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# COOPERATIVE AERIAL MANIPULATION WITH FORCE CONTROL AND ATTITUDE STABILIZATION 

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#### Abstract

Ranging from autonomous flying cars, fixed wing and rotorcraft UAVs, there has been a tremendous interest in aerial robotics over the last decade. This thesis presents contributions to the state-of-art in cooperative payload transport with force synthesis and dynamic interaction using quadcopter UAVs. In this report, we consider multiple quadcopter aerial robots and develop decentralized force controller for them to manipulate a payload. We use quadcopters with a rigid link attached to it to collaborately manipulate the payload. We develop a dynamic model of the payload for both point mass and rigid body cases. We model the contact force between the agents and the payload as a mass spring model. This assumption is valid when the vehicles are connected to the payload via elastic cables or when the payload is flexible or surrounded by elastic bumper materials. We also extend our aerial manipulation system to a multi-link arm attached to the quadcopter.


We develop an adaptive decentralized control law for transporting a payload of unknown mass without explicit communication between the agents. Our controller ensures that all quadcopters and the payload asymptotically converges to a constant reference velocity. It also ensures that all of the forces applied to the payload converges to desired set-points. Desired thrusts and attitude angles are computed from the control algorithms and a low-level PD controller is implemented to track the desired commands for each quadcopter. The sum of the estimates of the unknown mass from all the agents converge to the true mass. We also employ a consensus algorithm based on connected graphs to ensure that each agent gets an equal share of the payload mass. Furthermore, we develop an orientation control algorithm that guarantees attitude stabilization of the payload. In particular, we develop time varying force set-points to enforce attitude regulation without any moment inputs from the quadcopters.

## TABLE OF CONTENTS

Chapter Page
1 INTRODUCTION ..... 1
1.1 Motivation ..... 2
1.2 Related Work ..... 3
1.3 Contribution ..... 5
1.4 Thesis Organization ..... 5
2 PROBLEM FORMULATION ..... 7
2.1 Payload Transportation with Quadcopter with rigid extension ..... 7
2.2 Kinematics ..... 7
2.3 Dynamics ..... 9
2.3.1 Singularities ..... 10
2.4 Control Objective ..... 11
2.5 Quadcopter Attitude Controller ..... 11
2.6 Geometric Control ..... 12
2.6.1 Related Works ..... 12
2.6.2 Geometric Control of Cooperative UAVs ..... 13
2.6.3 Implementation ..... 15
2.6.4 Desired Rotation matrix $R_{i}^{d}$ ..... 15
2.6.5 Desired rate of Rotation $\dot{R}_{i}^{d}$ ..... 16
2.6.6 $\ddot{R}_{i}^{d}$ ..... 17
3 NOMINAL DESIGN ..... 19
3.1 Force Control with a Payload of Known Mass ..... 19
3.1.1 Computation of $T_{i}^{\text {des }}$ and $\Phi_{i}^{\text {des }}$ ..... 20
3.1.2 Stability Analysis ..... 21
3.2 Simulation Results ..... 22
3.2.1 Simulation Environment ..... 22
3.2.2 Known Payload with a Predesigned $v^{d}$ with a PD Controller ..... 23
3.2.3 Known Payload with a Predesigned $v^{d}$ with a Geometric Controller ..... 27
4 COOPERATIVE AERIAL MANIPULATION OF A PAYLOAD WITH UNKNOWN MASS ..... 33
4.1 Force Control with a Payload of Unknown Mass ..... 34
4.1.1 Computation of $T_{i}^{\text {des }}$ and $\Phi_{i}^{\text {des }}$ ..... 34
4.1.2 Stability Analysis ..... 34
4.1.3 Simulation Results: Unknown Payload with a Predesigned $v^{d}$ ..... 37
4.2 Lift Consensus During Cooperative Aerial Load Transport ..... 41
4.2.1 Adaptation Law ..... 41
4.2.2 Force Control Law ..... 42
4.2.3 Computation of $T_{i}^{d e s}$ and $\Phi_{i}^{d e s}$ for a quadcopter ..... 43
4.2.4 Main Result ..... 43
4.2.5 Simulation Results: Force control of a payload with an unknown mass for a constant $v^{d}$ ..... 46
4.2.6 Simulation Results: Force Control of a load with an unknown mass for a time-varying $v^{d}(t)$ ..... 51
5 COOPERATIVE AERIAL MANIPULATION WITH ATTITUDE STABILIZATION ..... 57
5.1 Related Work ..... 57
5.2 Problem Formulation ..... 58
5.2.1 Kinematics ..... 58
5.2.2 Dynamics ..... 58
5.2.3 Control objective ..... 60
5.3 Control Design ..... 60
5.3.1 Computation of $T_{i}^{d e s}$ and $\Phi_{i}^{\text {des }}$ for a quadcopter ..... 62
5.4 Stability Analysis ..... 62
5.4.1 Closed-loop system and equilibria ..... 62
5.4.2 Main result ..... 63
5.4.3 Computation of $\dot{f}_{i}^{d}$ and $\dot{z}_{i}^{d}$ ..... 66
5.4.4 Useful identities ..... 68
5.5 Simulation Results ..... 68
5.5.1 Simulation Environment ..... 68
5.5.2 Quadcopter Attitude Controller ..... 69
5.5.3 Numerical Examples ..... 69
6 FORCE CONTROL OF PAYLOAD WITH UNKNOWN MASS USING QUADCOPTERS WITH MULTI-LINK ROBOT ..... 77
6.1 Payload Transportation with an Aerial Manipulator ..... 77
6.1.1 Kinematics ..... 77
6.1.2 Dynamics ..... 81
6.2 Control Objective ..... 84
6.3 Payload Transport with an Unknown Mass ..... 84
6.4 Simulation Results ..... 85
6.4.1 Simulation Environment ..... 85
6.4.2 Unknown Payload with a Predesigned $v^{d}$ ..... 86
6.4.3 Unknown Payload with a Time Varying $v^{d}$ ..... 89
7 EXPERIMENTAL RESULTS ..... 93
7.1 Mechanical Design of the Test Bed ..... 93
7.1.1 Quadcopter ..... 93
7.1.2 Robotic Manipulator Design ..... 93
7.2 Computation Hardware ..... 94
7.2.1 Microprocessor ..... 94
7.2.2 Communications ..... 94
7.2.3 Motors ..... 95
7.3 Initial Integration and test ..... 95
8 CONCLUSION AND FUTURE WORK ..... 97
REFERENCES ..... 98

## LIST OF TABLES

Table Page
3.1 Attitude controller gains. ..... 23
5.1 Attitude controller gains ..... 69

## LIST OF FIGURES

## Figure

2.1 Three quadcopters transport a common flexible load. Note that $a_{i}$ is the initial position of agent $i$. ..... 8
2.2 $Z Y X$ Euler angle representation using tripleangle toolbox [1] ..... 10
2.3 ZYX Euler Angle Representation and Gimbal lock. Note $\theta_{i}=90^{\circ}$ ..... 10
3.1 Desired and actual roll angles ( $\phi$ ) for both agents. ..... 24
3.2 Desired and actual pitch angles $(\theta)$ for both agents. ..... 24
3.3 The thrusts for both agents are bounded and converge to 14.7 N . ..... 25
3.4 Linear velocities for both agents. The $x$ component of the velocity con- verges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 25
3.5 Linear velocity for the payload. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 26
3.6 Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload. ..... 26
3.7 Linear velocities for both agents. The $x$ component of the velocity con- verges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 28
3.8 Linear velocity for the payload. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 28
3.9 Desired thrust input to each quadrotor. ..... 29
3.10 Control input $M_{1}$ ..... 29
3.11 Control input $M_{2}$ ..... 30
3.12 Control input $M_{3}$ ..... 30
3.13 Attitude error of the quadrotors given by $\Psi_{i}=\frac{1}{2}\left\|R_{i}-R_{I}^{d}\right\|^{2}$ ..... 31
3.14 Euler Angles for each agent. ..... 31
3.15 Contact forces. ..... 32
4.1 Desired and actual roll angles ( $\phi$ ) for both agents (unknown $M_{c}$ case). ..... 38
4.2 Desired and actual pitch angles $(\theta)$ for both agents (unknown $M_{c}$ case). ..... 38
4.3 The thrusts for both agents converge to constant and remain bounded. ..... 39
4.4 Linear velocities for both agents. The $x$ component of the velocity con- verges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 39
4.5 Linear velocities for both agents. The $x$ component of the velocity con- verges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 40
4.6 Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload. ..... 40
4.7 Estimation of the unknown mass. The sum of the individual estimates converges to the actual mass of the load which is 1.5 kg . ..... 41
4.8 Communication topology between each agent. ..... 46
4.9 The magnitude of the velocity error $\left(v_{i}-v^{d}\right)$ goes to zero which concludesthat $v_{i}$ converges to predesigned $v^{d}$474.10 Linear velocity plot for the payload. The $x$ and $y$ velocity converges to thedesired $v^{d}$ i.e., $2.0 \mathrm{~m} / \mathrm{s}$ and $1.0 \mathrm{~m} / \mathrm{s}$ respectively and the $z$ component con-verges to zero.48
4.11 The magnitude of $f_{i}-f_{i}^{d}$ goes to zero. All the contact forces $f_{i}$ converge to
$f_{i}^{d}$. ..... 484.12 Plot of the estimates of the mass of the payload. Each individual estimatesconverge to $\frac{1}{3} M_{c}=0.5 \mathrm{~kg}$ and its sum adds up-to the true mass of the pay-load i.e., 1.5 kg .49
4.13 Plot for contact forces in all 3 directions for agent 1 . It converges to con- verges $0.0 \mathrm{~N},-2.0 \mathrm{~N}$ and 0.5 N for each direction respectively. Remaining agents have similar graphs ..... 49
4.14 Plot for desired and actual orientation angles $(\phi)$ for three agents in the $x$-direction. ..... 50
4.15 Plot for the desired and actual orientation angles $(\theta)$ for three agents in the $y$-direction ..... 50
4.16 Snapshot of 3D position trajectories for all the agents and the payload at different time instance for adaptive design with time varying velocity. ..... 52
4.17 Position trajectories for aerial robots and the payload in $x$ and $y$ direction. ..... 52
4.18 The magnitude of the velocity error $\left(v_{i}-v^{d}\right)$ for the agents. ..... 53
4.19 The magnitude of the velocity error $\left(v_{c}-v^{d}\right)$ for the payload ..... 53
4.20 Linear velocity for the payload in all 3 directions ..... 54
4.21 The magnitude of $f_{i}-f_{i}^{d}$ stays bounded and eventually converges to zero. ..... 54
4.22 Plot of squeeze forces in all 3 directions for agent 1 . Note that other agents have similar plot ..... 55
4.23 Plot of the estimates of the mass of the payload. Eachg individual esti-mates converges $\frac{1}{3} M_{c}=0.5 \mathrm{~kg}$. Notice that their sum adds up-to the truemass of the payload used in this simulation. The convergence of the esti-mates to $\frac{1}{3} M_{c}=0.5$ is because the $z$-component of the desired velocity $v^{d}$is still constant.55
4.24 Plot for desired and actual orientation angles $(\phi)$ for three agents in the $x$-direction. ..... 56
4.25 Plot for the desired and actual orientation angles $(\theta)$ for three agents in the $y$-direction. ..... 56
5.1 The magnitude of $\tilde{z}_{i}$ for the three agents. ..... 70
5.2 The magnitude of the velocity error $\xi_{i}$ for the three agents and the payload. ..... 71
5.3 The angular velocity of the load $\Omega_{c}$ goes to zero. ..... 71
5.4 The magnitude of the orientation error $\left\|q_{\nu}^{e}\right\|$ for the payload. ..... 72
5.5 3D position trajectory for three agents and the payload at different time steps for the first 20 seconds. The changes of the position of the red and the light blue dots confirm that the payload is able to achieve the desired rotation. ..... 73
5.6 The magnitude of $\tilde{z}_{i}$ goes to zero for all three agents. ..... 73
5.7 The magnitude of the velocity error $\xi_{i}$ for the three agents and the payload. ..... 74
5.8 The angular velocity of the load $\Omega_{c}$ goes to zero. ..... 74
5.9 The magnitude of the orientation error $\left\|q_{\nu}^{e}\right\|$ for the payload. ..... 75
5.10 Desired and actual roll angles ( $\phi$ ) for three agents. ..... 75
5.11 Desired and actual pitch angles $(\theta)$ for three agents. ..... 76
6.1 Two aerial manipulators transport a common flexible load. Note that body frame $O_{c, i}$ coincides with $O_{1, j}$ (not shown in the figure) and $a_{i}$ is the initial position of agent $i$. ..... 78
6.2 Linear velocities for both agents. The $x$ component of the velocity con- verges to $0.2 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 87
6.3 Linear velocity for the payload. The $x$ component of the velocity converges to $0.2 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero. ..... 87
6.4 Estimation of the unknown mass. The sum of the individual estimates converges to the actual mass of the load. ..... 88
6.5 Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload. . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 88
6.6 X and Y position trajectories for both agents (time varying $v^{d}$ ). Both agents travel in the $x$-direction until 2 seconds and make a smooth turn towards the $y$-direction after 2 seconds. ..... 89
6.7 3D position trajectory for both agents and the payload at different time interval. Note that all 3 agents stay at an equal distance and will never col- lide. They approach the $y$-axis by being aligned in a straight line as shown in Fig. 6.6. Notice that if the quadrotors are at the same height, they must keep a gap and the manipulator must be large enough to avoid their blades clashing. ..... 90
6.8 Velocities of the 2 agents in all 3 directions. The $x$ and $y$ velocities are smooth during transition and the $z$ velocity is zero. ..... 91
6.9 Desired and actual velocities of the load in all 3 directions. ..... 91
6.10 Estimation of the unknown mass. The sum of individual estimate con- verges to the actual mass of the load, which is 0.5 kg . ..... 92
6.11 Contact forces of both agents in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload. ..... 92
7.1 Custom made quadcopter in the lab. ..... 93
7.2 CAD model of the tow link robotic arm. ..... 94
7.3 Custom made aerial manipulator. ..... 95
7.4 Different instances of initial aerial manipulation test. ..... 96

## CHAPTER 1

## INTRODUCTION

Aerial robotics has a tremendous potential to impact the world in a positive way. In this report, we focus on cooperative aerial manipulation using multi-rotors. In particular, we focus on quadcopters because of their popularity and accessibility.

Quadcopters are a type of small unmanned rotorcraft with four rotors. It is an underactuated mechanical system and has six degrees of freedom with only four control inputs. Thus, quadcopters are dynamically unstable. This rich dynamics and mechanical simplicity of quadcopters can be exploited for agile and dynamic maneuvers in constrained spaces.

Quadcopter aerial robots have become increasingly popular because of their superior mobility in three dimensional space and high thrust-to-weight ratio. Because of their size and hovering properties, quadcopters can operate in constrained spaces such as damaged buildings, adversarial and nuclear environments. They can be used to collect information from the environment and can also dynamically interact with the environment using grippers. Possible applications of such robots include disaster response, transport, photography, search and rescue operations, construction and so on.

Multiple quadcopters can be employed together to increase the capability of a single quadcopter, which enables better manipulation and control of heavier payloads for various applications. Cooperative aerial manipulation of a payload requires motion coordination, regulation of forces applied to the payload by the quadcopters, attitude regulation of the payload and multirotors and analysis of stability of the combined system. This thesis presents a framework for force regulation, motion coordination and orienta-
tion control of the payload during cooperative aerial manipulation.

### 1.1 Motivation

Ranging from autonomous flying cars, fixed wing and rotor-craft UAVs, there has been tremendous interest in aerial robotics over the last decade. Application of aerial robotic systems range from military and civilian sectors. The ability to manipulate in 3D allows multiple dynamic tasks like aerial gripping, manipulation, photography and so on. Using quadcopter aerial robots dynamic interactions, vertical take off and landing, hovering in constraint spaces are possible. Recently, we have already seen the commercial applications of such aerial robots in commercial sectors. If such commercial applications like air package delivery [2] is to be realistic, it should be able to handle both rigid and flexible payloads and the dynamic interaction between the quadcopter and payload. It is clear that the aerial robotics community lacks a general method of force and dynamic motion synthesis during cooperative payload transport.

Motivated by the current challenges and their potential applications, in this work, we design decentralized adaptive force control algorithm for multiple quadcopters to manipulate a load of an unknown mass. We develop a complete motion control algorithm for cooperative payload transport to address both the orientation control of the payload and the force control during cooperative aerial manipulation [3-6]. The developed algorithm achieves asymptotic tracking of a desired velocity, controls the applied forces to the payload to predefined set-points and ensures that the individual estimate of the mass from each agent converges to the total mass divided by the number of agents. We develop an attitude control algorithm for the payload using time-varying force control without no any moments input from the rotors.

### 1.2 Related Work

Because of the complexity of the problem, cooperative aerial manipulation has recently become a popular research area (see for example [7-11]). Two distinct types of framework exist in the literature. Inspired by the traditional approach of helicopterbased towing via cables for load transport, most of the work in cooperative aerial load transport has been focused on suspended payload transport using cables [7, 8, 11-13]. Recently, quadcopters are quipped with robotic manipulators and control laws are designed for dynamic interactions, perching and aerial gripping [14-17].

Load Transport via towed cables has been extensively studied by researchers worldwide [7,11, 12, 18]. Reference [11] develops a geometric controller to follow a reference trajectory of a suspended payload with multiple quadcopters. The authors in [12] develop quasi-static models for payload transport via cables at a desired pose. Reference [7] develops a complete dynamic model of cooperative aerial manipulation and exploits the differential flatness property of the system to generate smooth trajectories for path planning. Reference [8] develops a software architecture for planning, control, collision avoidance and deployment of cable suspended payloads using multiple helicopters. Recently, [13] develops a passivity based formation control for multiple UAVs cooperatively carrying a suspended payload.

Aerial load transport via towed cables is a hybrid system and there are problems with swing load attenuation. So, different control laws have to be designed when the cable is tout and flexible. One possible solution to this problem is attach a robotic manipulator to a Unmanned Aerial Vehicle (UAV) and develop an aerial manipulator (AM). Since each AM is a coupled system, the control problem now becomes more challenging but it introduces more dynamic features to the system such as better gripping and control of payload as opposed to aerial towing. In one the earlier works, the authors in [19] develop gripping mechanisms and control laws for cooperative grasping and manipulation of a group of quadcopters. In some recent work, references $[16,17,20]$ explore the problem of
cooperative aerial load transport by using multirotors attached with multi-DOF robotic arms and develop robust control and motion planning algorithms.

There also exist control strategies based on port Hamiltonian framework [21,22] towards control of aerial manipulators. Acosta et al. propose an Interconnection and Damping Assignment Passivity Based Control (IDA-PBC) based on energy shaping methodology by exploiting the passivity of aerial manipulator. Nonlinear matching types of controller developed in [22] utilizes the cyclo-passivity property of AMs and combines the robust controllers for the quadcopter and the manipulator by exploiting their stability margins and optimizing different prescribed criteria in real time.

On the other hand, Wang and Schwager [23] develop a scalable, decentralized control strategy by which a large number of ground mobile robots can manipulate a comparatively large object through a desired trajectory without explicit communication network among the robots. Reference [24] considers a decentralized velocity controller to transport a common payload by using a team of robots in $\mathbb{R}^{2}$ rigidly attached to the payload. The authors in [25] develop decentralized adaptive controller to transport a rigid payload in $\mathbb{R}^{2}$ or $\mathbb{R}^{3}$ by controlling the spatial velocity of the payload.

Inspired by aerial towing and integration of robotic manipulator we suggested cooperative aerial manipulation with rigid extension. This approach allows better control of the payload and there is no problem of swing load attenuation. We can also extend the rigid link to a multi-link robot and develop similar manipulation system. In aerial manipulator there is a lot of redundancy in the system, this approach helps us reduce the redundancy of the system and develop a unified framework for cooperative aerial manipulation. As a possible example of our suggestions, recently, reference [26] develops a distributed wrench controller that enables a group of rigidly attached quadcopters to manipulate heavy payloads along a predefined trajectory without communication between the robots.

### 1.3 Contribution

Most of the aforementioned references have not considered dynamic interactions and force control between the payload and the quadcopters. In this paper, we develop an unified framework for cooperative aerial manipulation. Reference [27] considers a ring shaped flexible structure with multiple attached flying vehicles and develops a LQR controller for linearized dynamics to track the desired pose of the payload. It is assumed that the vehicles can exert both forces and moments on the payload. However, we assume that the inputs from the quadcopters to the payload are only contact forces with no moments. In particular, we develop force, motion and orientation control algorithm that grantees force regulation, attitude stabilization and translational motion of the quadcopters and the payload. This thesis makes the following contributions:

- We develop adaptive decentralized controller for transport of unknown payload without communication between agents. The sum of estimates of each payload converges to the true mass.
- We develop force consensus along the vertical direction such that the agents share the mass of the payload equally.
- We develop a force and orientation control algorithm that guarantees attitude stabilization of the payload and the convergence of the contact force and the velocity of each quadcopter to desired setpoints. We develop time varying force setpoints to enforce the orientation regulation of the payload.


