.1638E+05	
6	1	1	0.5800E+04		
6	2	1	0.5633E+04		
				0.1143E+05	
7	1	2	0.5176E+05		
7	2	2	0.5479E+05		
				0.1066E+06	
8	1	2	0.5339E+05		
8	2	1	0.1247E+05		
8	3	1	0.1448E+05		
				0.8034E+05	
9	1	1	0.2689E+04		
9	2	1	0.3120E+04		
9	3	1	0.4735E+04		
9	4	1	0.3017E+04		
				0.1356E+05	
10	1	1	0.8381E+04		
10	2	1	0.5163E+04		
10	3	1	0.4070E+04		
10	4	1	0.3906E+04		
10	5	1	0.5896E+04		
				0.2742E+05	
11	1	1	0.1989E+05		
11	2	1	0.1989E+05		
11	3	1	0.1989E+05		
11	4	1	0.1359E+05		
11	5	1	0.1946E+05		
				0.9271E+05	
					0.5275E+06

REAL TIME(HOURS)= 30000.0000

DIMENSIONLESS TIME= 0.1625

WELL	STR/L	NO. OF REGIONS	RECOVERY	WELL TOTAL	RESERVOIR TOTAL
1	1	1	0.4355E+05		
1	2	1	0.1569E+05		
1	3	1	0.6789E+05		
				0.1271E+06	
2	1	1	0.1960E+05		
2	2	1	0.1036E+05		
				0.2997E+05	
3	1	1	0.9572E+04		
3	2	1	0.1271E+05		
				0.2228E+05	
4	1	1	0.5662E+04		
4	2	1	0.6274E+04		
				0.1194E+05	
5	1	1	0.9165E+04		
5	2	1	0.7212E+04		
				0.1638E+05	
6	1	1	0.5800E+04		
6	2	1	0.5633E+04		
				0.1143E+05	
7	1	2	0.5176E+05		
7	2	2	0.5479E+05		
				0.1066E+06	
8	1	2	0.6412E+05		
8	2	1	0.1247E+05		
8	3	1	0.1448E+05		
				0.9106E+05	
9	1	1	0.2689E+04		
9	2	1	0.3120E+04		
9	3	1	0.4735E+04		
9	4	1	0.3017E+04		
				0.1356E+05	
10	1	1	0.8381E+04		
10	2	1	0.5163E+04		
10	3	1	0.4070E+04		
10	4	1	0.3906E+04		
10	5	1	0.5896E+04		
				0.2742E+05	
11	1	1	0.2402E+05		
11	2	1	0.2425E+05		
11	3	1	0.2253E+05		
11	4	1	0.1359E+05		
11	5	1	0.1946E+05		
				0.1038E+06	
					0.5616E+06

PREDICTED RECOVERY

TIME(DAYS)	RECOVERY(BBLS)
62.5000	0.1598E+05
125.0000	0.2940E+05
187.5000	0.3882E+05
250.0000	0.4596E+05
312.5000	0.5177E+05
375.0000	0.5682E+05
437.5000	0.6153E+05
500.0000	0.6613E+05
562.5000	0.7066E+05
625.0000	0.7503E+05
687.5000	0.7906E+05
750.0000	0.8277E+05
812.5000	0.8580E+05
875.0000	0.8868E+05
937.5000	0.9160E+05
1000.0000	0.9394E+05
1062.5000	0.9557E+05
1125.0000	0.9718E+05
1187.5000	0.9866E+05
1250.0000	0.1000E+06

APPENDIX O

LISTING OF COMPUTER PROGRAMS

C*****
 C PROGRAM FOR THE SOLUTION OF TWO DIMENSIONAL POTENTIAL
 C PROBLEMS BY THE BOUNDARY ELEMENT METHOD. APPLICABLE TO
 C HOMOGENEOUS AS WELL AS 2-REGION PIECEWISE HOMOGENEOUS RESERVOIR
 C NEUMANN AND DIRICHLET CONDITIONS SPECIFIED ON THE BOUNDARIES
 C*****
 C

```

0001      DIMENSION X1(120),Y1(120),XM1(120),YM1(120)
0002      DIMENSION FI1(120),DFI1(120),KUDE1(120),CX1(120)
0003      DIMENSION SOL1(120),H1(120,120),A(120,120),B(120)
0004      DIMENSION G1(120,120),CY1(120),SUM1(120),CX2(120)
0005      DIMENSION RO(120),X2(120),Y2(120),FI2(120),DFI2(120)
0006      DIMENSION CY2(120),SOL2(120),XSS1(120),YSS1(120)
0007      DIMENSION XSS2(120),YSS2(120),QSS2(120),SUM2(120)
0008      DIMENSION XM2(120),YM2(120),XI1(120),YI1(120)
0009      DIMENSION QSS1(120),KODE2(120),DYBLE1(120)
0010      DIMENSION DGDY1(120),DGDY1(120),DHDX1(120),DHDY1(120)
0011      DIMENSION AQSS2(120),DXBLE1(120),DXBLE2(120),DISMID(120)
0012      DIMENSION DGDY2(120),DGDY2(120),DHDX2(120),DHDY2(120)
0013      DIMENSION VT1(120),YI2(120),DYBLE2(120),H2(120,120)
0014      DIMENSION VX2(120),VY2(120),VT2(120),VX1(120),VY1(120)
0015      DIMENSION G2(120,120),XI2(120)
0016      DIMENSION DPHIDX(120),DPHIDY(120),VLCPI1(120),VLCPI2(120)
0017      DIMENSION DPDX(120),DPDY(120),XQ1(120),YQ1(120),XQ2(120)
0018      DIMENSION ANSL1(120),ANSL2(120),YQ2(120),A1(120,120)
0019      DIMENSION AQSS1(120),F(120),INFO(18),DIST(120),ELMNT1(120)
0020      DIMENSION SINK(120),STRML(120),STRMW(120),AVWIDT(120)
0021      DIMENSION MM(120),A2(120),LENTH(120,120),WIDTH(120,120)
0022      DIMENSION ATETA1(120),ATETA2(120),ADFI1(120),ADFI2(120)
0023      DIMENSION XIF(120),YIF(120),ELMNT2(120),VEL1(120)
0024      DIMENSION VEL2(120),BETA(120),FIPLT(120),DFIPLT(120)
0025      DIMENSION QSSS1(120),QSSS2(120)
0026      DIMENSION XAVE1(120),XAVE2(120),YAVE1(120),YAVE2(120)
0027      DIMENSION OLDL1(120),OLDL2(120),OLDW1(120),OLDW2(120)
0028      COMMON /BLK1/LEC,IMP,VICSTM,PI
0029      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0030      COMMON /BLK3/PRESSF,DPRESSF,PORFAC,SCLFAC
0031      INTEGER TITLE(18),TITLE1(18),TITLE2(18),TITLE3(18)

```

C

C.....

C CONVERSION FACTORS

C DISFAC=CONVERSION FROM (FT) TO (M)

C TRFAC=CONVERSION FROM (MD) TO (SQ M)

C RATEF=CONVERSION FROM (BBL/DAY) TO (CU. M/SEC)

C SCALE=SCALE FACTOR OF THE FIELD

C PRESSF=PRESSURE CONVERSION FROM (LBF/SQ IN) TO (NEWTON/SQ M)

C DPRESF=PRESSURE GRADIENT CONVERSION FROM ((LB/SQ IN)/FT)

C TO ((NEWTON/SQ M)/M)

C PORFAC=POROSITY CONVERSION FACTOR=1

C.....

C NTR=TOTAL NUMBER OF REGIONS

```

0032      NTR=2
0033      LEC=5
0034      IMP=6
0035      PI=3.1415926
0036      VICSTM=1.0

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C

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0037      DISFAC=.3048

```

```
0038      VICFAC=0.001
0039      TRFAC=9.869E-16
0040      RATEF=1.8401E-6
0041      SCALE=738.46
0042      PRESSF=6894.76
0043      DPRESF=PRESSF/DISFAC
0044      PORFAC=1.0

C
0045      COUNT=0.
0046      SCLFAC=DISFAC*SCALE

C
C WRITE HEADINGS OR TITLES
C
0047      WRITE(IMP,100)
0048      100  FORMAT(1H1,/,/,80('*'))
C READ THE NAME OF THE JOB
0049      READ(LEC,150)TITLE
0050      READ(LEC,150)TITLE1
0051      READ(LEC,150)TITLE2
0052      READ(LEC,150)TITLE3
0053      150  FORMAT(18A4)
0054      WRITE(IMP,250)TITLE
0055      WRITE(IMP,250)TITLE1
0056      WRITE(IMP,250)TITLE2
0057      WRITE(IMP,250)TITLE3
0058      250  FORMAT(/15X,18A4)

C
C READ INPUT DATA FOR FIRST REGION
C
0059      7    COUNT=COUNT+1.
0060      WRITE(IMP,153)COUNT
0061      153  FORMAT(/,30X,'REGION',F4.1,/)

C
0062      CALL INPUT(CX1,CY1,X1,Y1,KODE1,FI1,XSS1,YSS1,QSS1,NSS1,
& N1,NIF,L1,T1,THICK1,POR1,RI1,TOPP1)

C
C...CONVERT TO SI UNITS
C
0063      CALL CONVR5(CX1,CY1,X1,Y1,KODE1,FI1,XSS1,YSS1,QSS1,NSS1,
& N1,L1,T1,THICK1,POR1,RI1,TOPP1)

C
C CALCULATE THE ELEMENTS OF THE H AND G MATRICES AND ARRANGE
C THEM SO AS TO PUT ALL UNKNOWN IN G AND KNOWN IN H.
C THE ELEMENTS OF H AND THE SOURCE VECTOR FORM THE RHS VECTOR DFI
C
0064      CALL MATRX(X1,Y1,XM1,YM1,G1,H1,FI1,DFI1,KODE1,XSS1,
& YSS1,QSS1,NSS1,IJ1,NIF,SUM1,T1,N1,L1,CX1,CY1)

C
C..IF THERE ARE TWO REGIONS REPEAT THE ABOVE STEPS FOR SECOND
C..REGION
C
0065      IF(NTR .LT. 2)GO TO 155
0066      COUNT=COUNT+1.
0067      WRITE(IMP,154)COUNT
0068      154  FORMAT(1H1,/,/,30X,'REGION',F4.1,/)

C
0069      CALL INPUT(CX2,CY2,X2,Y2,KODE2,FI2,XSS2,YSS2,QSS2,NSS2,
& N2,NIF,L2,T2,THICK2,POR2,RI2,TOPP2)
```

```

C
0070      CALL CONVR5(CX2,CY2,X2,Y2,KODE2,FI2,XSS2,YSS2,QSS2,NSS2,
&      N2,L2,T2,THICK2,POR2,RI2,TOPP2)
C
0071      CALL MATRX(X2,Y2,XM2,YM2,G2,H2,FI2,DFI2,KODE2,XSS2,
&      YSS2,QSS2,NSS2,IJ2,NIF,SUM2,T2,N2,L2,CX2,CY2)
C
C ASSEMBLE THE G MATRICES(CONSIDERING CONTINUITY AND
C COMPATIBILITY ON THE INTERFACE) INTO THE GLOBAL MATRIX (A)
C THE RHS (B) NOW CONSISTS OF THE SUM OF THE DFI VECTOR AND THE
C VECTOR (SUM) DUE TO THE SOURCE/SINK TERMS.
C...NOTE ..IF THE BOUNDARY IS COMPLETELY CLOSED, DFI WOULD EQUAL
C ZERO
C
0072      CALL ASEMBL(G1,G2,DFI1,DFI2,N1,N2,A,B,NAA,NIF,NNN,
&      NOM,NNI,NNM,T1,T2)
0073      GO TO 156
C
C STORE THE A MATRIX AND B VECTOR FOR FUTURE CHECK OF THE SOLUTIO
C SOLVE AND CHECK THE SOLUTION FOR A HOMOGENEOUS SINGLE REGION
0074      155 DO 158 I=1,N1
0075           A2(I)=DFI1(I)
0076           DO 157 J=1,N1
0077              A1(I,J)=G1(I,J)
0078      157 CONTINUE
0079      158 CONTINUE
0080           N2=0
0081           CALL SLNPD(G1,DFI1,D,N1,N2,NJ)
0082           WRITE(IMP,159)
0083      159 FORMAT(1H1,///20X,'CHECKING THE SOLUTION',///,
&      9X,'RHS VECTOR',10X,'VALUE FROM SOLUTION VECTOR',///)
0084           DO 162 I=1,NJ
0085              F(I)=0.
0086           DO 161 J=1,NJ
0087              F(I)=F(I)+(A1(I,J)*DFI1(J))
0088      161 CONTINUE
0089           WRITE(IMP,163)A2(I),F(I)
0090      163 FORMAT(5X,E14.3,14X,E14.3)
0091      162 CONTINUE
0092           GO TO 164
C
C SOLVE AND CHECK THE SOLUTION IF A 2-REGION COMPOSITE
C
0093      156 DO 76 I=1,NAA
0094           A2(I)=B(I)
0095           DO 75 J=1,NAA
0096              A1(I,J)=A(I,J)
0097      75 CONTINUE
0098      76 CONTINUE
C
C SOLVE THE SYSTEM
C
0099      CALL SLNPD(A,B,D,N1,N2,NJ)
C
C CHECK THE SOLUTION
C
0100      WRITE(IMP,108)
0101      108 FORMAT(1H1,///20X,'CHECKING THE SOLUTION VECTOR'///,

