BEAM CONSTANTS BY HIGH SPEED COMPUTER

By

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PREFACE

The analysis of continuous structures of variable moment of inertia by moment distribution or slope deflection equations is a laborious process. The carry-over moment procedure (1) has certain advantages over these two methods in that no solutions of simultaneous equations are required and no distribution of moments is necessary. Analysis by this method, however, requires the evaluation of certain constants that are not readily available in existing tables.

In this thesis mathematical expressions and a computer program are developed to evaluate constants for beams with parabolic haunches.

The author wishes to express his indebtedness to Professor J. J. Tuma for his valuable guidance and assistance in the preparation of this thesis and for acting as the writer's major adviser. Acknowledgment is also due Professor William Granet for making the facilities of the Oklahoma State University Computing Center available; the staff of the School of Civil Engineering for their aid and instruction during the writer's years of graduate study.

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NOMENCLATURE

a,a Length Coefficients
b Width of Beam
f Angular Flexibility Coefficient
g
hoose and the main set of the seam of the
h Depth of Haunch
h_{x}, h_{z} Depth of Beam at x and z, Respectively
n Moving Load Position Coefficient
p Maximum Intensity of Haunch Load
q Specific Weight of Material
t_{AB}, t_{BA} Angular Function Coefficients
t_x, t_z Ratio of $\frac{x}{L\alpha}$ and $\frac{z}{L\beta}$, Respectively
W Intensity of Uniform Load
x, x^1, z, z^1 Coordinates of the Cross-Section
A,B,C,D,E Letters Designating Cross-Sections of Beam
(DL) Due to Dead Load
E Modulus of Elasticity
F_{AB}, F_{BA} Angular Flexibilities
G,G,G, G,
(HL)
I _O °°°°°°° Inertia
I,I,
L Length of Span

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(LL) Due to Live Load
P
Q_{k}, R_{k} Integral Functions
T Parabolic Function of Haunch
(UL) Due to Uniformly Distributed Load
X_{k} Sum of Integral Functions For Beam
X Sum of Integral Functions to n kn
$X_{kq}, X_{kq}^{i}, X_{kq}^{i}$ Integral Functions for Parabolic Haunch
$X_{kr}, X_{kr}^{i}, X_{kr}^{i}$ Integral Functions For Constant Depth
X_{kn}^{i},X_{kn}^{i} Integral Functions to n for Constant Depth
X ¹¹ Integral Functions to n for Parabolic Haunch kn
$lpha,\gamma$ Coefficients of Length of Constant Depth
eta Coefficient of Length of Haunch
ω Coefficient of Depth of Haunch
μ or $\frac{n-a}{\gamma}$
μ_1 · · · · · · · · · · · · · · · · · · ·
$\mathcal{T}_{AB}, \mathcal{T}_{BA}$ Angular Load Functions

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CHAPTER I

DEFINITION OF BEAM CONSTANTS

1. <u>General</u>.

The analysis of continuous beams by carry-over moments requires the evaluation of certain constants which are defined as either angular functions or moment functions of a simple beam. An isolated span of a continuous beam of variable cross-section is considered. The location of the cross-section is given by the ordinates x and x^{\dagger} (Fig. 1-1).

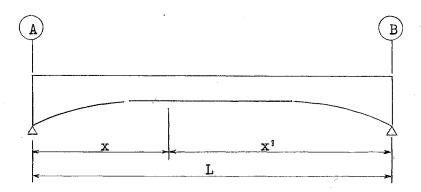


Fig. 1-1.-Isolated Span of Continuous Beam.

2. Angular Functions.

(a) <u>Angular Flexibilities</u>. The angular flexibility is the end slope of a simple beam due to a unit couple applied at that end. For the given beam (Fig. 1-2)

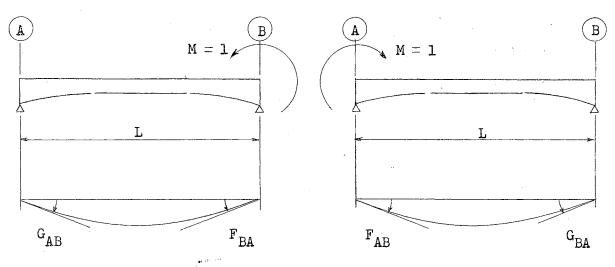
$$\mathbf{F}_{BA} = \frac{1}{L^2} \int_0^L \frac{\mathbf{x}^2 d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}}$$
(1-1a)

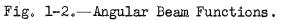
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$$F_{AB} = \frac{1}{L^2} \int_0^L \frac{x^{i^2} dx}{EI_x} \qquad (1-1b)$$

(b) <u>Carry-Over Values</u>. The carry-over value is the end slope of a simple beam due to a unit couple applied at the far end (Fig. 1-2).

$$G = G_{AB} = G_{BA} = \frac{1}{L^2} \int_0^L \frac{xx^{i} dx}{EI_x} \quad . \quad (1-2)$$





Substituting the identity

 $\mathbf{x}^{1} \equiv \mathbf{L} - \mathbf{x}$

and rearranging Eqs. (1-1b, 2) gives

$$G = \frac{1}{L} \int_{0}^{L} \frac{x dx}{E I_{x}} - F_{BA}$$
(1-3)

$$F_{AB} = \int_{0}^{L} \frac{dx}{EI_{x}} - 2G - F_{BA} \qquad (1-4)$$

(c) <u>Angular Load Functions</u>. The angular load functions are the end slopes of a simple beam due to applied transverse loads.

<u>Live Load</u>. If a unit load is applied to the beam (Fig. 1-3), the angular load functions are

$$\mathcal{T}_{BA}^{(LL)} = nLG + \frac{1}{L} \int_{0}^{Ln} \frac{x^2 dx}{EI_x} - n \int_{0}^{Ln} \frac{x dx}{EI_x}$$
(1-5)

$$\tau_{AB}^{(LL)} = nL \left[G + F_{AB} - \int_{0}^{Ln} \frac{dx}{EI_{x}} \right] + \int_{0}^{Ln} \frac{xdx}{EI_{x}} - \tau_{BA}^{(LL)} \quad . \quad (1-6)$$

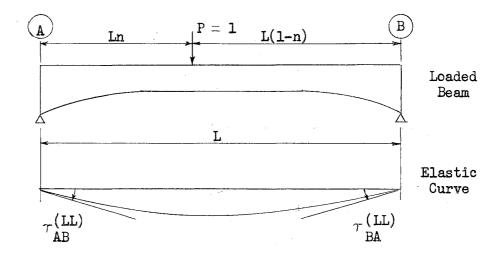


Fig. 1-3.-Angular Live Load Functions.

<u>Uniform Load</u>. If a uniform load of intensity w is applied to the beam (Fig. 1-4), the angular load functions are

$$\tau_{BA}^{(UL)} = \frac{\mathbf{w}}{2} \int_{0}^{L} \frac{\mathbf{x}^{2} d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}} - \frac{\mathbf{w}}{2\mathbf{L}} \int_{0}^{L} \frac{\mathbf{x}^{3} d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}}$$
(1-7a)
$$\tau_{AB}^{(UL)} = \frac{\mathbf{w}_{L}}{2} \int_{0}^{L} \frac{\mathbf{x} d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}} - \frac{\mathbf{w}}{2} \int_{0}^{L} \frac{\mathbf{x}^{2} d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}} - \tau_{BA}^{(UL)} .$$
(1-7b)

3

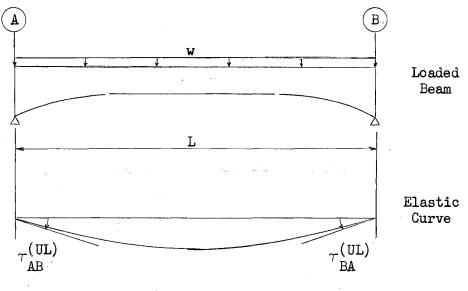


Fig. 1-4.-Angular Uniform Load Functions.

In terms of Eqs. (1-1a,2) Eqs. (1-7a,7b) become

$$\tau_{BA}^{(UL)} = \frac{wL^2}{2} F_{BA} - \frac{w}{2L} \int_{0}^{L} \frac{x^3 dx}{EI_x}$$
(1-8)

$$\tau_{AB}^{(UL)} = \frac{wL^2}{2} G - \tau_{BA}^{(UL)} . \qquad (1-9)$$

<u>Haunch Loads</u>. The angular load functions due to the dead load of the right haunch (Fig. 1-5) are

$$\tau_{BA}^{(HL)} = \frac{p\beta^2}{12} \int_0^L \frac{x^2 dx}{EI_x} - \frac{p}{12\beta^2 L^3} \int_L^L \frac{\left[\underline{x} - L(1-\beta)\right]^4 x dx}{EI_x}$$
(1-10a)

$$\tau_{AB}^{(HL)} = \frac{p\beta^2 L}{12} \int_0^L \frac{x dx}{EI_x} - \frac{p}{12(\beta L)^2} \int_L^L \frac{[x-L(1-\beta)]^4 x dx}{EI_x} - \tau_{BA}^{(HL)}.$$
 (1-10b)

where

p = The maximum intensity of the haunch load as shown in Fig. (1-5).

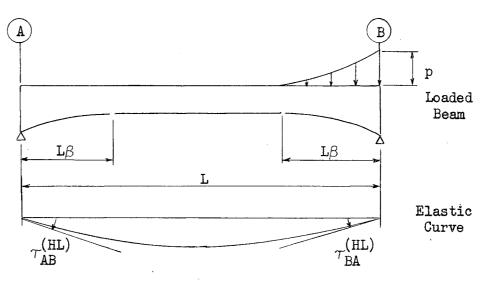


Fig. 1-5.-Angular Haunch Load Functions.