### 1.4 Thesis Organization

This thesis presents force, motion and orientation control algorithms during cooperative aerial manipulation. The remainder of the paper is structured as follows. In Chapter 2, we develop a complete dynamic model for payload transportation using point
mass model. In Chapter 3, we present nominal design for force control law while transporting a payload of known mass. In Chapter 4, we consider quadcopters with rigid extension and develop adaptive decentralize control laws. We also employ a consensus algorithm based on connected graphs to ensure that each agent gets an equal share of the payload mass. Desired thrusts and attitude angles are computed from the control algorithms and a low-level PD controller is implemented to track the desired commands for each quadcopter. Furthermore, in Chapter 5 we also develop orientation control algorithm that guarantees attitude stabilization of the payload. In particular, we develop time varying force setpoints to enforce attitude regulation. In Chapter 6, we developed complete dynamic model for payload transport using aerial manipulator. Using the inverse kinematics of the aerial manipulator, we implement the developed algorithm in Chapter 4 at the kinematic level for the aerial manipulators and demonstrate its effectiveness in simulations. We include the preliminary experimental results in Chapter 7. We present conclusion and future work in Chapter 8.

## CHAPTER 2

## PROBLEM FORMULATION

In this chapter we develop a dynamic model of the aerial manipulation system. We consider a group of quadcopters with a rigid extension and develop force control law without explicit communication between the agents.

### 2.1 Payload Transportation with Quadcopter with rigid extension

Consider $N$ agents holding a common flexible load as shown in Fig. $2.1(N=3)$. Each agent is a quadcopter with a rigid link extension. We model the payload as a point mass and assume that it is initially undeformed. Agent $i$ is attached to the load at the point $a_{i}$. Let $x_{i} \in \mathbb{R}^{3}$ be the position of the end-effector of agent $i$ in the inertial frame and $r_{i} \in \mathbb{R}^{3}$ be a fixed vector in the body frame of the load. Initially, $x_{i}(0)=a_{i}(0)=x_{c}(0)+r_{i}$, where $x_{c} \in \mathbb{R}^{3}$ is the position of the center of mass of the load in the inertial frame. Fig. 2.1 also shows the coordinate system defined to derive the kinematics of the system. $\Sigma_{I}$ is the world fixed inertial frame and $O_{c, i}$ is the body-fixed frame attached to each agent $i$.

### 2.2 Kinematics

The position of the attaching point by agent $i$ is given by

$$
\begin{equation*}
a_{i}(t):=x_{c}(t)+r_{i} \tag{2.1}
\end{equation*}
$$

The kinematics of this point is given by

$$
\begin{equation*}
\dot{a}_{i}=\dot{x}_{c} . \tag{2.2}
\end{equation*}
$$



Figure 2.1: Three quadcopters transport a common flexible load. Note that $a_{i}$ is the initial position of agent $i$.

As the quadcopters interact with the load, the load may be deformed. Therefore, $a_{i}(t) \neq$ $x_{i}(t)$ and the load experiences contact forces due to the deformation. As shown in Fig. 2.1, we approximate the deformation as

$$
\begin{equation*}
z_{i}=x_{i}-a_{i}, \quad i=1, \ldots, N . \tag{2.3}
\end{equation*}
$$

We consider that the reaction force $f_{i}$ exerted by each quadcopters is obtained by the gradient of some positive definite function $P_{i}\left(z_{i}\right)$, i.e.,

$$
\begin{equation*}
f_{i}=\nabla P_{i}\left(z_{i}\right) \tag{2.4}
\end{equation*}
$$

We also assume that $P_{i}\left(z_{i}\right)$ satisfies

$$
\begin{array}{r}
P_{i}\left(z_{i}\right)=0 \Longleftrightarrow z_{i}=0, \\
\nabla P_{i}\left(z_{i}\right)=0 \Longleftrightarrow z_{i}=0 . \tag{2.6}
\end{array}
$$

The contact force model in (2.4) includes the linear spring models used in [28,29] if the contact force is elastic.

### 2.3 Dynamics

The translational dynamics of the load can be expressed as

$$
\begin{equation*}
M_{c} \ddot{x}_{c}=\sum_{i=1}^{N} f_{i}-M_{c} g e_{3}, \tag{2.7}
\end{equation*}
$$

where $g$ is the gravitational acceleration, $M_{c}$ denotes the mass of the payload, and $e_{3}$ is the unit vector $\left[\begin{array}{lll}0 & 0 & 1\end{array}\right]^{T}$. In (2.7), we approximate the payload dynamics as a rigid body. This is valid when a load is surrounded by elastic or deformable materials and the deformations are small.

The translational dynamics for each agent $i$ can be modeled as

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=F_{i}-f_{i}-m_{i} g e_{3}, \quad i=1, \ldots, N, \tag{2.8}
\end{equation*}
$$

where $F_{i}$ denotes the force applied by agent $i, m_{i}$ denotes the mass of agent $i$, and $f_{i}$ is the force exerted to agent $i$ from the load. The applied force $F_{i}$ is given by

$$
\begin{equation*}
F_{i}=T_{i} R_{i}(\Phi) e_{3} \tag{2.9}
\end{equation*}
$$

where

$$
\begin{equation*}
T_{i}=\sum_{i=1}^{4} F_{\text {rotors }, i}, \tag{2.10}
\end{equation*}
$$

$\Phi=(\phi, \theta, \psi)$, which are the roll, pitch and yaw $(Z-Y-X)$ angles, respectively, and $R_{i}(\Phi) \in S O(3)$ is the rotation matrix that relates the body frame $O_{c, i}$ and the inertial frame $\Sigma_{i}$, which is given by

$$
R_{i}=\left[\begin{array}{ccc}
c \psi_{i} c \theta_{i} & c \psi_{i} s \phi_{i} s \theta_{i}-c \phi_{i} s \psi_{i} & s \phi_{i} s \psi_{i}+c \phi_{i} c \psi_{i} s \theta_{i}  \tag{2.11}\\
c \theta_{i} s \psi_{i} & c \phi_{i} c \psi_{i}+s \phi_{i} s \psi_{i} s \theta_{i} & c \phi_{i} s \psi_{i} s \theta_{i}-c \psi_{i} s \phi_{i} \\
-s \theta_{i} & c \theta_{i} s \phi_{i} & c \phi_{i} c \theta_{i}
\end{array}\right]
$$

in which $c \theta_{i}=\cos \theta_{i}$ and $s \theta_{i}=\sin \theta_{i}$.

### 2.3.1 Singularities

A fundamental problem with Euler angles is their singularity. This is also well known as gimbal lock [1]. This occurs when the rotational axis of the middle term in the sequence becomes parallel to the rotation axis of the first or third term. For $Z Y X$ Euler angles, singularity occurs at $\theta_{i}= \pm(2 k+1) \frac{\pi}{2}, k=0,1,2, \cdots$.


Figure 2.2: $Z Y X$ Euler angle representation using tripleangle toolbox [1]


Figure 2.3: $Z Y X$ Euler Angle Representation and Gimbal lock. Note $\theta_{i}=90^{\circ}$

### 2.4 Control Objective

The control objective is to design the desired thrusts $T_{i}$ and the desired attitude angles $\Phi_{i}$ in a decentralized way such that all the agents and the payload converge to a constant velocity $v^{d}$ and the contact force $f_{i}$ is regulated to a setpoint $f_{i}^{d}, i=1, \cdots, N$. Since the load converges to a constant velocity, the sum of $f_{i}^{d}$ 's in $x$ and $y$ direction is zero, respectively, and the sum of $f_{i}^{d}$ 's in the $z$ direction balances the weight of the load, i.e.,

$$
\begin{equation*}
\sum_{i=1}^{N} f_{i}^{d}=M_{c} g e_{3} \tag{2.12}
\end{equation*}
$$

We assume that for a given $f_{i}^{d}$ there exists a locally unique deformation $z_{i}^{d}$, such that

$$
\begin{equation*}
f_{i}^{d}=\nabla P_{i}\left(z_{i}^{d}\right) \quad \text { and } \quad \nabla^{2} P_{i}\left(z_{i}^{d}\right)>0 \tag{2.13}
\end{equation*}
$$

Assumption (2.13) is satisfied by linear spring-force model as well as certain classes of nonlinear models, such as $f_{i}=b_{i}\left|z_{i}\right|^{2} z_{i}$, with $b_{i}>0$.

### 2.5 Quadcopter Attitude Controller

We will use the following attitude control for rest of this report. We model the attitude dynamics of each quadcopter as [30]:

$$
\left[\begin{array}{c}
\ddot{\phi}_{i}  \tag{2.14}\\
\ddot{\theta}_{i} \\
\ddot{\psi}_{i}
\end{array}\right]=\left[\begin{array}{c}
\frac{I_{y}-I_{z}}{I_{x}} \dot{\theta}_{i} \dot{\psi}_{i}+\frac{1}{I_{x}} \tau_{\phi_{i}} \\
\frac{I_{z}-I_{x}}{I_{y}} \dot{\phi}_{i} \dot{\psi}_{i}+\frac{1}{I_{y}} \tau_{\theta_{i}} \\
\frac{I_{x}-I_{y}}{I_{z}} \dot{\theta}_{i} \dot{\psi}_{i}+\frac{1}{I_{x}} \tau_{\psi_{i}}
\end{array}\right]
$$

where $\phi_{i}, \theta_{i}$ and $\psi_{i}$ are the roll, pitch and yaw angles for the $i$-th quadcopter. $\tau_{\phi_{i}}, \tau_{\theta_{i}}$ and $\tau_{\psi_{i}}$ are the moments in $x, y$ and $z$ directions respectively and $I_{x}, I_{y}$ and $I_{z}$ are the moments of inertia of the quadcopter in its respective directions. The desired attitude $\Phi_{i}^{\text {des }}$ is given by,

$$
\Phi_{i}^{\text {des }}=\left[\begin{array}{c}
\phi_{i}^{d e s}  \tag{2.15}\\
\theta_{i}^{\text {des }} \\
\psi_{i}^{\text {des }}
\end{array}\right] .
$$

We set $\psi_{i}^{d e s}$ as constant and solve for $\theta_{i}^{d e s}$ and $\phi_{i}^{\text {des }}$ from (3.6) and (3.7) respectively for both cases presented in Section VI.

To track the desired attitude angles computed from our force controllers in Section 3.1 and 6.3, we follow a Proportional-derivative (PD) controller from [31] with some minor changes. The controller is modeled as:

$$
\left[\begin{array}{c}
T_{i}  \tag{2.16}\\
\tau_{\phi_{i}} \\
\tau_{\theta_{i}} \\
\tau_{\psi_{i}}
\end{array}\right]=\left[\begin{array}{c}
T_{i}^{d e s} \\
\frac{I_{x}}{L}\left(-K_{1}^{i} \dot{\phi}_{i}-\frac{I_{y}-I_{z}}{I_{x}} \dot{\theta}_{i} \dot{\psi}_{i}+K_{p, 1}^{i} \tilde{\phi}_{i}+K_{d, 1}^{i} \dot{\tilde{\phi}}_{i}\right) \\
\frac{I_{y}}{L}\left(-K_{2}^{i} \dot{\theta}_{i}-\frac{I_{z}-I_{x}}{I_{y}} \dot{\phi}_{i} \dot{\psi}_{i}+K_{p, 2}^{i} \tilde{\theta}_{i}+K_{d, 2}^{i} \dot{\tilde{\theta}}_{i}\right) \\
\frac{I_{z}}{L}\left(-K_{3}^{i} \dot{\psi}_{i}-\frac{I_{x}-I_{y}}{I_{z}} \dot{\phi}_{i} \dot{\theta}_{i}+K_{p, 3}^{i} \tilde{\psi}_{i}+K_{d, 3}^{i} \dot{\tilde{\psi}}_{i}\right)
\end{array}\right],
$$

where $K_{p, j}^{i}$ and $K_{d, j}^{i}, i=1,2, . ., N, j=1,2,3$ are the PD control gains, and $\tilde{\phi}_{i}, \tilde{\theta}_{i}$ and $\tilde{\psi}_{i}$ are the difference between the desired and actual roll, pitch and yaw angles, respectively, defined as

$$
\left[\begin{array}{c}
\tilde{\phi}_{i}  \tag{2.17}\\
\tilde{\theta}_{i} \\
\tilde{\psi}_{i}
\end{array}\right]=\left[\begin{array}{c}
\phi_{i}^{\text {des }}-\phi_{i} \\
\theta_{i}^{\text {des }}-\theta_{i} \\
\psi_{i}^{\text {des }}-\psi_{i}
\end{array}\right] .
$$

### 2.6 Geometric Control

Although we mainly use the PD controller in our simulation results, we have also implemented a geometric controller in the simulations.

### 2.6.1 Related Works

Geometric tracking control of a quadcopter UAV on $S E(3)$ is developed in [32], which uses the $S E(3)$ instead of the Euler angles. It's a nonlinear tracking controller which is similar to PD controller used in the literature for small angle approximations (see [33]) with two new important features. First, the orientation error is not based on Euler angles, which contain singularities. Second, the desired force is projected onto the actual $z$
body axis. Results shown in [32,34] verifies that this type of controller can be successfully applied for large attitude tracking error.

### 2.6.2 Geometric Control of Cooperative UAVs

Geometric control developed in [32] is extended to quadrotor with a cable suspended payload in [35] and the system is proved to be differentially flat. Similar paradigm was applied to multiple quadcopters transporting a cable-suspended regid body in [11, 18]. In a similar paper [7] developed a full dynamic model for cooperative manipulation of payloads (both point mass and rigid model) suspended by cables from multiple quadcopters for a hybrid model when the cables are tout and slack and utilized the differential flatness property to plan nominal trajectories in flat space. Recently [11] developed a controller based on [18] such that the payload asymptotically follows a given desired trajectory for its position and attitude in presence of uncertainties.

In [11] the force and moment acting on the payload are controlled by using a PID controller. These forces and moments are used to balance the tension along and normal to the cables. The force and moment controller for the payload is based on the error dynamics between the desired and actual trajectory. Similarly, the forces and moments for each quadrotor UAV is controlled using a PID controller.

Rewrite (2.12)

$$
\begin{equation*}
\sum_{i=1}^{N} f_{i}^{d}=M_{c} g e_{3} \tag{2.18}
\end{equation*}
$$

Also rewrite (3.1)

$$
\begin{equation*}
F_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+f_{i}^{d}+m_{i} g e_{3}, \tag{2.19}
\end{equation*}
$$

where $\Gamma_{i}=\Gamma_{i}^{T}>0$

$$
\begin{equation*}
F_{i}^{T} F_{i}=\left(T_{i} R_{i}\left(\Phi_{i}\right) e_{3}\right)^{T} T_{i} R_{i}\left(\Phi_{i}\right) e_{3} \Longrightarrow T_{i}^{d}=\sqrt{F_{i}^{T} F_{i}} \tag{2.20}
\end{equation*}
$$

The third body axis is always parallel to the control force $F_{i}$.

$$
\begin{equation*}
b_{3, i}^{d}=\frac{F_{i}}{\left\|F_{i}\right\|} \tag{2.21}
\end{equation*}
$$

The second body fixed axis $b_{1, i}^{d}$ is chosen such that it can form an orthonormal frame with $b_{3, i}^{d}$ Assuming $b_{1, i}^{d}$ is not parallel to $b_{3, i}^{d}$, the desired rotation matrix can be computed as

$$
R_{i}^{d}=\left[\begin{array}{lll}
b_{2, i}^{d} \times b_{3, i}^{d} & b_{2, i}^{d}, & b_{3, i}^{d} \tag{2.22}
\end{array}\right]
$$

The error vectors for attitude and angular velocity for $i$-th quadrotor is defined as:

$$
\begin{align*}
& e_{R, i}=\frac{1}{2}\left(R_{i}^{d^{T}} R_{i}-R_{i}^{T} R_{i}^{d}\right)^{\vee},  \tag{2.23}\\
& e_{\Omega, i}=\Omega_{i}-R_{i}^{T} R_{i}^{d} \Omega_{i}^{d} .
\end{align*}
$$

where vee map $\vee$ : so(3) $\rightarrow \mathbb{R}^{3}$.
We assume that the desired trajectory is given by constant velocity for each agent as

$$
\dot{x}^{d}=v^{d}=\left[\begin{array}{lll}
v_{x}^{d} & v_{y}^{d} & v_{z}^{d} \tag{2.24}
\end{array}\right], \quad \ddot{x}^{d}=a^{d}=0, \quad \dot{a}^{d}=0 .
$$

The kinematic equation for each quadrotor can be written as

$$
\begin{equation*}
\dot{R}_{i}=R_{i} \hat{\Omega}_{i} \tag{2.25}
\end{equation*}
$$

where $R_{i} \in S O(3)$ is the rotation matrix between the body frame $O_{c, i}$ and the inertial frame $\Sigma_{i}, \Omega_{i} \in \mathbb{R}^{3}$ is the angular velocity of $i$-th quadrotor in body-fixed frame and the operator converts a given vector $\omega=\left[\begin{array}{lll}\omega_{1} & \omega_{2} & \omega_{3}\end{array}\right]^{T}$, into a skew symmetric matrix as:

$$
\widehat{\omega}=\left[\begin{array}{ccc}
0 & -\omega_{3} & \omega_{2}  \tag{2.26}\\
\omega_{3} & 0 & -\omega_{1} \\
-\omega_{2} & \omega_{1} & 0
\end{array}\right] .
$$

The angular dynamics of the $i$-th quadrotor is given by

$$
\begin{equation*}
J_{i} \dot{\Omega}_{i}=M_{i}-\Omega_{i} \times J_{i} \Omega_{i} \tag{2.27}
\end{equation*}
$$

where $J_{i}$ is the inertia matrix and $M_{i} \in \mathbb{R}^{3}$ is the moment generated by the $i$-quadcopter. Based on [11], $M_{i}$ is designed as

$$
\begin{equation*}
M_{i}=-\frac{k_{R}}{\epsilon^{2}} e_{R, i}-\frac{K_{\Omega}}{\epsilon} e_{\Omega, i}+\Omega_{i} \times J_{i} \Omega_{i}-J_{i}\left(\hat{\Omega}_{i} R_{i}^{T} R_{i}^{d} \Omega_{i}^{d}-R_{i}^{T} R_{i}^{d} \dot{\Omega}_{i}^{d}\right), \tag{2.28}
\end{equation*}
$$

where $\epsilon, k_{R}, k_{\Omega}$ are positive constants. $\Omega_{i}^{\text {des }} \in \mathbb{R}^{3}$ is obtained from the attitude kinematics equation

$$
\begin{equation*}
\Omega_{i}^{d}=\left(R_{i}^{d^{T}} \dot{R}_{i}^{d}\right)^{\vee} \tag{2.29}
\end{equation*}
$$

### 2.6.3 Implementation

Rewrite the attitude kinematics as

$$
\begin{align*}
& \dot{R}_{i}^{d}=R_{i}^{d} \hat{\Omega}_{i}^{d}, \\
& \dot{R}_{i}^{d}=\left(\hat{\Omega}_{i}^{d}\right)^{I} R_{i}^{d}=\left(\Omega_{i}^{d}\right)^{I} \times R_{i}^{d}  \tag{2.30}\\
& \dot{R}_{i}^{d}=R_{i}^{d} \Omega_{i}^{d} \times R_{i}^{d}
\end{align*}
$$

Time derivative of (2.30) yields,

$$
\begin{align*}
\ddot{R}_{i}^{d} & =\frac{d}{d t}\left(R_{i}^{d} \hat{\Omega}_{i}^{d}\right) \\
\ddot{R}_{i}^{d} & =\dot{R}_{i}^{d} \hat{\Omega}_{i}^{d}+R_{i}^{d} \dot{\hat{\Omega}}_{i}^{d} \\
\ddot{R}_{i}^{d} & =R_{i}^{d}\left(\hat{\Omega}_{i}^{d}\right)^{2}+R_{i}^{d} \dot{\hat{\Omega}}_{i}^{d}  \tag{2.31}\\
\Longrightarrow \dot{\hat{\Omega}}_{i}^{d} & =\left(R_{i}^{d}\right)^{T} \ddot{R}_{i}^{d}-\left(\hat{\Omega}_{i}^{d}\right)^{2} \\
\Longrightarrow \dot{\Omega}_{i}^{d} & =\left(\left(R_{i}^{d}\right)^{T} \ddot{R}_{i}^{d}-\left(\hat{\Omega}_{i}^{d}\right)^{2}\right)^{v}
\end{align*}
$$

### 2.6.4 Desired Rotation matrix $R_{i}^{d}$

$$
R_{i}^{d}=\left[\begin{array}{lll}
b_{2, i}^{d} \times b_{3, i}^{d} & b_{2, i}^{d}, & b_{3, i}^{d} \tag{2.32}
\end{array}\right]
$$

Define

$$
b_{1, i}^{d}=\left[\begin{array}{lll}
\cos \pi t, & \sin \pi t, & 0 \tag{2.33}
\end{array}\right]
$$

Rewrite (2.21)

$$
\begin{gather*}
b_{3, i}^{d}=\frac{F_{i}}{\left\|F_{i}\right\|}=\frac{-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+f_{i}^{d}+m_{i} g e_{3}}{\left\|-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+f_{i}^{d}+m_{i} g e_{3}\right\|}  \tag{2.34}\\
b_{2, i}^{d}=\frac{b_{3, i}^{d} \times b_{1, i}^{d}}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|} \tag{2.35}
\end{gather*}
$$

Define

$$
\begin{align*}
& \beta=b_{3, i}^{d} \times b_{1, i}^{d}  \tag{2.36}\\
& \dot{\beta}=\dot{b}_{3, i}^{d} \times b_{1, i}^{d}+b_{3, i}^{d} \times \dot{b}_{1, i}^{d}  \tag{2.37}\\
& \ddot{\beta}=\ddot{b}_{3, i}^{d} \times b_{1, i}^{d}+2\left(\dot{b}_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right)+b_{3, i}^{d} \times \ddot{b}_{1, i}^{d} \tag{2.38}
\end{align*}
$$