```

```
      & 9X,'RHS VECTOR',10X,'VALUE FROM SOLUTION VECTOR',//)
0102      DO 72 I=1,NAA
0103      F(I)=0.
0104      DO 73 J=1,NAA
0105      F(I)=F(I)+(A1(I,J)*B(J))
0106      73 CONTINUE
0107      WRITE(IMP,74)A2(I),F(I)
0108      74 FORMAT(5X,E14.3,13X,E14.3)
0109      72 CONTINUE
C
C SPLIT THE SOLUTION VECTOR INTO ITS COMPONENT REGIONAL VALUES
C
0110      CALL SPLIT(N1,N2,NIF,NNN,NOM,NNI,NNM,NAA,KODE1,KODE2,
      & FI1,FI2,DFI1,DFI2,B,T1,T2)
0111      164 CONTINUE
0112      IF(NTR .GT. 1)GO TO 168
C
C REORDER THE FI AND DFI ARRAYS BY PUTTING ALL POTENTIALS IN FI
C AND ALL DERIVATIVES IN DFI
C
0113      DO 169 I=1,N1
0114      IF(KODE1(I))20,20,10
0115      10 CH=FI1(I)
0116      FI1(I)=DFI1(I)
0117      DFI1(I)=CH
0118      20 CONTINUE
0119      169 CONTINUE
0120      168 CONTINUE
C
C PRINT OUT THE BOUNDARY VALUES AND COMPUTE THE POTENTIAL
C VALUES AT DESIGNATED INTERNAL POINTS IF ANY FOR THE FIRST REGI
C
0121      CALL OUTPUT(X1,Y1,XM1,YM1,FI1,DFI1,CX1,CY1,SOL1,N1,
      & XSS1,YSS1,L1,NSS1,QSS1,TOPP1,THICK1,T1,QSSS1)
C
0122      IF(L1 .LE. 0)GO TO 173
C
C COMPUTE THE POTENTIAL VALUES AT INTERNAL POINTS FOR
C THE FIRST REGION
C
0123      CALL INTER(FI1,DFI1,L1,N1,CX1,CY1,X1,Y1,
      & XSS1,YSS1,QSS1,NSS1,SOL1)
C
0124      173 IF(NTR .LT. 2)GO TO 165
C
C PRINT OUT THE BOUNDARY VALUES FOR THE SECOND REGION AND
C COMPUTE THE POTENTIAL VALUES AT DESIGNATED INTERNAL POINTS
C IF ANY.
C
0125      CALL OUTPUT(X2,Y2,XM2,YM2,FI2,DFI2,CX2,CY2,SOL2,N2,
      & XSS2,YSS2,L2,NSS2,QSS2,TOPP2,THICK2,T2,QSSS2)
C
0126      IF(L2 .LE. 0)GO TO 165
C
C COMPUTE THE POTENTIAL VALUES AT INTERNAL POINTS FOR THE
C SECOND REGION
C
0127      CALL INTER(FI2,DFI2,L2,N2,CX2,CY2,X2,Y2,
```

```
      & XSS2,YSS2,QSS2,NSS2,SOL2)
C PLOT BOUNDARIES AND WELL LOCATIONS
C
0128      165      CONTINUE
0129              CALL BDRY(X1,Y1,XSS1,YSS1,NSS1,QSSS1,N1)
0130              IF(NTR .LT. 2)GO TO 166
0131              CALL BDRY(X2,Y2,XSS2,YSS2,NSS2,QSSS2,N2)
C
C GENERATE THE STREAMLINES
C
0132      166      CONTINUE
0133              CALL STRM(NSS1,QSS1,FI1,DFI1,X1,Y1,XSS1,YSS1,
      & N1,T1,KODE1,IP2,XO2,YO2,XM1,YM1,AQSS2,ANSL2,MM,
      & ATETA1,ADFI1,ELMNT1,THICK1,POR1,RI1,QSSS1,
      & OLDL1,OLDW1)
0134              IF(NTR .LT. 2)GO TO 873
C
0135              IF(IP2 .LE. 0)GO TO 878
0136              CALL COMPAT(NSS2,QSS2,FI2,DFI2,X2,Y2,XSS2,YSS2,
      & N2,T2,KODE2,IP2,XO2,YO2,XM2,YM2,AQSS2,ANSL2,MM,
      & ATETA1,ADFI1,ELMNT1,THICK2,POR2,RI2,T1,
      & OLDL1,OLDW1)
0137              GO TO 872
C
0138      878      WRITE(IMP,871)
0139      871      FORMAT(/5X,'NO STREAMLINE FROM REGION 1 CROSSES',
      & ' INTO REGION 2',/)
0140      872      CALL PLOT(0.0,0.0,3)
C
0141              CALL STRM(NSS2,QSS2,FI2,DFI2,X2,Y2,XSS2,YSS2,
      & N2,T2,KODE2,IP1,XO1,YO1,XM2,YM2,AQSS1,ANSL1,MM,
      & ATETA2,ADFI2,ELMNT2,THICK2,POR2,RI2,QSSS2,
      & OLDL2,OLDW2)
C
0142              IF(IP1 .LE. 0)GO TO 879
0143              CALL COMPAT(NSS1,QSS1,FI1,DFI1,X1,Y1,XSS1,YSS1,
      & N1,T1,KODE1,IP1,XO1,YO1,XM1,YM1,AQSS1,ANSL1,MM,
      & ATETA2,ADFI2,ELMNT2,THICK1,POR1,RI1,T2,
      & OLDL2,OLDW2)
0144              GO TO 873
C
0145      879      WRITE(IMP,874)
0146      874      FORMAT(/5X,'NO STREAMLINE FROM REGION 2 CROSSES',
      & ' INTO REGION 1',/)
0147      873      CALL PLOT(0.0,0.0,999)
0148      79      STOP
0149      END
```

```

0001      SUBROUTINE INPUT (CX,CY,X,Y,KODE,FI,XSS,YSS,QSS,NSS,N,
      & NIF,L,T,THICK,POR,RI,TOPP)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      DIMENSION CX(120),CY(120),X(120),Y(120),KODE(120)
0005      DIMENSION FI(120),XSS(120),YSS(120),QSS(120)
C
C THIS SUBROUTINE READS ALL THE INPUT DATA FOR A REGION
C
C N=THE NUMBER OF BOUNDARY ELEMENTS=NUMBER OF NODES
C L=THE NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS
C CALCULATED
C NSS= TOTAL NUMBER OF SOURCES AND SINKS
C QSS= THE STRENGTHS OF THE SOURCES AND SINKS
C NIF=NUMBER OF INTERFACE ELEMENTS OR NODES
C T=PERMEABILITY
C
C READ BASIC PARAMETERS
C
0006      READ(LEC,200)N,L,NSS,NIF
0007      200  FORMAT(4 I5)
0008      WRITE(IMP,300)N,L,NSS,NIF
0009      300  FORMAT(/,30X,'DATA '//8X,'NUMBER OF BOUNDARY ELEMENTS=',
      & I3/8X,'NUMBER OF INTERNAL POINTS WHERE THE FUNCTION IS '
      & ',CALCULATED=',I3/8X,'NUMBER OF SOURCES AND SINKS=',I5
      & '/8X,'NUMBER OF INTERFACE NODES=',I3,/)
C
C....READ RESERVOIR PROPERTIES
C
0010      READ(LEC,109)T,THICK,POR
0011      109  FORMAT(3F10.2)
0012      WRITE(IMP,107)T,THICK,POR
0013      107  FORMAT(20X,'PERMEABILITY (MD)=' ,F10.4, '//20X,
      & 'THICKNESS(FT)=' ,F10.4, '//20X, 'POROSITY=' ,F10.4, '//)
0014      WRITE(IMP,498)SCALE
0015      498  FORMAT(/20X,'SCALE: 1 INCH = ' ,F7.2, 'FEET',//)
C
0016      IF(L .LE. 0)GO TO 499
C
C READ INTERNAL POINTS CORDINATES
      DO 1 I=1,L
0017      1    READ(LEC,400)CX(I),CY(I)
0018      400  FORMAT(2F10.4)
C
C READ AND WRITE THE CORDINATES OF THE EXTREME POINTS
C OF THE BOUNDARY ELEMENTS INTO ARRAYS X AND Y
C
0020      499  WRITE(IMP,500)
0021      500  FORMAT(1H1,///8X,'THE COORDINATES OF THE EXTREME POINTS',
      & ' OF THE BOUNDARY ELEMENTS',//9X,'POINT',11X,'X(INCH)',
      & 12X,'Y(INCH)',/)
0022      DO 10 I=1,N
0023      READ(LEC,600)X(I),Y(I)
0024      600  FORMAT(2F10.4)
0025      10    WRITE(IMP,700)I,X(I),Y(I)
0026      700  FORMAT(9X,I3,2(9X,F10.4))

```

```
C
C READ THE BOUNDARY CONDITIONS INTO FI(I) ACCORDING TO A
C PRE-ASSIGNED CODE. IF THE POTENTIAL(PRESSURE) IS SPECIFIED,
C THE BOUNDARY IS ASSIGNED A CODE=0. IF THE POTENTIAL GRADIENT
C IS SPECIFIED, IT IS ASSIGNED A CODE=1. IF THE BOUNDARY IS ON
C AN INTERFACE BETWEEN TWO REGIONS(BOTH POTENTIAL AND POTENTIAL
C GRADIENT ARE UNKNOWN),IT IS ASSIGNED A CODE=2
C
0027      WRITE(IMP,800)
0028      800  FORMAT(1H1,///14X,'BOUNDARY CONDITIONS'//8X,'NODE',6X,
      &      'CODE',5X,'PRESCRIBED VALUE')
0029      DO 20 I=1,N
0030      READ(LEC,900)KCODE(I),FI(I)
0031      900  FORMAT(15,F10.4)
0032      20   WRITE(IMP,950)I,KCODE(I),FI(I)
0033      950  FORMAT(8X,I3,8X,I1,8X,E14.7)
C
C READ THE COORDINATES AND STRENGTHS OF THE SOURCES AND SINKS
C
0034      WRITE(IMP,556)
0035      556  FORMAT(1H1,///8X,'COORDINATES OF SOURCES AND SINKS',1X,
      &      'AND THEIR STRENGTHS',//)
0036      WRITE(IMP,557)
0037      557  FORMAT(20X,'X(INCH)',5X,'Y(INCH)',7X,'RATE(BBL/D)',/)
0038      DO 444 JJ=1,NSS
0039      READ(LEC,333)XSS(JJ),YSS(JJ),QSS(JJ)
0040      333  FORMAT( 3F10.4)
0041      WRITE(IMP,666)XSS(JJ),YSS(JJ),QSS(JJ)
0042      666  FORMAT(15X,2F12.4,6X,F12.4)
0043      444  CONTINUE
C READ THE AVERAGE RADIUS OF THE WELLS AND THE CHARACTERISTIC
C PRESSURE
C
0044      READ(LEC,448)RI,TOPP
0045      448  FORMAT(2F10.4)
0046      WRITE(IMP,449)RI,TOPP
0047      449  FORMAT(/20X,'THE AVERAGE RADIUS OF THE WELLS(INCH) IS=',
      &      F7.3,/20X,'THE CHARACTERISTIC PRESSURE(PSI) IS=',F10.4//)
0048      RETURN
0049      END
```

```
0001      SUBROUTINE CONVRS(CX,CY,X,Y,KODE,FI,XSS,YSS,QSS,NSS,N,  
      & L,T,THICK,PCR,RI,TOPP)  
C*****  
C  
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI  
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE  
0004      COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC  
0005      DIMENSION CX(120),CY(120),X(120),Y(120),KODE(120)  
0006      DIMENSION FI(120),XSS(120),YSS(120),QSS(120)  
C  
0007      T=T*TRFAC  
0008      RI=RI*DISFAC  
0009      THICK=THICK*DISFAC  
0010      PCR=PCR*PORFAC  
0011      TOPP=TOPP*PRESSF  
0012      IF(L .LE. 0)GO TO 15  
0013      DO 10 I=1,L  
0014      CX(I)=CX(I)*DISFAC*SCALE  
0015      CY(I)=CY(I)*DISFAC*SCALE  
0016      10 CONTINUE  
0017      15 DO 20 I=1,N  
0018      X(I)=X(I)*DISFAC*SCALE  
0019      Y(I)=Y(I)*DISFAC*SCALE  
0020      20 CONTINUE  
C  
C AFTER CONVERTING THE PRESSURE AND PRESSURE GRADIENTS TO  
C SI UNITS,DIVIDE THEM BY THE CHARACTERISTIC PRESSURE  
C  
0021      DO 40 I=1,N  
0022      IF(KODE(I) .EQ. 1)GO TO 30  
0023      FI(I)=FI(I)*PRESSF/TOPP  
0024      GO TO 40  
0025      30 FI(I)=FI(I)*DPRESF/TOPP  
0026      40 CONTINUE  
0027      DO 50 I=1,NSS  
0028      XSS(I)=XSS(I)*DISFAC*SCALE  
0029      YSS(I)=YSS(I)*DISFAC*SCALE  
0030      QSS(I)=QSS(I)*RATEF  
C  
C USE THE RATE TO DETERMINE THE STRENGTH  
0031      QSS(I)=QSS(I)*((VICSTM*VICFAC)/(THICK*T))  
C  
C DIVIDE THE STRENGTH BY THE CHARACTERISTIC PRESSURE  
0032      QSS(I)=QSS(I)/TOPP  
0033      50 CONTINUE  
0034      RETURN  
0035      END
```

```

0001      SUBROUTINE MATRX(X,Y,XM,YM,G,H,FI,DFI,KODE,XSS,YSS,
      &    QSS,NSS,IJ,NIF,SUM,T,N,L,CX,CY)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      DIMENSION X(120),Y(120),XM(120),YM(120),G(120,120)
0004      DIMENSION H(120,120),DYBLEN(120)
0005      DIMENSION FI(120),KODE(120),DFI(120),A(120,120),B(120)
0006      DIMENSION XSS(120),YSS(120),QSS(120),SUM(120),SIG(120)
0007      DIMENSION SUMC(120),CX(120),CY(120),DXBLEN(120)
0008      DIMENSION DGDY(120),DGDY(120),DHDY(120),DHDY(120)
0009      DIMENSION THETA(120),DIST(120),ELMNT(120)
C
C  THIS SUBROUTINE COMPUTES THE H AND G MATRIX ELEMENTS BY
C  THE METHOD OF GAUSSIAN QUADRATURE ALONG THE BOUNDARY
C  ELEMENTS. IT THEN FORMS THE SYSTEM AX=F
C  'INTE' CALCULATES ALL THE OFF-DIAGONAL ELEMENTS WHILE
C  'INLO' CALCULATES ONLY THE DIAGONAL ELEMENTS
C
C  COMPUTE THE MID-POINT COORDINATES AND STORE IN ARRAY
C  XM AND YM
C
0010      X(N+1)=X(1)
0011      Y(N+1)=Y(1)
0012      DO 10 I=1,N
0013      XM(I)=(X(I)+X(I+1))/2
0014      10  YM(I)=(Y(I)+Y(I+1))/2
C
C  COMPUTE THE G AND H MATRICES AND THE B VECTOR
C
0015      DO 31 I=1,N
0016      DO 30 J=1,N
0017      IF(I-J)20,25,20
0018      20  CALL INTE(XM(I),YM(I),X(J),Y(J),X(J+1),Y(J+1),H(I,J)
      &    ,G(I,J),DGDY(I),DGDY(I),DHDY(I),DHDY(I),DIST(J),
      &    ELMNT(J),THETA(J))
      GO TO 30
0019      25  CALL INLO(X(J),Y(J),X(J+1),Y(J+1),G(I,J))
0020      H(I,J)=PI
0021      30  CONTINUE
0022      31  CONTINUE
C  WRITE PORTIONS OF THE MATRICES OBTAINED
C  WRITE(IMP,34)
C  34  FORMAT(1H1,1X,'THE H MATRIX IS'//)
C  NNA=1
C  NA=0
C  102  NA=NA+12
C  DO 32 I=1,N
C  WRITE(IMP,33)(H(I,J),J=NNA,NA)
C  33  FORMAT(1X,12F10.4)
C  32  CONTINUE
C  IF(NA .GE. N)GO TO 101
C  NNA=NNA+12
C  GO TO 102
C  101  CONTINUE
C  WRITE(IMP,36)
C  36  FORMAT(//,1X,'THE G MATRIX IS'//)
C  NNA=1