In terms of Eqs. (1-1a,3), Eqs. (1-10a,b) become

$$\tau_{BA}^{(HL)} = \frac{p\beta^2 L^2}{12} F_{BA} - \frac{p}{12\beta^2 L^3} \int_{L(1-\beta)}^{L} \frac{\left[\frac{x-L(1-\beta)}{EI}\right]^4 x dx}{L(1-\beta)}$$
(1-11)

$$\tau_{AB}^{(HL)} = \frac{p_{\beta}^{2}L^{2}}{12} (G + F_{BA}) - \frac{p}{12\beta^{2}L^{2}} \int_{L(1-\beta)}^{L} \frac{\left[\underline{x} - (L-\beta)\right]^{4} x dx}{EI_{x}} - \tau_{BA}^{(HL)} \quad .(1-12)$$

If the beam is symmetrical with two haunches of length L β , Eqs. (1-11,12) become

$$\tau_{BA}^{(HL)} = \tau_{AB}^{(HL)} = \frac{p\beta^{2}L^{2}}{24} \int_{0}^{L} \frac{dx}{EI_{x}} - \frac{p}{12\beta^{2}L^{2}} \int_{L(1-\beta)}^{L} \frac{[x-L(1-\beta)]^{4}dx}{EI_{x}} .$$
 (1-13)

Noting that for the symmetrical beam

$$\mathbf{F}_{AB} = \mathbf{F}_{BA}$$

and writing Eq. (1-13) in terms of Eq. (1-4)

$$\tau_{BA}^{(HL)} = \tau_{AB}^{(HL)} = \frac{p\beta^2 L^2}{12} (F_{AB} + G) - \frac{p}{12\beta^2 L^2} \int_{L(1-\beta)}^{L} \frac{\left[\frac{x-L(1-\beta)}{EL}\right]^4 dx}{L(1-\beta)} .(1-14)$$

More complete derivations of Eqs. (1-1) to (1-14) may be found in (2). The apparent difference in the limits of integration for Eqs. (1-10a) to (1-14) in this work and the reference work is due to the selection of different reference axes.

CHAPTER II

INTEGRAL FUNCTIONS

1. General.

The expressions for angular functions in Chapter 1 contain the recurring integrals

$$\int_{0}^{L} \frac{x^{k} dx}{I_{x}}$$
(2-1)
$$\int_{0}^{Ln} \frac{x^{k} dx}{I_{x}}$$
(2-2)

The solutions of integrals (2-1,2) may be greatly facilitated by relocating the reference axis and expressing the integrals in terms of integral functions R, Q, X_{kr}, and X_{kq}. The angular functions may then be expressed in terms of the integral functions as coefficients.

2. Integral Functions - Constant Depth.

A typical haunched beam is considered Fig. (2-1). The cross-section of the beam is constant over the length $L\gamma$. The moment of inertia

$$I_x = I_0 = \frac{1}{12} bh_0^3$$
 (2-3)

where

b = constant width $h_0 = reference depth$.

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The location of the cross-section is given by

$$\mathbf{x} = \mathbf{L}\mathbf{a} + \mathbf{t}_{\mathbf{x}}\mathbf{L}\boldsymbol{\gamma} \tag{2-4}$$

in which t_x varies between the limits

$$t_x = 0, t_x = 1$$

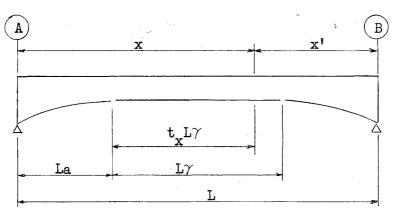


Fig. 2-1.-Typical Haunched Beam.

In terms of Eqs. (2-3,4) the integral (2-1) for constant crosssection becomes

$$\int_{La}^{L(a+\gamma)} \frac{\mathbf{x}^{k} d\mathbf{x}}{\mathbf{I}_{\mathbf{x}}} = \frac{\mathbf{L}^{k+1}}{\mathbf{I}_{0}} \int_{0}^{1} (a+\mathbf{t}_{\mathbf{x}}\gamma)^{k} \gamma d\mathbf{t}_{\mathbf{x}}$$

$$= \frac{\mathbf{L}^{k+1}}{\mathbf{I}_{0}} \mathbf{X}_{kr} \quad .$$
(2-5)

If the upper limit is x = Ln such that the corresponding upper limit

$$\mathbf{t}_{\mathbf{X}} = \boldsymbol{\mu} = \frac{\mathbf{n} - \mathbf{a}}{\gamma} , \qquad (2-6)$$

integral (2-2) becomes in terms of Eqs. (2-3,4,6)

$$\int_{La}^{Ln} \frac{\mathbf{x}^{\mathbf{k}} d\mathbf{x}}{\mathbf{I}_{\mathbf{x}}} = \frac{\mathbf{L}^{\mathbf{k}+\mathbf{l}}}{\mathbf{I}_{0}} \int_{0}^{\mu} (\mathbf{a} + \mathbf{t}_{\mathbf{x}} \gamma)^{\mathbf{k}} \gamma d\mathbf{t}_{\mathbf{x}}$$

$$= \frac{\mathbf{L}^{\mathbf{k}+\mathbf{l}}}{\mathbf{I}_{0}} \mathbf{X}^{*}_{\mathbf{kn}} \quad .$$

$$(2-7)$$

k=0, 1, 2, 3

in terms of the following integral functions and their equivalents.

$$R_{0} = \int_{0}^{1} dt_{x} = 1$$

$$R_{1} = \int_{0}^{1} t_{x} dt_{x} = \frac{1}{2}$$

$$R_{2} = \int_{0}^{1} t_{x}^{2} dt_{x} = \frac{1}{3}$$

$$R_{3} = \int_{0}^{1} t_{x}^{3} dt_{x} = \frac{1}{4}$$

$$R_{0n} = \int_{0}^{\mu} dt_{x} = \mu$$

$$R_{1n} = \int_{0}^{\mu} t_{x} dt_{x} = \frac{\mu^{2}}{2}$$

$$R_{2n} = \int_{0}^{\mu} t_{x}^{2} dt_{x} \frac{\mu^{3}}{3}$$

3. Integral Functions - Parabolic Haunch.

A typical haunch of length L β is considered (Fig. 2-2). The depth of haunch varies as a parabola of second degree and its depth at support B is

$$h_h = \omega h_0$$

The location of the cross-section is given by

$$\mathbf{x} = \mathbf{L}\mathbf{a}_{1} + \mathbf{z} = \mathbf{L}\mathbf{a}_{1} + \mathbf{t}_{\mathbf{z}}\mathbf{L}\boldsymbol{\beta} \qquad (2-9)$$

.

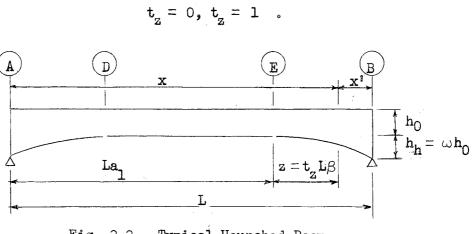


Fig. 2-2.-Typical Haunched Beam.

The depth of the cross-section at x is

$$h_{x} = h_{z} = h_{0}T_{z}$$
 (2-10)

where ${\rm T}_{_{\rm Z}}$ is the parabolic function

$$T_z = 1 + \omega t_z^2 \quad . \tag{2-11}$$

From Eq. (2-10) the moment of inertia at x is

$$I_x = I_z = I_0 T_z^3$$
 (2-12)

In terms of Eqs. (2-9,12) the integral (2-1) for the parabolic haunch

is

$$\int_{La_{1}}^{L(a_{1}+\beta)} \frac{x^{k} dx}{I_{x}} = \frac{L^{k+1}}{I_{0}} \int_{0}^{1} \frac{(a_{1}+t_{z}\beta)^{k} \beta dt_{z}}{T_{z}^{3}}$$
$$= \frac{L^{k+1}}{I_{0}} X_{kq} \quad .$$
(2-13)

If the upper limit x = Ln is such that the corresponding upper limit

$$t_z = \mu_1 = \frac{n-a_1}{\beta}$$
, (2-14)

integral (2-2) becomes in terms of Eqs. (2-9,12,14)

$$\int_{La}^{Ln} \frac{\mathbf{x}^{k} d\mathbf{x}}{\mathbf{I}_{\mathbf{x}}} = \frac{\mathbf{L}^{k+1}}{\mathbf{I}_{0}} \int_{0}^{\mu} \frac{(\mathbf{a}_{1} + \mathbf{t}_{z}\beta)^{k}\beta d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}}$$

$$= \frac{\mathbf{L}^{k+1}}{\mathbf{I}_{0}} \mathbf{x}^{**} \mathbf{k}_{n} \quad .$$

$$(2-15)$$

The following algebraic equivalents are introduced for the solution of Eqs. (2-13,15).

$$B = \frac{1}{4(1+\omega)^2}$$

$$B_n = \frac{\mu_1}{4(1+\omega\mu_1^2)^2}$$

$$C = \frac{1}{8(1+\omega)}$$

$$C_n = \frac{\mu_1}{8(1+\omega\mu_1^2)}$$

$$D = \frac{1}{8\sqrt{\omega}} \operatorname{Tan}^{-1} \sqrt{\omega}$$

$$D_n = \frac{1}{8\sqrt{\omega}} \operatorname{Tan}^{-1} \mu_1 \sqrt{\omega}$$

$$Lg = \frac{1}{2} \log (1+\omega)$$

$$(2-16a)$$

With notations (2-16a) Eqs. (2-13,15) will be evaluated in Chapters 3, 4, and 5 in terms of the following integral functions and their equivalents.

$$Q_{0} = \int_{0}^{1} \frac{dt_{z}}{T_{z}^{3}} = B + 3C + 3D$$

$$Q_{1} = \int_{0}^{1} \frac{t_{z}dt_{z}}{T_{z}^{3}} = B + 2C$$

$$Q_{2} = \int_{0}^{1} \frac{t_{z}^{2}dt_{z}}{T_{z}^{3}} = \frac{1}{\omega}(-B + C + D)$$

$$Q_{3} = \int_{0}^{1} \frac{t_{z}^{3}dt_{z}}{T_{z}^{3}} = B$$

(2**-**16b)

$$Q_{0n} = \int_{0}^{\mu} \frac{dt_{z}}{T_{z}^{3}} = B_{n} + 3C_{n} + 3D_{n}$$

$$Q_{1n} = \int_{0}^{\mu} \frac{t_{z}dt_{z}}{T_{z}^{3}} = \mu_{1}(B_{n} + 2C_{n}) \qquad (2-16b)$$
(Cont.)

$$Q_{2n} = \int_{0}^{\mu_{1}} \frac{t_{z}^{2} dt_{z}}{T_{z}^{3}} = \frac{1}{\omega} (-B_{n} + C_{n} + D_{n})$$
.