Rewrite the rotation matrix as

$$
R_{i}^{d}=\left[\begin{array}{lll}
b_{2, i}^{d} \times b_{3, i}^{d}, & \frac{\beta}{\|\beta\|}, & b_{3, i}^{d} \tag{2.39}
\end{array}\right]
$$

### 2.6.5 Desired rate of Rotation $\dot{R}_{i}^{d}$

Derivative of (2.32) yields

$$
\begin{equation*}
\dot{R}_{i}^{d}=\left[\dot{b}_{2, i}^{d} \times b_{3, i}^{d}+b_{2, i}^{d} \times \dot{b}_{3, i}^{d}, \quad \dot{b}_{2, i}^{d}, \quad \dot{b}_{3, i}^{d}\right] \tag{2.40}
\end{equation*}
$$

Derivative of (2.34) yields

$$
\begin{gather*}
\dot{b}_{3, i}^{d}=\frac{d}{d t}\left(\frac{F_{i}}{\left\|F_{i}\right\|}\right) \\
\dot{b}_{3, i}^{d}=\frac{\dot{F}_{i}\left\|F_{i}\right\|-F_{i} \frac{F_{i}^{T} \cdot \dot{F}_{i}}{\left\|F_{i}\right\|}}{\left\|F_{i}\right\|^{2}}  \tag{2.41}\\
\dot{b}_{3, i}^{d}=\frac{\dot{F}_{i}}{\left\|F_{i}\right\|}-\frac{F_{i} F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{3}} \\
\dot{b}_{1, i}^{d}=\left[\begin{array}{lll}
-\pi \sin \pi t, \quad \pi \cos \pi t, & 0
\end{array}\right] \tag{2.42}
\end{gather*}
$$

Calculate $\dot{b}_{2, i}^{d}$

$$
\begin{align*}
\dot{b}_{2, i}^{d} & =\frac{d}{d t} \frac{b_{3, i}^{d} \times b_{1, i}^{d}}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|} \\
& =\frac{\frac{d}{d t}\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|-\left(b_{3, i}^{d} \times b_{1, i}^{d}\right) \frac{d}{d t}\left(\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|^{2}}  \tag{2.43}\\
& =\frac{\frac{d}{d t}\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|}-\frac{\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)^{T} \frac{d t}{d t}\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|^{3}}
\end{align*}
$$

where

$$
\begin{equation*}
\frac{d}{d t}\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)=\left(\dot{b}_{3, i}^{d} \times b_{1, i}^{d}\right)+\left(b_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right) \tag{2.44}
\end{equation*}
$$

Rewrite (2.45) as

$$
\begin{align*}
\dot{b}_{2, i}^{d} & =\frac{\left(\dot{b}_{3, i}^{d} \times b_{1, i}^{d}\right)+\left(b_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|}-\frac{\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)^{T}\left(\left(\dot{b}_{3, i}^{d} \times b_{1, i}^{d}\right)+\left(b_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right)\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|^{3}} \\
\dot{b}_{2, i}^{d} & =\frac{\dot{\beta}}{\|\beta\|}-\frac{\beta \beta^{T} \dot{\beta}}{\|\beta\|^{3}} \tag{2.45}
\end{align*}
$$

Rewrite $\dot{R}_{i}^{d}$ as

$$
\dot{R}_{i}^{d}=\left[\begin{array}{ll}
\dot{\beta}, & \frac{\dot{\beta}}{\|\beta\|}-\frac{\beta \beta^{T} \dot{\beta}}{\|\beta\|^{3}}, \quad \frac{\dot{F}_{i}}{\left\|F_{i}\right\|}-\frac{F_{i} F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{3}} \tag{2.46}
\end{array}\right]
$$

### 2.6.6 $\ddot{R}_{i}^{d}$

Derivative of (2.40) yields

$$
\begin{equation*}
\ddot{R}_{i}^{d}=\left[\ddot{b}_{2, i}^{d} \times b_{3, i}^{d}+2\left(\dot{b}_{2, i}^{d} \times \dot{b}_{3, i}^{d}\right)+b_{2, i}^{d} \times \ddot{b}_{3, i}^{d}, \quad \ddot{b}_{2, i}^{d}, \quad \ddot{b}_{3, i}^{d}\right] \tag{2.47}
\end{equation*}
$$

Differentiate (2.41)

$$
\begin{align*}
\ddot{b}_{3, i}^{d} & =\frac{d}{d t}\left\{\frac{\dot{F}_{i}\left\|F_{i}\right\|^{2}}{\left\|F_{i}\right\|^{3}}-\frac{F_{i} F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{3}}\right\} \\
& =\frac{d}{d t}\left\{\frac{\dot{F}_{i}}{\left\|F_{i}\right\|}-\frac{F_{i} F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{3}}\right\} \\
& =\frac{\ddot{F}_{i}\left\|F_{i}\right\|-\dot{F}_{i}\left\|\dot{F}_{i}\right\|}{\left\|F_{i}\right\|^{2}}-\frac{\frac{d t}{d t}\left(F_{I} F_{i}^{T} \dot{F}_{i}\right)\left\|F_{i}\right\|^{3}-\left(F_{i} F_{i}^{T} \dot{F}_{i}\right) \frac{d}{d t}\left(\left\|F_{i}\right\|^{3}\right)}{\left\|F_{i}\right\|^{6}}  \tag{2.48}\\
& =\frac{\ddot{F}_{i}}{\left\|F_{i}\right\|}-\frac{\dot{F}_{i} F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{3}}-\frac{\left(\dot{F}_{i} F_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} \dot{F}_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} F_{i}^{T} \ddot{F}_{i}\right)}{\left\|F_{i}\right\|^{3}}+\frac{3\left(F_{i} F_{i}^{T} \dot{F}_{i}\right) F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{5}} \\
& =\frac{\ddot{F}_{i}}{\left\|F_{i}\right\|}-\frac{2\left(\dot{F}_{i} F_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} \dot{F}_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} F_{i}^{T} \ddot{F}_{i}\right)}{\left\|F_{i}\right\|^{3}}+\frac{3\left(F_{i} F_{i}^{T} \dot{F}_{i}\right) F_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{5}}
\end{align*}
$$

Compute the second derivative of $b_{2, i}^{d}$ as

$$
\begin{align*}
\ddot{b}_{2, i}^{d} & =\frac{d}{d t}\left\{\frac{\left(\dot{b}_{3, i}^{d} \times b_{1, i}^{d}\right)+\left(b_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|}-\frac{\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)\left(b_{3, i}^{d} \times b_{1, i}^{d}\right)^{T}\left(\left(\dot{b}_{3, i}^{d} \times b_{1, i}^{d}\right)+\left(b_{3, i}^{d} \times \dot{b}_{1, i}^{d}\right)\right)}{\left\|b_{3, i}^{d} \times b_{1, i}^{d}\right\|^{3}}\right\} \\
\ddot{b}_{2, i}^{d} & =\frac{d}{d t}\left\{\frac{\dot{\beta}}{\|\beta\|}-\frac{\beta \beta^{T} \dot{\beta}}{\|\beta\|^{3}}\right\} \\
& =\frac{\ddot{\beta}\|\beta\|-\dot{\beta}\|\dot{\beta}\|}{\|\beta\|^{2}}-\frac{\frac{d t}{d t}\left(\beta \beta^{T} \dot{\beta}\right)\|\beta\|^{3}-\left(\beta \beta^{T} \dot{\beta}\right) \frac{d}{d t}\left(\|\beta\|^{3}\right)}{\|\beta\|^{6}} \\
& =\frac{\ddot{\beta}}{\|\beta\|}-\frac{\dot{\beta} \beta^{T} \dot{\beta}}{\|\beta\|^{3}}-\frac{\left(\dot{\beta} \beta^{T} \dot{\beta}\right)+\left(\beta \dot{\beta}^{T} \dot{\beta}\right)+\left(\beta \beta^{T} \ddot{\beta}\right)}{\|\beta\|^{3}}+\frac{3\left(\beta \beta^{T} \dot{\beta}\right) \beta^{T} \dot{\beta}}{\|\beta\|^{5}} \\
& =\frac{\ddot{\beta}}{\|\beta\|}-\frac{2\left(\dot{\beta} \beta^{T} \dot{\beta}\right)+\left(\beta \dot{\beta}^{T} \dot{\beta}\right)+\left(\beta \beta^{T} \ddot{\beta}\right)}{\|\beta\|^{3}}+\frac{3\left(\beta \beta^{T} \dot{\beta}\right) \beta^{T} \dot{\beta}}{\|\beta\|^{5}} \tag{2.49}
\end{align*}
$$

Rewrite $\ddot{R}_{i}^{d}$ as

$$
\ddot{R}_{i}^{d}=\left[\begin{array}{lll}
\ddot{\beta}, & \frac{\ddot{\beta}}{\|\beta\|}-\frac{2\left(\dot{\beta} \beta^{T} \dot{\beta}\right)+\left(\beta \dot{\beta}^{T} \dot{\beta}\right)+\left(\beta \beta^{T} \ddot{\beta}\right)}{\|\beta\|^{3}}+\frac{3\left(\beta \beta^{T}\right)}{\|\beta\|^{5}} \boldsymbol{\beta}^{T} \dot{\beta}  \tag{2.50}\\
& \frac{\ddot{F}_{i}}{\left\|F_{i}\right\|}-\frac{2\left(\dot{F}_{i} F_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} \dot{F}_{i}^{T} \dot{F}_{i}\right)+\left(F_{i} F_{i}^{T} \ddot{F}_{i}\right)}{\left\|F_{i}\right\|^{3}}+\frac{3\left(F_{i} F_{i}^{T} \dot{F}_{i}\right)_{i}^{T} \dot{F}_{i}}{\left\|F_{i}\right\|^{5}}
\end{array}\right]
$$

## CHAPTER 3

## NOMINAL DESIGN

In this chapter, we review the main results in [36] and extend it to address cooperative transportation of a payload in 3D.

### 3.1 Force Control with a Payload of Known Mass

We assume that the orientation of the load is controlled separately and consider only the translational part of the motion in this chapter. We first choose $F_{i}$ as

$$
\begin{equation*}
F_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+f_{i}^{d}+m_{i} g e_{3}, \tag{3.1}
\end{equation*}
$$

where $\Gamma_{i}=\Gamma_{i}^{T}>0$, which component wise becomes

$$
\left[\begin{array}{c}
F_{i}^{x}  \tag{3.2}\\
F_{i}^{y} \\
F_{i}^{z}
\end{array}\right]=-\Gamma_{i}\left(\begin{array}{c}
\dot{x}_{i, x}-v_{x}^{d} \\
\dot{x}_{i, y}-v_{y}^{d} \\
\dot{x}_{i, z}-v_{z}^{d}
\end{array}\right)+\left(\begin{array}{c}
f_{i, x}^{d} \\
f_{i, y}^{d} \\
\frac{1}{N} M_{c} g+m_{i} g
\end{array}\right)
$$

Note that in (3.2) we have chosen $f_{i, z}^{d}=\frac{1}{N} M_{c} g$ and the $x$ and $y$ components of $f_{i}^{d}$ are pre-designed to satisfy the following constraints

$$
\begin{equation*}
\sum_{i=1}^{N} f_{i, x}^{d}=0 \quad \text { and } \quad \sum_{i=1}^{N} f_{i, y}^{d}=0 \tag{3.3}
\end{equation*}
$$

### 3.1.1 Computation of $T_{i}^{d e s}$ and $\Phi_{i}^{d e s}$

For a coplanar quadcopter and away from singularities, we next establish the following map (for any yaw $\psi_{i}$ )

$$
\left[\begin{array}{c}
F_{i}^{x}  \tag{3.4}\\
F_{i}^{y} \\
F_{i}^{z}
\end{array}\right] \rightarrow\left[\begin{array}{c}
\phi_{i}^{d e s}\left(\psi_{i}\right) \\
\theta_{i}^{d e s}\left(\psi_{i}\right) \\
T_{i}^{d e s}
\end{array}\right] .
$$

Starting from

$$
\begin{equation*}
T_{i}^{d e s} R_{i}\left(\Phi_{i}^{d e s}\right) e_{3}=F_{i} \tag{3.5}
\end{equation*}
$$

and premultiplying by $R_{i}^{T}\left(\Phi_{i}^{\text {des }}\right)$, we get

$$
T_{i}^{d e s} e_{3}=R_{i}^{T}\left(\Phi_{i}^{d e s}\right)\left[\begin{array}{c}
F_{i}^{x} \\
F_{i}^{y} \\
F_{i}^{z}
\end{array}\right],
$$

which are three scalar equations. From the first we get

$$
\begin{equation*}
\theta_{i}^{d e s}=\operatorname{atan}\left(\frac{F_{i}^{x} c \psi_{i}^{d e s}+F_{i}^{y} s \psi_{i}^{d e s}}{F_{i}^{z}}\right) \tag{3.6}
\end{equation*}
$$

and from the other two equations

$$
\begin{equation*}
\phi_{i}^{d e s}=\operatorname{asin}\left(\frac{F_{i}^{x} s \psi_{i}^{d e s}-F_{i}^{y} c \psi_{i}^{d e s}}{\sqrt{F_{i}^{T} F_{i}}}\right), \tag{3.7}
\end{equation*}
$$

where in (3.7) we have used (3.5) to obtain the relationship

$$
\begin{equation*}
T_{i}^{d e s}=\sqrt{F_{i}^{T} F_{i}} . \tag{3.8}
\end{equation*}
$$

Remark 3.1 Define the velocity error vector in (3.1) as

$$
\begin{equation*}
\xi_{i}=\dot{x}_{i}-v^{d} \tag{3.9}
\end{equation*}
$$

and recall from (3.2) that $v^{d}=\left[\begin{array}{lll}v_{x}^{d} & v_{y}^{d} & v_{z}^{d}\end{array}\right]^{T}$. Suppose that $\left|\xi_{i}\right| \leq \eta_{i}^{m}$ and $\left|f_{i, y}^{d}\right|+\left|f_{i, x}^{d}\right|$ $\leq f_{c}^{m}$ for some positive bounds $\eta_{i}^{m}$ and $f_{c}^{m}$, and fix $\Gamma_{i}$ as a positive definite and diagonal
matrix. Then, it is straightforward to see that $T_{i}^{\text {des }}$ can be upper bounded as follows

$$
\begin{align*}
\left|T_{i}^{d e s}\right| & =\sqrt{F_{i}^{T} F_{i}} \\
& \leq\left\|\Gamma_{i}\right\|\left|\xi_{i}\right|+\left|f_{i, x}^{d}\right|+\left|f_{i, y}^{d}\right|+\frac{1}{N} M_{c} g+m_{i} g  \tag{3.10}\\
& \leq\left\|\Gamma_{i}\right\| \eta_{i}^{m}+f_{c}^{m}+\frac{1}{N} M_{c} g+m_{i} g .
\end{align*}
$$

### 3.1.2 Stability Analysis

The dynamics (2.8) with the proposed control law (3.1) takes the following form:

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=-\Gamma_{i} \xi_{i}+f_{i}^{d}-f_{i} . \tag{3.11}
\end{equation*}
$$

Theorem 3.1 The decentralized control law (3.1) ensures that the closed-loop equilibrium of (3.11), given by,

$$
\begin{equation*}
\mathscr{E}=\left\{\left(\dot{x}_{i}, \dot{x}_{c}, f_{i}\right) \mid \dot{x}_{i}=v_{i}^{d}, \quad \dot{x}_{c}=v_{i}^{d}, \quad \text { and } \quad f_{i}=f_{i}^{d}\right\} \tag{3.12}
\end{equation*}
$$

is asymptotically stable.

Proof. Consider the following energy-motivated positive definite candidate Lyapunov function

$$
\begin{align*}
V & =\sum_{i=1}^{N}\left[P_{i}\left(z_{i}\right)-P_{i}\left(z_{i}^{d}\right)-\left(f_{i}^{d}\right)^{T}\left(z_{i}-z_{i}^{d}\right)\right] \\
& +\frac{1}{2}\left(\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \xi_{i}+\xi_{c}^{T} M_{c} \xi_{c}\right) . \tag{3.13}
\end{align*}
$$

From (2.2) and (2.3), the kinematics of $z_{i}$ is given by

$$
\begin{equation*}
\dot{z}_{i}=\dot{x}_{i}-\dot{a}_{i}=\dot{x}_{i}-\dot{x}_{c} . \tag{3.14}
\end{equation*}
$$

Define $\xi_{c}=\dot{x}_{c}-v^{d}$. The time derivative of $V$ yields

$$
\begin{equation*}
\dot{V}=\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \dot{z}_{i}+\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \ddot{x}_{i}+\xi_{c}^{T} M_{c} \ddot{x}_{c} . \tag{3.15}
\end{equation*}
$$

We rewrite (3.15) from (2.7), (2.8), (2.12) and (3.14) as:

$$
\begin{align*}
\dot{V}= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\dot{x}_{i}-\dot{x}_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(F_{i}-f_{i}\right) \\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g_{3}, \\
= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\xi_{i}-\xi_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(-\Gamma_{i} \xi_{i}+f_{i}^{d}-f_{i}\right)  \tag{3.16}\\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g e_{3} \\
= & -\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i} \leq 0
\end{align*}
$$

which implies the stability of the equilibrium $\mathscr{E}$.
To conclude the asymptotic stability of the system, we apply the LaSalle invariance principle [37] to investigate the largest invariant set $M$. On $M, \xi_{i}=0$ which implies from (3.9) that $\dot{x}_{i}=v^{d}$. Since $\dot{x}_{i}$ is constant, $\ddot{x}_{i}=0$, which implies from (3.11) that $f_{i}^{d}=f_{i}$. Since $f_{i}$ is constant and equal to $f_{i}^{d}, z_{i}$ is constant, which leads to $\dot{x}_{c}=\dot{x}_{i}=v^{d}$ and $\ddot{x}_{c}=0$.

Remark 3.2 Since $\xi_{i}$ is bounded, the existence of a bounded thrust $T_{i}$ is always guaranteed from the result of Remark 3.1.

Remark 3.3 The proof of proposition 3.1 is similar to [38, Chapter 8] except that the dynamics are extended to incorporate the gravity of the agents and the payload.

### 3.2 Simulation Results

### 3.2.1 Simulation Environment

In this section, we present simulation results for two quadcopters with a rigid extension transporting a load. The mass and inertia of each quadcopter are $m_{i}=0.75 \mathrm{~kg}$ and $I_{i}=\operatorname{diag}\{0.0820,0.0845,0.1377\} \mathrm{kg} \mathrm{m}^{2}$ respectively. The mass of the payload is $M_{c}=1.5$ kg with a radius of 15 cm . We use a linear spring-force model to compute the contact
force $f_{i}$ in the simulation,

$$
\begin{equation*}
f_{i}=k z_{i} \tag{3.17}
\end{equation*}
$$

where $k$ is the spring constant of the payload. If the contact between agents and payload is elastic at a constant pressure, the contact force model in (3.17) is equivalent to the models presented in [28,29]. For all of the results presented in this section we set $\psi_{i}^{d e s}=$ 0 and $k=2.5 \times 10^{4} \mathrm{~N} / \mathrm{m}$.

### 3.2.2 Known Payload with a Predesigned $v^{d}$ with a PD Controller

We choose

$$
\begin{aligned}
v^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
5.0 & 0 & 0
\end{array}\right]^{T}, \\
f_{1}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
0 & 2.0 & 0.5 M_{c} g
\end{array}\right]^{T}, \\
f_{2}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
0 & -2.0 & 0.5 M_{c} g
\end{array}\right]^{T}, \\
\Gamma_{i} & =\left[\begin{array}{lll}
10 & 0 & 0 \\
0 & 10 & 0 \\
0 & 0 & 10
\end{array}\right] .
\end{aligned}
$$

Note that $v^{d}(t), f_{1}^{d}(t), f_{2}^{d}(t)$ converge to constants exponentially. The parameters for the gains for the attitude controller are given in Table 3.1.

| $K_{1}^{i}, K_{2}^{i}$ | $K_{3}$ | $K_{p 1}^{i}, K_{p 2}^{i}$ | $K_{p 3}^{i}$ | $K_{d 1}^{i}, K_{d 2}^{i}$ | $K_{d 3}^{i}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 21.93 | 18 | 4.65 | 3.76 | 0.19 | 0.15 |

Table 3.1: Attitude controller gains.

Fig. 3.1 and 3.2 show that the attitude controller implemented in Section 3.2 is successful in tracking the desired attitude angles, $\phi_{i}^{\text {des }}$ and $\theta_{i}^{\text {des }}$ generated from the force controller (3.6) and (3.7). It can be observed that the actual pitch $\left(\theta_{i}\right)$ and roll $\left(\phi_{i}\right)$ angles for each quadcopter converge to steady states within the first 6 seconds of the simulation. Fig. 3.3 shows that the desired thrust force ( $T_{i}$ ) converges to a constant value within


Figure 3.1: Desired and actual roll angles $(\phi)$ for both agents.


Figure 3.2: Desired and actual pitch angles $(\theta)$ for both agents.


Figure 3.3: The thrusts for both agents are bounded and converge to 14.7 N .


Figure 3.4: Linear velocities for both agents. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 3.5: Linear velocity for the payload. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 3.6: Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload.
first 2 seconds and stays constant throughout the simulation. We observe from Fig. 3.4 and 3.5 that the velocities of the agents and the payload converge to $v^{d}$. We see from Fig. 3.6 that the contact forces are regulated to the setpoints $f_{i}^{d}$.