```

```

C      NA=0
C 104    NA=NA+12
C      DO 38 I=1,N
C      WRITE(IMP,39)(G(I,J),J=NNA,NA)
C 39    FORMAT(1X,12E10.2)
C 38    CONTINUE
C      IF(NA .GE. N)GO TO 103
C      NNA=NNA+12
C      GO TO 104
C 103    CONTINUE
C
C ARRANGE THE SYSTEM TO PUT ALL THE UNKNOWNNS IN G(I,J) ON THE LHS
C

```

```

0024      IJ=0
0025      NEIF=N+NIF
0026      DO 51 J=1,N
0027      IF(KODE(J) .EQ. 1)GO TO 41
0028      IF(KODE(J) .EQ. 0)GO TO 51
0029      IJ=IJ+1
0030      DO 50 I=1,N
0031      CH=H(I,J)
0032      G(I,N+IJ )=-CH
0033      H(I,J)=0.
0034      50    CONTINUE
0035      40    GO TO 51
0036      41    DO 52 I=1,N
0037      CH=G(I,J)
0038      G(I,J)=-H(I,J)
0039      H(I,J)=-CH
0040      52    CONTINUE
0041      51    CONTINUE
C      WRITE(IMP,46)
C 46    FORMAT(1H1,1X,'THE REARRANGED G MATRIX IS'//)
C      NNA=1
C      NAA=0
C 106    NAA=NAA+12
C      DO 48 I=1,N
C      WRITE(IMP,47)(G(I,J),J=NNA,NAA)
C 47    FORMAT(1X,12E10.2)
C 48    CONTINUE
C      IF(NAA .GE. N)GO TO 105
C      NNA=NNA+12
C      GO TO 106
C 105    CONTINUE
C      WRITE(IMP,42)
C 42    FORMAT(//,1X,'THE REARRANGED H MATRIX IS'//)
C      NNA=1
C      NAA=0
C 110    NAA=NAA+12
C      DO 43 I=1,N
C      WRITE(IMP,44)(H(I,J),J=NNA,NAA)
C 44    FORMAT(1X,12E10.2)
C 43    CONTINUE
C      IF(NAA .GE. N)GO TO 109
C      NNA=NNA+12
C      GO TO 110
C 109    CONTINUE
C

```

C CALCULATE THE CONTRIBUTIONS OF THE SOURCES AND SINKS

C

```

0042      DO 55 I=1,N
0043      SUM(I)=0.
0044      DO 54 JJ=1,NSS
0045      CALL SOURCE(XM(I),YM(I),XSS(JJ),YSS(JJ),QSS(JJ),NSS,
      & BLENT,DXSURS,DYSURS)
0046      SUM(I)=SUM(I)+BLENT
0047      54 CONTINUE
0048      55 CONTINUE
C      WRITE(IMP,56)
C 56      FORMAT(//,1X,'THE VALUES DUE TO SOURCES AND SINKS ARE'//
C      DO 58 I=1,N
C      WRITE(IMP,57)SUM(I)
C 57      FORMAT(30X,F12.4)
C 58      CONTINUE
C
C COMPUTE THE R.H.S. VECTOR
C
C      WRITE(IMP,71)
C 71      FORMAT(1H1,//10X,'THE RHS WITHOUT EFFECTS OF ',
C      & 'SOURCES AND SINKS'//)
0049      DO 61 I=1,N
0050      DFI(I)=0.
0051      DO 60 J=1,N
0052      DFI(I)=DFI(I)+H(I,J)*FI(J)
0053      60 CONTINUE
C      WRITE(IMP,72)DFI(I)
C 72      FORMAT(10X,E14.7)
0054      DFI(I)=DFI(I)-SUM(I)
0055      61 CONTINUE
0056      RETURN
0057      END

```

```

0001      SUBROUTINE INTE(XP,YP,X1,Y1,X2,Y2,H,G,DGDX,DGDY,
        & DHDX,DHDY,DIST,ELMNT,THETA)
C*****
C
0002      DIMENSION XCO(4),YCO(4),GI(4),OME(4)
0003      COMMON /BLK1/LEC,IMP,VICSTM,PI
C
C ALL THE OFF-DIAGONAL ELEMENTS OF MATRICES H AND G ARE
C CALCULATED BY THIS SUBROUTINE
C DIST=PERPENDICULAR DISTANCE FROM THE NODE POINT UNDER
C CONSIDERATION TO THE BOUNDARY ELEMENTS
C RA=DISTANCE FROM THE POINT UNDER CONSIDERATION TO THE
C INTEGRATION POINTS IN THE BOUNDARY ELEMENTS
C
0004      GI(1)=0.86113631
0005      GI(2)=-GI(1)
0006      GI(3)=0.33998104
0007      GI(4)=-GI(3)
0008      OME(1)=0.34785485
0009      OME(2)=OME(1)
0010      OME(3)=0.65214515
0011      OME(4)=OME(3)
0012      AX=(X2-X1)/2
0013      BX=(X2+X1)/2
0014      AY=(Y2-Y1)/2
0015      BY=(Y2+Y1)/2
0016      ELMNT=SQRT((X1-X2)**2+(Y1-Y2)**2)
0017      IF(AX)10,20,10
0018      10    TA=AY/AX
0019           THETA=ATAN(TA)
0020           XLEN=(TA*XP-YP+Y1-TA*X1)/SQRT(TA**2+1.0)
0021           DIST=ABS(XLEN)
0022           DDISTX=(TA/SQRT(TA**2+1.0))
0023           DDISTY=(-1.0/SQRT(TA**2+1.0))
0024           IF(XLEN.GT.0.0)GO TO 30
0025      36    DDISTX=-DDISTX
0026           DDISTY=-DDISTY
0027           THETA=-THETA
0028           GO TO 30
0029      20    ALEN=XP-X1
0030           DIST=ABS(ALEN)
0031           THETA=PI/2.0
0032           DDISTX=1.0
0033           DDISTY=0.0
0034           IF(ALEN.GT.0.0)GO TO 30
0035           DDISTX=-DDISTX
0036           DDISTY=-DDISTY
0037           THETA=-THETA
0038      30    SIG=(X1-XP)*(Y2-YP)-(X2-XP)*(Y1-YP)
0039           IF(SIG)31,32,32
0040      31    DIST=-DIST
0041           DDISTX=-DDISTX
0042           DDISTY=-DDISTY
0043           THETA=-THETA
0044      32    G=0.
0045           H=0.
0046           DGDX=0.
0047           DGDY=0.

```

```

0048          DHDX=0.
0049          DHDY=0.
0050          DO 40 I=1,4
0051          XCD(I)=AX*GI(I)+BX
0052          YCD(I)=AY*GI(I)+BY
0053          RA=SQRT((XP-XCD(I))**2+(YP-YCD(I))**2)
0054          G=G+ALOG(1/RA)*OME(I)*SQRT(AX**2+AY**2)
0055          DGDY=DGDY+(-1.0*(XP-XCD(I))/RA**2)*OME(I)*SQRT(AX**2+
&            AY**2)
0056          DGDY=DGDY+(-1.0*(YP-YCD(I))/RA**2)*OME(I)*SQRT(AX**2+
&            AY**2)
0057          H=H+(-1.0*(DIST*OME(I)*SQRT(AX**2+AY**2))/RA**2)
0058          DHDX=DHDX+((-1.0)*(((1.0/RA**2)*DDISTX)+(DIST*(-2.0)
&            *(XP-XCD(I))/RA**4)))*OME(I)*SQRT(AX**2+AY**2)
0059          DHDY=DHDY+((-1.0)*(((1.0/RA**2)*DDISTY)+(DIST*(-2.0)
&            *(YP-YCD(I))/RA**4)))*OME(I)*SQRT(AX**2+AY**2)
0060          40 CONTINUE
0061          RETURN
0062          END

```

0001

SUBROUTINE INLG(X1,Y1,X2,Y2,G)

C*****

C

C THIS SUBROUTINE COMPUTES THE DIAGONAL ELEMENTS OF THE
C 'G' MATRIX. THE DIAGONAL ELEMENTS OF THE 'H' MATRIX
C ARE ALL ZERO.

C

0002

AX=(X2-X1)/2

0003

AY=(Y2-Y1)/2

0004

SR=SQRT(AX**2+AY**2)

0005

G=2*SR*(ALOG(1/SR)+1)

0006

ELMNT=SQRT((X1-X2)**2+(Y1-Y2)**2)

0007

RETURN

0008

END

```
0001      SUBROUTINE SOURCE(XM,YM,XSS,YSS,QSS,NSS,BLENT,DXSURS,  
        & DYSURS)  
C*** ****  
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI  
0003      DENT=SQRT((XM-XSS)**2+(YM-YSS)**2)  
0004      BLENT=QSS*ALOG(1.0/DENT)  
0005      DXSURS=QSS*(-1.0)*(XM-XSS)/DENT**2  
0006      DYSURS=QSS*(-1.0)*(YM-YSS)/DENT**2  
0007      RETURN  
0008      END
```

```
0001      SUBROUTINE ASEMBL(G1,G2,DFI1,DFI2,N1,N2,A,B,NAA,NIF,
      &   NNN,NOM,NNI,NNM,T1,T2)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      DIMENSION G1(120,120),G2(120,120),DFI1(120),DFI2(120)
0004      DIMENSION A(120,120),B(120)
C
C STORE THE MATRIX G(I,J) INTO A(I,J) AND DFI(I,J) INTO B(I,J)
C
C INITIALIZE THE ARRAYS
C
0005      NAA=N2+N1
0006      NNN=N1-NIF
0007      NDIF=N2-NIF
0008      NOM=N1+NIF
0009      NNI=NIF+1
0010      NNM=N2+NIF+1
0011      DO 6 I=1,NAA
0012      DO 7 J=1,NAA
0013      A(I,J)=0.0
0014      7 CONTINUE
0015      B(I)=0.0
0016      6 CONTINUE
C
C ASSEMBLE G1 INTO A
C
0017      806 DO 801 J=1,NOM
0018      DO 800 I=1,N1
0019      A(I,J)=G1(I,J)
0020      800 CONTINUE
0021      801 CONTINUE
0022      DO 805 J=1,NDIF
0023      DO 804 I=1,N1
0024      A(I,NOM+J)=0.
0025      804 CONTINUE
0026      805 CONTINUE
0027      DO 808 J=1,N1
0028      808 B(J)=DFI1(J)
C
C ASSEMBLE G2 INTO A
C
0029      810 DO 812 J=1,NNN
0030      DO 812 I=1,N2
0031      A(N1+I,J)=0.
0032      812 CONTINUE
0033      DO 814 J=1,NIF
0034      DO 813 I=1,N2
0035      A(N1+I,NNN+J)=-G2(I,NNI-J)*(T1/T2)
0036      813 CONTINUE
0037      814 CONTINUE
0038      DO 816 J=1,NIF
0039      DO 815 I=1,N2
0040      A(N1+I,N1+J)=G2(I,NNM-J)
0041      815 CONTINUE
0042      816 CONTINUE
0043      DO 818 J=1,NDIF
0044      DO 817 I=1,N2
```

```
0045      A(N1+I,NOM+J)=G2(I,NIF+J)
0046      817  CONTINUE
0047      818  CONTINUE
0048      DO 830 I=1,N2
0049      B(N1+I)=DFI2(I)
0050      830  CONTINUE
C
C  WRITE OUT THE FINAL RESULTING MATRIX
C
C  821  WRITE(IMP,62)
C  62   FORMAT(1H1,1X,'THE FINAL ASSEMBLED MATRIX IS :')
C      NB=1
C      NBB=0
C  740  NBB=NBB+12
C      DO 1101 I=1,NAA
C      WRITE(IMP,1100)(A(I,J),J=NB,NBB),B(I)
C 1100  FORMAT(1X,12F9.3,4X,F9.4)
C 1101  CONTINUE
C      IF(NB .GE. 48)GO TO 741
C      NB=NB+12
C      GO TO 740
C  741  CONTINUE
0051      RETURN
0052      END
```

```
0001      SUBROUTINE SLNPD(A,B,D,N,NC,NJ)
C *****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      DIMENSION A(120,120),B(120)
0004      DOUBLE PRECISION AA(120,120),BB(120),C
C
C THIS IS A SUBROUTINE THAT CAN SOLVE THE EQUATIONS BY
C GAUSSIAN ELIMINATION
C
C A=THE SYSTEM MATRIX
C B=ORIGINALLY IT CONTAINS THE RHS COEFFICIENTS. AFTER
C SOLUTION IT CONTAINS THE VALUES OF THE SOLUTION VECTOR
C N=ACTUAL NUMBER OF UNKNOWNNS
C NX=ROW AND COLUMN DIMENSION OF MATRIX 'A'
C
0005      NJ=N+NC
0006      N11=NJ-1
C
C CHANGE TO DOUBLE PRECISION VARIABLES
C
0007      DO 11 I=1,NJ
0008      DO 12 J=1,NJ
0009      AA(I,J)=A(I,J)
0010      12 CONTINUE
0011      BB(I)=B(I)
0012      11 CONTINUE
C
C INTERCHANGE ROWS TO GET NON ZERO DIAGONAL COEFFICIENT
C
0013      DO 100 K=1,N11
0014      K1=K+1
0015      C=A(K,K)
0016      IF(DABS(C)-0.000001)1,1,3
0017      1 DO 7 J=K1,NJ
0018      IF(DABS(AA(J,K))-0.000001)7,7,5
0019      5 DO 6 L=K,NJ
0020      C=AA(K,L)
0021      AA(K,L)=AA(J,L)
0022      6 AA(J,L)=C
0023      C=BB(K)
0024      BB(K)=BB(J)
0025      BB(J)=C
0026      C=AA(K,K)
0027      GO TO 3
0028      7 CONTINUE
0029      8 WRITE(IMP,2)K
0030      2 FORMAT('***** SINGULARITY IN ROW',I5)
0031      D=0.
0032      GO TO 300
C
C DIVIDE ROW BY DIAGONAL COEFFICIENT
C
0033      3 C=AA(K,K)
0034      DO 4 J=K1,NJ
0035      4 AA(K,J)=AA(K,J)/C
0036      BB(K)=BB(K)/C
C
```

```
C  ELIMINATE UNKNOWN X(K) FROM ROW 1
C
0037      DO 10 I=K1,NJ
0038      C=AA(I,K)
0039      DO 9 J=K1,NJ
0040      9  AA(I,J)=AA(I,J)-C*AA(K,J)
0041      10  BB(I)=BB(I)-C*BB(K)
0042      100  CONTINUE
C
C
C  COMPUTE LAST UNKNOWN
C
0043      IF(DABS(AA(NJ,NJ))-0.000001)14,14,101
0044      101  BB(NJ)=BB(NJ)/AA(NJ,NJ)
C
C  APPLY BACK SUBSTITUTION PROCESS TO COMPUTE REMAINING UNKNOWNNS
C
0045      DO 200 L=1,N11
0046      K=NJ-L
0047      K1=K+1
0048      DO 200 J=K1,NJ
0049      200  BB(K)=BB(K)-AA(K,J)*BB(J)
0050      WRITE(IMP,1307)
0051      1307  FORMAT(1H1,/,/,2X,'THE SOLUTION VECTOR IS'//)
0052      DO 1300 I=1,NJ
0053      WRITE(IMP,1301)BB(I)
0054      1301  FORMAT(10X,E14.7)
0055      1300  CONTINUE
C
C  RETURN TO SINGLE PRECISION VARIABLES
C
0056      255  DO 13 J=1,NJ
0057      B(J)=BB(J)
0058      13  CONTINUE
C
C  COMPUTE VALUE OF DETERMINANT
C
C      D=1.
C      DO 250 I=1,NJ
C  250  D=D*AA(I,I)
0059      GO TO 300
0060      14  WRITE(IMP,15)K
0061      15  FORMAT('****SINGULARITY IN ROW',I5)
0062      D=0.
0063      300  RETURN
0064      END
```

```

0001      SUBROUTINE SPLIT(N1,N2,NIF,NNN,NOM,NNI,NNM,NAA,KODE1,
      & KODE2,F11,F12,DFI1,DFI2,B,T1,T2)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      DIMENSION KODE1(120),KODE2(120),F11(120),F12(120)
0004      DIMENSION DFI1(120),DFI2(120),B(120)
C
0005      DO 50 I=1,N1
0006      IF(KODE1(I) .EQ. 2)GO TO 50
0007      IF(KODE1(I) .EQ. 0)GO TO 52
0008      CH=F11(I)
0009      DFI1(I)=CH
0010      F11(I)=B(I)
0011      GO TO 50
0012      52 DFI1(I)=B(I)
0013      50 CONTINUE
0014      54 DO 60 I=1,NIF
0015      DFI1(NNN+I)=B(NNN+I)
0016      DFI2(I)=-B(N1+1-I)*(T1/T2)
0017      F12(I)=B(NOM+1-I)
0018      F11(NNN+I)=B(N1+I)
0019      60 CONTINUE
0020      NDIF=N2-NIF
0021      DO 70 I=1,N2
0022      IF(KODE2(I) .EQ. 2)GO TO 70
0023      IF(KODE2(I) .EQ. 0)GO TO 61
0024      CH=F12(I)
0025      F12(I)=B(N1+I)
0026      DFI2(I)=CH
0027      GO TO 70
0028      61 DFI2(I)=B(N1+I)
0029      70 CONTINUE
0030      51 CONTINUE
0031      RETURN
0032      END