Expressed in terms of t_z and T_z , the integrals contained in the expressions for haunch load angular functions (Eqs. 1-11,12,13,14) are

$$\int_{0}^{L\beta} \frac{z^{4}dz}{I_{z}} = \frac{L}{I_{0}}^{5} \int_{0}^{1} \frac{\beta^{5}t_{z}^{4}dt_{z}}{T_{z}^{3}}$$
(2-17)
$$= \frac{L}{I_{0}}^{5} X_{4q} .$$
$$\int_{0}^{L\beta} \frac{z^{4}(La_{1}+z)dz}{I_{z}} = \frac{L}{I_{0}}^{6} \int_{0}^{1} \frac{\beta^{5}t_{z}^{4}(a_{1}+t_{z}\beta)dt_{z}}{T_{z}^{3}}$$
(2-18)
$$= \frac{L}{I_{0}}^{6} X_{5q} .$$

With notations (2-16a,b) Eqs. (2-17,19) will be evaluated in Chapters 3, 4, and 5 in terms of the following integral functions and algebraic equivalents.

$$Q_{4} = \int_{0}^{1} \frac{t_{z}^{4} dt_{z}}{T_{z}^{3}} = \frac{1}{\omega^{2}} (Q_{0} - 8C)$$

$$Q_{5} = \int_{0}^{1} \frac{t_{z}^{5} dt_{z}}{T_{z}^{3}} = \frac{1}{\omega^{3}} (Lg - 2\omega^{2}Q_{3} - Q_{1})$$
(2-19)

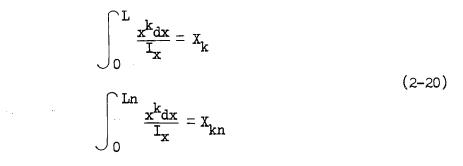
4. Integral Substitution And Numerical Evaluation.

Substituting the lower limits

$$t_{x,z} = 0$$

into Eqs. (2-8,16a) reduces all expressions to zero. Hence, Eqs. (2-8, 16b,19) need be evaluated only at the upper limit to obtain the desired solutions.

The integrals (2-1,2) will be replaced by the integral functions



and the second second

in which the symbols X_k , X_{kn} are the sum of the X-functions as indicated by the limits.

The algebraic expressions for R-functions (Eqs. 2-8) and Q-functions (Eqs. 2-16b) will be evaluated numerically on the IBM 650 Electronic Computer. From these values, the numerical solutions for the X-functions (Eqs. 2-20) and finally the angular function coefficients to be derived in Chapters 4 and 5 will be obtained.

CHAPTER III

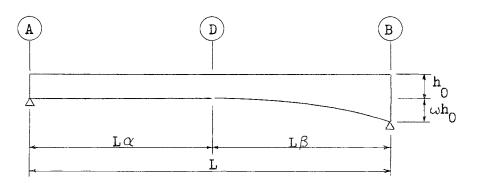
ANGULAR COEFFICIENTS - BEAMS WITH ONE HAUNCH

1. <u>General</u>.

A beam of length L with one parabolic haunch is considered Fig. (3-1). The beam is of constant depth over the length L α and the reference moment of inertia is Eq. (2-3)

$$I_0 = \frac{1}{12}bh_0^3$$
 (3-1)

The length of the haunch is $L\beta$ and the moment of inertia at any crosssection within the haunch is Eq. (2-11)



 $I_x = I_0 T_z^3$ (3-2)

Fig. 3-1. - Beam With One Haunch.

2. Integral Functions X_{kr} $(x = 0 \rightarrow L\alpha)$.

The location of the cross-section is given by (Fig. 3-2)

$$\mathbf{x} = \mathbf{t}_{\mathbf{x}} \mathbf{L} \boldsymbol{\alpha}$$
.

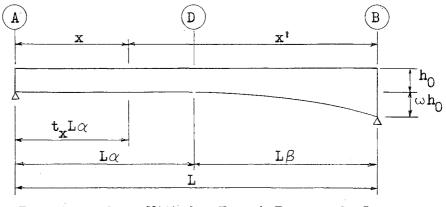


Fig. 3-2.—Beam With One Haunch For $x = 0 \rightarrow L\alpha$.

Substituting

into Eqs. (2-5,7) gives

$$\mathbf{X}_{\mathbf{kr}} = \alpha^{\mathbf{k}+\mathbf{l}} \int_{0}^{\mathbf{l}} \mathbf{t}_{\mathbf{x}}^{\mathbf{k}} d\mathbf{t}_{\mathbf{x}}$$
(3-4)

$$X_{kn}^{*} = \alpha^{k+1} \int_{0}^{\mu} t_{x}^{k} dt_{x} . \qquad (3-5)$$

In terms of functions R_k , Eqs. (2-7; 3-4,5) are

$$X_{0r} = \alpha \int_{0}^{1} dt_{x} = \alpha R_{0}$$

$$X_{1r} = \alpha^{2} \int_{0}^{1} t_{x} dt_{x} = \alpha^{2} R_{1}$$

$$X_{1r} = \alpha^{2} \int_{0}^{1} t_{x} dt_{x} = \alpha^{2} R_{1}$$

$$X_{1r} = \alpha^{2} \int_{0}^{\mu} t_{x} dt_{x} = \alpha^{2} R_{1n}$$

$$X_{2r} = \alpha^{3} \int_{0}^{1} t_{x}^{2} dt_{x} = \alpha^{3} R_{2}$$

$$X_{2r} = \alpha^{3} \int_{0}^{\mu} t_{x}^{2} dt_{x} = \alpha^{3} R_{2n}$$

$$X_{2r} = \alpha^{3} \int_{0}^{\mu} t_{x}^{2} dt_{x} = \alpha^{3} R_{2n}$$

$$\mathbf{x}_{3r} = \alpha^{4} \int_{0}^{1} \mathbf{t}_{\mathbf{x}}^{3} d\mathbf{t}_{\mathbf{x}} = \alpha^{4} \mathbf{R}_{3} \quad . \tag{3-6}$$
(Cont.)

If the upper limit x = Ln (Eqs. 2-6,7) is such that

 $0 \leq n \leq \alpha$,

with notation (2-20)

$$X_{kn} = X^{*}_{kn}$$
 (3-7)

3. Integral Functions \underline{X}_{kq} (x = L $\alpha \rightarrow L$).

The location of the cross-section is given by (Fig. 3-3)

$$\mathbf{x} = \mathbf{L}\alpha + \mathbf{t}_{\mathbf{z}}\mathbf{L}\beta . \tag{3-8}$$

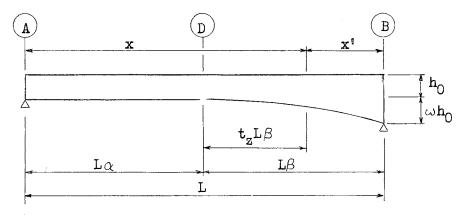


Fig. 3-3.—Beam With One Haunch For $\mathbf{x} = \mathbf{L}\alpha \rightarrow \mathbf{L}$.

Substituting

$$\mathbf{a} = \boldsymbol{\alpha} = \mathbf{1} - \boldsymbol{\beta} \tag{3-9}$$

into Eqs. (2-12,13) yields

$$X_{kq} = \int_{0}^{1} \frac{\left[1 - \beta(1 - t_z)\right]^k \beta dt_z}{T_z^3}$$
(3-10)

$$X^{**}_{kn} = \int_{0}^{\mathcal{H}_{1}} \frac{\left[1-\beta(1-t_{z})\right]^{k}\beta dt_{z}}{T_{z}^{3}} \quad . \tag{3-11}$$

The bracketed terms in Eqs. (3-10,11) may be replaced by the identities

$$Z = 1 - \beta(1 - t_z)$$

$$Z^0 = 1$$

$$Z^1 = Z$$

$$Z^2 = \beta^2 (1 - t_z)^2 + 2Z - 1$$

$$Z^3 = \beta^3 (1 - t_z)^3 + 3Z^2 - 3Z + 1$$
.
(3-12)

Thus, each succeeding power of Z is written in terms of preceeding powers.

Expressed in terms of Q_k functions (Eqs. 2-18) and with notations (3-12) Eqs. (3-10,11) are

$$\begin{split} \mathbf{X}_{0q} &= \int_{0}^{1} \frac{\mathbf{Z}_{\beta}^{0} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = \beta \mathbf{Q}_{0} \\ \mathbf{X}_{1q} &= \int_{0}^{1} \frac{\mathbf{Z}_{\beta}^{1} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = -\beta^{2}(\mathbf{Q}_{0} - \mathbf{Q}_{1}) + \mathbf{X}_{0q} \\ \mathbf{X}_{2q} &= \int_{0}^{1} \frac{\mathbf{Z}_{\beta}^{2} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = \beta^{3}(\mathbf{Q}_{0} - 2\mathbf{Q}_{1} + \mathbf{Q}_{2}) + 2\mathbf{X}_{1q} - \mathbf{X}_{0q} \\ \mathbf{X}_{3q} &= \int_{0}^{1} \frac{\mathbf{Z}_{\beta}^{3} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = \beta^{4}(-\mathbf{Q}_{0} + 3\mathbf{Q}_{1} - 3\mathbf{Q}_{2} + \mathbf{Q}_{3}) + 3\mathbf{X}_{2q} - 3\mathbf{X}_{1q} + \mathbf{X}_{0q} \quad . \end{split}$$

The X functions for uniform and dead loads (Eqs. 2-17,18) expressed in terms of notations (2-19; 3-8) are

$$x_{4q} = \int_{0}^{1} \frac{\beta^{5} t_{z}^{4} dt_{z}}{T_{z}^{3}} = \beta^{5} Q_{4}$$

$$x_{5q} = \int_{0}^{1} \frac{\beta^{5} t_{z}^{4} Z dt_{z}}{T_{z}^{3}} = x_{4q} + \beta^{6} (Q_{5} - Q_{4}) \quad .$$
(3-14)

Similarly as for functions X

$$\begin{split} \mathbf{X}_{0n}^{iii} &= \int_{0}^{\mu_{1}} \frac{\mathbf{Z}_{\beta}^{0} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = \beta \mathbf{Q}_{0n} \\ \mathbf{X}_{1n}^{iii} &= \int_{0}^{\mu_{1}} \frac{\mathbf{Z}_{\beta}^{1} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = -\beta^{2} (\mathbf{Q}_{0} - \mathbf{Q}_{1}) + \mathbf{X}_{0n}^{iii} \\ \mathbf{X}_{2n}^{iii} &= \int_{0}^{\mu_{1}} \frac{\mathbf{Z}_{\beta}^{2} d\mathbf{t}_{z}}{\mathbf{T}_{z}^{3}} = \beta^{3} (\mathbf{Q}_{0n} - \mathbf{Q}_{1n} + \mathbf{Q}_{2n}) + 2\mathbf{X}_{1n}^{ii} - \mathbf{X}_{0n}^{iii} \end{split}$$
(3-15)

If the upper limit x = Ln (Eqs. 2-14,15) is such that

$$\alpha \leq n \leq 1,$$

$$X_{kn} = X_{kr} + X_{kn}^{**} \quad . \tag{3-16}$$

4. Angular Functions - Coefficients.

The angular function coefficients are introduced as noted. With notations (2-5,7,13,15) the angular beam functions become:

(a) <u>Angular Flexibility</u>, $\underline{F}_{BA}(\underline{Eq. 1-la})$. $F_{BA} = \frac{1}{L^2} \int_{0}^{L} \frac{x^2 dx}{\underline{EI}_x} = \frac{L}{\underline{EI}_0} X_2$ $= \frac{L}{\underline{EI}_0} f_{BA}$ (3-17)

hence,

$$f_{BA} = X_{2r} + X_{2q}$$
 (3-18)

(b) Angular Carry Over Value, G (Eq. 1-3).