### 3.2.3 Known Payload with a Predesigned $v^{d}$ with a Geometric Controller

## Simulation Parameters:

$$
\begin{gathered}
b_{i}^{d}=\left[\begin{array}{lll}
1 & 0 & 0
\end{array}\right] \\
v^{d}=\left[\begin{array}{lll}
2.0 & 0 & 0
\end{array}\right] \\
f_{i}^{d}=\left[\begin{array}{lll}
0 & \pm 2.0 & 0.5 M_{c} g
\end{array}\right]
\end{gathered}
$$

Gains: $k_{R}=1.5, \epsilon=1.0, k_{\Omega}=1.0$

$$
\Gamma_{i}=\left[\begin{array}{lll}
5 & 0 & 0 \\
0 & 5 & 0 \\
0 & 0 & 5
\end{array}\right]
$$

spring constant $k=2.5 \times 10^{4} \mathrm{~N} / \mathrm{m}$.


Figure 3.7: Linear velocities for both agents. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 3.8: Linear velocity for the payload. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 3.9: Desired thrust input to each quadrotor.


Figure 3.10: Control input $M_{1}$


Figure 3.11: Control input $M_{2}$


Figure 3.12: Control input $M_{3}$


Figure 3.13: Attitude error of the quadrotors given by $\Psi_{i}=\frac{1}{2}\left\|R_{i}-R_{I}^{d}\right\|^{2}$


Figure 3.14: Euler Angles for each agent.


Figure 3.15: Contact forces.

## CHAPTER 4

## COOPERATIVE AERIAL MANIPULATION OF A PAYLOAD WITH UNKNOWN MASS

In Chapter 3, the mass of the payload is assumed available so that $f_{i}^{d}$ in (2.12) can be specified. However, in practical scenarios, an exact knowledge of the mass of the payload may not be available. To address such a scenario, we develop an adaptive controller to estimate $M_{c}$.

In this chapter, we consider a group quadcopters with a single rigid link attached to it and develop an adaptive decentralized control law for transporting a payload with an unknown mass without explicit communication between the agents. The algorithm provides desired thrust and attitude angles required for each quadcopter to cooperatively transport a payload. It also guarantees that all the agents converge to a desired velocity and regulates the contact force. The sum of the estimates of the unknown mass from all the agents converge to the true mass. We model the contact force between the agents and the payload as a mass spring model. This assumption is valid when the vehicles are connected to the payload via elastic cables or when the payload is flexible or surrounded by elastic bumper materials.

This chapter also introduces a new force consensus algorithm to ensure that each quadcopter gets an equal share of the payload mass. Such an algorithm is important because without it, some of the quadcopters may provide more lift to the unknown payload than others, and therefore consume more power, which can lead to a shorter life span of the entire team. We provide extended stability analysis for the new algorithm and additional simulation results. To test the performance of the controllers for different conditions, we include simulation results for a time-varying reference velocity.

### 4.1 Force Control with a Payload of Unknown Mass

In Section 3.1, the mass of the payload is assumed available so that $f_{i}^{d}$ in (2.12) can be specified. However, in practical scenarios, an exact knowledge of the mass of the payload may not be available. To address such a scenario, we develop an adaptive controller to estimate $M_{c}$.

We define $\widehat{M}_{c, i}$ as an estimate of mass of payload for each agent $i=1, \ldots, N$ and propose the following update law for $\widehat{M}_{c, i}$

$$
\begin{equation*}
\dot{\widehat{M}}_{c, i}=\gamma_{i}\left(\dot{x}_{i}-v^{d}\right)^{T} e_{3} \tag{4.1}
\end{equation*}
$$

where $\gamma_{i}>0$. Note that $\widehat{M}_{c, i}$ stops updating when each agent reaches the predesigned velocity $\nu^{d}$.

### 4.1.1 Computation of $T_{i}^{\text {des }}$ and $\Phi_{i}^{\text {des }}$

We modify the design in (3.1) as:

$$
\begin{equation*}
F_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+\hat{f}_{i}^{d}+m_{i} g e_{3}, \tag{4.2}
\end{equation*}
$$

and estimate the $z$-component of $f_{i}^{d}$ as:

$$
\begin{equation*}
\hat{f}_{i, z}^{d}=\widehat{M}_{c, i} g \tag{4.3}
\end{equation*}
$$

The $x$ and $y$ components of $\hat{f}_{i}^{d}$ are the steady state values of the contact forces that are pre-designed to satisfy the following constraints

$$
\begin{equation*}
\sum_{i=1}^{N} \hat{f}_{i, x}^{d}=0 \quad \text { and } \quad \sum_{i=1}^{N} \hat{f}_{i, y}^{d}=0 \tag{4.4}
\end{equation*}
$$

Given (4.2), the corresponding $T_{i}$ and $\Phi_{i}$ are found based on (3.6)-(3.8).

### 4.1.2 Stability Analysis

The dynamics (2.8) with the proposed control law (4.1)- (4.3) takes the following form:

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=-\Gamma_{i} \xi_{i}+\hat{f}_{i}^{d}-f_{i} . \tag{4.5}
\end{equation*}
$$

Theorem 4.1 The decentralized control law (4.2) ensures that the equilibrium of (2.7), (4.1), and (4.5), given by

$$
\begin{gather*}
\mathscr{E}^{*}=\left\{\left(\dot{x}_{i}, f_{i}, v_{i}^{d}, \widehat{M}_{c, i}\right) \mid \dot{x}_{i}=v_{i}^{d}, \quad \dot{x}_{c}=v_{i}^{d}\right. \\
\left.\sum_{i=1}^{N} \widehat{M}_{c, i}=M_{c}, \quad \text { and } \quad \hat{f}_{i}^{d}=f_{i}\right\} \tag{4.6}
\end{gather*}
$$

is asymptotically stable.

Proof. Consider the following energy-motivated positive definite candidate Lyapunov function

$$
\begin{align*}
V_{1} & =\sum_{i=1}^{N}\left[P_{i}\left(z_{i}\right)-P_{i}\left(z_{i}^{d}\right)-\left(f_{i}^{d}\right)^{T}\left(z_{i}-z_{i}^{d}\right)\right]  \tag{4.7}\\
& +\frac{1}{2}\left(\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \xi_{i}+\xi_{c}^{T} M_{c} \xi_{c}\right) .
\end{align*}
$$

The time derivative of $V_{1}$ yields

$$
\begin{equation*}
\dot{V}_{1}=\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \dot{z}_{i}+\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \ddot{x}_{i}+\xi_{c}^{T} M_{c} \ddot{x}_{c} \tag{4.8}
\end{equation*}
$$

We rewrite (4.8) from (2.7), (2.8), (2.12), (3.14) and (4.2) as:

$$
\begin{align*}
\dot{V}_{1}= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\dot{x}_{i}-\dot{x}_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(F_{i}-f_{i}\right) \\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g_{3}, \\
= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\xi_{i}-\xi_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(-\Gamma_{i} \xi_{i}+\hat{f}_{i}^{d}-f_{i}\right) \\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g e_{3}  \tag{4.9}\\
= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \xi_{i}-\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \xi_{c}-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i} \\
& +\sum_{i=1}^{N} \xi_{i}^{T}\left(\hat{f}_{i}^{d}-f_{i}\right)+\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g e_{3} \\
= & -\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{f}_{i}^{d}
\end{align*}
$$

where $\tilde{f}_{i}^{d}=\hat{f}_{i}^{d}-f_{i}^{d}$, i.e.,

$$
\tilde{f}_{i}^{d}=\left(\begin{array}{c}
0  \tag{4.10}\\
0 \\
\left(\widehat{M}_{c, i}-M_{c, i}\right) g
\end{array}\right)=\left(\begin{array}{c}
0 \\
0 \\
\tilde{M}_{c, i} g
\end{array}\right)=\tilde{M}_{c, i} g e_{3},
$$

where $M_{c, i}, i=1, \cdots, N$, are arbitrary constants that satisfy $\sum_{i=1}^{N} M_{c, i}=M_{c}$. From (4.9) and (4.10),

$$
\begin{equation*}
\dot{V}_{1}=-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3} . \tag{4.11}
\end{equation*}
$$

Consider another positive definite candidate Lyapunov function:

$$
\begin{equation*}
V_{2}=\frac{1}{2} \sum_{i=1}^{N} \Lambda_{i}\left(\tilde{M}_{c, i}\right)^{2} . \tag{4.12}
\end{equation*}
$$

The time derivative of $V_{2}$ yields

$$
\begin{equation*}
\dot{V}_{2}=\sum_{i=1}^{N} \Lambda_{i} \tilde{M}_{c, i} \dot{\tilde{M}}_{c, i} \tag{4.13}
\end{equation*}
$$

Let $W=V_{1}+V_{2}$. The time derivative of $W$ is given by,

$$
\begin{align*}
\dot{W} & =\dot{V}_{1}+\dot{V}_{2}, \\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}+\sum_{i=1}^{N} \tilde{M}_{c, i} \dot{\tilde{M}}_{c, i},  \tag{4.14}\\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}-\sum_{i=1}^{N} \tilde{M}_{c, i} \Lambda_{i}\left(\gamma_{i} \xi_{i}^{T} e_{3}\right) .
\end{align*}
$$

Now from (4.1) and (4.14) and choosing $\Lambda_{i} \gamma_{i}=g$, we obtain

$$
\begin{equation*}
\dot{W}=-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i} \leq 0, \tag{4.15}
\end{equation*}
$$

which implies the stability of the equilibrium $\mathscr{E}^{*}$.
To conclude the asymptotic stability of the system, we apply the LaSalle invariance principle [37] to investigate the largest invariant set $M$. On $M, \xi_{i}=0$ and (3.9) leads to $\dot{x}_{i}=v^{d}$. Since $\dot{x}_{i}$ is constant, $\ddot{x}_{i}=0$, which implies from (4.25) that $\hat{f}_{i}^{d}=f_{i}$. Because $\dot{x}_{i}=v^{d}, \dot{\vec{M}}_{c, i}=0$, which means that $\widehat{M}_{c, i}$ is a constant. It follows from (4.3) that $\hat{f}_{i}^{d}$ is constant and equal to $f_{i}$. Since $f_{i}$ is constant, $z_{i}$ is constant, which leads to $\dot{x}_{c}=\dot{x}_{i}=v^{d}$ and $\ddot{x}_{c}=0$. From (2.7), we get $\sum_{i=1}^{N} f_{i}=M_{c} g e_{3}$, which, together with $f_{i}=\hat{f}_{i}^{d}$, results in $\sum_{i=1}^{N} \hat{f}_{i}^{d}=M_{c} g$. Given (4.3), we conclude $\sum_{i=1}^{N} \widehat{M}_{c, i}=M_{c}$.

### 4.1.3 Simulation Results: Unknown Payload with a Predesigned $v^{d}$

In this case, we assume that the mass of the payload is unknown. We keep the other parameters to be the same as in the previous chapter and set $\widehat{M}_{c, 1}(0)=0.1 \mathrm{~kg}, \widehat{M}_{c, 2}(0)=$ 0.4 kg ,

$$
\begin{aligned}
\hat{f}_{1}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
0 & 2.0 & \widehat{M}_{c, 1} g
\end{array}\right]^{T}, \\
\hat{f}_{2}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
0 & -2.0 & \widehat{M}_{c, 2} g
\end{array}\right]^{T} . \\
\Gamma_{i} & =\left[\begin{array}{ccc}
10 & 0 & 0 \\
0 & 10 & 0 \\
0 & 0 & 25
\end{array}\right] .
\end{aligned}
$$

Similar to the known mass case, Fig. 4.1 shows that actual roll angle ( $\phi_{i}$ ) converges to the steady state $\left(\phi_{i}^{\text {des }}\right.$ ) computed from the force controller (using (3.6) for Section. 6.3) within the first 6 seconds of simulation. Likewise, Fig. 4.2 shows that the actual pitch angle $\left(\theta_{i}\right)$ converges to its steady state $\left(\theta_{i}^{d e s}\right)$ for each quadrotor within first 6 seconds of the simulation. Fig. 4.3 shows that the desired thrust ( $T_{i}$ ) converges to some constant value within first 4 seconds and stays constant throughout the simulation. We observe from Fig. 4.4 and 4.5 that the velocities of the agents and the payload converge to desired $v^{d}$. The contact forces is also regulated at the setpoints $f_{i}^{d}$ as shown in Fig. 4.6. The update law from (4.1) successfully recovers the actual mass of the payload as seen in Fig. 4.7.


Figure 4.1: Desired and actual roll angles $(\phi)$ for both agents (unknown $M_{c}$ case).


Figure 4.2: Desired and actual pitch angles $(\theta)$ for both agents (unknown $M_{c}$ case).


Figure 4.3: The thrusts for both agents converge to constant and remain bounded.


Figure 4.4: Linear velocities for both agents. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 4.5: Linear velocities for both agents. The $x$ component of the velocity converges to $5.0 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 4.6: Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload.


Figure 4.7: Estimation of the unknown mass. The sum of the individual estimates converges to the actual mass of the load which is 1.5 kg .

### 4.2 Lift Consensus During Cooperative Aerial Load Transport

In Section 4.1, we have developed an adaptive control to estimate the payload when the mass of the payload is not available to the controller. However, the estimates of the mass from the adaptive controller may not converge to the same value. In this section, we enhance the adaptive controller developed in Section 4.1 by incorporating a lift consensus term to ensure that the mass estimate of each agent converges to the same value, i.e., the total mass divided by the number of agents. In particular, we develop a lift consensus algorithm in the $z$-direction such that convergence of the estimates is achieved.

### 4.2.1 Adaptation Law

We denote by $\widehat{M}_{c, i}$ the estimate of the payload mass for agent $i, i=1, \ldots, N$. We propose to update $\widehat{M}_{c, i}$ based on the following equation

$$
\begin{equation*}
\dot{\vec{M}}_{c, i}=\underbrace{-\lambda_{i}\left(\dot{x}_{i}-v^{d}\right)^{T} e_{3}}_{\widehat{\beta}_{i}}+\sum_{j \in \mathcal{N}_{i}} \lambda_{i}\left(\widehat{M}_{c, j}-\widehat{M}_{c, i}\right) \tag{4.16}
\end{equation*}
$$

where $\lambda_{i}>0$ and $\mathscr{N}_{i}$ is the set of neighboring agents that can share their estimates of the payload mass with agent $i$. Note that $\dot{\widehat{M}}_{c, i}$ becomes zero if each agent attains $v^{d}$ in the $e_{3}$ direction and $\widehat{M}_{c, j}=\widehat{M}_{c, i}, \forall j \in \mathscr{N}_{i}$.

Let $1_{N}$ be the $N \times 1$ vectors of ones. Define

$$
\begin{gather*}
\widehat{M}_{c}=\left[\begin{array}{cccc}
\widehat{M}_{c, 1}, & \widehat{M}_{c, 2}, & \widehat{M}_{c, 3}, & \cdots
\end{array} \widehat{M}_{c . N}\right]^{T},  \tag{4.17}\\
\tilde{M}_{c}=\widehat{M}_{c}-\frac{1}{N} M_{c} 1_{N} \tag{4.18}
\end{gather*}
$$

and

$$
\widehat{\beta}=\left[\begin{array}{lllll}
\widehat{\beta}_{1}, & \widehat{\beta}_{2}, & \widehat{\beta}_{3}, & \cdots, & \widehat{\beta}_{N} \tag{4.19}
\end{array}\right]^{T}
$$

We can rewrite (4.1) as

$$
\begin{equation*}
\dot{\vec{M}}_{c}=\widehat{\beta}-\Lambda L \tilde{M}_{c}, \tag{4.20}
\end{equation*}
$$

where $\Lambda=\operatorname{diag}\left\{\lambda_{1} \cdot \lambda_{2}, \cdots, \lambda_{n}\right\}$ and $L \in R^{N \times N}$ is the Laplacian matrix defined as

$$
\ell_{i j}=\left\{\begin{array}{l}
\left|\mathscr{N}_{i}\right| \quad \text { if } i=j  \tag{4.21}\\
-1 \quad \text { if } i \in \mathscr{N}_{i} \\
0 \quad \text { otherwise }
\end{array}\right.
$$

where $\left|\mathscr{N}_{i}\right|$ is the cardinality of the set $\mathscr{N}_{i}$. We assume that the communication topology is bidirectional and connected. Therefore, $L$ is symmetric and positive semidefinite and has the only nontrivial null space spanned by $1_{N}[39]$.

### 4.2.2 Force Control Law

Rewrite (4.2)

$$
\begin{equation*}
F_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+\hat{f}_{i}^{d}+m_{i} g e_{3} \tag{4.22}
\end{equation*}
$$

and construct the $z$-component of $\hat{f}_{i}^{d}$ as:

$$
\begin{equation*}
\hat{f}_{i, z}^{d}=\widehat{M}_{c, i} g \tag{4.23}
\end{equation*}
$$

For the $x$ and $y$ components of $\hat{f}_{i}^{d}$, we assume that they are predefined constant values satisfying

$$
\begin{equation*}
\sum_{i=1}^{N} \hat{f}_{i, x}^{d}=0 \quad \text { and } \quad \sum_{i=1}^{N} \hat{f}_{i, y}^{d}=0 . \tag{4.24}
\end{equation*}
$$

### 4.2.3 Computation of $T_{i}^{d e s}$ and $\Phi_{i}^{d e s}$ for a quadcopter

Given (4.22), the corresponding $T_{i}$ and $\Phi_{i}$ are found based on (3.6)-(3.8).

### 4.2.4 Main Result

The closed-loop dynamics (2.8) with (4.16), (4.22) and (4.23) can be expressed as

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=-\Gamma_{i} \xi_{i}+\hat{f}_{i}^{d}-f_{i} . \tag{4.25}
\end{equation*}
$$

Theorem 4.2 The decentralized adaptive control law in (4.22), (4.23) and (4.16) asymptotically stabilizes the equilibrium of (2.7), (4.16), and (4.25), given by

$$
\begin{gather*}
\mathscr{E}^{*}=\left\{\left(\dot{x}_{i}, \dot{x}_{c}, f_{i}, \widehat{M}_{c, i}\right) \mid \dot{x}_{i}=v^{d}, \quad \dot{x}_{c}=v^{d},\right. \\
\left.\widehat{M}_{c, i}=\frac{1}{N} M_{c}, \quad \text { and } \quad f_{i}=\hat{f}_{i}^{d}\right\} . \tag{4.26}
\end{gather*}
$$

Proof. We first consider the following Lyapunov function

$$
\begin{align*}
V_{1} & =\sum_{i=1}^{N}\left[P_{i}\left(z_{i}\right)-P_{i}\left(z_{i}^{d}\right)-\left(f_{i}^{d}\right)^{T}\left(z_{i}-z_{i}^{d}\right)\right] \\
& +\frac{1}{2}\left(\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \xi_{i}+\xi_{c}^{T} M_{c} \xi_{c}\right) . \tag{4.27}
\end{align*}
$$

From (2.2) and (2.3), the kinematics of $z_{i}$ is given by

$$
\begin{equation*}
\dot{z}_{i}=\dot{x}_{i}-\dot{a}_{i}=\left(\dot{x}_{i}-\dot{x}_{c}\right) \tag{4.28}
\end{equation*}
$$

Define

$$
\begin{equation*}
\xi_{i}=\dot{x}_{i}-v^{d}, \quad \xi_{c}=\dot{x}_{c}-v^{d} \tag{4.29}
\end{equation*}
$$

We compute the time derivative of $V_{1}$ as

$$
\begin{equation*}
\dot{V}_{1}=\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \dot{z}_{i}+\sum_{i=1}^{N} \xi_{i}^{T} m_{i} \ddot{x}_{i}+\xi_{c}^{T} M_{c} \ddot{x}_{c} \tag{4.30}
\end{equation*}
$$

Rewrite (4.30) from (2.7), (2.8), (2.12), (4.22) and (4.28) as:

$$
\begin{align*}
\dot{V}_{1}= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\dot{x}_{i}-\dot{x}_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(F_{i}-f_{i}\right) \\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g_{3}, \\
= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T}\left(\xi_{i}-\xi_{c}\right)+\sum_{i=1}^{N} \xi_{i}^{T}\left(-\Gamma_{i} \xi_{i}+\hat{f}_{i}^{d}-f_{i}\right) \\
& +\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g e_{3},  \tag{4.31}\\
= & \sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \xi_{i}-\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)^{T} \xi_{c}-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i} \\
& +\sum_{i=1}^{N} \xi_{i}^{T}\left(\hat{f}_{i}^{d}-f_{i}\right)+\xi_{c}^{T} \sum_{i=1}^{N} f_{i}-\xi_{c}^{T} M_{c} g e_{3}, \\
= & -\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{f}_{i}^{d},
\end{align*}
$$

where $\tilde{f}_{i}^{d}=\hat{f}_{i}^{d}-f_{i}^{d}$, i.e.,

$$
\tilde{f}_{i}^{d}=\left(\begin{array}{c}
0  \tag{4.32}\\
0 \\
\left(\widehat{M}_{c, i}-\frac{1}{N} M_{c} g\right.
\end{array}\right)=\left(\begin{array}{c}
0 \\
0 \\
\tilde{M}_{c, i} g
\end{array}\right)=\tilde{M}_{c, i} g e_{3}
$$

where $\tilde{M}_{c, i}=\widehat{M}_{c, i}-\frac{1}{N} M_{c}$. From (4.31) and (4.32),

$$
\begin{equation*}
\dot{V}_{1}=-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3} \tag{4.33}
\end{equation*}
$$