```

```

0001          SUBROUTINE INTER(FI,DFI,L,N,CX,CY,X,Y,
&  XSS,YSS,QSS,NSS,SOL)
C*****
C
0002          COMMON /BLK1/LEC,IMP,VICSTM,PI
0003          COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004          COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC
0005          DIMENSION FI(120),DFI(120),CX(120),CY(120),X(120)
0006          DIMENSION Y(120),DYBLEN(120),ELMNT(120)
0007          DIMENSION SOL(120),SUM(120),DGDY(120),DGDY(120)
0008          DIMENSION DHDX(120),DHDY(120),SIG(120),DXBLEN(120)
0009          DIMENSION THETA(120),XSS(120),YSS(120),QSS(120)
0010          DIMENSION XIF(120),YIF(120),FILOT(120),DIST(120)
C
C  THIS SUBROUTINE COMPUTES POTENTIAL VALUES FOR INTERNAL POINTS
C
0011          DO 40 K=1,L
0012          SOL(K)=0.
0013          DO 30 J=1,N
0014          CALL INTE(CX(K),CY(K),X(J),Y(J),X(J+1),Y(J+1),HI,
&  GI,DGDY(K),DGDY(K),DHDX(K),DHDY(K),DIST(J),ELMNT(J),
&  THETA(J))
0015          SOL(K)=SOL(K)+DFI(J)*GI-FI(J)*HI
0016          30  CONTINUE
C  CALCULATE THE VALUES DUE TO SOURCES AND SINKS AT INTERNAL POINT
0017          SUM(K)=0.
0018          DO 90 JJ=1,NSS
0019          CALL SOURCE(CX(K),CY(K),XSS(JJ),YSS(JJ),QSS(JJ),NSS,
&  BLENT,DXSURS,DYSURS)
0020          SUM(K)=SUM(K)+BLENT
0021          90  CONTINUE
0022          SOL(K)=(SOL(K)+SUM(K))/(2.0*PI)
0023          40  CONTINUE
0024          WRITE(IMP,300)
0025          300  FORMAT(//,2X,' INTERNAL POINTS',//11X,'X',18X,'Y',
&  14X,'POTENTIAL',/)
0026          DO 20 K=1,L
0027          XIF(K)=CX(K)/SCLFAC
0028          YIF(K)=CY(K)/SCLFAC
0029          FILOT(K)=SOL(K)/PRESSF
0030          20  WRITE(IMP,400)XIF(K),YIF(K),FILOT(K)
0031          400  FORMAT(3(5X,E14.7))
0032          RETURN
0033          END

```

```

0001      SUBROUTINE OUTPUT(X,Y,XM,YM,FI,DFI,CX,CY,SOL,N,XSS,
      &  YSS,L,NSS,QSS,TOPP,THICK,T,QSSS)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC
0005      DIMENSION XM(120),YM(120),FI(120),DFI(120),CX(120)
0006      DIMENSION SOL(120),XSS(120),YSS(120),XIF(120)
0007      DIMENSION CY(120),QSSS(120),X(120),Y(120),QSS(120)
0008      DIMENSION YIF(120),FIPLDT(120),DFIPLT(120)
C
0009      WRITE(IMP,100)
0010      100  FORMAT(1H1,80(' '))//30X,'RESULTS'//25X,'BOUNDARY NODES'
      &  //11X,'X(INCH)',13X,'Y(INCH)',10X,'PRESSURE(PSI)',7X,
      &  'NORMAL GRADIENT'//)
C
0011      DO 10 I=1,N
0012      FI(I)=FI(I)*TOPP
0013      DFI(I)=DFI(I)*TOPP
C
C FOR SAKE OF OUTPUT ONLY RETURN THE FOLLOWING TO FPS UNITS
0014      XIF(I)=XM(I)/SCLFAC
0015      YIF(I)=YM(I)/SCLFAC
0016      FIPLDT(I)=FI(I)/PRESSF
0017      DFIPLT(I)=DFI(I)/DPRESF
0018      WRITE(IMP,200)XIF(I),YIF(I),FIPLDT(I),DFIPLT(I)
0019      200  FORMAT(4(8X,E12.4))
0020      10  CONTINUE
0021      DO 30 J=1,NSS
0022      QSS(J)=QSS(J)*TOPP
C
C DETERMINE THE RATE FROM THE STRENGTH
0023      QSSS(J)=QSS(J)/(VICSTM*VICFAC/(THICK*T))
0024      30  CONTINUE
0025      WRITE(IMP,500)
0026      500  FORMAT(' ',80(' '))
0027      RETURN
0028      END

```

```
0001      SUBROUTINE BDRY(X,Y,XSS,YSS,NSS,QSSS,N)
C*****
C
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC
0005      DIMENSION X(120),Y(120),XSS(120),YSS(120),QSS(120)
0006      DIMENSION QSSS(120)
C
C SUBROUTINE BDRY PLOTS THE BOUNDARY AND WELL LOCATIONS
C
0007      FACPLO=2.0
0008      N1=N+1
0009      CALL FACTOR(FACPLO)
0010      CALL PLOT(X(N1)/SCLFAC,Y(N1)/SCLFAC,3)
0011      DO 741 I=1,N1
0012      CALL PLOT(X(I)/SCLFAC,Y(I)/SCLFAC,2)
0013      741 CONTINUE
0014      DO 742 I=1,NSS
0015      FACPLO=2.0
0016      XSSI=XSS(I)/SCLFAC
0017      YSSI=YSS(I)/SCLFAC
0018      QSSI=QSSS(I)/RATEF
0019      PN1=XSSI*FACPLO
0020      PN2=YSSI*FACPLO
0021      FACPLO=1.0
0022      CALL FACTOR(FACPLO)
0023      CALL NUMBER(PN1-0.15,PN2-0.15,0.08,QSSI,0.0,2)
0024      IF(QSSI .GE. 0.0) GO TO 744
0025      CALL SYMBOL(PN1,PN2,0.08,11,0,-1)
0026      GO TO 742
0027      744 CALL SYMBOL(PN1,PN2,0.08,10,0,-1)
0028      742 CONTINUE
0029      RETURN
0030      END
```

```

0001      SUBROUTINE STRM(NSS,QSS,FI,DFI,X,Y,XSS,YSS,
      & N,T,KODE, IPE,XOE,YOE,XM,YM,AQSS,ANSL,MM,ATETA,
      & ADFI,ELMNT,THICK,POR,RI,QSSS,OLDL,OLDW)
C *****
C
0002      COMMON /ELK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC
0005      DIMENSION X(120),Y(120),FI(120),DFI(120),XSS(120)
0006      DIMENSION YSS(120),G(120,120),DHDY(120),RD(120)
0007      DIMENSION QSS(120),XI(120),YI(120),KODE(120)
0008      DIMENSION MM(120),VX(120),VY(120),VT(120),H(120,120)
0009      DIMENSION DIST(120),XM(120),YM(120),SUM(120),DISMID(120)
0010      DIMENSION ANSL(120),DGDY(120),DGDY(120),DHDX(120)
0011      DIMENSION DXBLEN(120),DPDX(120),DPDY(120),DYBLEN(120)
0012      DIMENSION DPHIDX(120),DPHIDY(120),VLCP(120),XOE(120)
0013      DIMENSION SINK(120),STRML(120),STRMW(120),AVWIDT(120)
0014      DIMENSION AQSS(120),ELMNT(120),LENTH(120,120)
0015      DIMENSION QSSS(120),CQSS(120,120),WIDTH(120,120)
0016      DIMENSION ADFI(120),ATETA(120),THETA(120),VEL(120)
0017      DIMENSION YOE(120),ZETA(120),XIF(120),YIF(120)
0018      DIMENSION OLDL(120),OLDW(120)
0019      INTEGER SINK
0020      REAL LENTH
C
C FACT=THE CONVERSION FACTOR CHANGING (ATM/CM) TO
C (NEWTON/SQ M)/M
0021      FACT=1.0333E+7
C CALCULATE THE VELOCITY NEAR THE LOWEST PRODUCING WELL
C
0022      CALL LOWVEL(NSS,QSS,FI,DFI,X,Y,XSS,YSS,N,T,VLOW,RI,
      & QMNI,RAD,POR,THICK)
C
C INITIALIZE ARRAYS
C
0023      NST=20
0024      DO 206 I=1,NSS
0025      SINK(I)=0.0
0026      206 CONTINUE
0027      IPE=0
0028      DO 220 I=1,NSS
0029      IF(QSSS(I) .LE. 0.0) GO TO 220
0030      NSL=ABS((-FLOAT(NST)*QSSS(I)/QMNI)+0.5)
0031      IF(NSL .LT. 10) NSL=10
0032      IF(NSL .GT. 20) NSL=20
C
C CALCULATE THE STARTING POSITIONS ON A STREAMLINE
C
0033      DO 210 K=1,NSL
0034      ICANT=0
0035      XI(K)=XSS(I)+RAD*COS(1.0+(K*2.0*PI)/NSL)
0036      YI(K)=YSS(I)+RAD*SIN(1.0+(K*2.0*PI)/NSL)
0037      STRML(K)=0.0
0038      STRMW(K)=(2.0*PI*RI)/NSL
0039      AVWIDT(K)=STRML(K)
0040      FACPLD=2.0
0041      CALL FACTOR(FACPLD)
0042      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,3)