~ 7

$$G = \frac{1}{L} \int_{0}^{L} \frac{x dx}{E I_{x}} - F_{BA} = \frac{L}{E I_{0}} (X_{1} - f_{BA})$$

$$= \frac{L}{E I_{0}} g$$
(3-19)

18

in which

$$g = X_{lr} + X_{lq} - f_{BA} \quad . \tag{3-20}$$

(c) <u>Angular Flexibility</u>, \underline{F}_{AB} (Eq. 1-4). $F_{AB} = \int_{0}^{L} \frac{dx}{EI_{x}} - 2G - F_{BA} = \frac{L}{EI_{0}} (X_{0} - 2g - f_{BA})$ $= \frac{L}{EI_{0}} f_{AB}$ (3-21)

and

$$f_{AB} = X_{0r} + X_{0q} - 2g - f_{BA}$$
 (3-22)

Expressed in terms of X_{kq} functions (Eqs. 2-5,13) and angular beam function coefficients (Eqs. 3-18,20,22), the uniform load angular coefficients are:

(a) Angular Load Function
$$T_{BA}^{(UL)}$$
 (Eq. 1-8) .

$$\tau_{BA}^{(UL)} = \frac{wL^2}{2} F_{BA} - \frac{w}{2L} \int_0^L \frac{x^3 dx}{EI_x} = \frac{wL^3}{EI_0} \frac{1}{2} (f_{BA} - X_3)$$

$$= \frac{wL^3}{EI_0} t_{BA}^{(UL)}$$
(3-23)

in which

$$t_{BA}^{(UL)} = \frac{1}{2} (f_{BA} - X_{3r} - X_{3q}) . \qquad (3-24)$$
(b) Angular Load Function $T_{AB}^{(UL)} (Eq. 1-9) .$

$$\tau_{AB}^{(UL)} = \frac{wL^{2}}{2} G - \tau_{BA}^{(UL)} = \frac{wL^{3}}{EI_{0}} (\frac{1}{2} g - t_{BA}^{(UL)}) \qquad (3-25)$$

$$= \frac{wL^{3}}{EI_{0}} t_{AB}^{(UL)}$$

in which

$$t_{AB}^{(UL)} = \frac{1}{2} g - t_{BA}^{(UL)}$$
 (3-26)

The haunch dead load angular coefficients expressed in terms of X_{kq} functions (Eqs. 2-17,18; 3-14) and angular beam function coefficients (Eqs. 3-18,20,22) are

(a) <u>Angular Load Function</u> $T_{BA}^{(HL)}$ (Eq. 1-11).

$$\tau_{BA}^{(HL)} = \frac{p_{\beta}^{2}L^{2}}{L^{2}} F_{BA} - \frac{p}{12\beta^{2}L^{3}} \int_{0}^{L} \frac{z^{4}(La_{1}+z)dz}{EI_{z}}$$
$$= \frac{pL^{3}}{EI_{0}} \left(\frac{\beta^{2}}{12} f_{BA} - \frac{1}{12\beta^{2}} X_{5q}\right)$$
$$= \frac{pL^{3}}{EI_{0}} t_{BA}^{(HL)}$$
(3-27)

since

$$\mathbf{p} = \mathbf{q}\mathbf{h}_0\omega$$

$$\mathbf{t}_{\mathrm{BA}}^{(\mathrm{HL})} = \frac{\omega\beta^2}{12} \mathbf{f}_{\mathrm{BA}} - \frac{\omega}{12}\beta^2 \mathbf{X}_{5q} \quad . \tag{3-28}$$

(b) <u>Angular Load Function</u> $T_{AB}^{(HL)}$ (Eq. 1-12).

$$\tau_{AB}^{(HL)} = \frac{p\beta^{2}L^{2}}{12} (F_{BA} + G) - \frac{p}{12E\beta^{2}L^{2}} \int_{0}^{L\beta} \frac{z^{4}dz}{I_{z}}$$
$$= \frac{pL^{3}}{EI_{0}} (\frac{\beta^{2}}{12} f_{BA} + \frac{\beta^{2}}{12} g - \frac{1}{12\beta^{2}} X_{4q})$$
$$= \frac{pL^{3}}{EI_{0}} t_{AB}^{(HL)}$$
(3-29)

in which

$$t_{AB}^{(HL)} = \frac{\beta^2}{12} (f_{BA} + g) - \frac{\omega}{12\beta^2} X_{4q}$$
 (3-30)

The live load angular coefficients expressed in terms of X functions (Eqs. 3-6,12) and angular beam function coefficients (Eqs. 3-18, 20,22) are

(a) Angular Load Function
$$\underline{\tau}_{BA}^{(LL)}$$
 (Eq. 1-5).
 $\tau_{BA}^{(LL)} = nLG + \frac{1}{L} \int_{0}^{Ln} \frac{\mathbf{x}^2 d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}} - n \int_{0}^{Ln} \frac{\mathbf{x} d\mathbf{x}}{\mathbf{EI}_{\mathbf{x}}}$
 $= \frac{L^2}{\mathbf{EI}_0} (ng - nX_{1n} + X_{2n})$ (3-31)
 $= \frac{L^2}{\mathbf{EI}_0} t_{BA}^{(LL)}$

where

$$t_{BA}^{(LL)} = n(g - X_{ln}) + X_{2n} \qquad (3-32)$$
(b) Angular Load Function $T_{AB}^{(LL)}$ (Eq. 1-6).
$$\tau_{AB}^{(LL)} = nL(G + F_{AB} - \frac{1}{R} \int_{-\infty}^{-1} \frac{dx}{T} + \frac{1}{R} \int_{-\infty}^{-1} \frac{xdx}{T} - \tau_{BA}^{(LL)}$$

$$\begin{array}{l} \mathcal{A}_{AB} = \operatorname{HI}(\mathbf{G} + \mathbf{I}_{AB} - \mathbf{E}_{B}) \\ = \frac{\mathbf{L}^{2}}{\mathbf{E}\mathbf{I}_{0}} \left(\operatorname{ng} + \operatorname{nf}_{AB} - \operatorname{nX}_{0n} + \mathbf{X}_{1n} - \mathbf{t}_{BA}^{(\mathrm{LL})} \right) \\ = \frac{\mathbf{L}^{2}}{\mathbf{E}\mathbf{I}_{0}} \mathbf{t}_{AB}^{(\mathrm{LL})} \end{array}$$
(3-33)

and

$$\mathbf{t}_{AB}^{(LL)} = \mathbf{n}(\mathbf{g} + \mathbf{f}_{AB} - \mathbf{X}_{On}) + \mathbf{X}_{ln} - \mathbf{t}_{BA}^{(LL)} . \qquad (3-34)$$

CHAPTER IV

ANGULAR COEFFICIENTS - SYMMETRICAL BEAMS

1. <u>General</u>.

A beam of length L with two symmetrical parabolic haunches is considered (Fig. 4-1). The beam is of constant depth over the length $L\gamma$ and from Eq. (2-3)

$$I_x = I_0 = \frac{1}{12} bh_0^3$$
. (4-1)

The length of each haunch is $L\beta$ and for any cross-section within either haunch from Eqs. (2-11,12)

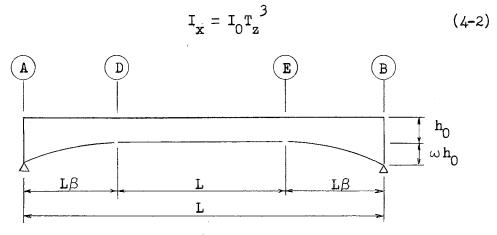


Fig. 4-1.—Symmetrical Beam.

2. Integral Functions X_{kq} (x = 0 -> L β).

The location of the cross-section is given by (Fig. 4-2)

$$\mathbf{x} = \mathbf{L}\boldsymbol{\beta} - \mathbf{t}_{\mathbf{z}}\mathbf{L}\boldsymbol{\beta}$$
 .

For

$$a_1 = \beta$$

and, noting the change in sign, from Eq. (2-13)

$$\mathbf{X}_{\mathbf{kq}}^{*} = \frac{\mathbf{L}^{\mathbf{k}+1}}{\mathbf{I}_{0}} \int_{0}^{1} \frac{\beta^{\mathbf{k}}(1-\mathbf{t}_{\mathbf{z}})^{\mathbf{k}}\beta d\mathbf{t}_{\mathbf{z}}}{\frac{\mathbf{T}_{\mathbf{z}}^{3}}{\mathbf{T}_{\mathbf{z}}}} \quad .$$
(4-3)

With notations (3-9,10) the integral functions X_{kq}^i may be expressed in terms of integral functions X_{kq}^i . Thus,

$$X_{kq}^{*} = \frac{\underline{L}^{k+1}}{I_{0}} \int_{0}^{1} \frac{\left[1 - \{1 - \beta(1 - t_{z})\}\right]^{k} \beta dt_{z}}{T_{z}^{3}}$$

$$= \frac{\underline{L}^{k+1}}{I_{0}} \int_{0}^{1} \frac{(1 - Z)^{k} \beta dt_{z}}{T_{z}^{3}} .$$
(4-4)

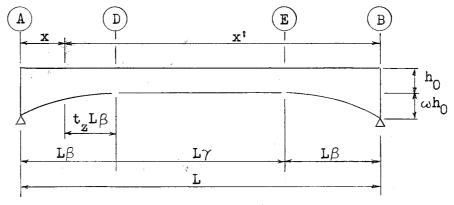


Fig. 4-2.—Symmetrical Beam For $x = 0 \rightarrow L\beta$.