We next take another Lyapunov function:

$$
\begin{equation*}
V_{2}=\frac{g}{2} \tilde{M}_{c}^{T} \Lambda^{-1} \tilde{M}_{c} \tag{4.34}
\end{equation*}
$$

Let $W=V_{1}+V_{2}$. The time derivative of $W$ is given by,

$$
\begin{align*}
\dot{W} & =\dot{V}_{1}+\dot{V}_{2}, \\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}+g \tilde{M}_{c}^{T} \Lambda^{-1} \dot{\tilde{M}}_{c}, \\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}+g \tilde{M}_{c}^{T} \Lambda^{-1}\left(\widehat{\beta}-\Lambda L \tilde{M}_{c}\right),  \tag{4.35}\\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}+g \tilde{M}_{c}^{T} \Lambda^{-1} \widehat{\beta}-g \tilde{M}_{c}^{T} \Lambda^{-1} \Lambda L \tilde{M}_{c}, \\
& =-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}+\sum_{i=1}^{N} \xi_{i}^{T} \tilde{M}_{c, i} g e_{3}-\Lambda^{-1} \Lambda g \sum_{i=1}^{N} \tilde{M}_{c, i} \xi_{i}^{T} e_{3}-g \tilde{M}_{c}^{T} L \tilde{M}_{c} .
\end{align*}
$$

Now from (4.16) and (4.35), we obtain

$$
\begin{equation*}
\dot{W}=-\sum_{i=1}^{N} \xi_{i}^{T} \Gamma_{i} \xi_{i}-g \tilde{M}_{c}^{T} L \tilde{M}_{c} \leq 0, \tag{4.36}
\end{equation*}
$$

Since $\dot{W}$ is negative semi-definite, we conclude that the equilibrium $\mathscr{E}^{*}$ is stable.
We employ the LaSalle invariance principle [37] to complete the analysis for the asymptotic stability of $\mathscr{E}^{*}$. Let the largest invariant set where $\dot{W}=0$ be $M$. On $M, \xi_{i}=0$ and $L \tilde{M}_{c}=0$. Since the only null space of $L$ is spanned by $1_{N}, L \tilde{M}_{c}=0$ implies $\tilde{M}_{c}=1_{N} \alpha$, for some scalar $\alpha$. Because $\tilde{M}_{c}=\widehat{M}_{c}-\frac{1}{N} M_{c} 1_{N}=1_{N} \alpha$, we conclude that $\widehat{M}_{c, i}=\widehat{M}_{c, j}, \forall$ $i, j$. We further note that $\xi_{i}=0$ leads to $\dot{x}_{i}=v^{d}$, which means that $\ddot{x}_{i}=0$ and $\hat{f}_{i}^{d}=f_{i}$ from (4.25). Because $\dot{x}_{i}=v^{d}, \dot{\widehat{M}}_{c, i}=0$, which indicates that $\widehat{M}_{c, i}$ is constant. Thus, $f_{i}$ and $z_{i}$ are constant, which leads to $\dot{x}_{c}=\dot{x}_{i}=v^{d}$ and $\ddot{x}_{c}=0$. From (2.7), we obtain $\sum_{i=1}^{N} f_{i}=M_{c} g e_{3}$, which, together with $f_{i}=\hat{f}_{i}^{d}$, yields $\sum_{i=1}^{N} \hat{f}_{i}^{d}=M_{c} g$. Given (4.23), it follows that $\sum_{i=1}^{N} \widehat{M}_{c, i}=M_{c}$ and $\widehat{M}_{c, i}=\widehat{M}_{c, j}=\frac{1}{N} M_{c}$.

Remark 4.1 Theorem 4.2 guarantees that the agents share the weight of the payload equally. If the second term in (4.16) is removed, we still ensure that $\sum_{i=1}^{N} \widehat{M}_{c, i} \rightarrow M_{c}, \dot{x}_{i} \rightarrow v^{d}$, $\dot{x}_{c} \rightarrow v^{d}$ and $f_{i} \rightarrow \hat{f}_{i}^{d}$. However, the estimates $\widehat{M}_{c, i}$ 's may not converge to the same value, which in turn implies that the agents provide different vertical forces to balance the payload mass.

### 4.2.5 Simulation Results: Force control of a payload with an unknown mass for a constant $v^{d}$

We consider three quadcopters with rigid extensions transporting a payload. The mass and inertia are kept identical to our previous chapter 2, [4]. We also employ the same PD control as in Chapter 2, to track the desired orientation for each quadcopter.

Throughout this simulation, we assume that agent 1 communicates with agent 2 and 3 but agent 2 and 3 do not communicate with each other. We choose $\widehat{M}_{c, 1}(0)=0.1 \mathrm{~kg}$,


Figure 4.8: Communication topology between each agent.
$\widehat{M}_{c, 2}(0)=0.5 \mathrm{~kg}$ and $\widehat{M}_{c, 3}(0)=1.2 \mathrm{~kg}$,

$$
\begin{aligned}
v^{d} & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
2.0 & 1.0 & 0
\end{array}\right]^{T} \mathrm{~m} / \mathrm{s}, \\
\hat{f}_{1}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
0.0 & -2.0 & \widehat{M}_{c, 1} g
\end{array}\right]^{T}, \\
\hat{f}_{2}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
2.0 & 1.0 & \widehat{M}_{c, 2} g
\end{array}\right]^{T}, \\
\hat{f}_{3}^{d}(t) & =\left(1-e^{-t}\right)\left[\begin{array}{lll}
-2.0 & 1.0 & \widehat{M}_{c, 3} g
\end{array}\right]^{T}, \\
\Gamma_{i} & =\left[\begin{array}{ccc}
10 & 0 & 0 \\
0 & 10 & 0 \\
0 & 0 & 25
\end{array}\right], \\
L & =\left[\begin{array}{ccc}
2 & -1 & -1 \\
-1 & 1 & 0 \\
-1 & 0 & 1
\end{array}\right] .
\end{aligned}
$$

Note that $v^{d}(t), f_{1}^{d}(t)$, and $f_{2}^{d}(t)$ exponentially converge to constants. The term $\left(1-e^{-t}\right)$ is added to improve performance of the controller during the transition.

Fig. 4.9 and 4.10 illustrate that the velocities of the three agents and the load asymptotically approach $v^{d}$. The contact forces $f_{i}$ 's also converge to the desired values $f_{i}^{d}$ 's as observed in Fig. 4.11. The adaptation law from (4.16) estimates the true mass of the load as illustrated in Fig. 4.12 and the consensus algorithm guarantees that each agent estimates converge to $\frac{1}{3} M_{c}=0.5 \mathrm{~kg}$.


Figure 4.9: The magnitude of the velocity error $\left(v_{i}-v^{d}\right)$ goes to zero which concludes that $v_{i}$ converges to predesigned $v^{d}$


Figure 4.10: Linear velocity plot for the payload. The $x$ and $y$ velocity converges to the desired $v^{d}$ i.e., $2.0 \mathrm{~m} / \mathrm{s}$ and $1.0 \mathrm{~m} / \mathrm{s}$ respectively and the $z$ component converges to zero.


Figure 4.11: The magnitude of $f_{i}-f_{i}^{d}$ goes to zero. All the contact forces $f_{i}$ converge to $f_{i}^{d}$.


Figure 4.12: Plot of the estimates of the mass of the payload. Each individual estimates converge to $\frac{1}{3} M_{c}=0.5 \mathrm{~kg}$ and its sum adds up-to the true mass of the payload i.e., 1.5 kg .


Figure 4.13: Plot for contact forces in all 3 directions for agent 1. It converges to converges $0.0 \mathrm{~N},-2.0 \mathrm{~N}$ and 0.5 N for each direction respectively. Remaining agents have similar graphs.


Figure 4.14: Plot for desired and actual orientation angles $(\phi)$ for three agents in the $x$-direction.


Figure 4.15: Plot for the desired and actual orientation angles $(\theta)$ for three agents in the $y$-direction.

### 4.2.6 Simulation Results: Force Control of a load with an unknown mass for a timevarying $v^{d}(t)$

To test the performance of the controllers for different conditions, we illustrate a simulation for a pre-designed time-varying desired velocity. We keep the other parameters identical to the previous constant velocity case and change $v^{d}$ to the following:

$$
v^{d}(t)= \begin{cases}{\left[\begin{array}{lll}
2.0, & 0, & 0
\end{array}\right]^{T},} & \mathrm{t} \leq 1.57 \mathrm{~s} \\
{\left[\begin{array}{ll}
2 \sin (t), & -2 \cos (t),
\end{array}\right]^{T},} & 2<\mathrm{t} \leq 9.42 \mathrm{~s} \\
{\left[\begin{array}{lll}
0, & 2.0, & 0
\end{array}\right]^{T},} & \mathrm{t}>9.42 \mathrm{~s}\end{cases}
$$

Fig. 4.16 and 4.17 illustrate the 2D and 3D trajectories for all three agents and the load respectively. Notice from Fig. 4.18, 4.19 and 4.20 that the velocity errors of the three agents and the load are bounded. The contact force errors are also bounded and stay close to zero as illustrated in Fig. 4.21 and 4.22. The update law in (4.16) recovers the true mass of the load as illustrated in Fig. 4.23 and the consensus algorithm guarantees that each agent estimates converge to $\frac{1}{3} M_{c}=0.5 \mathrm{k}$ even with a time varying velocity in $x$ and $y$ direction. Note that the convergence of the estimates to $\frac{1}{3} M_{c}=0.5$ is because the $z$-component of the desired velocity $v^{d}$ is still constant. Fig. 4.24 and 4.25 illustrate that the low level PD controller closely tracks the attitude of the quadcopter during the transient and as well as steady state.


Figure 4.16: Snapshot of 3D position trajectories for all the agents and the payload at different time instance for adaptive design with time varying velocity.


Figure 4.17: Position trajectories for aerial robots and the payload in $x$ and $y$ direction.


Figure 4.18: The magnitude of the velocity error $\left(v_{i}-v^{d}\right)$ for the agents.


Figure 4.19: The magnitude of the velocity error $\left(v_{c}-v^{d}\right)$ for the payload.


Figure 4.20: Linear velocity for the payload in all 3 directions.


Figure 4.21: The magnitude of $f_{i}-f_{i}^{d}$ stays bounded and eventually converges to zero.


Figure 4.22: Plot of squeeze forces in all 3 directions for agent 1 . Note that other agents have similar plot.


Figure 4.23: Plot of the estimates of the mass of the payload. Eachg individual estimates converges $\frac{1}{3} M_{c}=0.5 \mathrm{~kg}$. Notice that their sum adds up-to the true mass of the payload used in this simulation. The convergence of the estimates to $\frac{1}{3} M_{c}=0.5$ is because the $z$-component of the desired velocity $\nu^{d}$ is still constant.


Figure 4.24: Plot for desired and actual orientation angles $(\phi)$ for three agents in the $x$-direction.


Figure 4.25: Plot for the desired and actual orientation angles $(\theta)$ for three agents in the $y$-direction.

## CHAPTER 5

## COOPERATIVE AERIAL MANIPULATION WITH ATTITUDE STABILIZATION

In our previous Chapters 3 and 4 we have developed force and motion control algorithms that guarantee force regulation and translational motion of the quadcopters and the payload. For complete motion coordination and path planning, orientation control is important. In this chapter, we extend our work in Chapters 3 and 4, to address both the orientation control of the payload and the force control during cooperative aerial manipulation. In particular, we develop time varying force setpoints to enforce attitude regulation. The algorithm provides desired thrust and attitude angles required for each quadcopter to cooperatively transport the payload. We analyze the stability of the system using singular perturbation theory. We demonstrate the effectiveness of the algorithms in numerical simulations.

### 5.1 Related Work

Reference [18] uses three or more quadcopters for payload transport to balance the attitude of the payload. Reference [27] considers a ring shaped flexible structure with multiple attached flying vehicles and develops a LQR controller for linearized dynamics to track the desired pose of the payload. It is assumed that the vehicles can exert both forces and moments on the payload. However, we assume that the inputs from the quadcopters to the payload are only contact forces with no moments.

### 5.2 Problem Formulation

### 5.2.1 Kinematics

Let the orientation of the load with respect to $\Sigma_{I}$ be $R_{c} \in S O(3)$. Let $q=\left[\begin{array}{l}q_{0} \\ q_{\nu}\end{array}\right]$ be the unit quaternion representation parameterizing $R_{c}$. Then $R_{c}$ is related to $q$ through the Rodriguez formula [40]:

$$
\begin{equation*}
R_{c}=I_{3}+2\left(\widehat{q_{v}}\right)^{2}+2 q_{0} \widehat{q_{v}} \tag{5.1}
\end{equation*}
$$

where the operator ${ }^{\wedge}$ converts a given vector $\omega=\left[\begin{array}{lll}\omega_{1} & \omega_{2} & \omega_{3}\end{array}\right]^{T}$ into a skew symmetric matrix as

$$
\widehat{\omega}=\left[\begin{array}{ccc}
0 & -\omega_{3} & \omega_{2}  \tag{5.2}\\
\omega_{3} & 0 & -\omega_{1} \\
-\omega_{2} & \omega_{1} & 0
\end{array}\right]
$$

Note that $\widehat{x} y=x \times y$ for all $x, y \in \mathbb{R}^{3}$. The kinematic equations of $q$ are given by

$$
\begin{gather*}
\frac{d q_{0}}{d t}=-\frac{1}{2}\left(\Omega_{c}\right)^{T} q_{v}  \tag{5.3}\\
\frac{d q_{v}}{d t}=\frac{1}{2} q_{0} \Omega_{c}+\frac{1}{2} \widehat{\Omega}_{c} q_{v} \tag{5.4}
\end{gather*}
$$

where $\Omega_{c}$ is the angular velocity of the payload in the inertial frame.
Because the contact points $a_{i}$ satisfy

$$
\begin{equation*}
a_{i}(t):=x_{c}(t)+R_{c} r_{i}, \tag{5.5}
\end{equation*}
$$

the kinematics of the contact points are given by

$$
\begin{equation*}
\dot{a}_{i}=\dot{x}_{c}+\left(R_{c} r_{i}\right)=\dot{x}_{c}+\widehat{\Omega}_{c} R_{c} r_{i} \tag{5.6}
\end{equation*}
$$

### 5.2.2 Dynamics

We model the contact force between the agents and the payload as a mass spring model. This assumption is valid when the vehicles are connected to the payload via
elastic cables or when the payload is flexible or surrounded by elastic bumper materials. When the agents move, the payload experiences tension or compression. Therefore, $a_{i}(t) \neq x_{i}(t)$. We approximate the deformation as

$$
\begin{equation*}
z_{i}=x_{i}-a_{i}, \quad i=1, \cdots, N . \tag{5.7}
\end{equation*}
$$

In this paper, we consider a linear spring model for the contact force, i.e., $f_{i}=k z_{i}$, where $k$ is the spring constant (possibly unknown). If the contact between the agents and the payload is elastic at a constant pressure, the linear spring model is a valid representation of the contact force as shown in [29].

Rewrite the translational dynamics of the load as

$$
\begin{equation*}
M_{c} \ddot{x}_{c}=\sum_{i=1}^{N} f_{i}-M_{c} g e_{3} \tag{5.8}
\end{equation*}
$$

where $M_{c}$ is the mass of the load, $g$ is the gravitational constant, and $e_{3}$ is the unit vector $\left[\begin{array}{lll}0 & 0 & 1\end{array}\right]^{T}$. The rotational dynamics of the payload is given by

$$
\begin{equation*}
I_{c} \dot{\Omega}_{c}=\sum_{i=1}^{N} \widehat{R_{c} r_{i}} f_{i}-\widehat{\Omega}_{c} I_{c} \Omega_{c}, \tag{5.9}
\end{equation*}
$$

where the $\sum_{i=1}^{N} \widehat{R_{c} r_{i}} f_{i}$ term represents the total moment exerted on the payload due to the contact forces. We assume that the agents cannot directly exert moments on the payload.

Rewrite the translational dynamics of the $N$ agents as

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=F_{i}-f_{i}-m_{i} g e_{3}, \quad i=1, \ldots, N, \tag{5.10}
\end{equation*}
$$

where $m_{i}$ is the mass of the $i$-th agent, $F_{i}$ is the force applied to agent $i$ and $f_{i}$ is the contact force to agent $i$. For a quadcopter, the applied force $F_{i}$ is given by

$$
\begin{equation*}
F_{i}=T_{i} R_{i}(\Phi) e_{3}, \tag{5.11}
\end{equation*}
$$

where $T_{i}$ is the total thrust from the four rotors, given by

$$
\begin{equation*}
T_{i}=\sum_{i=1}^{4} F_{\text {rotors }, i}, \tag{5.12}
\end{equation*}
$$

$\Phi=\left[\begin{array}{lll}\phi & \theta & \psi\end{array}\right]^{T}$ are the roll, pitch and yaw angles and $R_{i}(\Phi) \in S O(3)$ is the rotation matrix between the body frame $O_{c, i}$ and the inertial frame $\Sigma_{i}$ given by

$$
R_{i}=\left[\begin{array}{ccc}
c \psi_{i} c \theta_{i} & c \psi_{i} s \phi_{i} s \theta_{i}-c \phi_{i} s \psi_{i} & s \phi_{i} s \psi_{i}+c \phi_{i} c \psi_{i} s \theta_{i}  \tag{5.13}\\
c \theta_{i} s \psi_{i} & c \phi_{i} c \psi_{i}+s \phi_{i} s \psi_{i} s \theta_{i} & c \phi_{i} s \psi_{i} s \theta_{i}-c \psi_{i} s \phi_{i} \\
-s \theta_{i} & c \theta_{i} s \phi_{i} & c \phi_{i} c \theta_{i}
\end{array}\right]
$$

in which $c \theta_{i}=\cos \theta_{i}$ and $s \theta_{i}=\sin \theta_{i}$. The rotational dynamics of the $i$-th quadcopter is given by

$$
\begin{equation*}
I_{i} \dot{\Omega}_{i}=M_{i}-\widehat{\Omega}_{i} I_{i} \Omega_{i},+\tau_{i, f}, \tag{5.14}
\end{equation*}
$$

where $I_{i}$ is the inertia matrix, $M_{i} \in \mathbb{R}^{3}$ is the moment generated by the $i$ th-quadcopter and $\tau_{i, f}$ is the torque generated due to the contact force $f_{i}$.

When $f_{i}$ is measured, e.g., based on a force sensor or an accelerometer, we may cancel $\tau_{i, f}$ by choosing $M_{i}=\bar{M}_{i}-\tau_{i, f}$. Then we can design $\bar{M}_{i}$ to control the attitude dynamics of the $i$-th quadcopter.

### 5.2.3 Control objective

Our control objective is to design $F_{i}$ such that all the agents and the payload translate with a constant velocity $v^{d}$ while the orientation of the payload $R_{c}$ is driven to a desired setpoint $R_{c}^{d}$, i.e., $\dot{x}_{i} \rightarrow v^{d}, \forall i, \dot{x}_{c} \rightarrow v^{d}, \Omega_{c} \rightarrow 0$, and $R_{c} \rightarrow R_{c}^{d}$. Meanwhile, we also regulate the contact force $f_{i}$ to desired squeezing force $f_{i}^{d}$, i.e., $f_{i} \rightarrow f_{i}^{d}, \forall i$.

### 5.3 Control Design

Compared with Chapters 4 and 6, our control objective consists of an additional goal of stabilizing $R_{c}$ to $R_{c}^{d}$. To achieve this goal, we design time-varying force setpoints $f_{i}^{d}(t)$ such that when substituted in (5.9), the $f_{i}^{d}(t)$ 's achieve the orientation control. Specifically, we choose $f_{i}^{d}(t)$ dynamically to satisfy

$$
\begin{equation*}
\sum_{i=1}^{N} f_{i}^{d}=M_{c} g e_{3} \tag{5.15}
\end{equation*}
$$

$$
\begin{equation*}
\sum_{i=1}^{N} \widehat{R_{c} r_{i}} f_{i}^{d}=-k_{\Omega_{c}} \Omega_{c}-k_{q_{v}} q_{v}^{e} \tag{5.16}
\end{equation*}
$$

where $k_{\Omega_{c}}$ and $k_{q_{v}}$ are positive scalars, and $q_{\nu}^{e}$ is the vector part of the quaternion parameterizing $\left(R_{c}^{d}\right)^{T} R_{c}$. Note that the right hand side of (5.16) is an exponentially stable attitude control that ensures $\Omega_{c} \rightarrow 0$ and $q_{v}^{e} \rightarrow 0[40,41]$. Therefore, if $f_{i}=f_{i}^{d}$, it follows from (5.8) and (5.9) and the constraints (5.15) and (5.16) that the payload will move at a constant velocity with a stabilized orientation.

Define

$$
\begin{equation*}
u_{i}=f_{i}^{d}(t)-f_{i} \tag{5.17}
\end{equation*}
$$

and

$$
\begin{equation*}
\xi_{i}:=\dot{x}_{i}-v^{d}-\Gamma_{i} u_{i} \quad \Gamma_{i}=\Gamma_{i}^{T}>0 \tag{5.18}
\end{equation*}
$$

We next design $F_{i}$ as

$$
\begin{equation*}
F_{i}=-K_{i} \xi_{i}+f_{i}^{d}(t)+m_{i} \Gamma_{i} \dot{u}_{i}+m_{i} g e_{3} \tag{5.19}
\end{equation*}
$$

where $K_{i}=K_{i}^{T}>0$. As we shall see in Section 5.4, the $\dot{u}_{i}$ term allows us to establish stability of the combined orientation and force control system.