```

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C
C CALCULATE THE POTENTIAL AND VELOCITY AT THIS POSITION
C
C IPE IS THE COUNTER FOR STREAMLINES THAT CROSS INTO NEW REGIONS
C ICANT IS THE COUNTER FOR THE NUMBER OF POINTS ON A STREAMLINE
C   WRITE(IMP,10)
C 10   FORMAT(6X,'XI',8X,'YI',7X,'DPDX',5X,'DPDY',5X,
C   $ 'DXBLEN',5X,'DYBLEN',7X,'VT',7X,'DT',7X,'DSX',7X,'DSY')
0043   171 ICANT=ICANT+1
0044       IF(ICANT .GT. 100) GO TO 210
0045       VX(K)=0.
0046       VY(K)=0.
0047       VT(K)=0.
0048       DPDX(K)=0.
0049       DPDY(K)=0.
0050       DO 172 J=1,N
0051       DISMID(J)=SQRT((XI(K)-XM(J))**2+(YI(K)-YM(J))**2)
0052       CALL INTE(XI(K),YI(K),X(J),Y(J),X(J+1),Y(J+1),H(K,J),
           & G(K,J),DGDY(J),DGDY(J),DHDX(J),DHDY(J),DIST(J),
           & ELMNT(J),THETA(J))
0053       ELMNT2=0.5*ELMNT(J)
C....CHECK IF POINT IS NEAR A BOUNDARY, IF SO GO TO 347
C...OTHERWISE FIND THE VELOCITY AT THE CURRENT POSITION
0054       IF(DISMID(J) .LT. ELMNT2)GO TO 347
0055       GO TO 748
0056   347 IF(ICANT .GT. 1)GO TO 258
0057       WRITE(IMP,259)I
0058   259 FORMAT(/8X,' INJECTOR NUMEER',I2,' IS TOO CLOSE TO',
           & ' A BOUNDARY',/)
0059       GO TO 220
0060   258 IF(KODE(J) .GT. 1) GO TO 749
0061       GO TO 210
0062   749 IF(DFI(J) .LT. 0.0)GO TO 348
0063       GO TO 210
0064   748 DPDX(K)=DPDX(K)+DFI(J)*DGDY(J)-FI(J)*DHDX(J)
0065       DPDY(K)=DPDY(K)+DFI(J)*DGDY(J)-FI(J)*DHDY(J)
0066   172 CONTINUE
C
C CALCULATE THE CONTRIBUTIONS OF THE SOURCES AND SINKS
C
0067       DXBLEN(K)=0.
0068       DYBLEN(K)=0.
0069       DO 54 JJ=1,NSS
0070       CALL SOURCE(XI(K),YI(K),XSS(JJ),YSS(JJ),QSS(JJ),NSS,
           & BLEN,DXSURS,DYSURS)
0071       DXBLEN(K)=DXBLEN(K)+DXSURS
0072       DYBLEN(K)=DYBLEN(K)+DYSURS
0073   54 CONTINUE
0074       DPHDX(K)=(DPDX(K)+DXBLEN(K))/(2.0*PI)*FACT
0075       DPHDY(K)=(DPDY(K)+DYBLEN(K))/(2.0*PI)*FACT
0076       VX(K)=-1.0*(T/VICSTM*VICFAC)*DPHDX(K)
0077       VY(K)=-1.0*(T/VICSTM*VICFAC)*DPHDX(K)
0078       VT(K)=SQRT(VX(K)**2+VY(K)**2)
0079       DT=50.0*RI/ABS(VT(K))
0080       DSX=DT*VX(K)
0081       DSY=DT*VY(K)
0082       DS=DT*VT(K)
0083       XIF(K)=XI(K)/SCLFAC

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0084      YIF(K)=YI(K)/SCLFAC
0085      IF(ICANT .GT. 5)GO TO 3005
C        WRITE(IMP,11)XIF(K),YIF(K),DPDX(K),DPDY(K),DXBLEN(K),
C      &  DYBLEN(K),VT(K),DT,DSX,DSY
C 11      FORMAT(2X,10E10.3)
C...CHECK IF POINT IS NEAR A PRODUCER, IF SO BRANCH TO 333
C...OTHERWISE CALCULATE THE LENGTH AND WIDTH AT THE CURRENT
C...POSITION. PLOT THE CURRENT POSITION
0086      3005  CONTINUE
0087      IF(VT(K) .GT. VLOW) GO TO 333
0088      111    STRML(K)=STRML(K)+(VT(K)*DT)
0089      STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT(K))
0090      AVWIDT(K)=STRMW(K)/ICANT
0091      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
C...CALCULATE NEW POSITION
0092      XI(K)=XI(K)+VX(K)*DT
0093      YI(K)=YI(K)+VY(K)*DT
0094      XIF(K)=XI(K)/SCLFAC
0095      YIF(K)=YI(K)/SCLFAC
0096      GO TO 171
0097      333    DO 200 LL=1,NSS
0098      IF(QSS(LL) .GE. 0.0)GO TO 200
0099      RO(LL)=SQRT((XI(K)-XSS(LL))**2+(YI(K)-YSS(LL))**2)
0100      IF(RAD .GE. RO(LL)) GO TO 208
0101      200    CONTINUE
C...POINT IS NOT NEAR PRODUCER, BUT RATHER, IT IS NEAR INJECTOR
0102      GO TO 111
C IDENTIFY THE PRODUCER AT WHICH THE STREAMLINE TERMINATED
0103      208    SINK(LL)=SINK(LL)+1
0104      LAST=SQRT((XI(K)-XSS(LL))**2+(YI(K)-YSS(LL))**2)
0105      WLAST=(2.0*PI*RI)/FLOAT(NSL)
0106      STRMW(K)=STRMW(K)+WLAST
0107      AVWIDT(K)=STRMW(K)/(FLOAT(ICANT)+1.0)
0108      LENTH(LL,SINK(LL))=STRML(K)+LAST
0109      WIDTH(LL,SINK(LL))=AVWIDT(K)
0110      QSS(LL,SINK(LL))=QSSS(I)/FLOAT(NSL)
0111      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0112      GO TO 210
C
C...POINT IS NEAR AN INTERFACE BOUNDARY
C
0113      348    IPE=IPE+1
0114      MM(IPE)=J
0115      EX=XI(K)
0116      EY=YI(K)
C
C...CHECK FOR HORIZONTAL BOUNDARY
C
0117      IF(ABS(Y(J+1)-Y(J)) .GT. 0.001)GO TO 244
0118      243    XI(K)=XI(K)
0119      YI(K)=Y(J)
0120      XOE(IPE)=XI(K)
0121      YOE(IPE)=YI(K)
0122      AQSS(IPE)=QSSS(I)
0123      ANSL(IPE)=NSL
0124      ADFI(IPE)=DFI(J)
C
C CALCULATE THE ANGLE THE OUTWARD NORMAL MAKES WITH THE

```

C X-DIRECTION

C

```

0125      ATETA(IPE)=0.5*PI
0126      IF(Y(J) .LT. EY) ATETA(IPE)=-ATETA(IPE)
0127      VT(K)=-1.0*(T/VICFAC*VICSTM)*DFI(J)
0128      VEL(IPE)=VT(K)
0129      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0130      STRML(K)=STRML(K)+DIST(J)
0131      STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT(K))
0132      AVWIDT(K)=STRMW(K)/ICANT
0133      OLDL(IPE)=STRML(K)
0134      OLDW(IPE)=AVWIDT(K)
0135      EX=EX/SCLFAC
0136      EY=EY/SCLFAC
0137      XIF(K)=XI(K)/SCLFAC
0138      YIF(K)=YI(K)/SCLFAC
0139      714      GO TO 2121

```

C

C...CHECK FOR VERTICAL BOUNDARY

C

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0140      244      IF(ABS(X(J+1)-X(J)) .GT. 0.001) GO TO 245
0141      246      XI(K)=X(J)
0142              YI(K)=YI(K)
0143              XOE(IPE)=XI(K)
0144              YOE(IPE)=YI(K)
0145              AQSS(IPE)=QSSS(I)
0146              ANSL(IPE)=NSL
0147              ADFI(IPE)=DFI(J)

```

C

C THETA IS THE ANGLE THAT THE BOUNDARY ELEMENT MAKES WITH
C THE X-AXIS

C ATETA IS THE ANGLE THE OUTWARD NORMAL MAKES WITH THE
C X-AXIS

C

```

0148      ATETA(IPE)=0.0
0149      IF(EX .GT. X(J)) GO TO 3001
0150      ATETA(IPE)=-ATETA(IPE)
0151      ADFI(IPE)=-ADFI(IPE)
0152      3001      VT(K)=-1.0*(T/VICFAC*VICSTM)*DFI(J)
0153              VEL(IPE)=VT(K)
0154              CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0155              STRML(K)=STRML(K)+DIST(J)
0156              STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT(K))
0157              AVWIDT(K)=STRMW(K)/ICANT
0158              OLDL(IPE)=STRML(K)
0159              OLDW(IPE)=AVWIDT(K)
0160              EX=EX/SCLFAC
0161              EY=EY/SCLFAC
0162              XIF(K)=XI(K)/SCLFAC
0163              YIF(K)=YI(K)/SCLFAC
0164              GO TO 2121
0165      245      XI(K)=(((X(J)-X(J+1))/(Y(J+1)-Y(J)))*EX-((Y(J+1)-Y(J))/
&      (X(J+1)-X(J)))*X(J)+Y(J)-EY)/((X(J)-X(J+1))/(Y(J+1)-Y(J))
&      )-((Y(J+1)-Y(J))/(X(J+1)-X(J))))
0166              YI(K)=(((X(J)-X(J+1))/(Y(J+1)-Y(J)))*XI(K)-((X(J)-X(J+1))/
&      (Y(J+1)-Y(J)))*EX+EY
0167              XOE(IPE)=XI(K)
0168              YOE(IPE)=YI(K)

```

C X-DIRECTION

C

```

0125      ATETA(IPE)=0.5*PI
0126      IF(Y(J) .LT. EY) ATETA(IPE)=-ATETA(IPE)
0127      VT(K)=-1.0*(T/VICFAC*VICSTM)*DFI(J)
0128      VEL(IPE)=VT(K)
0129      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0130      STRML(K)=STRML(K)+DIST(J)
0131      STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT(K))
0132      AVWIDT(K)=STRMW(K)/ICANT
0133      OLDL(IPE)=STRML(K)
0134      OLDW(IPE)=AVWIDT(K)
0135      EX=EX/SCLFAC
0136      EY=EY/SCLFAC
0137      XIF(K)=XI(K)/SCLFAC
0138      YIF(K)=YI(K)/SCLFAC
0139      714      GO TO 2121

```

C

C...CHECK FOR VERTICAL BOUNDARY

C

```

0140      244      IF(ABS(X(J+1)-X(J)) .GT. 1.0) GO TO 245
0141      246      XI(K)=X(J)
0142      YI(K)=YI(K)
0143      XOE(IPE)=XI(K)
0144      YOE(IPE)=YI(K)
0145      AQSS(IPE)=QSSS(I)
0146      ANSL(IPE)=ANSL(I)
0147      ADFI(IPE)=ADFI(I)

```

C

C THETA IS

C THE X-AXIS

C ATETA IS

C X-AXIS

C

```

0148      ATETA(IPE)=ATETA(IPE)+THETA
0149      IF(ATETA(IPE) .GT. 3.14159) ATETA(IPE)=ATETA(IPE)-3.14159
0150      ATETA(IPE)=ATETA(IPE)
0151      ADI=ADI+1
0152      3001      VT=VEL(IPE)
0153      VEL(IPE)=VEL(IPE)+VT
0154      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0155      STRML(K)=STRML(K)+DIST(J)
0156      STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT)
0157      AVWIDT(K)=STRMW(K)/ICANT
0158      OLDL(IPE)=STRML(K)
0159      OLDW(IPE)=AVWIDT(K)
0160      EX=EX/SCLFAC
0161      EY=EY/SCLFAC
0162      XIF(K)=XI(K)/SCLFAC
0163      YIF(K)=YI(K)/SCLFAC
0164      GO TO 2121
0165      245      XI(K)=((X(J)-X(J+1))/(Y(J+1)-Y(J)))*EX-((Y(J+1)-Y(J))/
&      (X(J+1)-X(J)))*X(J)+Y(J)-EY)/((X(J)-X(J+1))/(Y(J+1)-Y(J))
&      )-((Y(J+1)-Y(J))/(X(J+1)-X(J)))
0166      YI(K)=((X(J)-X(J+1))/(Y(J+1)-Y(J)))*XI(K)-((X(J)-X(J+1))/
&      (Y(J+1)-Y(J)))*EX+EY
0167      XOE(IPE)=XI(K)
0168      YOE(IPE)=YI(K)

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```

0169      AQSS(IPE)=QSSS(I)
0170      ANSL(IPE)=NSL
0171      ADFI(IPE)=DFI(J)
0172      VT(K)=-1.0*(T/VICFAC*VICSTM)*DFI(J)

C
C...ADJUST FOR STREAMLINES APPROACHING FROM THE LEFT
0173      IF(XI(K) .GT. EX)ADFI(IPE)=-ADFI(IPE)
0174      VEL(IPE)=VT(K)
0175      STRML(K)=STRML(K)+(VT(K)*DT)
0176      STRMW(K)=STRMW(K)+QSSS(I)/(NSL*THICK*POR*VT(K))
0177      AVWIDT(K)=STRMW(K)/ICANT
0178      OLDL(IPE)=STRML(K)
0179      OLDW(IPE)=AVWIDT(K)

C
C CALCULATE THE ANGLE THE OUTWARD NORMAL MAKES WITH THE
C X-AXIS
C
0180      ZETA(J)=-1.0/TAN(THETA(J))
0181      ATETA(IPE)=ATAN(ZETA(J))
0182      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0183      XIF(K)=XI(K)/SCLFAC
0184      YIF(K)=YI(K)/SCLFAC
0185      EX=EX/SCLFAC
0186      EY=EY/SCLFAC
0187      XIF(K)=XI(K)/SCLFAC
0188      YIF(K)=YI(K)/SCLFAC
0189      2121 WRITE(IMP,2111)IPE,XIF(K),YIF(K),EX,EY
0190      2111 FORMAT(/,8X,'STREAMLINE INTERFACE BEGINING POINTS ARE',
& 15,2F9.3,/,8X,'PRECEEDING INTERIOR POINT IS',2F10.3)
0191      WRITE(IMP,605)VEL(IPE),ADFI(IPE),ATETA(IPE),AQSS(IPE)
& ,ANSL(IPE)
0192      605 FORMAT(/,2X,5F10.4)
0193      210 CONTINUE
0194      220 CONTINUE
C...WRITE OUT THE ANSWERS
0195      DO 440 I=1,NSS
0196      IF(QSSS(I) .GE. 0.0) GO TO 440
0197      WRITE(IMP,441)I,SINK(I)
0198      441 FORMAT(1H1,///,15X,'PRODUCER NUMBER',I3,/,15X,
& 'NUMBER OF STREAMLINES=',I3,/)
0199      WRITE(IMP,443)
0200      443 FORMAT(8X,'STREAMLINE NUMBER',3X,'TOTAL LENGTH',3X,
& 'AVERAGE WIDTH',6X,'INJ. RATE',/)
0201      ITEMP=SINK(I)
0202      IF(ITEMP .LE. 0) GO TO 440
0203      DO 444 J=1,ITEMP
0204      LENTH(I,J)=LENTH(I,J)/DISFAC
0205      WIDTH(I,J)=WIDTH(I,J)/DISFAC
0206      CQSS(I,J)=CQSS(I,J)/RATEF
0207      WRITE(IMP,445)J,LENTH(I,J),WIDTH(I,J),CQSS(I,J)
0208      445 FORMAT(14X,I3,10X,E12.4,4X,E12.4,5X,E12.4)
0209      444 CONTINUE
0210      440 CONTINUE
0211      RETURN
0212      END