For values

$$k = 0, 1, 2$$

$$x_{0q} = \int_{0}^{1} \frac{(1-Z)^{0} \beta dt_{z}}{T_{z}^{3}} = x_{0q}$$
 (4-5)

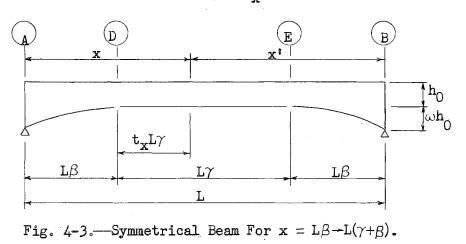
$$X_{lq}^{i} = \int_{0}^{1} \frac{(1-Z)^{l}\beta dt_{z}}{T_{z}^{3}} = X_{0q} - X_{lq}$$

$$X_{2q}^{i} = \int_{0}^{1} \frac{(1-Z)^{2}\beta dt_{z}}{T_{z}^{3}} = X_{0q} - 2X_{lq} + X_{2q} .$$
(4-5)
(Cont.)

Since the beam is symmetrical the angular live load coefficients need not be computed for the unit load in positions that require the computation of X_{kn} functions for the left haunch.

3. Integral Functions $\underline{X}_{kr} (\underline{x} = \underline{L}\beta \rightarrow \underline{L}(\beta + \gamma))_{\bullet}$

The location of the cross-section is given by (Fig. 4-3)



$$\mathbf{x} = \mathbf{L}\boldsymbol{\beta} + \mathbf{t}_{\mathbf{v}}\mathbf{L}\boldsymbol{\alpha}$$

Substituting into Eq. (2-4)

 $\mathbf{a} = \beta$

gives for functions X_{kr}^{\dagger} (Eq. 2-5)

$$\mathbf{X}_{\mathbf{kr}}^{*} = \int_{0}^{1} \left(\beta + \mathbf{t}_{\mathbf{x}}^{\gamma}\right)^{\mathbf{k}} \gamma d\mathbf{t}_{\mathbf{x}} \quad . \tag{4-6}$$

Expressed in terms of notations (3-6), Eq. (4-6) for values

k = 0, 1

$$\begin{aligned} \mathbf{x}_{0\mathbf{r}}^{i} &= \int_{0}^{1} (\beta + \gamma \mathbf{t}_{\mathbf{x}})^{0} \gamma d\mathbf{t}_{\mathbf{x}} = \mathbf{x}_{0\mathbf{r}} \\ \mathbf{x}_{\mathbf{l}\mathbf{r}}^{i} &= \int_{0}^{1} (\beta + \gamma \mathbf{t}_{\mathbf{x}})^{1} \gamma d\mathbf{t}_{\mathbf{x}} = \beta \mathbf{x}_{0\mathbf{r}} + \mathbf{x}_{\mathbf{l}\mathbf{r}} \end{aligned}$$
(4-7)

Evaluating Eq. (4-6) in terms of Eqs. (4-7) and with notations (3-6) for k = 2 gives

$$\mathbf{X}_{2\mathbf{r}}^{i} = \int_{0}^{1} (\beta + \gamma \mathbf{t}_{\mathbf{x}})^{2} \gamma d\mathbf{t}_{\mathbf{x}} = -\beta^{2} \mathbf{X}_{0\mathbf{r}}^{i} + 2\beta \mathbf{X}_{1\mathbf{r}}^{i} + \mathbf{X}_{2\mathbf{r}} \quad .$$
 (4-8)

Thus, each function X_{kr}^{i} is written in terms of previously evaluated functions X_{kr}^{i} and notations (3-6).

For uniform load angular functions (Eqs. 1-7a,b) the integral (2-1) must be evaluated for

$$\alpha^{1} = \frac{\gamma}{2}$$

Substituting this value into notations (3-6) gives for uniform load

$$\begin{aligned} \mathbf{x}_{0\mathbf{r}}^{\text{ii}} &= \frac{\gamma}{2} \int_{0}^{1} d\mathbf{t}_{\mathbf{x}} = \frac{\gamma}{2} \mathbf{R}_{0} \\ \mathbf{x}_{2\mathbf{r}}^{\text{ii}} &= (\frac{\gamma}{2})^{3} \int_{0}^{1} \mathbf{t}_{\mathbf{x}}^{2} d\mathbf{t}_{\mathbf{x}} = (\frac{\gamma}{2})^{3} \mathbf{R}_{2} \quad . \end{aligned}$$
(4-9)

No other X_{kr}^{ii} functions are necessary.

Similarly as for functions X_{kr}^{i}

$$X_{0n}^{iii} = \int_{0}^{\mu} (\beta + \gamma t_{x})^{0} \gamma dt_{x} = X_{0n}^{i}$$

$$X_{ln}^{iii} = \int_{0}^{\mu} (\beta + \gamma t_{x})^{1} \gamma dt_{x} = X_{0n}^{i} + X_{ln}^{i}$$
(4-10)

$$\mathbf{X}_{2n}^{\mathbf{i}\mathbf{i}\mathbf{i}} = \int_{0}^{\mu} (\beta + \gamma \mathbf{t}_{\mathbf{x}})^{2} \gamma d\mathbf{t}_{\mathbf{x}} = -\beta^{2} \mathbf{X}_{0n}^{\mathbf{i}} + 2 \mathbf{X}_{1n}^{\mathbf{i}\mathbf{i}\mathbf{i}} + \mathbf{X}_{2n}^{\mathbf{i}} . \quad (4-10)$$
(Cont.)

If the upper limit x = Ln (Eqs. 2-6,7) is such that

$$\beta \leq n \leq \beta + \gamma$$

$$X_{kn} = X_{kq}^{i} + X_{kn}^{i+i} . \qquad (4-11)$$

4. Integral Functions $\underline{X}_{kq} \xrightarrow{(\mathbf{x} = \mathbf{L}(\beta + \gamma) \rightarrow \mathbf{L})}$.

The location of the cross-section is given by (Fig. 4-4)

$$\mathbf{x} = \mathbf{L}(\boldsymbol{\beta} + \boldsymbol{\gamma}) + \mathbf{t}_{\mathbf{z}}\mathbf{L}\boldsymbol{\beta}$$

and

$$\mathbf{a_1} = \beta + \gamma = \mathbf{1} - \beta$$

which is identical to the value for a_1 in Eq. (3-8), therefore with the exception of uniform load (Eqs. 1-7a,b), the solutions for functions X_{kq} and X_{kn} are given by Eqs. (3-10,11,12,13,15,16).

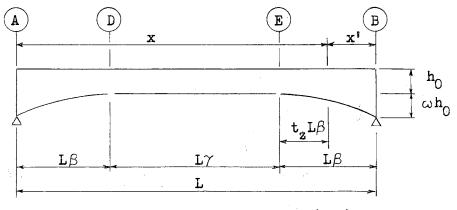


Fig. 4-4.—Symmetrical Beam For $x = L(\beta + \gamma) \rightarrow L$.

For uniform load the cross-section is measured from the center of the beam (Fig. 4-5)

$$\mathbf{x} = \mathbf{L} \frac{\gamma}{2} + \mathbf{t_z} \mathbf{L} \boldsymbol{\beta}$$
 .

Substituting

$$a_1 = \frac{\gamma}{2} = \frac{1}{2} (1 - 2\beta)$$

into Eq. (2-12) and considering notation (3-9) gives for k = 2

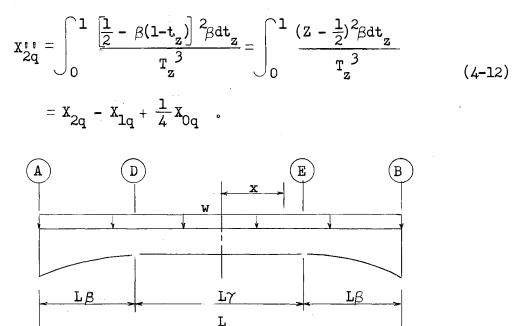


Fig. 4-5.-Symmetrical Beam For Uniform Load.

5. Angular Functions - Coefficients.

Similarly as for the unsymmetrical beam, the angular functions will be expressed in terms of coefficients as follows:

(a) <u>Angular Flexibility Coefficient</u>, f_{BA} (Eq. 3-17).

$$f_{BA} = X_2 = X_{2q}^{i} + X_{2r}^{i} + X_{2q}^{i}$$
 (4-13)

(b) <u>Angular Carry-Over Value Coefficient</u>, g (Eq. 3-19).

$$g = X_1 - f_{BA} = X_{1q}^{i} + X_{1r}^{i} + X_{1q}$$
 (4-14)

(c) Angular Flexibility Coefficient,
$$f_{AB} \xrightarrow{(Eq. 3-21)}$$
.
 $f_{AB} = X_0 - 2g - f_{BA} = X_{0q}^{!} + X_{0r}^{!} + X_{0q} - 2g - f_{BA}$. (4-15)

For symmetrical beams the uniform load angular functions may be written (3)

$$\tau_{BA}^{(UL)} = \tau_{AB}^{(UL)} = \frac{wL^2}{8E} \int_0^{\frac{L}{2}} \frac{dx}{I_x} - \frac{w}{2E} \int_0^{\frac{L}{2}} \frac{x^2 dx}{I_x} . \qquad (4-16)$$

Proceeding as in Eqs. (3-23,24) and with notations (3-13; 4-9)

$$t_{BA_{c}}^{(UL)} = \frac{1}{8} \left(X_{0r}^{ii} + X_{0q}^{i} \right) - \frac{1}{2} \left(X_{2r}^{ii} + X_{2q}^{ii} \right) \quad . \tag{4-17}$$

The haunch load angular functions are (Eq. 1-13)

$$\tau_{BA}^{(HL)} = \tau_{AB}^{(HL)} = \frac{p\beta^2}{12E} \int_0^{\frac{L}{2}} \frac{dx}{I_x} - \frac{p}{12\beta^2 L^2 E} \int_0^{L\beta} \frac{z \, dz}{I_z} \, . \quad (4-18)$$

Proceeding as in Eqs. (3-27,28) and with notations (3-13,14; 4-9)

$$t_{BA}^{(HL)} = \frac{\beta^2 \omega}{12} (X_{0r}^{**} + X_{0q}) - \frac{\omega}{12\beta^2} X_{4q} \quad .$$
 (4-19)

The coefficients for the angular live load functions are given by Eqs. (3-32,34).

CHAPTER V

PROGRAM FOR THE IBM 650 ELECTRONIC COMPUTER

1. <u>General</u>

Programming for the solution of any problem on a digital computer is usually accomplished in two steps. First, a drawing showing each phase of the problem and the sequence of operations is made. Second, from the schematic drawing, or flow chart, a series of instructions for the computer is established.