Remark 5.1 From (5.15) and (5.16), we obtain

$$
\underbrace{\left[\begin{array}{cccc}
I & I & \ldots & I  \tag{5.20}\\
\widehat{R_{c} r_{1}} & \widehat{R_{c} r_{2}} & \ldots & \widehat{R_{c} r_{N}}
\end{array}\right]}_{\Lambda} \underbrace{\left[\begin{array}{c}
f_{1}^{d} \\
f_{2}^{d} \\
\vdots \\
f_{N}^{d}
\end{array}\right]}_{f^{d}}=\underbrace{\left[\begin{array}{c}
M_{c} g e_{3} \\
-k_{\Omega_{c}} \Omega_{c}-k_{q_{v}} q_{v}^{e}
\end{array}\right]}_{h\left(\Omega_{c}, q_{v}^{e}\right)}
$$

Thus, $f^{d}$ can be solved as

$$
\begin{equation*}
f^{d}=\Lambda^{+} h\left(\Omega_{c}, q_{v}^{e}\right)+f_{N(\Lambda)}^{d} \tag{5.21}
\end{equation*}
$$

where $\Lambda^{+}=\Lambda^{T}\left(\Lambda \Lambda^{T}\right)^{-1}$ is the Moore-Penrose generalized inverse of $\Lambda$ and $f_{N(\Lambda)}^{d}$ belongs to the null space of $\Lambda$, i.e., $\Lambda f_{N(\Lambda)}^{d}=0$.

Note that $f^{d}$ consists of two components: the first one solving the constraints (5.15) and (5.16) and the second one aligned with the squeezing direction (the null space of $\Lambda$ ). The $f^{d}$ is a function of $\left(\Omega_{c}, q_{\nu}^{e}\right)$ since its two components depend only on $\left(\Omega_{c}, q_{\nu}^{e}\right)$. At the desired equilibrium where $\Omega_{c}=0$ and $q_{v}^{e}=0$, the sum of $f_{i}^{d}$ 's in the $z$ direction balances the weight of the load and the $f_{i}^{d}$ 's contribute no moments to the payload.

The $\dot{f}_{i}^{d}$ used in $\dot{u}_{i}$ can be computed in a similar fashion. The expression of $\dot{f}_{i}^{d}$ is given in the Appendix.

### 5.3.1 Computation of $T_{i}^{d e s}$ and $\Phi_{i}^{\text {des }}$ for a quadcopter

Given (5.11), the corresponding $T_{i}$ and $\Phi_{i}$ are found based on (3.6)-(3.8). Once $\theta_{i}^{\text {des }}$, $\phi_{i}^{d e s}$ and $T_{i}^{d e s}$ are computed, low-level attitude and thrust tracking controllers (e.g., a PD controller) can be implemented to track these desire commands for the $i$-th quadcopter.

### 5.4 Stability Analysis

### 5.4.1 Closed-loop system and equilibria

In this section, we derive the closed-loop dynamics. We note that $u_{i}=k\left(z_{i}^{d}-z_{i}\right)$, where $z_{i}^{d}=f_{i}^{d} / k$. Define $\tilde{z}_{i}=z_{i}^{d}-z_{i}, \forall i$. Then $u_{i}=0$ is equivalent to $\tilde{z}_{i}=0$. Therefore, we derive the dynamics for $\tilde{z}_{i}$. Differentiating (5.7) leads to

$$
\begin{equation*}
\dot{z}_{i}=\dot{x}_{i}-\dot{a}_{i}=\xi_{i}+v^{d}+\Gamma_{i} u_{i}-\dot{x}_{c}-\widehat{\Omega}_{c} R_{c} r_{i} \tag{5.22}
\end{equation*}
$$

which can be further written as

$$
\begin{align*}
\dot{z}_{i} & =\xi_{i}-\xi_{c}+\Gamma_{i} u_{i}-\widehat{\Omega}_{c} R_{c} r_{i}  \tag{5.23}\\
& =\xi_{i}-\xi_{c}+K_{z} \tilde{z}_{i}-\widehat{\Omega}_{c} R_{c} r_{i},
\end{align*}
$$

where $\xi_{c}=\dot{x}_{c}-v^{d}$ and $K_{z}=\Gamma_{i} k$. Thus, we obtain

$$
\begin{equation*}
\dot{\tilde{z}}_{i}=-\xi_{i}+\xi_{c}-K_{z} \tilde{z}_{i}+\widehat{\Omega}_{c} R_{c} r_{i}+\dot{z}_{i}^{d} . \tag{5.24}
\end{equation*}
$$

Note from (5.46) that $\dot{z}_{i}^{d}$ is a function of $\left(\Omega_{c}, q_{\nu}^{e}, \tilde{z}_{i}\right)$.

For the translational motion of the payload, we obtain from (5.8) and (5.15) that

$$
\begin{equation*}
M_{c} \dot{\xi}_{c}=\sum_{i=1}^{N}\left(f_{i}-f_{i}^{d}\right)=-k \sum_{i=1}^{N} \tilde{z}_{i} . \tag{5.25}
\end{equation*}
$$

The derivative of $\xi_{i}$ defined in (5.18) is given by

$$
\begin{equation*}
m_{i} \dot{\xi}_{i}=m_{i} \ddot{x}_{i}-m_{i} \Gamma_{i} \dot{u}_{i}=F_{i}-f_{i}-m_{i} g e_{3}-m_{i} \Gamma_{i} \dot{u}_{i}, \tag{5.26}
\end{equation*}
$$

which, together with the design of $F_{i}$ in (5.19), leads to

$$
\begin{align*}
m_{i} \dot{\xi}_{i} & =-K_{i} \xi_{i}+f_{i}^{d}+m_{i} \Gamma_{i} \dot{u}_{i}+m_{i} g e_{3} \\
& -f_{i}-m_{i} g e_{3}-m_{i} \Gamma_{i} \dot{u}_{i}  \tag{5.27}\\
m_{i} \dot{\xi}_{i} & =-K_{i} \xi_{i}+k \tilde{z}_{i} .
\end{align*}
$$

The rotational dynamics in (5.9) with $f_{i}^{d}$ in (5.16) takes the following form:

$$
\begin{equation*}
I_{c} \dot{\Omega}_{c}=-\widehat{\Omega}_{c} I_{c} \Omega_{c}-k \sum_{i=1}^{N} \widehat{R_{c} r_{i}} \tilde{z}_{i}-k_{\Omega_{c}} \Omega_{c}-k_{q_{v}} q_{v}^{e}, \tag{5.28}
\end{equation*}
$$

where the dynamics of the unit quaternion $\left(q_{\nu}^{0}, q_{\nu}^{e}\right)$ is given by

$$
\begin{gather*}
\dot{q}_{0}^{e}=-\frac{1}{2}\left(\left(R_{c}^{d}\right)^{T} \Omega_{c}\right)^{T} q_{v}^{e},  \tag{5.29}\\
\dot{q}_{v}^{e}=\frac{1}{2} q_{0}^{e}\left(\left(R_{c}^{d}\right)^{T} \Omega_{c}\right)+\frac{1}{2} \widehat{\left(R_{c}^{d}\right)^{T} \Omega_{c}} q_{v}^{e} . \tag{5.30}
\end{gather*}
$$

The desired equilibrium of (5.24), (5.25), (5.27), (5.28), and (5.30) is given by the origin of $\left\{\Xi, \xi_{c}, \tilde{z}, \Omega_{c}, q_{\nu}^{e}\right\}$, where $\Xi=\left(\xi_{1}^{T}, \cdots, \xi_{N}^{T}\right)^{T}$ and $\tilde{z}=\left(\tilde{z}_{1}^{T}, \cdots, \tilde{z}_{N}^{T}\right)$.

### 5.4.2 Main result

Theorem 5.1 Suppose that $K_{z}=\Gamma_{i} k=\lambda I>0$ and $\lambda$ is sufficiently large. Define $\zeta_{i}=\dot{\tilde{z}}_{i}$. The control law (5.19) ensures that the equilibria in (5.31) below are exponentially stable.

$$
\begin{gather*}
\mathscr{E}=\left\{\left(\xi_{i}, \zeta_{i}, \tilde{z}_{i}, \Omega_{c}, q_{v}^{e}\right) \mid \xi_{i}=0, \zeta_{i}=0,\right.  \tag{5.31}\\
\left.\tilde{z}_{i}=0, \Omega_{c}=0, \text { and } q_{\nu}^{e}=0, \forall i\right\}
\end{gather*}
$$

In addition, convergence to $\mathscr{E}$ is equivalent to convergence to the origin of $\left\{\Xi, \xi_{c}, \tilde{z}, \Omega_{c}, q_{v}^{e}\right\}$.

Proof. From $\zeta_{i}=\dot{\tilde{z}}_{i}$ and (5.24), we obtain

$$
\begin{equation*}
\zeta_{i}=\dot{\tilde{z}}_{i}=-\xi_{i}+\xi_{c}-K_{z} \tilde{z}_{i}+\widehat{\Omega}_{c} R_{c} r_{i}+\dot{z}_{i}^{d}, \forall i . \tag{5.32}
\end{equation*}
$$

We first divide both sides of (5.32) by $K_{z}$ and obtain

$$
\begin{equation*}
\epsilon \dot{\tilde{z}}_{i}=-\epsilon \xi_{i}+\epsilon \xi_{c}-\tilde{z}_{i}+\epsilon \widehat{\Omega}_{c} R_{c} r_{i}+\epsilon \dot{z}_{i}^{d} \tag{5.33}
\end{equation*}
$$

where $\epsilon=1 / \lambda$ and $\dot{z}_{i}^{d}$ is given in (5.46).
We next derive the dynamics of $\zeta_{i}$. We obtain from (5.32), (5.25), (5.27), and (5.28) that

$$
\begin{align*}
& \dot{\zeta}_{i}=-\dot{\xi}_{i}+\dot{\xi}_{c}-K_{z} \dot{\tilde{z}}_{i}+\dot{\widehat{\Omega}}_{c} R_{c} r_{i}+\widehat{\Omega}_{c} \widehat{\Omega}_{c} R_{c} r_{i}+\ddot{z}_{i}^{d} \\
& \dot{\zeta}_{i}=\frac{1}{m_{i}} K_{i} \xi_{i}-\frac{k}{m_{i}} \tilde{z}_{i}-\frac{k}{M_{c}} \sum_{i=1}^{N} \tilde{z}_{i}-K_{z} \zeta_{i} \\
& -\left(I_{c}^{-1}\left(\sum_{i=1}^{N} \widehat{R_{c} r_{i}} k \tilde{z}_{i}+\widehat{\Omega}_{c} I_{c} \Omega_{c}+k_{\Omega_{c}} \Omega_{c}+k_{q_{v}} q_{v}^{e}\right)\right) \times R_{c} r_{i}  \tag{5.34}\\
& +\left(\widehat{\Omega}_{c}\right)^{2} R_{c} r_{i}+\ddot{z}_{i}^{d} .
\end{align*}
$$

Dividing both sides by $K_{z}$ yields

$$
\begin{align*}
\epsilon \dot{\zeta}_{i} & =\epsilon \frac{1}{m_{i}} K_{i} \xi_{i}-\frac{1}{m_{i}} \Gamma_{i}^{-1} \tilde{z}_{i}-\frac{1}{M_{c}} \Gamma_{i}^{-1} \sum_{i=1}^{N} \tilde{z}_{i} \\
& -\zeta_{i}+\epsilon\left(\widehat{\Omega}_{c}\right)^{2} R_{c} r_{i}+\epsilon \ddot{z}_{i}^{d}+\Gamma_{i}^{-1} \widehat{R_{c} r_{i}}\left(I_{c}^{-1} \sum_{i=1}^{N} \widehat{R_{c} r_{i}} \tilde{z}_{i}\right)  \tag{5.35}\\
& +\epsilon \widehat{R_{c} r_{i}}\left(I_{c}^{-1} \widehat{\Omega}_{c} I_{c} \Omega_{c}\right)+\epsilon \widehat{R_{c} r_{i}}\left(I_{c}^{-1} k_{\Omega_{c}} \Omega_{c}\right) \\
& +\epsilon \widehat{R_{c} r_{i}}\left(I_{c}^{-1} k_{q_{v}} q_{v}^{e}\right)
\end{align*}
$$

Further computation from (5.46) shows that $\ddot{z}_{i}^{d}$ is a function of $\left(\Omega_{c}, \tilde{z}_{i}, q_{v}^{e}, \zeta_{i}\right)$ and does not contain terms multiplied by $K_{z}$ or $k$.

For a sufficiently small $\epsilon$, there exists a time-scale separation between the fast dynamics (5.33) and (5.35) and the slow dynamics of $\xi_{i}, \Omega_{c}$ and $q_{\nu}^{e}$ in (5.27), (5.28) and (5.30), respectively. To analyze such a singularly perturbed system, we first set $\epsilon=0$ and equate the right hand side of (5.33) and (5.35) to zero to obtain

$$
\begin{equation*}
\tilde{z}_{i}=0 \quad \text { and } \quad \zeta_{i}=0, \forall i . \tag{5.36}
\end{equation*}
$$

Substituting (5.36) into (5.27) and (5.28) yields the reduced order system, which consists of

$$
\begin{align*}
\dot{\xi}_{i} & =-\frac{K_{i}}{m_{i}} \xi_{i}  \tag{5.37}\\
\dot{\Omega}_{c} & =-I_{c}^{-1}\left(k_{\Omega_{c}} \Omega_{c}+\widehat{\Omega}_{c} I_{c} \Omega_{c}+k_{q v} q_{v}^{e}\right) \tag{5.38}
\end{align*}
$$

and (5.30). The origin of the reduced system is exponentially stable because the $\xi_{i}$ dynamics are exponentially stable and the dynamics of $\Omega_{c}$ and $q_{v}^{e}$ are also exponentially stable according to $[40,41]$.

The boundary layer system of (5.33) and (5.35) is given by

$$
\begin{align*}
& \frac{d \tilde{z}_{i}}{d \tau}=-\tilde{z}_{i}  \tag{5.39}\\
& \frac{d \zeta_{i}}{d \tau}=-\zeta_{i}- \\
& \quad \Gamma_{i}^{-1}\left(\frac{1}{m_{i}} \tilde{z}_{i}+\frac{1}{M_{c}} \sum_{i=1}^{N} \tilde{z}_{i}-\widehat{R_{c} r_{i}}\left(I_{c}^{-1} \sum_{i=1}^{N} \widehat{R_{c} r_{i}} \tilde{z}_{i}\right)\right) . \tag{5.40}
\end{align*}
$$

Note that (5.39) and (5.40) is a cascaded system, where the coupling term in (5.40) is linear in $\tilde{z}_{i}$ with a bounded growth rate. Without the coupling term, (5.39) and (5.40) are both exponentially stable linear systems. Therefore, it is straightforward to employ a Lyapunov function $V=\frac{1}{2} \sum_{i=1}^{N} \tilde{z}_{i}^{T} P \tilde{z}_{i}+\zeta_{i}^{T} \zeta_{i}$ with a sufficiently large $P$ to show that the origin of $(\tilde{z}, \zeta)$ is exponentially stable, where $\zeta=\left(\zeta_{1}^{T}, \cdots, \zeta_{N}^{T}\right)$. We then conclude from [37, Theorem 11.4], the origin of the combined system (5.27), (5.28), (5.30), (5.33), and (5.35) given in (5.31) is exponentially stable for a sufficiently large $\Gamma_{i} k$.

We now show that convergence to $\mathscr{E}$ is equivalent to convergence to the origin of $\left\{\Xi, \xi_{c}, \tilde{z}, \Omega_{c}, q_{\nu}^{e}\right\}$. Specifically, we establish that convergence to $\mathscr{E}$ leads to $\xi_{c} \rightarrow 0$. From (5.24), $\dot{\tilde{z}}_{i}=\zeta_{i} \rightarrow 0, \xi_{i} \rightarrow 0, \Omega_{c} \rightarrow 0$, and $\tilde{z} \rightarrow 0$, we conclude that $\xi_{c}+\dot{z}_{i}^{d} \rightarrow 0$. The expression of $\dot{z}^{d}$ in (5.46) indicates that $\dot{z}^{d}-\frac{1}{k} \dot{f}_{N(\Lambda)}^{d} \rightarrow 0$. As $\Omega_{c} \rightarrow 0$ and $q_{v}^{e} \rightarrow 0, R_{c}$ approaches $R_{c}^{d}$, which means that $f_{N(\Lambda)}^{d}$ can be chosen to converge to a constant vector and thus $\dot{f}_{N(\Lambda)}^{d} \rightarrow 0$. Therefore, it follows that $\dot{z}_{i}^{d} \rightarrow 0$ and $\xi_{c} \rightarrow 0$.

The proof for Theorem 5.1 illustrates that with a sufficiently large $\Gamma_{i} k$, the force error dynamics evolves at the fast time-scale while the velocity dynamics of the agents and the rotational dynamics of the payload evolve at the slow time-scale. If $K_{i}$ in (5.27) is also chosen sufficiently large, the velocity dynamics of the agents $\xi_{i}$ becomes part of the fast dynamics. We can follow a similar singular perturbation analysis to obtain the same stability and convergence results.

The value of $\Gamma_{i} k$ depends on both $\Gamma_{i}$ and $k$. For a stiff spring, where $k$ is large, a moderate $\Gamma_{i}$ will be sufficient. For a soft spring where $k$ is small, a large $\Gamma_{i}$ is needed. Although our analysis assumes a uniform $k$ for all the agents, the analysis can be easily extended to non-uniform cases.

The current implementation of $F_{i}$ requires $f_{i}, \dot{f}_{i}, f_{i}^{d}$ and $\dot{f}_{i}^{d}$. The $\dot{f}_{i}$ may be estimated using a derivative estimator, such as [42], based on the $f_{i}$ measurements. As shown in the below, the expression of $\dot{f}_{i}^{d}$ is complicated and may be difficult to implement. We are currently investigating the sensitivity and robustness of the proposed algorithm with respect to $\dot{f}_{i}^{d}$.

### 5.4.3 Computation of $\dot{f}_{i}^{d}$ and $\dot{z}_{i}^{d}$

The time derivative of (2.12) yields

$$
\begin{equation*}
\sum_{i=1}^{N} \dot{f}_{i}^{d}=0 \tag{5.41}
\end{equation*}
$$

The time derivative of (5.16) yields

$$
\begin{align*}
\sum_{i=1}^{N} \widehat{\left(R_{c} r_{i}\right)} \dot{f}_{i}^{d} & =-k_{\Omega_{c}} I_{c}^{-1}\left(-\widehat{\Omega}_{c} I_{c} \Omega_{c}-\sum_{i=1}^{N} \widehat{R_{c} r_{i}} u_{i}\right. \\
& \left.-k_{\Omega_{c}} \Omega_{c}-k_{q_{v}} q_{v}^{e}\right) \\
& -k_{q_{v}} \dot{q}_{v}^{e}-\sum_{i=1}^{N} \widehat{\left({ }_{c} R_{c} r_{i}\right.} f_{i}^{d}  \tag{5.42}\\
& =k_{\Omega_{c}} I_{c}^{-1} \widehat{\Omega}_{c} I_{c} \Omega_{c}+k_{\Omega_{c}} \sum_{i=1}^{N} \widehat{R_{c} r_{i}} u_{i} \\
& -k_{\Omega_{c}} I_{c}^{-1} k_{q_{v}} q_{v}^{e}-k_{q_{v}} \dot{q}_{v}^{e}+k_{\Omega_{c}}^{2} I_{c}^{-1} \Omega_{c} \\
& +k_{q v} \widehat{\Omega}_{c} q_{v}^{e}+\sum_{i=1}^{N} \widehat{R_{c} r_{i}} \widehat{\Omega}_{c} f_{i}^{d} .
\end{align*}
$$