```

```

0001      SUBROUTINE LOWVEL(NSS,QSS,FI,DFI,X,Y,XSS,YSS,
&      N,T,VLOW,RI,QMNI,RAD,POR,THICK)
C
C*****
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      COMMON /BLK3/PRESSF,DPRESF,PORFAC,SCLFAC
0005      DIMENSION X(120),Y(120),FI(120),DFI(120),XSS(120)
0006      DIMENSION YSS(120),DXBLEN(120),ELMNT(120)
0007      DIMENSION QSS(120),XI(120),YI(120),DPDX(120),DPDY(120)
0008      DIMENSION VX(120),VY(120),VT(120),A(120,120),G(120,120)
0009      DIMENSION DGDY(120),DGDY(120),DHDY(120),DHDY(120)
0010      DIMENSION DPHDX(120),DPHDY(120),VLCP(120),DYBLEN(120)
0011      DIMENSION RO(120),THETA(120),SUM(120),DIST(120)
C
C
0012      NST=20
0013      FACT=1.0333E+7
C  DEFINE A CAPTURE RADIUS (RAD)
0014      RAD=RI*50.0
C
C  CALCULATE THE VELOCITY NEAR THE LOWEST PRODUCTION WELL
C
0015      QMNI=1.0E+10
0016      DO 110 KK=1,NSS
0017      IF(QSS(KK) .GT. 0.0) GO TO 110
0018      IF(ABS(QSS(KK)) .GT. ABS(QMNI))GO TO 110
0019      QMNI=QSS(KK)
0020      K=KK
0021      110  CONTINUE
0022      VX(K)=0.
0023      VY(K)=0.
0024      VT(K)=0.
0025      DPDX(K)=0.
0026      DPDY(K)=0.
0027      XI(K)=XSS(K)+RAD
0028      YI(K)=YSS(K)+RAD
0029      DO 172 J=1,N
0030      CALL INTE(XI(K),YI(K),X(J),Y(J),X(J+1),Y(J+1),A(K,J),
&      G(K,J),DGDY(J),DGDY(J),DHDY(J),DHDY(J),DIST(J),
&      ELMNT(J),THETA(J))
0031      DPDX(K)=DPDX(K)+DFI(J)*DGDY(J)-FI(J)*DHDY(J)
0032      DPDY(K)=DPDY(K)+DFI(J)*DGDY(J)-FI(J)*DHDY(J)
0033      172  CONTINUE
C
C  CALCULATE THE CONTRIBUTIONS OF THE SOURCES AND SINKS
C
0034      DXBLEN(K)=0.
0035      DYBLEN(K)=0.
0036      SUM(K)=0.
0037      DO 54 JJ=1,NSS
0038      CALL SOURCE(XI(K),YI(K),XSS(JJ),YSS(JJ),QSS(JJ),NSS
&      ,BLENT,DXSURS,DYSURS)
0039      DXBLEN(K)=DXBLEN(K)+DXSURS
0040      DYBLEN(K)=DYBLEN(K)+DYSURS
0041      54  CONTINUE
0042      DPHDX(K)=(DPDX(K)+DXBLEN(K))/(2.0*PI)*FACT
0043      DPHDY(K)=(DPDY(K)+DYBLEN(K))/(2.0*PI)*FACT

```

```
0044      VX(K)=-1.0*(T/VICSTM*VICFAC)*DPHIDX(K)
0045      VY(K)=-1.0*(T/VICSTM*VICFAC)*DPHIDY(K)
0046      VLCP(K)=SQRT(VX(K)**2+VY(K)**2)
0047      VLOW=VLCP(K)
0048      QMNI=QMNI/(VICSTM*VICFAC/(THICK*T))
0049      QMNIW=QMNI/RATEF
0050      VELO=VLOW/DISFAC
0051      WRITE(IMP,113)QMNIW,VELO
0052      113  FORMAT(1H1,///3X,'THE RATE OF THE LOWEST PRODUCER(BBL/D)=
           & F12.3/3X,'THE VELOCITY NEAR THE LOWEST PRODUCER(FT/SEC)='
           & E12.3,/)
0053      RETURN
0054      END
```

```

0001      SUBROUTINE COMPAT(NSS,QSS,FI,DFI,X,Y,XSS,YSS,N,T,
&      KODE,IPC,XOC,YOC,XM,YM,BQSS,BNSL,MMX,BETA,BDFI,ELT,
&      THICK,POR,RI,TADJ,OLDL,OLDW)
C
C*****
0002      COMMON /BLK1/LEC,IMP,VICSTM,PI
0003      COMMON /BLK2/DISFAC,VICFAC,DENFAC,TRFAC,RATEF,SCALE
0004      COMMON /BLK3/PRESSF,DPRE SF,PORFAC,SCLFAC
0005      DIMENSION X(120),Y(120),FI(120),DFI(120),XSS(120)
0006      DIMENSION XI(120),YI(120),QSS(120),KODE(120)
0007      DIMENSION YSS(120),G(120,120),DISMID(120)
0008      DIMENSION VX(120),VY(120),VT(120),H(120,120)
0009      DIMENSION DIST(120),XM(120),YM(120),SUM(120)
0010      DIMENSION BNSL(120),DGDX(120),DGDY(120),DHDX(120)
0011      DIMENSION DXBLEN(120),DPDX(120),DPDY(120),DYBLEN(120)
0012      DIMENSION DHDY(120),RO(120),YOC(120),AVWIDT(120)
0013      DIMENSION DPHIDX(120),DPHIDY(120),VLCP(120),XOC(120)
0014      DIMENSION MMX(120),SINK(120),STRML(120),STRMW(120)
0015      DIMENSION BQSS(120),ELMNT(120),LENTH(120,120)
0016      DIMENSION XIF(120),YIF(120),THETA(120),BDFI(120)
0017      DIMENSION WIDTH(120,120),VEL(120),ELT(120),BETA(120)
0018      DIMENSION OLDL(120),OLDW(120),CQSS(120,120)
0019      DIMENSION QLENTH(120,120),QWIDTH(120,120)
0020      INTEGER SINK
0021      REAL LENTH
C
C...THIS SUBROUTINE CALCULATES THE STREAMLINES THAT ORIGINATED
C....IN THE ADJACENT REGION INTO THE REGION UNDER CONSIDERATION
C....THEIR STARTING POINTS ARE ON THE INTERFACE BOUNDARY
C
0022      FACT=1.0333E+7
0023      WRITE(IMP,107)THICK,POR
0024      107  FORMAT(1H1,/,20X,'CONTINUATION IN ADJACENT REGION',/,20X,
&      'THICKNESS=',F10.4,/,20X,'POROSITY=',F10.4,/,/)
C
0025      WRITE(IMP,64)
0026      64  FORMAT(30X,'INTERFACE POINTS',/,2X,'COORDINATES',
&      3X,'VELOCITY',2X,'BOUNDARY CONDS.',5X,'ANGLE',
&      2X,'FLOW RATE AND NO OF STRMLNS'//)
0027      DO 66 I=1,IPC
0028      BDFI(I)=-1.0*(TADJ/T)*BDFI(I)
0029      VEL(I)=-1.0*(T/(VICSTM*VICFAC))*BDFI(I)*FACT
0030      WRITE(IMP,63)XOC(I),YOC(I),VEL(I),BDFI(I),BETA(I),
&      BQSS(I),BNSL(I)
0031      63  FORMAT(7F10.4)
0032      66  CONTINUE
C
C CALCULATE THE VELOCITIES NEAR THE PRODUCING WELLS
C
0033      CALL LOWVEL(NSS,QSS,FI,DFI,X,Y,XSS,YSS,N,T,VLOW,RI,
&      QMNI,RAD,POR,THICK)
C
C INITIALIZE ARRAYS
C
0034      DO 206 I=1,NSS
0035      SINK(I)=0.0
0036      206  CONTINUE
C

```

C...XOC,YOC ARE THE INTERFACE BOUNDARY STARTING COORDINATES
 C...XI,YI ARE UNSCALED INTERNAL CALCULATION COORDINATES
 C...XIF,YIF ARE SCALED COORDINATES USED ONLY FOR OUTPUT
 C

```

0037      DO 210 K=1,IPC
0038      ICANT=0
0039      FACPLD=2.0
0040      CALL FACTOR(FACPLD)
0041      CALL PLOT(XOC(K)/SCLFAC,YOC(K)/SCLFAC,3)
0042      STRML(K)=0.0
0043      STRMW(K)=0.0
0044      AVWIDT(K)=0.0
0045      171 ICANT=ICANT+1
0046      IF(ICANT .GT. 100) GO TO 210
0047      VX(K)=0.
0048      VY(K)=0.
0049      VT(K)=0.
0050      DPDX(K)=0.
0051      DPDY(K)=0.
0052      IF(ICANT .GT. 1)GO TO 501
0053      ANGLE=BETA(K)
0054      VX(K)=VEL(K)*COS(ANGLE)
0055      VY(K)=VEL(K)*SIN(ANGLE)
0056      VT(K)=SQRT(VX(K)**2+VY(K)**2)
0057      STRML(K)=0.0
0058      STRMW(K)=STRMW(K)+BQSS(K)/(BNSL(K)*THICK*POR*VT(K))
0059      AVWIDT(K)=STRMW(K)/ICANT
0060      DT=50.0*RI/ABS(VT(K))
0061      DO 502 L=1,1000
0062      DTD=DT+(L*0.005*DT)
0063      XI(K)=XOC(K)+(VX(K)*DTD)
0064      YI(K)=YOC(K)+(VY(K)*DTD)
0065      XIF(K)=XI(K)/SCLFAC
0066      YIF(K)=YI(K)/SCLFAC
0067      IF(L .GT. 10)GO TO 503
0068      WRITE(IMP,646)VX(K),VY(K),XIF(K),YIF(K)
0069      646 FORMAT(/2X,4F10.4)
0070      503 CONTINUE
0071      BDIST=SQRT((XI(K)-XOC(K))**2+(YI(K)-YOC(K))**2)
0072      IF(BDIST .GT. ELT(K))GO TO 608
0073      502 CONTINUE
0074      WRITE(IMP,609)XIF(K),YIF(K)
0075      609 FORMAT(/2X,'POSITION AFTER 100 TRIALS IS',2F10.4,/)
0076      GO TO 210
0077      608 CONTINUE
0078      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0079      XIF(K)=XI(K)/SCLFAC
0080      YIF(K)=YI(K)/SCLFAC
C      WRITE(IMP,606)XIF(K),YIF(K)
C 606 FORMAT(/2X,'XNEW=',F10.4,5X,'YNEW=',F10.4,/)
0081      501 DO 172 J=1,N
0082      DISMID(J)=SQRT((XI(K)-XM(J))**2+(YI(K)-YM(J))**2)
0083      CALL INTE(XI(K),YI(K),X(J),Y(J),X(J+1),Y(J+1),H(K,J),
& G(K,J),DGDY(J),DGDY(J),DHDX(J),DHDY(J),DIST(J),
& ELMNT(J),THETA(J))
0084      ELMNT2=0.5*ELMNT(J)
C..IF POINT IS NEAR ANY BOUNDARY, STOP AND START NEW
C..STREAMLINE

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0085      IF(DISMID(J) .LT. ELMNT2)GO TO 210
0086      360  DPDX(K)=DPDX(K)+DFI(J)*DGDY(J)-FI(J)*DHDX(J)
0087      DPDY(K)=DPDY(K)+DFI(J)*DGDY(J)-FI(J)*DHDY(J)
0088      172  CONTINUE
C
C CALCULATE THE CONTRIBUTIONS OF THE SOURCES AND SINKS
C
0089      DXBLEN(K)=0.
0090      DYBLEN(K)=0.
0091      DO 54 JJ=1,NSS
0092      CALL SOURCE(XI(K),YI(K),XSS(JJ),YSS(JJ),QSS(JJ),NSS,
&      BLEN,DXSURS,DYSURS)
0093      DXBLEN(K)=DXBLEN(K)+DXSURS
0094      DYBLEN(K)=DYBLEN(K)+DYSURS
0095      54  CONTINUE
0096      DPHIDX(K)=(DPDX(K)+DXBLEN(K))/(2.0*PI)*FACT
0097      DPHIDY(K)=(DPDY(K)+DYBLEN(K))/(2.0*PI)*FACT
0098      VX(K)=-1.0*(T/VICSTM*VICFAC)*DPHIDX(K)
0099      VY(K)=-1.0*(T/VICSTM*VICFAC)*DPHIDY(K)
0100      VT(K)=SQRT(VX(K)**2+VY(K)**2)
C...CHECK IF POINT IS NEAR A PRODUCER, IF SO BRANCH TO 333
C...OTHERWISE CALCULATE THE LENGTH AND WIDTH AT THE CURRENT
C...POSITION. PLOT THE CURRENT POSITION
0101      DT=50.0*RI/ABS(VT(K))
0102      IF(VT(K) .GT. VLOW) GO TO 333
0103      111  STRML(K)=STRML(K)+(VT(K)*DT)
0104      STRMW(K)=STRMW(K)+BQSS(IPC)/(BNSL(IPC)*THICK*POR*VT(K))
0105      AVWIDT(K)=STRMW(K)/ICANT
C...CALCULATE NEW POSITION
0106      XI(K)=XI(K)+VX(K)*DT
0107      YI(K)=YI(K)+VY(K)*DT
0108      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0109      XIF(K)=XI(K)/SCLFAC
0110      YIF(K)=YI(K)/SCLFAC
0111      WRITE(IMP,332)K,XIF(K),YIF(K),VX(K),VY(K),VT(K)
0112      332  FORMAT(6X,I3,2X,5E10.3)
0113      GO TO 171
0114      333  DO 200 LL=1,NSS
0115      IF(QSS(LL) .GE. 0.0)GO TO 200
0116      RO(LL)=SQRT((XI(K)-XSS(LL))**2+(YI(K)-YSS(LL))**2)
0117      IF(RAD .GE. RO(LL)) GO TO 208
0118      200  CONTINUE
0119      GO TO 111
C...IDENTIFY THE PRODUCER AT WHICH THE STREAMLINE TERMINATED
0120      208  SINK(LL)=SINK(LL)+1
0121      LAST=SQRT((XI(K)-XSS(LL))**2+(YI(K)-YSS(LL))**2)
0122      LENTH(LL,SINK(LL))=STRML(K)+LAST
0123      WLAST=(2.0*PI*RI)/BNSL(IPC)
0124      STRMW(K)=STRMW(K)+WLAST
0125      AVWIDT(K)=STRMW(K)/(FLOAT(ICANT)+1.0)
0126      WIDTH(LL,SINK(LL))=AVWIDT(K)
0127      OLENTH(LL,SINK(LL))=OLDL(K)
0128      ODWIDTH(LL,SINK(LL))=OLDW(K)
0129      CQSS(LL,SINK(LL))=BQSS(K)/BNSL(K)
0130      CALL PLOT(XI(K)/SCLFAC,YI(K)/SCLFAC,2)
0131      210  CONTINUE
C...WRITE OUT THE ANSWERS
0132      DO 440 I=1,NSS