The program in this chapter was prepared in floating decimal arithmetic for the IEM 650 Electronic Computer at Oklahoma State University's Computing Center. The coding form used is that of IEM's Symbolic Optimum Assembly Programming, Type II. Storage locations have been re-used as soon as possible in order to take advantage of the available sixty high-speed storage locations.

2. Functional Evaluation.

The subroutines for the evaluation of the square roots, arc-tangents, and logarithms required for the solution of Eqs. (2-16a) are an integral part of the program, therefore no library subroutines are necessary.

The square roots are obtained by Newton's method which is as follows:

$$R_{i} = \frac{1}{2} (R_{i-1} + \frac{\omega}{R_{i+1}})$$
 (5-1)

in which

 ω = number for which the square root is desired,

 $R_{i-1} =$ preceding approximation of the square root.

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For all values of ω the initial approximation is one and fifteen successive approximations are made.

With notation (2-10) the arc-tangents are evaluated by the infinite series

$$\operatorname{Tan}^{-1} \mathbf{t}_{\mathbf{z}} / \overline{\omega} = \mathbf{V} = \frac{\mathbf{R}_{\mathbf{t}}}{\mathbf{t}_{\mathbf{z}} / \overline{\omega}} (1 + \frac{2}{3} \mathbf{R}_{\mathbf{t}} + \frac{2 \cdot 4}{3 \cdot 5} \mathbf{R}_{\mathbf{t}}^{2} + \frac{2 \cdot 4 \cdot 6}{3 \cdot 5 \cdot 7} \mathbf{R}_{\mathbf{t}}^{3} \cdots (5 - 2)$$

in which

$$R_{t} = \frac{\omega t_{z}^{2}}{1 + \omega t_{z}^{2}} = \frac{T_{z} - 1}{T_{z}}$$

The series converges for all values of

$$\omega t_z^2 < \infty$$
 .

For values

$$\omega t_z^2 \leq 0.1$$

computations are limited to eleven terms to avoid overflow in the computer. For larger values twenty-five terms are evaluated to obtain the desired accuracy.

With notation (2-10) the logarithms are evaluated by the infinite series

$$L_{g} = \frac{1}{2} L_{0} g_{e} T_{z} = R_{f} + \frac{1}{3} R_{f}^{3} + \frac{1}{5} R_{f}^{5} + \frac{1}{7} R_{f}^{7} \dots$$
(5-3)

in which

$$R_{l} = \frac{T_{z}-1}{T_{z}+1}$$

The above series converges for all values

and nineteen terms are used.

The accuracy of Eqs. (5-1,2,3) depends upon the magnitude of the value for which the function is required. Using the number of terms indicated and rounding off the angular function coefficients to four decimal places gives results that agree with those published in other works for

 $\omega_{\max} \leq 2$.

3. Input Card Format.

The description of the beam for which constants are desired is introduced into the computer with seven words. Fig. (5-1) shows the arrangement of input data.

Word	Card Columns Inclusive	Data
1	1 - 10	. W
2	11 - 20	β
3	21 - 30	Δω
4	31 - 40	Δβ
5	41 - 50	
6	51 - 60	β_{max}
7	61 - 70	Beam Type
8	71 - 80	zeros

Fig. 5-1.-Input Card Data.

The meanings of ω and β have already been established. The symbols

 $\Delta \omega$ and $\Delta \beta$ are the increments by which the dimension coefficients are to be increased. These two words must have some positive value even though computations may be required for only one beam. The fifth and sixth words are the maximum values the dimension coefficients may attain. The beam type number is zero for unsymmetrical beams and one for symmetrical.

The first six words of the input card must be in floating decimal form. The position of the decimal is obtained by subtracting 50 from the last two digits. If the result is zero the decimal immediately precedes the first digit. If the result is negative or positive the decimal is shifted to the left or right the indicated number of places. (Fig. 5-2).

The beam type number must be entered as a fixed point number, either ten zeros for unsymmetrical beams or one preceeded by nine zeros for symmetrical beams.

Number	Floating Decimal Form
345.6	3456000053
0.3456	3456000050
0.03456	3456000049

Fig. 5-2.—Examples of Floating Decimals,

Example 1.

Beam constants are to be computed for the symmetrical beam in Fig. (5-3).

Since computations are required for only one beam the maximum values of dimension coefficients are equal to the initial values and the incre-

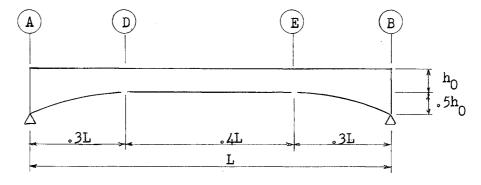


Fig. 5-3.—Symmetrical Beam.

The beam type number is one and the data is entered as in Fig. (5-4).

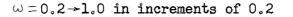
Word	Data Entered
1	500000050
2	300000050
3	100000050
4	100000050
5	500000050
6	300000050
7	00000001
8	Not Used

Fig. 5-4.-Input Card For One Symmetrical Beam.

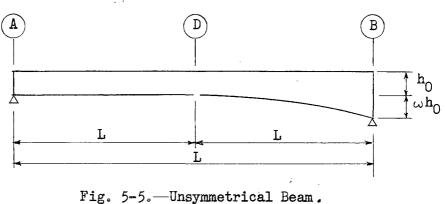
Example 2.

Beam constants are required for the beam shown in Fig. (5-5) for all combinations of

ments will be unimportant as long as they have some positive value.



 $\beta=0.1\!\rightarrow\!1.0$ in increments of 0.1 .



The input data is entered as in Fig. (5-6) and the computer will calculate the constants for the required 50 beams.

Word	Data Entered
1	200000050
2	100000050
3	200000050
4	100000050
5	100000051
6	100000051
7	00000000
8	(Not Used)

Fig. 5-6.—Input Card For Series of Unsymmetrical Beams.

4. Output Card Format.

The angular function coefficients will be in floating decimal form

on either three or four cards depending on the type of beam. The first word of each card will be an identification number. The first two digits are ten times the value ω , the fourth and fifth digits are ten times the value β_{g} and the last digit is the beam type number. The identification number for angular live load functions will have, in addition, ten times the last computed value of n as the seventh digit. Thus, the identification number

05 003 00 001

will appear on the first output card for the symmetrical beam for which

$\omega = 0.5$ $\beta = 0.3,$

and the number

05 003 09 001.

will appear on the card containing influence coefficients for

$$n = 0.7, 0.8, 0.9$$
.

The first output card for each beam will be arranged as follows:

Word	Information
1	Identification
2	f _{BA}
3	g
4	$\mathtt{f}_{\mathtt{A}\mathtt{B}}$
5	$\mathtt{t}_{\mathtt{BA}}^{(\mathtt{UL})}$
6	t ^(UL) AB
7	$t_{BA}^{(DL)}$
8	$t_{AB}^{(DL)}$

Fig. 5-7.-First Output Card.

The angular live load coefficients will appear as in Fig. (5-8). In the case of symmetrical beams the second and third words are not used for cards bearing identification numbers of the form.

xx 0xx 06 001 .

Word	Information	Position Of Load
1	Identification	
2	t ^(LL) BA	n – 2
3	$t_{AB}^{(LL)}$	n - 2
4	$t_{BA}^{(LL)}$	n - 1
5	$t_{AB}^{(LL)}$	n – 1
6	$t_{\rm BA}^{({ m LL})}$	n
7	$t_{AB}^{(LL)}$	n
8	(Not Used)	1

Fig. 5-8.—Output Card For Live Load Function Coefficients.

5. Flow Chart.

The flow chart in Fig. (5-9) was prepared as an aid to programming the solutions for angular function coefficients for beams with one haunch and symmetrical beams.

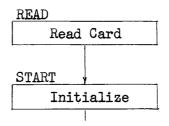


Fig. 5-9.—Flow Chart.