From (5.41) and (5.42), we obtain

$$
\begin{align*}
& {\left[\begin{array}{cccc}
I & I & \ldots & I \\
\widehat{R_{c} r_{1}} & \widehat{R_{c} r_{2}} & \ldots & \widehat{R_{c} r_{N}}
\end{array}\right]\left[\begin{array}{c}
\dot{f}_{1}^{d} \\
\dot{f}_{2}^{d} \\
\vdots \\
\dot{f}_{N}^{d}
\end{array}\right]} \\
& =k_{\Omega_{c}} k\left[\begin{array}{cccc}
0 & 0 & \ldots & 0 \\
\widehat{R_{c} r_{1}} & \widehat{R_{c} r_{2}} & \ldots & \widehat{R_{c} r_{N}}
\end{array}\right]\left[\begin{array}{c}
\tilde{z}_{1} \\
\tilde{z}_{2} \\
\vdots \\
\tilde{z}_{N}
\end{array}\right] \\
& +\left[\begin{array}{cccc}
0 & 0 & \ldots & 0 \\
\widehat{R_{c} r_{1}} \widehat{\Omega}_{c} & \widehat{R_{c} r_{2}} \widehat{\Omega}_{c} & \ldots & \widehat{R_{c} r_{N}} \widehat{\Omega}_{c}
\end{array}\right]\left[\begin{array}{c}
f_{1}^{d} \\
f_{2}^{d} \\
\vdots \\
f_{3}^{d}
\end{array}\right]  \tag{5.43}\\
& +\left[\begin{array}{c}
O \\
k_{\Omega_{c}} I_{c}^{-1} \widehat{\Omega}_{c} I_{c} \Omega_{c}+k_{\Omega_{c}}^{2} I_{c}^{-1} \Omega_{c}
\end{array}\right] \\
& +\left[\begin{array}{c}
O \\
-k_{\Omega_{c}} I_{c}^{-1} k_{q_{v}} q_{v}^{e}-k_{q_{v}} \dot{q}_{v}^{e}+k_{q_{v}} \widehat{\Omega}_{c} q_{v}^{e}
\end{array}\right],
\end{align*}
$$

which can be written as

$$
\begin{equation*}
\Lambda \dot{f^{d}}=g\left(\Omega_{c}, f^{d}, \tilde{z}, q_{v}^{e}\right)=\bar{g}\left(\Omega_{c}, \tilde{z}, q_{v}^{e}\right) \tag{5.44}
\end{equation*}
$$

since $f^{d}$ in (5.21) is only a function of $\Omega_{c}$ and $q_{\nu}^{e}$. Then $\dot{f}^{d}$ can be obtained as

$$
\begin{equation*}
\dot{f}^{d}=\Lambda^{+} \bar{g}\left(\Omega_{c}, \tilde{z}, q_{v}^{e}\right)+\dot{f}_{N(\Lambda)}^{d} \tag{5.45}
\end{equation*}
$$

where $\Lambda^{+}$is the Moore-Penrose generalized inverse of $\Lambda$ and $\dot{f}_{N(\Lambda)}^{d}$ belongs to the null space of $\Lambda$. Because $\dot{f}^{d}=k \dot{z}^{d}$, we obtain

$$
\begin{equation*}
\dot{z}^{d}=\frac{1}{k}\left(\Lambda^{+} \bar{g}\left(\Omega_{c}, \tilde{z}, q_{v}^{e}\right)+\dot{f}_{N(\Lambda)}^{d}\right) \tag{5.46}
\end{equation*}
$$

### 5.4.4 Useful identities

$$
\begin{align*}
\hat{\omega}^{T} & =-\hat{\omega}  \tag{5.47}\\
\hat{\omega} \omega & =0  \tag{5.48}\\
\hat{v} \omega & =-\hat{\omega} \nu \Longrightarrow \vec{v} \times \vec{\omega}=-\vec{\omega} \times \vec{v}  \tag{5.49}\\
\dot{R} & =R \hat{\omega}  \tag{5.50}\\
\frac{d}{d t}\|b\| & =\frac{b^{T} \cdot \dot{b}}{\|b\|}  \tag{5.51}\\
\frac{d\left(r_{1} \times r_{2}\right)}{d t} & =\dot{r}_{1} \times r_{2}+r_{1} \times \dot{r}_{2}  \tag{5.52}\\
\text { Product Rule: } \frac{d}{d t}(f \cdot g) & =\dot{f} g+f \dot{g}  \tag{5.53}\\
\text { Quotient Rule: } \frac{d}{d t} \frac{f}{g} & =\frac{\dot{f} g-f \dot{g}}{g^{2}} \tag{5.54}
\end{align*}
$$

### 5.5 Simulation Results

### 5.5.1 Simulation Environment

In this section, we present a numerical example for three agents transporting a load. The control objective is transport the payload such that it will rotate by 90 degrees and
translate with the constant velocity $v^{d}$. The mass and inertia of each quadcopter are $m_{i}=0.75 \mathrm{~kg}$ and $I_{i}=\operatorname{diag}\{0.0820,0.0845,0.1377\} \mathrm{kgm}^{2}$ respectively. The mass and inertia of the payload are $M_{c}=1.5$ and $I_{c}=\operatorname{diag}\{0.0135,0.0135,0.0135\} \mathrm{kgm}^{2} \mathrm{~kg}$ with a radius of 15 cm . For all of the results presented in this section we set $\psi_{i}^{\text {des }}=0$ and $k=1000 \mathrm{~N} / \mathrm{m}$. We set $r_{1}=[0,0.15,0]^{T}, r_{2}=[0.15 \sin (\pi / 3),-0.15 \cos (\pi / 3), 0]^{T}$ and $r_{3}=$ $[-0.15 \sin (\pi / 3),-0.15 \cos (\pi / 3), 0]^{T}$. The $f_{i}^{d}(t)$ 's are chosen online based on (5.21), where $f_{N(\Lambda)}^{d}$ for each agent is given by $R_{c}[0.0,-2.0,0]^{T} \mathrm{~N}, R_{c}[2.0,1.0,0]^{T} \mathrm{~N}$, and $R_{c}[-2.0,1.0,0]^{T}$ N , respectively. As $R_{c} \rightarrow R_{d}$ and $\Omega_{c} \rightarrow 0, f_{N(\Lambda)}^{d}$ approaches a constant vector and $\dot{f}_{N(\Lambda)}^{d}$ approaches zero.

### 5.5.2 Quadcopter Attitude Controller

The gains for the attitude controller are given in Table 5.1.

| $K_{1}^{i}, K_{2}^{i}$ | $K_{3}$ | $K_{p 1}^{i}, K_{p 2}^{i}$ | $K_{p 3}^{i}$ | $K_{d 1}^{i}, K_{d 2}^{i}$ | $K_{d 3}^{i}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 21.93 | 18 | 4.65 | 3.76 | 0.19 | 0.15 |

Table 5.1: Attitude controller gains.

### 5.5.3 Numerical Examples

In this example we illustrate the time-scale separation between the fast dynamics (5.33) and (5.35) and the slow dynamics of $\xi_{i}, \Omega_{c}$ and $q_{\nu}^{e}$. We choose

$$
\begin{align*}
& v^{d}=\left[\begin{array}{lll}
2.0, & 0.0, & 0
\end{array}\right]^{T}, q_{v}^{d}=\left[\begin{array}{ccc}
0, & 0, & \frac{1}{\sqrt{2}}
\end{array}\right]^{T}, \\
& K_{i}=\left[\begin{array}{ccc}
0.2 & 0 & 0 \\
0 & 0.2 & 0 \\
0 & 0 & 0.2
\end{array}\right], \Gamma_{i}=\left[\begin{array}{ccc}
0.01 & 0 & 0 \\
0 & 0.01 & 0 \\
0 & 0 & 0.01
\end{array}\right],  \tag{5.55}\\
& k_{q_{v}}=0.1, k_{\Omega_{c}}=0.2 .
\end{align*}
$$

Note that $q_{v}^{d}$ corresponds to a desired vector quaternion for 90 degrees rotation on the $x-y$ plane. The product $\Gamma_{i} k$ is large compared with other control gains.

As shown in Fig. 5.1, 5.2, 5.3 and 5.4, the convergence of $\tilde{z}_{i}, \xi_{i}, \Omega_{c}$ and the attitude error $q_{v}^{e}$ verifies that there exists a time-scale separation between the fast dynamics (5.33) and (5.35) and the slow dynamics of $\xi_{i}, \Omega_{c}$ and $q_{v}^{e}$.


Figure 5.1: The magnitude of $\tilde{z}_{i}$ for the three agents.

Since singular perturbation and Lyapunov analysis can be conservative, we further tune the gains of the controllers to achieve improved transient performance. We increase $K_{i}, k_{q_{v}}$ and $k_{\Omega_{c}}$ to

$$
K_{i}=\left[\begin{array}{lll}
1 & 0 & 0  \tag{5.56}\\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right], k_{q_{v}}=2, k_{\Omega_{c}}=1
$$

while keeping other parameters the same as in the previous simulation.
Fig. 5.5 shows that the 3D trajectories of the agents and the payload. It can be observed that the payload is able to make a 90 degree turn and maintain the desired orientation throughout the simulation. Fig. 5.6 shows that the contact force is regulated to


Figure 5.2: The magnitude of the velocity error $\xi_{i}$ for the three agents and the payload.


Figure 5.3: The angular velocity of the load $\Omega_{c}$ goes to zero.


Figure 5.4: The magnitude of the orientation error $\left\|q_{\nu}^{e}\right\|$ for the payload.
the time varying setpoints. Fig. 5.7 shows that the velocity errors of agents and the payload converge to zero. Fig. 5.8 shows that the angular velocity of the payload converges to zero and Fig. 5.9 shows that $\left\|q_{v}^{e}\right\| \rightarrow 0$, which means that $R_{c}$ converges to $R_{c}^{d}$. Fig. 5.10 and Fig. 5.11 show that the attitude controller tracks the desired attitude angles $\phi_{i}^{\text {des }}$ and $\theta_{i}^{\text {des }}$ generated from the force controller (3.6) and (3.7). Note that the convergence speed for all the variables are improved due to the increased gain parameters in (5.56).

Time: 20.00 sec


Figure 5.5: 3D position trajectory for three agents and the payload at different time steps for the first 20 seconds. The changes of the position of the red and the light blue dots confirm that the payload is able to achieve the desired rotation.


Figure 5.6: The magnitude of $\tilde{z}_{i}$ goes to zero for all three agents.


Figure 5.7: The magnitude of the velocity error $\xi_{i}$ for the three agents and the payload.


Figure 5.8: The angular velocity of the load $\Omega_{c}$ goes to zero.


Figure 5.9: The magnitude of the orientation error $\left\|q_{\nu}^{e}\right\|$ for the payload.


Figure 5.10: Desired and actual roll angles $(\phi)$ for three agents.


Figure 5.11: Desired and actual pitch angles $(\theta)$ for three agents.

Because of the complexity in computation of $\dot{f}_{i}^{d}$, we have also experimented with $\dot{f}_{i}^{d}$ set to 0 in (5.19). Our simulation results show that $\xi_{i} \rightarrow 0, \xi_{c} \rightarrow 0, \Omega_{c} \rightarrow 0, q_{v}^{e} \rightarrow 0$ and $\tilde{z}_{i} \rightarrow 0$ are still achieved. The simulation results are not included in this report and we plan to investigate more in our future work.

## CHAPTER 6

## FORCE CONTROL OF PAYLOAD WITH UNKNOWN MASS USING QUADCOPTERS WITH MULTI-LINK ROBOT

In this chapter, we tested our force control algorithms from Chapter 4 for a group of aerial manipulators (AM) transporting a flexible load. Each AM is a combination of an Unmanned Aerial Vehicle (UAV) with a two-degree-of-freedom robotic manipulator (RM) attached to it. This work is similar to the past work in [36] and we extend it to address cooperative transportation of a payload with an unknown mass. Specifically, we develop decentralized control and estimation algorithms that guarantees force regulation, velocity convergence and convergence of the estimate of the mass of the payload without explicit communication between the AMs. The algorithms are implemented at the kinematic level using the inverse kinematics of the AMs.

### 6.1 Payload Transportation with an Aerial Manipulator

### 6.1.1 Kinematics

Each aerial manipulator is a 2-DOF robotic arm mounted at the bottom of a quadrotor. Similar derivations for single aerial manipulator can be found in [43, 44]. Position of the center of mass of the quadrotor in the inertial frame $\Sigma_{I}$ is given by, $p_{c, i}^{I}=$ $\left[\begin{array}{lll}x_{c, i} & y_{c, i} & z_{c, i}\end{array}\right]^{T}$, orientation of the quadrotor is described by the triple of ZYX (yaw-pitch-roll) Euler angles, $\Phi_{i}=\left[\begin{array}{lll}\psi_{i} & \theta_{i} & \phi_{i}\end{array}\right]^{T}$ and joint angles of the two-DOF manipulator are given by, $\eta_{i j}=\left[\begin{array}{ll}\eta_{i 1} & \eta_{i 2}\end{array}\right]^{T}$ (see Fig 2.1). Both of the joint angles are defined about positive $x_{c, i}$ axis. The origin of frame for the load and each link of the manipula-


Figure 6.1: Two aerial manipulators transport a common flexible load. Note that body frame $O_{c, i}$ coincides with $O_{1, j}$ (not shown in the figure) and $a_{i}$ is the initial position of agent $i$.
tor is placed at its center of mass making the axis coincident with the inertial axes. The vector containing all the generalized coordinates for agent $i$ is given by

$$
q_{i}=\left[\begin{array}{lll}
\left(p_{c, i}^{I}\right)^{T} & \Phi_{i}^{T} & \eta_{i}^{T} \tag{6.1}
\end{array}\right]^{T}
$$

Let $\omega_{c, i}^{I}$ and $\omega_{i, c}^{b}$ be the angular velocity of the quadrotor in the inertial frame and the body-fixed frame respectively. We can map the time derivative of Euler angles $\dot{\Phi}_{i}$ to $\omega_{c, i}^{I}$ by a transformation matrix $T_{i}$. The following equalities hold:

$$
\begin{align*}
& \omega_{c, i}^{I}=R_{c, i}^{I} \omega_{c, i}^{b}  \tag{6.2}\\
& \omega_{c, i}^{I}=T_{i} \dot{\Phi}_{i}  \tag{6.3}\\
& \omega_{c, i}^{b}=\left(R_{c, i}^{I}\right)^{T} T_{i} \dot{\Phi}_{i}=Q_{i} \dot{\Phi}_{i} \tag{6.4}
\end{align*}
$$

where $R_{c, i}^{I} \in S O(3)$ is the rotation matrix between the body frame $O_{c, i}$ and the inertial frame $\Sigma_{I}$.

Let $p_{i, j}^{b}$ be the position of center of mass of link $j=1,2$ of the $i$-th aerial manipulator
in the body-fixed frame $O_{c, i}$. The position $p_{i j}^{I}$ of center of mass of link $j$ is given by

$$
\begin{equation*}
p_{i, j}^{I}=p_{c, i}^{I}+R_{c, i}^{I} p_{i, j}^{b} \tag{6.5}
\end{equation*}
$$

The linear and angular velocity of each link of manipulator are related to the time derivatives of joint angles $\eta_{i, j}$ by

$$
\begin{align*}
& \dot{p}_{i, j}^{b}=J_{v, i j} \dot{\eta}_{i, j}  \tag{6.6}\\
& \omega_{i, j}^{b}=J_{\omega, i j} \dot{\eta}_{i, j} \tag{6.7}
\end{align*}
$$

where $J_{\nu, i j} \in \mathbb{R}^{3 \times 2}$ and $J_{\omega, i j} \in \mathbb{R}^{3 \times 2}$ are the linear and angular Jacobian of the manipulator in the body-fixed frame [45]. We can compute the linear and angular velocity of each link in the inertial frame from the following equations,

$$
\begin{align*}
& \dot{p}_{i, j}^{I}=\dot{p}_{c, i}^{I}+\dot{R}_{c, i}^{I} p_{i, j}^{b}+R_{c, i}^{I} \dot{p}_{i, j}^{b}  \tag{6.8}\\
& \dot{p}_{i, j}^{I}=\dot{p}_{c, i}^{I}+\widehat{\omega}_{c, i}^{I} R_{c, i}^{I} p_{i, j}^{b}+R_{c, i}^{I} J_{v, i} \dot{\eta}_{i, j}  \tag{6.9}\\
& \omega_{i, j}^{I}=\omega_{c, i}^{I}+R_{c, i}^{I} J_{\omega, i} \dot{\eta}_{i, j} \tag{6.10}
\end{align*}
$$

where the operator^converts a given vector $\omega=\left[\begin{array}{lll}\omega_{1} & \omega_{2} & \omega_{3}\end{array}\right]^{T}$, into a skew symmetric matrix as:

$$
\widehat{\omega}=\left[\begin{array}{ccc}
0 & -\omega_{3} & \omega_{2}  \tag{6.11}\\
\omega_{3} & 0 & -\omega_{1} \\
-\omega_{2} & \omega_{1} & 0
\end{array}\right]
$$

Let $x_{i}^{b}$ be the position of the end-effector in the body-fixed frame $O_{c, i}$, which is given by

$$
\begin{equation*}
x_{i}^{b}=P_{c 1}^{C} R_{c 1}^{C} P_{12}+R_{c 2}^{C} P_{2 T} \tag{6.12}
\end{equation*}
$$

where $P_{c 1}^{C}$ is the vector representing the initial distance between the body frame of UAV and the body frame of the 1 st link of the manipulator expressed in frame C. $R_{c 1}^{C}$, and $R_{c 2}^{C}$ are the rotation matrices between the frames of the manipulator and its base frame. $P_{12}$ is the initial distance between the first and the second link, $P_{2 T}$ is the initial distance between the second link and the end-effector both expressed in its body-frame.

The position of the end effector $x_{i}^{I}$ in the inertial frame for the $i$-th aerial manipulator is given by,

$$
\begin{equation*}
x_{i}^{I}=p_{c, i}^{I}+R_{c, i}^{I} x_{i}^{b} \tag{6.13}
\end{equation*}
$$

The linear and angular velocities of the end effector can be computed as:

$$
\begin{align*}
& \dot{x}_{i}^{I}=\dot{p}_{c, i}^{I}+\dot{R}_{c, i}^{I} x_{i}^{b}+R_{c, i}^{I} \dot{x}_{i}^{b}  \tag{6.14}\\
& \dot{x}_{i}^{I}=\dot{p}_{c, i}^{I}+\widehat{\omega}_{c, i}^{I} R_{c, i}^{I} x_{i}^{b}+R_{c, i}^{I} J_{v, i j} \dot{\eta}_{i, j}  \tag{6.15}\\
& \omega_{i}^{I}=\omega_{c, i}^{I}+R_{c, i}^{I} J_{\omega, i j} \dot{\eta}_{i, j} \tag{6.16}
\end{align*}
$$

Equations (6.13), (6.15) and (6.16) can be simplified as:

$$
\begin{align*}
\dot{p}_{c, i}^{I} & =\left[\begin{array}{lll}
I_{3 \times 3} & 0_{3 \times 3} & 0_{3 \times 2}
\end{array}\right] \dot{q}_{i}=: J_{t, b} \dot{q}_{i},  \tag{6.17}\\
\omega_{c, i}^{I} & =\left[\begin{array}{lll}
0_{3 \times 3} & T_{i} & 0_{3 \times 2}
\end{array}\right] \dot{q}_{i}=: J_{r, b} \dot{q}_{i},  \tag{6.18}\\
\dot{x}_{i}^{I} & =\left[\begin{array}{lll}
I_{3 \times 3} & \left.-\widehat{\left(R_{c, i}^{I} x_{i}^{b}\right.}\right) T & R_{c, i}^{I} J_{v, i j}
\end{array}\right] \dot{q}_{i}=: J_{t, i} \dot{q}_{i},  \tag{6.19}\\
\omega_{i}^{I} & =\left[\begin{array}{lll}
0_{3 \times 3} & T_{i} & R_{c, i}^{I} J_{w, i j}
\end{array}\right] \dot{q}_{i}=: J_{r, i} \dot{q}_{i} . \tag{6.20}
\end{align*}
$$

Define $\nu_{i}=\left[\begin{array}{llll}\left(\dot{p}_{c, i}^{I}\right)^{T} & \left(\omega_{c, i}^{I}\right)^{T} & \left(\dot{x}_{i}^{I}\right)^{T} & \left(\omega_{i}^{I}\right)^{T}\end{array}\right]^{T}$, which contains the linear and angular velocities of the UAV and the end-effector of the $i$-th aerial manipulator. The time derivative of generalized joints vector $q_{i}$ can be mapped to $v_{i}$ by the following equation,

$$
\begin{equation*}
v_{i}=J_{i} \dot{q}_{i} \tag{6.21}
\end{equation*}
$$

where $J_{i} \in \mathbb{R}^{12 \times q_{i}}$ is the geometric Jacobian of the system ([45]) whose elements is computed as:

$$
J_{i}=\left[\begin{array}{llll}
J_{t, b} & J_{r, b} & J_{t, i} & J_{r, i} \tag{6.22}
\end{array}\right]^{T}
$$

We can represent the position/orientation of the end-effector as the minimal representation ( [45]) and express the forward kinematics as:

$$
\begin{equation*}
x_{i}=k\left(q_{i}\right) \tag{6.23}
\end{equation*}
$$

where $k(\cdot)$ is an $m \times 1$ vector function based on the system configuration and $x_{i}$ is an $m \times 1$ state vector representing the system in minimal configuration. The time derivative of $x_{i}$ and $q_{i}$ can be related by,

$$
\begin{equation*}
\dot{x}_{i}=J_{a, i} \dot{q}_{i} \tag{6.24}
\end{equation*}
$$

where $J_{a, i} \in \mathbb{R}^{m \times n_{q_{i}}}$ is the analytical Jacobian of the system which can be derived via differentiation of $k(\cdot)$ ([45]). In this paper, we use the Cartesian position of the endeffector as a minimal representation, which gives us $J_{a, i}=J_{t, i}$.

### 6.1.2 Dynamics

We derive the dynamic model based on Euler-Lagrange formulation.

$$
\begin{gather*}
\frac{d}{d t} \frac{\partial \mathscr{L}}{\partial \dot{q}_{i}}-\frac{\partial \mathscr{L}}{\partial q}=\tau+\tau_{\text {ext }}  \tag{6.25}\\
\mathscr{L}=\mathscr{K}-\mathscr{U} \tag{6.26}
\end{gather*}
$$

where $\mathscr{L}$ is the Lagrangian of the system. $\mathscr{K}$ and $\mathscr{U}$ are the total kinetic and potential energy of the $i$-th aerial manipulator, $\tau$ represents generalized forces and $\tau_{\text {ext }}$ refers to any external disturbance applied to the system.