```

```
0133      IF(QSS(I) .GE. 0.0) GO TO 440
0134      WRITE(IMP,441)I,SINK(I)
0135      441  FORMAT(1H1,15X,'PRODUCER NUMBER',I3,///15X,
           &  'NUMBER OF STREAMLINES',I3,///)
0136      WRITE(IMP,443)
0137      443  FORMAT(8X,'STREAMLINE NUMBER',3X,'TOTAL LENGTH',3X,
           &  'AVERAGE WIDTH',6X,'INJ. RATE',///)
0138      ITEMP=SINK(I)
0139      IF(ITEMP .LE. 0) GO TO 440
0140      DO 444 J=1,ITEMP
0141      LENTH(I,J)=LENTH(I,J)/DISFAC
0142      WIDTH(I,J)=WIDTH(I,J)/DISFAC
0143      OLENTH(I,J)=OLENTH(I,J)/DISFAC
0144      OWIDTH(I,J)=OWIDTH(I,J)/DISFAC
0145      CQSS(I,J)=CQSS(I,J)/RATEF
0146      WRITE(IMP,445)J,OLENTH(I,J),OWIDTH(I,J),CQSS(I,J)
0147      WRITE(IMP,445)J,LENTH(I,J),WIDTH(I,J),CQSS(I,J)
0148      445  FORMAT(14X,I3,8X,E12.4,6X,E12.4,5X,E12.4)
0149      444  CONTINUE
0150      440  CONTINUE
0151      272  CONTINUE
0152      RETURN
0153      END
```

```

0001 DIMENSION RL1(30,30),TID(30),TD(10),WR1(30,30)
0002 DIMENSION RHOCR(30),PERM(4),WIDTH(30),FB(30,30)
0003 DIMENSION AVSATW(4),AVSATG(4),AVSATG(4),FA(30,30)
0004 DIMENSION DENR(4),CR(4),POR(4),THICKR(4),LAMBDA(4)
0005 DIMENSION VA(30,30),VB(30,30),INJECA(30,30)
0006 DIMENSION DELTD(30),RTIME(30),RL2(30,30),WR2(30,30)
0007 DIMENSION DIST1(30,30),CUMREC(30),QST(30)
0008 DIMENSION INJECB(30,30),NSTLINE(30),OILREC(30,30)
0009 DIMENSION TOTREC(30),ENDTM1(30,30),ENDTM2(30,30)
0010 DIMENSION TZTD1(30,30),DIST2(30,30),CODE(30,30)
0011 DIMENSION TZTD2(30,30),SWI(30),SDI(30),SGI(30)
0012 DIMENSION QSTRM1(30,30),QSTRM2(30,30),END1(30,30)
0013 DIMENSION END2(30,30)
0014 INTEGER TITLE1(18),TITLE2(18),TITLE3(18),TITLE4(18)
0015 INTEGER TITLES(18),TITLE6(18),TITLE7(18)
0016 REAL LAMBDA,LAMBDA1,LAMBDA2,INJEC,INJEC1,INJEC2,INJEC3
0017 REAL INJECA,INJECA,INJECB
0018 INTEGER CODE
C*****
C..OBJECTIVE: TO CALCULATE OIL RECOVERY BY STEAMFLOODING
C IN STREAMTUBES THAT CONSIST OF ONE OR TWO
C PERMEABILITY REGIONS.
C*****
C
C..DEFINE STATEMENT FUNCTIONS AND SUBROUTINE FUNCTIONS USED
C
0019 RTD(A,B)=(2.0*SQRT((A**2)*B)/1.7725)-1.0
0020 VELTD(A,B)=EXP((A**2)*B)*ERFC(SQRT((A**2)*B))
0021 ZTD(A,B,C)=C*(VELTD(A,B)+RTD(A,B))
0022 VTD(A,B,D)=D*VELTD(A,B)
C
C..READ THE TITLES FOR THE JOB
C
0023 READ(5,1005)TITLE1
0024 READ(5,1005)TITLE2
0025 READ(5,1005)TITLE3
0026 READ(5,1005)TITLE4
0027 READ(5,1005)TITLE5
0028 READ(5,1005)TITLE6
0029 READ(5,1005)TITLE7
0030 FORMAT(18A4)
0031 WRITE(6,1006)TITLE1
0032 WRITE(6,1007)TITLE2
0033 WRITE(6,1007)TITLE3
0034 WRITE(6,1007)TITLE4
0035 WRITE(6,1007)TITLE5
0036 WRITE(6,1007)TITLE6
0037 WRITE(6,1007)TITLE7
0038 FORMAT(///10X,18A4)
0039 FORMAT(//10X,18A4)
C
C..... READ AND WRITE INITIAL AND FIXED VALUES
C
0040 READ(5,5)NPROD,PST,TST,TI,X,EPS,THICK,NREG
0041 FORMAT(15,6F10.4,11)
0042 WRITE(6,6)NPROD,TST,X,EPS,THICK,PST,TI
0043 FORMAT(///40X,'DATA',/39X,'-----',/6X,
6 'TOTAL NUMBER OF PRODUCERS=',15X,13,/6X,
6

```

```

& 'STEAM TEMPERATURE (DEG. F)=' ,14X,F10.4,/6X,
& 'STEAM QUALITY=' ,25X,F10.4,/6X,'CONVERGENCE LIMIT=' ,
& 21X,F10.4,/6X,'RESERVOIR THICKNESS (FT)=' ,15X,F10.4/
& 6X,'STEAM PRESSURE (PSI)=' ,20X,F10.4,/6X,
& 'INITIAL RESERVOIR TEMPEFATURE(DEG F)=' ,4X,F10.4,/)

```

C

C..... READ AND WRITE FLUID PROPERTIES

C

```

0044      READ(5,10)DENW1,DENO1,DENW2,DENO2
0045      10      FORMAT(4F10.4)
0046      READ(5,12)CW,CG,CG
0047      12      FORMAT(3F10.4)

```

C

C CALCULATE THE DENSITY AND LATENT HEAT OF STEAM USING
C CORRELATIONS

C

```

0048      AAA=-0.9588
0049      BBB=-0.08774
0050      DENST=1.0/(363.9*(PST**AAA))
0051      STLAT=1318.0*(PST**BBB)

```

C

```

0052      WRITE(6,20)DENW1,DENO1,DENW2,DENO2,DENST,CW,CG,CG,STLAT
0053      20      FORMAT(/40X,'FLUID PROPERTIES',/39X,'-----'
& //,44X,'WATER',8X,'OIL',8X,'STEAM',/43X,'-----',6X,
& '-----',
& 6X,'-----',/6X,'DENSITY(LB/CU FT) AT STD. TEMP.',2X,
& F10.4,2X,F10.4,/6X,'DENSITY(LB/CU FT) AT STEAM TEMP.',1X
& F10.4,2X,F10.4,2X,F10.4,/6X,
& 'SPECIFIC HEAT(BTU/LB *F)',8X,F10.4,2X,F10.4,3X,F10.4,/6X,
& 'LATENT HEAT(BTU/LB)',40X,F10.4,/)

```

C

C.....READ AND WRITE THE PROPERTIES OF THE CAP AND BASE ROCK

C

```

0054      READ(5,70)DENCE,CCB,CBK
0055      70      FORMAT(3F10.4)
0056      WRITE(6,72)DENCE,CCB,CBK
0057      72      FORMAT(30X,'PROPERTIES OF THE CAP AND BASE ROCK',/29X,
& '-----',/6X,
& 'DENSITY(LB/CU. FT)',
& 8X,F10.4,/6X,'SPECIFIC HEAT(BTU/LB*F)',F10.4,/6X,
& 'THERM. COND.(BTU/HR-FT*F)',F10.4/)

```

C

C.....READ AND WRITE RESERVOIR ROCK PROPERTIES

C

```

0058      DO 43 I=1,NREG
0059      READ(5,44)POR(I),DENR(I),CR(I),PERM(I)
0060      44      FORMAT(4F10.4)
0061      43      CONTINUE
0062      WRITE(6,50)
0063      50      FORMAT(1H1,////,30X,'RESERVOIR ROCK PROPERTIES',/29X,
& '-----',/31X,'REGION 1',7X,
& 'REGION 2',/30X,'-----',5X,'-----',/)
0064      WRITE(6,60)(POR(I),I=1,NREG),(DENR(I),I=1,NREG),
& (CR(I),I=1,NREG),(PERM(I),I=1,NREG)
0065      60      FORMAT(/6X,'POROSITY',13X,2(F10.4,7X),/6X,
& 'DEN.(LB/CU FT)',8X,2(F10.4,7X),/6X,'SPEC. HEAT(BTU/LB*F',
& ,2X,2(F10.4,7X),/6X,'PERMEABILITY(MD)',6X,2(F10.4,7X),/)

```

C

C READ AND WRITE INITIAL FLUID SATURATIONS

C

```

0066      READ(5,38)(SWI(I),I=1,2)
0067      READ(5,38)(SOI(I),I=1,2)
0068      READ(5,38)(SGI(I),I=1,2)
0069      38  FORMAT(2F10.4)
0070      WRITE(6,39)(SWI(I),I=1,NREG),(SOI(I),I=1,NREG),
&      (SGI(I),I=1,NREG)
0071      39  FORMAT(/20X,'INITIAL FLUID SATURATIONS',/19X,
&      '-----',
&      //,23X,'REGION 1',3X,'REGION 2',/22X,'-----',1X,
&      '-----',
&      //16X,'WATER',2(F10.4),/16X,'OIL',2X,2(F10.4),/16X,
&      'STEAM',2(F10.4),//)

```

C

C.....READ AND WRITE THE AVERAGE RESIDUAL SATURATIONS AFTER
C STEAM FLOODING FOR EACH REGION

C

```

0072      READ(5,82)(AVSATW(I),I=1,2)
0073      READ(5,82)(AVSATO(I),I=1,2)
0074      READ(5,82)(AVSATG(I),I=1,2)
0075      82  FORMAT(2F10.4)
0076      WRITE(6,84)(AVSATW(I),I=1,NREG),(AVSATO(I),I=1,NREG),
&      (AVSATG(I),I=1,NREG)
0077      84  FORMAT(/20X,'AVERAGE RESIDUAL SATURATIONS BEHIND THE '
&      'FRONT',/19X,'-----'
&      '-----',//,23X,'REGION 1',3X,'REGION 2',/22X,
&      '-----',1X'-----',
&      //16X,'WATER',2(F10.4),/16X,'OIL',2X,2(F10.4),/16X,
&      'STEAM',2(F10.4),//)