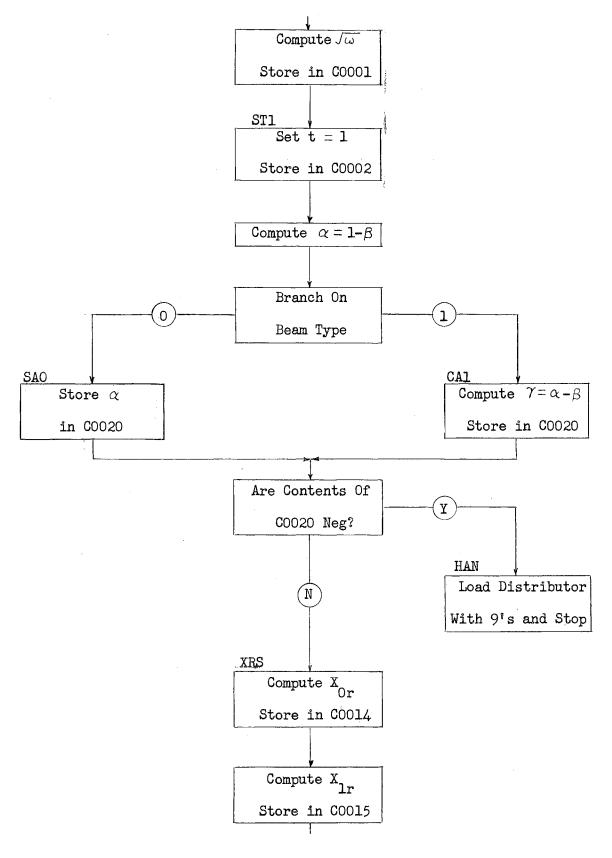
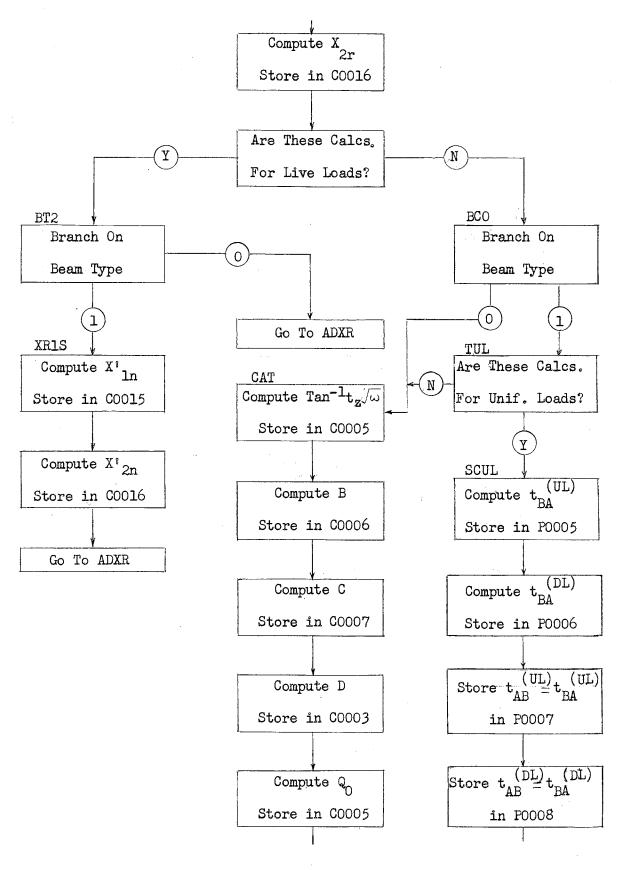
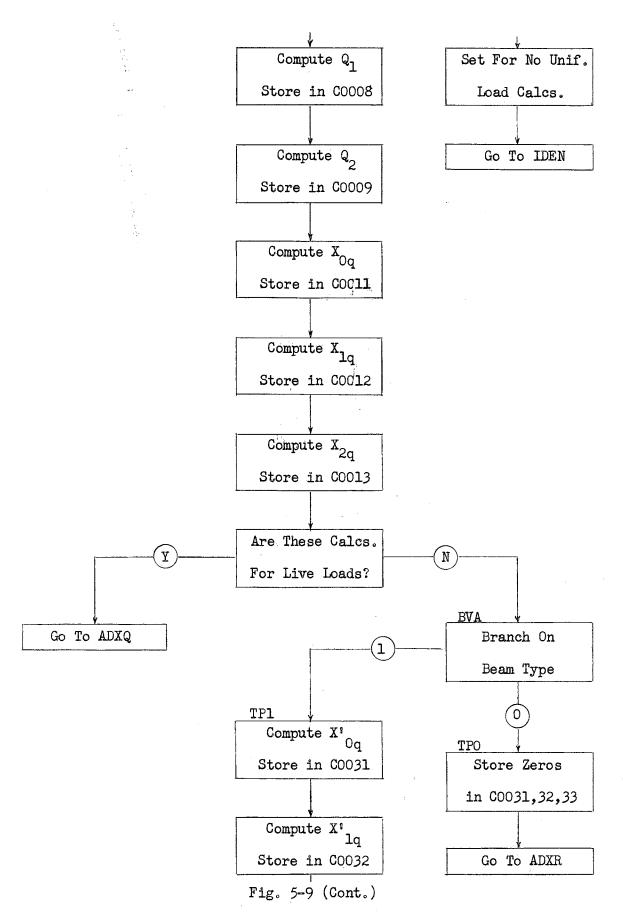
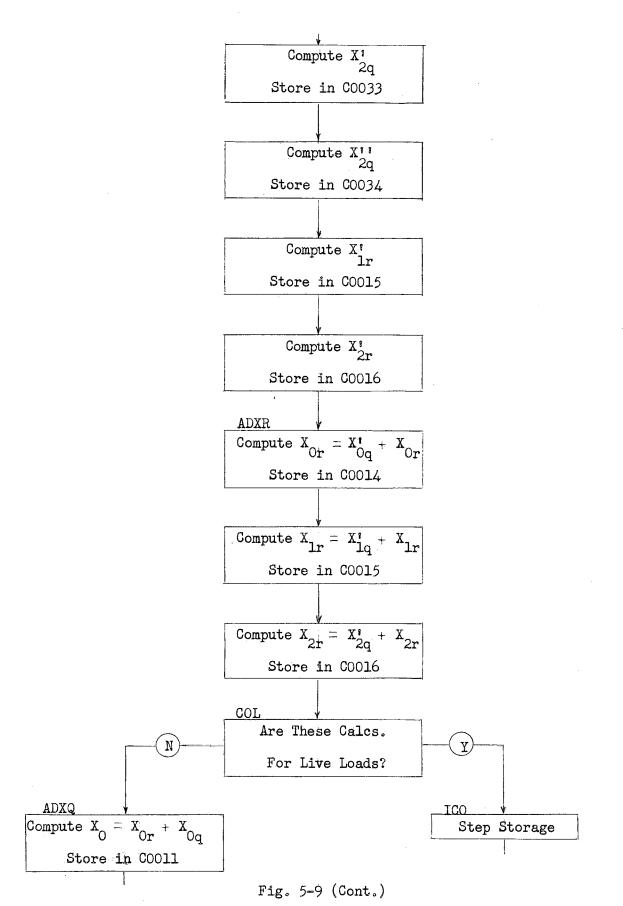


Fig. 5-9 (Cont.)







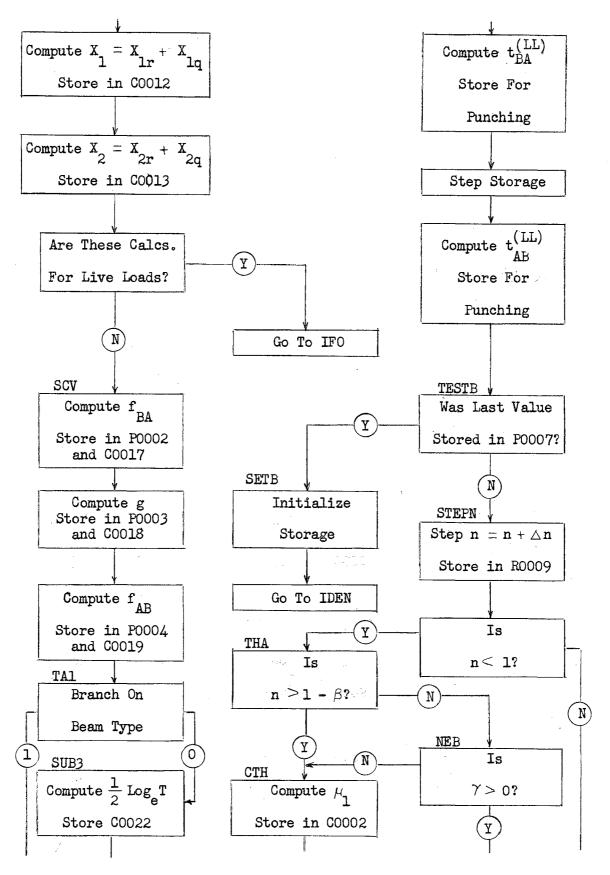
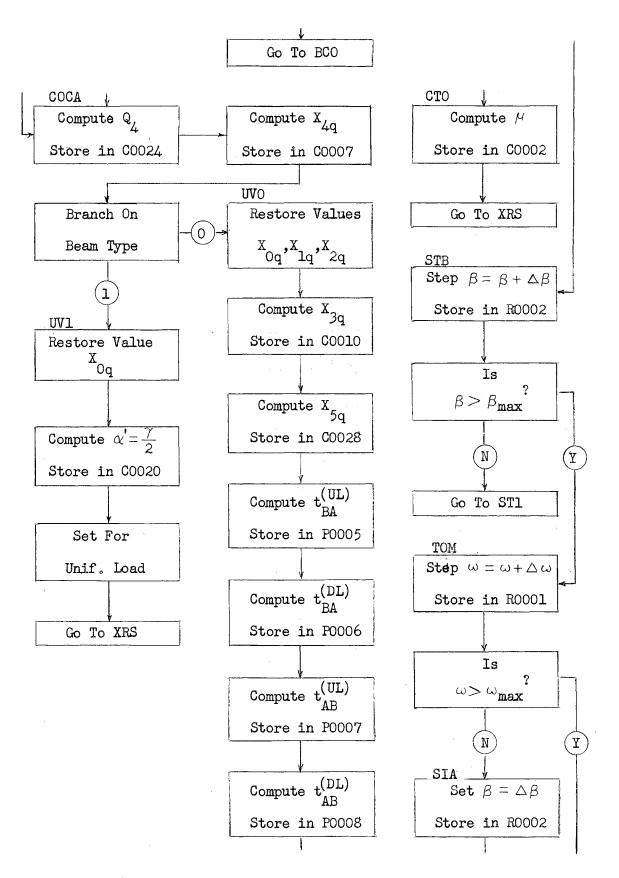
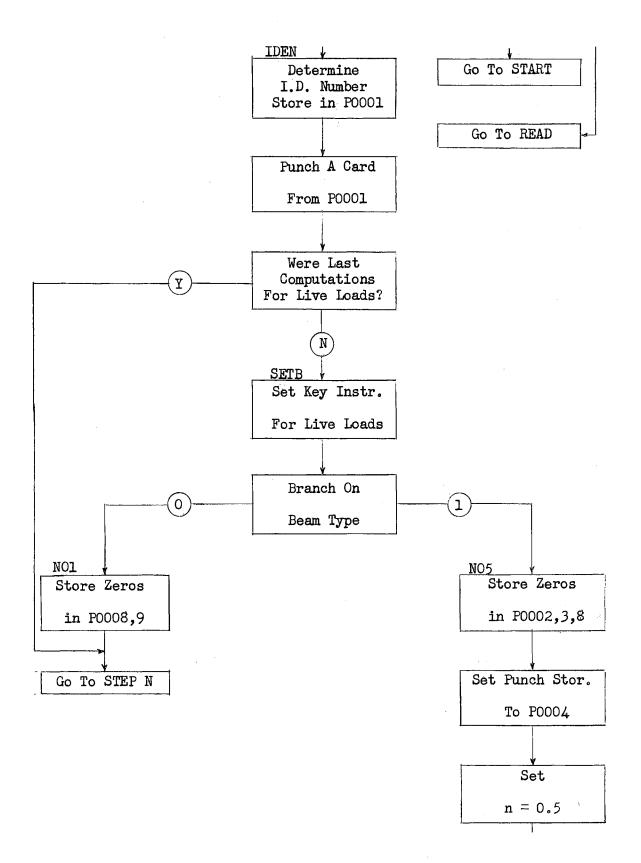
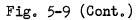


Fig. 5-9 (Cont.)







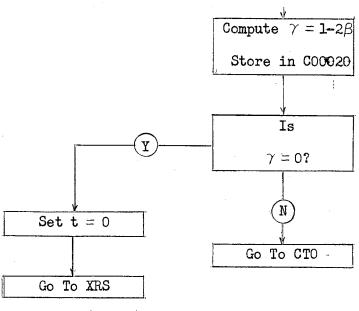


Fig. 5-9 (Cont.)

6. IBM 650 Program.

With the program in Fig. (5-10) the computer will calculate beam constants including 10-point influence coefficients for beams of either type for which is expressed as a multiple of one-tenth. Beam constants are available for beams with constant moment of inertia (2). No provision is made for them in this program and entry of ω or β equal to zero will result in an attempt to divide by zero. The computer will stop with ten nines in the distributor if α or γ become negative as would be the case if for a symmetrical beam β would be entered as 0.6. No other programmed stops are incorporated.