The total kinetic energy for the $i$-th aerial manipulator is given by the sum of individual contributions of the aerial vehicle and two links of the manipulator.

$$
\begin{equation*}
\mathscr{K}_{i}=\mathscr{K}_{b, i}+\sum_{i}^{2} \mathscr{K}_{i j} \tag{6.27}
\end{equation*}
$$

where $\mathcal{K}_{b, i}$ is the kinetic energy of the quadrotor and $\mathbb{K}_{i j}$ is the kinetic energy of the link $j$ of the $i$-th aerial manipulator.

The kinetic energy for the quadrotor $\mathcal{K}_{b, i}$ is given by

$$
\begin{equation*}
\mathscr{K}_{b, i}=\frac{1}{2} \dot{p}_{b, i}^{T} m_{b, i} \dot{p}_{b, i}+\frac{1}{2} \dot{\Phi}_{i}^{T} T^{T} R_{b, i} I_{b, i} R_{b, i}^{T} T \dot{\Phi}_{i} \tag{6.28}
\end{equation*}
$$

where $m_{b, i}$ and $I_{b, i}$ is the mass and inertia matrix of the quadrotor respectively of the $i$-th aerial manipulator. Kinetic energy of each link can be expressed as follows

$$
\begin{equation*}
\mathscr{K}_{i j}=\frac{1}{2} \dot{p}_{i j}^{T} m_{i j} \dot{p}_{i j}+\frac{1}{2} w_{i j}^{T}\left(R_{b, i} R_{i j}^{b}\right) I_{i j}\left(R_{b, i} R_{i j}^{b}\right)^{T} w_{i j} \tag{6.29}
\end{equation*}
$$

where $R_{i j}^{b}$ is the rotation matrix between bodyframe of center of mass of $j$-th link and $O_{b i}$, while $m_{i j}$ and $I_{i j}$ is the mass and constant inertia matrix of link $j$ respectively.

By considering (6.4), (6.8), (6.9), (6.10), (6.28) and (6.29) total kinetic energy of $i$-th aerial manipulator (6.27) can be expressed into following form

$$
\begin{equation*}
\mathscr{K}_{i}=\frac{1}{2} \dot{q}_{i}^{T} M_{i}\left(q_{i}\right) \dot{q}_{i} \tag{6.30}
\end{equation*}
$$

where $M_{i} \in \mathbb{R}^{8 \times 8}$ is a symmetric and positive definite inertia matrix of aerial manipulator $i$, which can be defined as a $3 \times 3$ block matrix with individual elements as follows:

$$
\begin{align*}
& M_{i_{11}}=\left(m_{b, i}+\sum_{i}^{2} m_{i j}\right) I_{3 \times 3}  \tag{6.31}\\
& M_{i_{22}}= Q_{i}^{T} I_{b, i} Q_{i}+\sum_{i}^{2}\left(T_{i}^{T}\left(\widehat{R_{b, i} p_{i j}^{b}}\right)^{T} m_{i j}\left(\widehat{R_{b, i} p_{i j}^{b}}\right) T_{i}+Q_{i}^{T} R_{i j}^{b} I_{i j}\left(R_{i j}^{b}\right)^{T} Q_{i}\right)  \tag{6.32}\\
& M_{i_{33}}=\sum_{i}^{2}\left(J_{\nu, i}^{T} m_{i j} J_{v, i}+J_{w, i}^{T} R_{i j}^{b} I_{i j}\left(R_{i j}^{b}\right)^{T} J_{w, i}\right)  \tag{6.33}\\
& M_{i_{12}}=M_{i_{21}}^{T}=-\sum_{i}^{2}\left(m_{i j}\left(\widehat{R_{b, i} p_{i, j}^{b}}\right) T_{i}\right)  \tag{6.34}\\
& M_{i_{13}}=M_{i_{13}}^{T}=\sum_{i}^{2}\left(m_{i j} R_{b, i} J_{v, i}\right)  \tag{6.35}\\
& M_{i_{23}}=M_{i_{23}}^{T}=\sum_{i}^{2}\left(Q_{i}^{T} R_{i j}^{b} I_{i j}\left(R_{i j}^{b}\right)^{T} J_{w, i}-m_{i j} T_{i}^{T}\left(\widehat{R_{b, i} p_{i j}^{b}}\right)^{T} R_{b, i} J_{v, i}\right) \tag{6.3}
\end{align*}
$$

The total potential energy for the $i$-th aerial manipulator is given by the sum of individual contributions of the aerial vehicle and two links of the manipulator.

$$
\begin{equation*}
\mathscr{U}_{i}=\mathscr{U}_{b, i}+\sum_{i}^{2} \mathscr{U}_{i j} \tag{6.3}
\end{equation*}
$$

The potential energy for the quadrotor $\mathscr{U}_{b, i}$ is given by

$$
\begin{equation*}
\mathscr{U}_{b, i}=m_{b, i} g e_{3}^{T} p_{b, i} \tag{6.38}
\end{equation*}
$$

where $g$ is acceleration due to gravity and $e_{3}$ is the unit vector $\left[\begin{array}{lll}0 & 0 & 1\end{array}\right]^{T}$. Potential energy of each links is given by

$$
\begin{equation*}
\mathscr{U}_{i j}=m_{i j} g e_{3}^{T}\left(p_{b, i}+R_{b, i} p_{i j}^{b}\right) \tag{6.39}
\end{equation*}
$$

The total potential energy for $i$-th aerial manipulator can be computed as

$$
\begin{equation*}
\mathscr{U}_{i}=m_{b, i} g e_{3}^{T} p_{b, i}+\sum_{i}^{2}\left[m_{i j} g e_{3}^{T}\left(p_{b, i}+R_{b, i} p_{i j}^{b}\right)\right] \tag{6.40}
\end{equation*}
$$

By substituting equations (6.26), (6.30), (6.40), into (6.25), the full dynamics of $i$-th aerial manipulator can be computed as

$$
\begin{equation*}
M_{i}\left(q_{i}\right) \ddot{q}_{i}+C_{i}\left(q_{i}, \dot{q}_{i}\right) \dot{q}_{i}+G_{i}\left(q_{i}\right)=\tau_{i}+\tau_{e x t, i} \tag{6.41}
\end{equation*}
$$

where $C_{i}\left(q_{i}, \dot{q}_{i}\right)$ is the Coriolis matrix and $G_{i}\left(q_{i}\right)$ is the vector containing gravity effects. The generic element of the Coriolis matrix $C_{i}\left(q_{i}, \dot{q}_{i}\right)$ is given by [45, 46],

$$
\begin{equation*}
c_{i j}=\sum_{1}^{8} c_{i j k} \dot{q}_{k} \tag{6.42}
\end{equation*}
$$

where the coefficients $c_{i j k}$ is given by

$$
\begin{equation*}
c_{i j k}=\frac{1}{2}\left(\frac{\partial m_{i j}}{\partial q_{k}}+\frac{\partial m_{i k}}{\partial q_{j}}-\frac{\partial m_{j k}}{\partial q_{i}}\right) \tag{6.43}
\end{equation*}
$$

which are also termed as Christoffel symbols of the first type. $G_{i}\left(q_{i}\right)$ is computed as

$$
\begin{equation*}
G_{i}\left(q_{i}\right)=\frac{\partial \mathscr{U}_{i}}{\partial q_{i}} \tag{6.44}
\end{equation*}
$$

If we neglect the aerodynamic effects and consider low speed displacements [47] for the quadrotor model, the input $\tau$ mapped to actuation command as:

$$
\begin{equation*}
\tau_{i}=R_{f, i} N_{i} f_{i}=\Xi_{i} f_{i} \tag{6.45}
\end{equation*}
$$

where $f_{i}=\left[\begin{array}{ll}f_{q, i}^{T} & \tau_{i}^{T}\end{array}\right]^{T}$ includes $f_{q, i}=\left[\begin{array}{llll}f_{q, 1} & f_{q, 2} & f_{q, 3} & f_{q, 4}\end{array}\right]^{T}$ as input forces actuated by the quadrotor motors and $\tau_{i}=\left[\begin{array}{cc}\tau_{i, 1} & \tau_{i, 2}\end{array}\right]^{T}$ for manipulator joint torques. $R_{f, i}=$ $\operatorname{diag}\left(R_{b, i}, Q^{T}, I_{n_{i} \times n_{i}}\right) \in \mathbb{R}^{6+n_{i} \times 6+n_{i}}$ and $N_{i}=\operatorname{diag}\left(\Omega_{i}, I_{n_{i} \times n_{i}}\right) \in \mathbb{R}^{8 \times 4+n_{i}}$, in which

$$
\Omega=\left[\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
1 & 1 & 1 & 1 \\
0 & d & 0 & -d \\
-d & 0 & d & 0 \\
c & -c & c & -c
\end{array}\right]
$$

where $d$ is the distance from a motor to the center of the vehicle and $c=k_{m} / k_{f}>0$, $k_{m}$ the drag coefficient and $k_{f}$ the motor thrust coefficient which are obtained from the motor test data using a 6 -axis force-torque sensor [17].

### 6.2 Control Objective

The control objective is to design $F_{i}$ in a decentralized way such that all the agents and the payload converge to a constant velocity $v^{d}$ and the contact force $f_{i}$ is regulated to a setpoint $f_{i}^{d}, i=1, \cdots, N$.

### 6.3 Payload Transport with an Unknown Mass

Rewrite the update law from Section 4.1

$$
\begin{equation*}
\dot{\widehat{M}}_{c, i}=\gamma_{i}\left(\dot{x}_{i}-v^{d}\right)^{T} e_{3} \tag{6.46}
\end{equation*}
$$

We modify the design in (2.9) as:

$$
\begin{equation*}
F_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+\hat{f}_{i}^{d} \tag{6.47}
\end{equation*}
$$

and estimate the $z$-component of $f_{i}^{d}$ as:

$$
\begin{equation*}
\hat{f}_{i, z}^{d}=\widehat{M}_{c, i} g \tag{6.48}
\end{equation*}
$$

The $x$ and $y$ components of $\hat{f}_{i}^{d}$ are the steady state values of the squeeze forces that are pre-designed to satisfy the following constraints

$$
\begin{equation*}
\sum_{i=1}^{N} \hat{f}_{i, x}^{d}=0 \quad \text { and } \quad \sum_{i=1}^{N} \hat{f}_{i, y}^{d}=0 \tag{6.49}
\end{equation*}
$$

The dynamics (2.8) with the proposed control law (6.46), (6.47), takes the following form:

$$
\begin{equation*}
m_{i} \ddot{x}_{i}=-\Gamma_{i}\left(\dot{x}_{i}-v^{d}\right)+\hat{f}_{i}^{d}-f_{i} . \tag{6.50}
\end{equation*}
$$

Theorem 6.1 Consider the decentralized control law in (6.47) together with (6.46) and
(6.48). The equilibrium of (2.7), (6.46), and (3.11) given by

$$
\begin{gather*}
\mathscr{E}=\left\{\left(\dot{x}_{i}, f_{i}, v^{d}, \widehat{M}_{c, i}\right) \mid \dot{x}_{i}=v^{d}, \quad \dot{x}_{c}=v^{d},\right. \\
\left.\sum_{i=1}^{N} \widehat{M}_{c, i}=M_{c}, \quad \text { and } \quad f_{i}=f_{i}^{d}\right\}, \tag{6.51}
\end{gather*}
$$

is asymptotically stable.

Proof. The proof of the closed loop system is exactly same as in the one presented in Section 4.1.

### 6.4 Simulation Results

### 6.4.1 Simulation Environment

In this section, we present simulation results for two AMs transporting a load. Each quadrotor has a mass of 5 kg . Each link of the robotic manipulator has a length of 20 cm and a mass of 0.5 kg . The mass of the payload is 0.5 kg with a radius of 15 cm . All the joints of the manipulator are assumed to be spherical. The first joint rotates about the $x$ axis and the second joint rotates about the $y$ axis as shown in Fig. 2.1. We select $x_{c, i}, y_{c, i}$ and $\phi_{i}$ as state variables to be controlled for all the results presented in this section. We use a linear spring-force model to compute the contact force $f_{i}$ in this simulation.

As a preliminary investigation, we implement the control in Chapter 4 at the kinematic level. This is equivalent to assuming that there exists a sufficiently-fast velocity tracking controller for each AM. To do so, we set right-hand side of (3.11) to 0 and solve for $\dot{x}_{i}$ :

$$
\begin{equation*}
\dot{x}_{i}=K_{i}\left(\hat{f}_{i}^{d}-f_{i}\right)+v^{d}, \quad i=1, \ldots, N, \tag{6.52}
\end{equation*}
$$

where, $K_{i}>0$ and $x_{i}$ is the position of the end-effector of $i$-th the aerial manipulator.
The end-effector velocity $\dot{x}_{i}$ is computed from (6.52) and transformed to the generalized velocities using (6.19). Because of the redundancy in the system, we use the partial-inverse of the Jacobian matrix to calculate the joint velocities for 3 desired states.

An example to calculate the joint velocities for 3 states, $x_{c, i}, y_{c, i}$ and $\phi_{i}$ from the partialinverse of the Jacobian is shown below. Let

$$
\begin{equation*}
J_{(t, i) s}=\left[J_{t, i}(:, 1), \quad J_{t, i}(:, 2), \quad J_{t, i}(:, 6)\right] \tag{6.53}
\end{equation*}
$$

be the partial Jacobian matrix for the $x_{c, i}, y_{c, i}$ and $\phi_{i}$ states. Then the corresponding joint velocities are computed as

$$
\begin{align*}
\dot{q}_{(i) p r t l} & =\left(J_{(t, i) s}\right)^{-1} \dot{x}_{i}, \\
\dot{q}_{i} & =\left[\begin{array}{lllllll}
\dot{q}_{(i) p r t l}(1), & \dot{q}_{(i) p r t l}(2), & 0, & 0, & 0, & \dot{q}_{(i) p r t l}(3), & 0,
\end{array} 0^{T} .\right. \tag{6.54}
\end{align*}
$$

Note that $k$ is not used in the controller (6.52).

### 6.4.2 Unknown Payload with a Predesigned $v^{d}$

We set $v^{d}=\left[\begin{array}{lll}0.2 & 0 & 0\end{array}\right]^{T}, \widehat{M}_{c, i}=0.1 \mathrm{~kg}, K_{i}=0.1, k=100 \mathrm{~N} / \mathrm{m}, f_{1}^{d}=\left[\begin{array}{lll}0 & 2.0 & 0.5 \hat{M}_{c} g\end{array}\right]^{T}$ and $f_{2}^{d}=\left[\begin{array}{lll}0 & -2.0 & 0.5 \hat{M}_{c} g\end{array}\right]^{T}$. The initial conditions of the two AMs are set as:

$$
\begin{gathered}
q_{1}=\left[\begin{array}{llllllll}
0.1 & -0.35 & 0.7 & 0 & 0 & -0.1 & 0 & -\pi / 2
\end{array}\right]^{T}, \\
q_{2}=\left[\begin{array}{llllllll}
0 & 0.35 & 0.68 & 0 & 0 & 0.1 & 0 & \pi / 2
\end{array}\right]^{T}
\end{gathered}
$$

We observe from Fig. 6.2 and 6.3 that the velocities of the agents and the payload converge to $v^{d}$. The update law from (6.46) successfully recovers the actual mass of the payload as seen in Fig. 6.4. We observe from Fig. 6.5 that the squeeze force is regulated to the setpoint $f_{i}^{d}$.


Figure 6.2: Linear velocities for both agents. The $x$ component of the velocity converges to $0.2 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 6.3: Linear velocity for the payload. The $x$ component of the velocity converges to $0.2 \mathrm{~m} / \mathrm{s}$ and the rest converge to zero.


Figure 6.4: Estimation of the unknown mass. The sum of the individual estimates converges to the actual mass of the load.


Figure 6.5: Contact forces for both AMs in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload.

### 6.4.3 Unknown Payload with a Time Varying $v^{d}$

In this case study, we test the robustness of our controllers to time varying velocities. We keep the other parameters to be the same as in the previous case and change the desired velocity of the agents to

$$
v^{d}(t)= \begin{cases}{\left[\begin{array}{lll}
2.0, & 0, & 0
\end{array}\right]^{T},} & \mathrm{t} \leq 2 \mathrm{~s} \\
{[2-2 \sin (t-2),} & 2 \sin (t-2), \\
{\left[\begin{array}{ll}
2
\end{array}\right]^{T},} & 2<\mathrm{t} \leq 4 \mathrm{~s} \\
{\left[\begin{array}{lll}
0, & 2.0, & 0
\end{array}\right]^{T},} & \mathrm{t}>4.0 \mathrm{~s}\end{cases}
$$

Fig. 6.6 and 6.7 show the trajectories of the agents and the payload in 2D and 3D, respectively. We also observe from Fig. 6.8 and 6.9 that the velocities of the agents and the payload in both $x$ and $y$ directions closely track $v^{d}$. The contact force is also regulated close to the setpoint $f_{i}^{d}$ as shown in Fig. 6.11.


Figure 6.6: X and Y position trajectories for both agents (time varying $v^{d}$ ). Both agents travel in the $x$-direction until 2 seconds and make a smooth turn towards the $y$-direction after 2 seconds.


Figure 6.7: 3D position trajectory for both agents and the payload at different time interval. Note that all 3 agents stay at an equal distance and will never collide. They approach the $y$-axis by being aligned in a straight line as shown in Fig. 6.6. Notice that if the quadrotors are at the same height, they must keep a gap and the manipulator must be large enough to avoid their blades clashing.


Figure 6.8: Velocities of the 2 agents in all 3 directions. The $x$ and $y$ velocities are smooth during transition and the $z$ velocity is zero.


Figure 6.9: Desired and actual velocities of the load in all 3 directions.


Figure 6.10: Estimation of the unknown mass. The sum of individual estimate converges to the actual mass of the load, which is 0.5 kg .


Figure 6.11: Contact forces of both agents in all 3 direction. The forces in the $x$ direction converge to zero. In the $y$-direction, they converge to 2.0 N and -2.0 N , respectively. In the $z$ direction, the sum of the force is equal to the weight of the payload.

## CHAPTER 7

## EXPERIMENTAL RESULTS

In this chapter we present our preliminary experimental results.

### 7.1 Mechanical Design of the Test Bed

There are plenty of quadcopters and robotic arm available. However due to expensive SDK and propriety controllers, we develop our own test bed from custom made quadcopter and 3D printed light weight robotic arm.

### 7.1.1 Quadcopter



Figure 7.1: Custom made quadcopter in the lab.

### 7.1.2 Robotic Manipulator Design

We decided to use [48] file and significanty modified it to fit our needs for the quadcopter. This design is simple and can be printed using 3D printers. The motors and
other parts are also inexpensive. The robotic arm is composed a two links which can


Figure 7.2: CAD model of the tow link robotic arm.
be actuated using high torque servo motors that can be mounted on the bottom of the arm. Another light weight servo motor is mounted close the gripper to enable dynamic gripping.

### 7.2 Computation Hardware

### 7.2.1 Microprocessor

We use an Arduino Nano for the processing. The Arduino platform excels at basic input/output operations, such as inputting sensor information and outputting servo commands. It is also has enough computational power to enable the implementation of low-level algorithms. The Arduino is only used as a low-level controller. It inputs user command via serial input, processes it with various sensory information, and outputs the desired angular position of each servo.

### 7.2.2 Communications

The microprocessor communicates with high-level controllers via serial input. Upon detecting a new character in the serial buffer, the program enters a while() loop to con-
tinuously add incoming characters to char input[]. Once the end of the stream is detected, or a newline character is read, the program proceeds to process input[]. Each character is checked against a menu of options in a switch-case statement. If the character is found, the statement either the remaining characters of input[] are parsed to integers and return for further processing.

### 7.2.3 Motors

We used two Hitec HS-77BB servo motor in the linkage. The Hitec HS-77BB is a 3-pole brushed servo. At only 35.0 grams and dimensions of $44.0 \mathrm{~mm} \times 23.0 \mathrm{~mm} x$ 25.0 mm , the HS-77BB operates at $4.8-6.0 \mathrm{~V}$ DC and has a maximum torque range of 4.4 âĂŞ $5.5 \mathrm{~kg} / \mathrm{cm}$ [49]. When compared with servos of similar volume and mass (such as the SPMSA6050), the HS-77BB has above average torque.

### 7.3 Initial Integration and test



Figure 7.3: Custom made aerial manipulator.

We also tested the viability of our system by manfully flying the aerial manipulator in
the lab. Currently we are exploring the aerial gripping with a single AM and eventually plan to extend this to multiple AMs.

(a)
(b)

(c)

Figure 7.4: Different instances of initial aerial manipulation test.

## CHAPTER 8

## CONCLUSION AND FUTURE WORK

In this thesis, we considered multiple quadcopters transporting a payload. We developed a force and orientation control algorithm to transport the payload. We developed adaptive decentralized control law for transporting the payload of unknown mass without explicit communication between the agents. Our controller ensures that all quadcopters asymptotically converge to a constant reference velocity. It also ensures that all of the forces applied to the payload converges to desired set-points. Desired thrusts and attitude angles are computed from the control algorithms and a low-level PD controller is implemented to track the desired commands for each quadcopter. The sum of the estimates of the unknown mass from all the agents converges to the true mass.

We also developed a consensus algorithm based on connected graphs to ensure that each agent gets an equal share of the payload mass. Furthermore, we developed orientation control algorithm that guarantees attitude stabilization of the payload. In particular, we developed time varying force setpoints to enforce attitude regulation.

Future work will involve relaxation of the required information in the algorithm and experimental validation of the proposed design. We plan to design and implement robust controllers for disturbance rejection in the presence of wind gusts. Furthermore, we plan to develop a complete aerial manipulation system in the lab.

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