```

C

C...READ AND WRITE INPUT VALUES FROM THE STREAMLINE PROGRAM

C

```

0078      WRITE(6,18)
0079      18  FORMAT(1H1,///20X,'INPUT DATA FROM STREAMLINE PROGRAM',
&      /19X,'-----',//32X,'REGION 1',
&      20X,'REGION 2'/31X,'-----',19X,'-----',/5X,
&      'WELL NO.',2X,'S/L NO.',2X,'CODE',2X,'LENGTH',
&      2X,'WIDTH',2X,'RATE',6X,'LENGTH',2X,'WIDTH',2X,'RATE'//)
0080      DO 45 M=1,NPROD
0081      READ(5,13)NSTLNE(M),QST(M)
0082      13  FORMAT(13,F10.4)
0083      NTEMP=NSTLNE(M)
0084      DO 15 K=1,NTEMP
0085      READ(5,16)CODE(M,K),RL1(M,K),WR1(M,K),QSTRM1(M,K)
0086      16  FORMAT(15,3F10.4)
0087      IF(CODE(M,K).EQ.1)GO TO 131
0088      READ(5,129)RL2(M,K),WR2(M,K)
0089      129  FORMAT(2F10.4)
0090      QSTRM2(M,K)=QSTRM1(M,K)
0091      GO TO 415
0092      131  RL2(M,K)=0.0
0093      WR2(M,K)=0.0
0094      QSTRM2(M,K)=0.0
0095      415  WRITE(6,17)M,K,CODE(M,K),RL1(M,K),WR1(M,K),QSTRM1(M,K)
&      ,RL2(M,K),WR2(M,K),QSTRM2(M,K)
0096      17  FORMAT(5X,15,3X,15,3X,15,3X,3F7.2,4X,3F7.2)
0097      15  CONTINUE

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```

0098          45      CONTINUE
C
C.....CALCULATE THE DIMENSIONLESS CONSTANTS
C
0099          WRITE(6,98)
0100          98      FORMAT(1H1,/,30X,'CALCULATED RESULTS',/30X,'-----'
&              ,',-----',/10X,'REGION',7X,'LAMDA',5X,
&              'DIMENSIONLESS TIME FACTOR',/10X,'-----',5X,'-----',
&              3X,'-----',/))
0101          RHOCB=DENCB*CCB
0102          TDFAC=4.0*CBK/(DENCB*CCB*THICK**2)
0103          DELTST=TST-TI
0104          DO 94 I=1,2
0105          RHOCR(I)=DENR(I)*CR(I)*(1-POR(I))+(POR(I)*DENW2*CW*
&              AVSATW(I))+(POR(I)*DENO2*CO*AVSATD(I))+(POR(I)*DENST
&              *STLAT*AVSATG(I)/DELTST)
0106          LAMDA(I)=RHOCB/RHOCR(I)
0107          94      CONTINUE
0108          DO 99 I=1,2
0109          WRITE(6,96)I,LAMDA(I),TDFAC
0110          96      FORMAT(12X,I2,6X,E10.4,7X,E10.4)
0111          99      CONTINUE
C
C...INITIALIZE ARRAYS
0112          DO 122 M=1,NPROD
0113          CUMREC(M)=0.0
0114          NTEMP=NSTLNE(M)
0115          DO 123 K=1,NTEMP
0116          OILREC(M,K)=(1.0/5.615)*0.0
0117          ENDTM1(M,K)=0.0
0118          ENDTM2(M,K)=0.0
C
C CONVERT THE INJECTION RATE FROM BARRELS/DAY
C (WATER EQUIVALENT) TO CUBIC FT PER HOUR
C
0119          QSTRM1(M,K)=QSTRM1(M,K)*5.615*(64.5/DENST)/24.0
0120          QSTRM2(M,K)=QSTRM2(M,K)*5.615*(64.5/DENST)/24.0
0121          123      CONTINUE
0122          122      CONTINUE
C
C.....FOR EACH STREAMTUBE,CALCULATE THE TIMES TO TRAVERSE THE
C.....FIRSTND SECOND REGIONS AND STORE IN ENDTM1 AND ENDTM2
C.....RESPECTIVELY
C...TTDI=DIMENSIONLESS TIME VALUES OBTAINED DURING ITERATION
C.....TDI IS AN INITIAL ARBITRARY GUESS
C.....TTDI ARE IMPROVED ITERATES.
C
C...ASSUME AN INITIAL TIME GUESS OF T=1.0
0123          TIME=1.0
0124          TDI=TIME*TDFAC
0125          ITER=30
0126          126      DO 124 M=1,NPROD
0127          NTEMP=NSTLNE(M)
0128          DO 125 K=1,NTEMP
0129          FA(M,K)=X*DENST*(STLAT+(CW*DELTST))+(1.0-X)*DENW2
&              *CW*DELTST
0130          INJCA(M,K)=FA(M,K)*QSTRM1(M,K)*THICK/(WR1(M,K)
&              *LAMDA(I)*4.0*CBK*DELTST)

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0131      VA(M,K)=FA(M,K)*QSTRM1(M,K)*LAMDA(1)/(WR1(M,K)*THICK*
      & RHOCB*DELTST)
C
C SINCE FORTGCLG COMPILER DOES NOT ALLOW ARRAY NAMES IN THE
C STATEMENT FUNCTION PARAMETER LIST, REASSIGN ARRAYS TO
C VARIABLES
C
0132      LAMDA1=LAMDA(1)
0133      INJEC1=INJEC(M,K)
0134      INJEC3=VA(M,K)
0135      102 TTDI=0.0
0136      TTDI=RL1(M,K)/VTD(LAMDA1,TDI,INJEC3)
C      QQ=VTD(LAMDA1,TDI,INJEC3)
C      WRITE(6,723)FA(M,K),LAMDA1,QSTRM1(M,K),INJEC1,WR1(M,K)
C      & ,INJEC3,TTDI,QQ
C 723      FORMAT(/10X,'FA(M,K)=',F12.4,/10X,'LAMDA1=',F10.4,/10X,
C      & 'QSTRM1=',F12.8,/10X,'INJEC1=',F12.8,/10X,'WR1(M,K)=',
C      & F12.8,/10X,'INJEC3=',F12.8,/10X,'TTDI=',E12.4,/10X,
C      & 'QQ=',F12.8//)
0137      TTDI=TTDI*TDFAC
0138      DO 116 J=1,ITER
0139      IF(TTDI .LT. 10.0)GO TO 724
0140      WRITE(6,505)
0141      505      FORMAT(/10X,'VELOCITY WILL BE ZERO BEFORE TUBE END'//)
0142      GO TO 304
0143      724      DELTD(J)=-1.0*(ZTD(LAMDA1,TTDI,INJEC1)-RL1(M,K))/
      & VTD(LAMDA1,TTDI,INJEC3)
0144      725      DELTD(J)=DELTD(J)*TDFAC
0145      TTDI=TTDI+DELTD(J)
0146      IF(ABS(DELTD(J)/TTDI) .LT. EPS) GO TO 118
0147      116      CONTINUE
0148      WRITE(6,117)
0149      117      FORMAT(/2X,'NO CONVERGENCE'//)
0150      GO TO 304
0151      118      ENDTM1(M,K)=TTDI
0152      TZTD1(M,K)=ZTD(LAMDA1,TTDI,INJEC1)
0153      324      IF(CUDE(M,K) .LE. 1)GO TO 125
C...CALCULATE THE TIME TO TRAVERSE THE SECOND REGION(ENDTM2)
0154      FB(M,K)=X*DENST*(STLAT+(CW*DELTST))+(1.0-X)*DENW2*CW*
      & DELTST
0155      INJECB(M,K)=FB(M,K)*QSTRM2(M,K)*THICK/(WR2(M,K)*
      & LAMDA(2)*4.0*CBK*DELTST)
0156      VB(M,K)=FB(M,K)*QSTRM2(M,K)*LAMDA(2)/(WR2(M,K)*
      & THICK*RHCCB*DELTST)
0157      LAMDA2=LAMDA(2)
0158      INJEC2=INJECB(M,K)
0159      INJEC4=VB(M,K)
0160      TDEL=RL2(M,K)/VTD(LAMDA2,TTDI,INJEC4)
0161      TDEL=TDEL*TDFAC
0162      TTDI=ENDTM1(M,K)+TDEL
0163      TINIT=ENDTM1(M,K)
0164      DO 662 J=1,ITER
0165      IF(TTDI .LT. 10.0)GO TO 726
0166      WRITE(6,506)
0167      506      FORMAT(/10X,'VELOCITY WILL BE ZERO BEFORE TUBE END'//)
0168      GO TO 304
0169      726      DELTD(J)=-1.0*(ZTD(LAMDA2,TTDI,INJEC2)
      & -RL2(M,K))/VTD(LAMDA2,TTDI,INJEC4)

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0170      727   TTDI=TTDI+DELTD(J)
0171      IF(ABS(DELTD(J)/TTDI) .LT. EPS)GO TO 663
0172      662   CONTINUE
0173      WRITE(6,664)
0174      664   FORMAT(//2X,'NO CONVERGENCE',//)
0175      GO TO 304
0176      663   ENDTM2(M,K)=ENDTM1(M,K)+TTDI
0177      TZTD2(M,K)=ZTD(LAMDA2,TTDI,INJEC2)
0178      125   CONTINUE
0179      124   CONTINUE
0180      WRITE(6,36)
0181      36     FORMAT(1H1///10X,'CALCULATED BREAKTHROUGH TIMES(HOURS)'
&          ' FOR EACH STREAMTUBE',2)
0182      WRITE(6,149)
0183      149   FURMAT(//.6X,'PROD. NO.',3X,'STREAMLINE NO.',3X,'CODE',
&          3X,'ENDTIME(1)',4X,'ENDTIME(2)')//)
0184      DO 134 M=1,NPROD
0185      NTEMP=NSTLNE(M)
0186      DO 135 K=1,NTEMP
0187      END1(M,K)=ENDTM1(M,K)/TDFAC
0188      END2(M,K)=ENDTM2(M,K)/TDFAC
0189      136   WRITE(6,137)M,K,CODE(M,K),END1(M,K),END2(M,K)
0190      137   FORMAT(4X,I5,12X,I5,6X,I5,4X,2(F10.4,3X))
0191      135   CONTINUE
0192      134   CONTINUE

```

C

C...CALCULATE THE VELOCITY,DISTANCE,AND OIL RECOVERY AT
C...VARIOUS TIMES

```

0193      READ(5,14)TIME
0194      14     FORMAT(F10.2)
0195      CTIME=TIME/20.0
0196      DO 923 KK=1,20
0197      TIME=CTIME*KK
0198      TIMED=TIME*TDFAC
0199      WRITE(6,323)TIME,TIMED
0200      323   FORMAT(1H1///25X,'REAL TIME(HOURS)=',F10.4,//
&          25X,'DIMENSIONLESS TIME=',F10.4,//2X,'WELL',2X,
&          'STREAMLINE',2X,'NO. OF REGIONS',2X,'RECOVERY',5X,
&          'WELL TOTAL',5X,'RESERVOIR TOTAL'//)
0201      TOTREC(KK)=0.0
0202      DO 432 M=1,NPROD
0203      NTEMP=NSTLNE(M)
0204      CUMREC(M)=0.0
0205      DO 433 K=1,NTEMP
0206      IF(TIMED .GT. ENDTM1(M,K))GO TO 260
0207      TDEL=TIMED
0208      INJEC1=INJEC(M,K)
0209      DIST1(M,K)=ZTD(LAMDA1,TDEL,INJEC1)
0210      DIST2(M,K)=0.0
0211      OILREC(M,K)=(1.0/5.615)*DIST1(M,K)*WR1(M,K)
&          *THICK*POR(1)*(SOI(1)-AVSATO(1))
0212      CUMREC(M)=CUMREC(M)+OILREC(M,K)
0213      GO TO 777

```

C...CALCULATE THE BREAKTHROUGH RECOVERY IN FIRST REGION

```

0214      260   TDEL=ENDTM1(M,K)
0215      INJEC1=INJEC(M,K)
0216      DIST1(M,K)=RL1(M,K)
0217      OILREC(M,K)=(1.0/5.615)*DIST1(M,K)*WR1(M,K)

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      6 *THICK*POR(1)*(SOI(1)-AVSATO(1))
0218      IF(CODE(M,K) .GT. 1.0)GO TO 602
0219      DIST2(M,K)=0.0
0220      CUMREC(M)=CUMREC(M)+OILREC(M,K)
0221      GO TO 777
0222      602 TEMREC=OILREC(M,K)
0223      IF(TIMED .GT. ENDTM2(M,K))GO TO 665
0224      TIME2=TIMED
0225      TINIT=ENDTM1(M,K)
0226      INJEC2=INJECB(M,K)
      C...CALCULATE THE RECOVERY FROM SECOND REGION AND ADD TO
      C...FIRST
0227      DIST2(M,K)=ZTD(LAMDA2,TIME2,INJEC2)
      6 -ZTD(LAMDA2,TINIT,INJEC2)
0228      OILREC(M,K)=(1.0/5.615)*DIST2(M,K)*WR2(M,K)
      6 *THICK*POR(2)*(SOI(2)-AVSATO(2))+TEMREC
0229      CUMREC(M)=CUMREC(M)+OILREC(M,K)
0230      GO TO 777
      C... THERE IS BREAKTHROUGH IN THE SECOND REGION OF A TWO
      C...PERMEABILITY REGION STREAMTUBE, CALCULATE RECOVERY .
      C...FROM BOTH REGIONS
0231      665 TIME2=ENDTM2(M,K)
0232      INJEC2=INJECB(M,K)
0233      DIST2(M,K)=RL2(M,K)
0234      OILREC(M,K)=(1.0/5.615)*DIST2(M,K)*WR2(M,K)
      6 *THICK*POR(2)*(SOI(2)-AVSATO(2))+TEMREC
0235      CUMREC(M)=CUMREC(M)+OILREC(M,K)
0236      777 WRITE(6,307)M,K, CODE(M,K),OILREC(M,K)
0237      307 FORMAT(15,I8,13,6X,E13.4)
0238      433 CONTINUE
0239      TOTREC(KK)=TOTREC(KK)+CUMREC(M)
0240      WRITE(6,671)CUMREC(M)
0241      671 FORMAT(45X,E13.4)
0242      432 CONTINUE
0243      WRITE(6,672)TOTREC(KK)
0244      672 FORMAT(60X,E13.4)
0245      923 CONTINUE
0246      WRITE(6,675)
0247      675 FORMAT(1H1,///,17X,'PREDICTED RECOVERY',///,10X,
      6 'TIME(DAYS)',10X,'RECOVERY(BBLS)'///)
0248      DO 674 KK=1,20
0249      TIME=CTIME*KK
0250      RTIME(KK)=CTIME*KK
      C...CONVERT THE TIME TO DAYS AND RECOVERY TO BARRELS
0251      RTIME(KK)=RTIME(KK)/24.0
0252      TOTREC(KK)=TOTREC(KK)/5.615
0253      WRITE(6,676)RTIME(KK),TOTREC(KK)
0254      676 FORMAT(9X,F10.4,10X,E13.4)
0255      674 CONTINUE
      C
      C....PLOT THE CALCULATED RESULTS
      C
0256      CALL PLOT(1.0,1.0,-3)
      C NP=NO OF INTERVALS OF THE SCALED PLOT
      C K=EVERY K-TH DATA WILL BE USED TO GENERATE THE SCALE
0257      NP=20
0258      K=1
0259      AXLEN=4.0

```

```
0260      CALL SCALE(RTIME,AXLEN,NP,K)
0261      XS=0.0
0262      YS=0.0
0263      D=0.0
0264      CALL AXIS(XS,YS,'TIME(DAYS)',-10,AXLEN,D,RTIME(21),
& RTIME(22))
0265      NP=20
0266      K=1
0267      AXLEN=4.0
0268      CALL SCALE(TOTREC,AXLEN,NP,K)
0269      D=90.0
0270      CALL AXIS(XS,YS,'CUMULATIVE OIL RECOVERY(BBLS)',23,
& AXLEN,D,TOTREC(21),TOTREC(22))
0271      CALL PLOT(0.0,0.0,3)
0272      DO 32 K=1,20
0273      CALL PLOT(RTIME(K)/RTIME(22),TOTREC(K)/TOTREC(22),2)
0274      32  CONTINUE
0275      CALL PLOT(0.0,0.0,999)
0276      304 STOP
0277      END
```

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