T P	S N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
1			BE	AM CO	N	STANT	S	
-			REG	R9000		9009		READ AREA

Fig. 5-10.-IBM 650 Program.

					T		Т	
Ţ ₽	<u>s</u> N	Location	Oper. Code	Data Address	A	Instr. Address	A G	Remarks
ı			REG REG SYN	C9010 P9050 READ		9049 9059 0000		CALC AREA PCH AREA BEGIN 0000
		READ START	RD1 LDD STD STD RAA RAB	R0001 ZER0 R0009 P0009 R0007 0000		START		READ CARD INITIALIZE N NO DL OP BM TYPE NO LL OP
1			RAC RAU FAD FMP STU	0015 R0001 ONE HALF C0001		SUB1		TRMS SQ RT STORE INITIAL APPROX SQ RT
		SUBL	RAU FDV FAD FMP STU SXC BMC	R0001 C0001 HALF C0001 0001 ST1		SUB1		COMPUTE AND STORE SQ RT
1		ST1 CA1 SAO	RAU STU FSB NZA FSB STU	ONE C0002 R0002 CA1 R0002 C0020		SAO SAO		SET T TO ONE AND COMPUTE ALPHA OR GAMMA AND STORE
1		HAN	BMI LDD HLT	HAN NINES		XRS		IF NEG STOP
		XRS	RAU FMP STU FMP STU FMP FMP FMP	C0020 C0012 C0014 C0014 HALF C0015 C0014 FRAC				COMPUTE AND STORE XOR XIR AND
1			STU NZB	COO16 BT2		BCO		X2R

T P	<u>s</u> N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
		BCO TP9	NZA RAU NZU	TP9 P0009 SCUL		CAT CAT		IS ARC TANGENT REQD?
1		CAT ATSV NAT	RAU FMP STU FMP FAD STU FSB FSB BMI RAC RAC	C0001 C0002 C0003 ONE C0004 ONE DELTN ATSV 0010 0024		NAT ATC ATC		COMP AND STORE TANGENT SQUARE AND SET NUMBER OF TERMS BY MAG
1		ATC	LDD STD STD STD STD	ONE C0022 C0023 C0024 C0025		SUM		INITIALIZE FOR SERIES
-		SUM	RAU FAD FAD STU FSB FMP FDV STU RAU FSB FMP FDV STU FMP	C0022 C0022 ONE C0026 ONE C0024 C0024 C0024 C0024 C0023 C0004 C0023 C0004				COMPUTE SERIES TERMS
1			FAD STU BMC	C0025 C0025 ARTAN		STEP		SUM AND STORE
-		STEP	SXC RAU FAD STU	0001 C0002 ONE C0022		SUM		STEP FOR NEXT TERM AND REPEAT
1		ARTAN	RAU	C0025				COMPUTE

Fig. 5-10 (Cont.)

	T P	<u>S</u> N	Location	Oper. Code	Data Address	T A G	Intr. Address	T A G	Remarks
			Docation	FMP FDV	C0003 C0004		AUUI 655		ARC TAN AND
	l			STU	C0005				STORE
				FDV FDV FDV	FOUR C0004 C0004				AND STORE
				STU FMP FMP	C0004 C0004 HALF				В
				STU RAU FMP	C0007 C0002 C0005				C
•,*		-		FDV FDV STU	000 00 00 00 00 00 00 00 00 00 00 00 00				AND D
	1			FAD FMP	COOO7 THREE				COMPUTE AND STORE
				FAD STU RAU	COOO6 COO05 COO07				ଭୁଠ
				FMP FAD FMP	TWO COOO6 COOO2				
				STU RAU FAD	C0008 C0007 C000 <u>3</u>				Ql
	_			FSB FDV STU	COOO6 ROOO1 COOO9				AND Q2
	1			RAU FMP	C0005 R0002				COMPUTE AND STORE
				STU RAU FSB	C0011 C0008 C0005				XQQ
				FMP FMP FAD	R0002 R0002 C0011				AND
				STU RAU FMP	C0012 R0002 R0002				X1Q COMPUTE AND STORE
				FMP STU	R0002 P0010		- -		3 POWRS OMEGA

m	q	· · · · · · · · · · · · · · · · · · ·	0	Data	T	Tueta	T	
T P	<u>s</u> N	Location	Oper. Code	Data Address	A G	Instr. Address	A G	Remarks
			RAU FAD FSB FSB FMP FAD FAD FSB STU	C0005 C0009 C0008 C0008 P0010 C0012 C0012 C0012 C0011 C0013				COMPUTE AND STORE X2Q
1			NZB	ADXQ		BVA		BRN LL OP
1		BVA	NZA	TPl		TPO		BRN BM TP
		TPO	LDD STD STD STD	ZERO COO31 COO32 COO33		ADXR		LT HAUNCH XOQ X1Q AND X2Q
1		TPl	LDD STD RAU FSB STU	C0011 C0031 C0011 C0012 C0032				LT HAUNCH XOQ XIQ
			RAU FAD FSB STU RAU FDV FAD FSB STU	C0013 C0011 C0012 C0012 C0033 C0011 FOUR C0013 C0012 C0034		XRIS		X2Q AND MOD X2Q
1		XRIS	RAU FMP FAD STU FAD FDV FSB FMP FMP FAD STU	R0002 C0014 C0015 C0015 R0002 C0014 R0002 R0002 C0016 C0016		ADXR		COMPUTE AND STORE FOR TYPE 1 X1R AND X2R
1		ADXR	RAU	C0031			 	SUM LT XKQ

T P	s N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks						
			FAD STU RAU FAD	C0014 C0014 C0032 C0015				AND XKR XOR						
			STU RAU	C0015 C0033				Xlr						
-			FAD STU	C0016 C0016		COL		AND X2R						
1		COL	NZB	ICO		ADXQ		BRN LL OP						
		ADXQ	RAU FAD STU RAU	C0011 C0014 C0011 C0012				SUM XKR AND XKQ XOQ						
			FAD STU RAU	COO15 COO12 COO13				XlQ						
_			FAD STU	C0016 C0013				AND X2Q						
	1	1							NZB	IFO		SCV		BRN LL OP
-				LDD STD STD RAU FSB	C0013 P0002 C0017 C0012 C0017				COMPUTE AND STORE FBA					
			STU STU RAU FSB FSB	P0003 C0018 C0011 C0018 C0018			4	G						
1			FSB STU STU	C0017 P0004 C0019		TAl		AND FAB						
1		TAL	NZA	COCA		SUB3		BRN BM TP						
-		SUB3	RAC LDD STD LDD STD	0018 ZERO C0022 ONE C0023				SET TERMS INITALIZE FOR LOG SERIES						
			STD RAU	C0024 C0004				COMPUTE						

T P	<u>S</u> N	Tassta	Oper.	Data	T A	Instr.	T A	
P	N	Location	Code FAD STU	Address ONE C0025	G	Address	G	Remarks RATIO FOR SERIES
l			FSB FDV STU FMP STU	TW0 C0025 C0026 C0026 C0028		LOOP		
1		LOOP	RAU FMP FAD STU BMC	C0024 C0026 C0022 C0022 C0CA		STA		COMPUTE SUM OF TERMS STORE LOG
		STA	RAU FAD STU RAU FMP FDV STU LDD STD SXC	C0023 TW0 C0027 C0023 C0024 C0028 C0027 C0024 C0027 C0023 0001		LOOP		STEP FOR NEXT TERM AND REPEAT
		COCA	RAU FMP STU RAU FDV FSB FMP FDV STU RAU FMP FMP	R0001 R0001 C0023 Q D005 OCT C0007 OCT C0023 C0024 P0010 R0002 R0002		Ð		COMPUTE AND STORE Q4 5 POWRS OF
1			STU FMP STU NZA	P0010 C0024 C0007 UV1		υvo		BETA AND X4Q BRN BM TP
		UVO	RAU FSB STU	C0011 C0014 C0011				RESTORE XOQ

T P	S N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
			RAU FSB STU RAU	C0012 C0015 C0012 C0013				XIQ
			FSB STU RAU FSB FMP FAD	C0016 C0013 C0008 C0009 THREE C0006				X2Q COMPUTE AND STORE
			FSB FMP FDV STU RAU FSB FMP FAD FAD FAD FMP FMP FAD FMP FMP FAD	C0005 P0010 R0002 C0010 C0013 C0012 THREE C0010 C0010 C0010 C0006 R0001 TW0 C0008 R0001 C0025				X3Q
			RAU FSB FDV FSB FMP FMP FAD STU	C0022 C0025 C0023 R0001 C0024 P0010 R0002 C0007 C0028				Q5 AND X5Q
1			RAU NZU	COO2O ANO		AEQO		BRN IF ALPH EQ O
1		AEQO	RAU	C0017		SX4R		do not
		ANO	RAU FDV FDV FSB FMP	C0017 C0014 DEC C0016 C0014				DIVIDE BY ZERO IN THESE CALCS

T	S		Oper.	Data	T A	Instr.	T A	
T P	<u>S</u> N	Location	Code	Address	1	Address	G	Remarks
1			FMP	DEC		SX4R	r r	
_ _		SX4R	FSB FMP STU RAU	C0010 HALF P0005 P0010				COMPUTE AND STORE UL TBA
			FDV STU FMP FSB FMP FDV FDV	R0002 P0010 C0017 C0028 R0001 TWELV R0002				FOUR POWRS OF BETA
			FDV FAD STU RAU FMP FSB	R0002 P0005 P0006 C0018 HALF P0005				DL TBA
			STU RAU FAD FMP FSB FMP FDV FDV FDV FSB FAD	P0007 C0017 C0018 P0010 C0007 R0002 TWELV R0002 R0002 P0006 P0007				UL TAB
l		UVI	STU RAU FSB STU RAU	P0008 C0011 C0014 C0011 C0020		IDEN		DL TAB RESTORE XOQ COMPUTE
1			FMP STU RAU STU	HALF COO2O ONE POOO9		XRS		GAMMA AND SET FOR DL OP
-		SCUL	RAU FAD FDV FSB FSB	C0014 C0011 FOUR C0016 C0034				COMPUTE AND STORE

T P	<u>S</u> N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
				1	ļ.,		-	
			FMP STU	HALF POOO5				UL TBA
			RAU	P0010				UI IDA
			FDV	R0002				
			STU	P0010				
			RAU	C0014				
			FAD	C0011				
			FMP	P0010				
			FSB	C0007				
			$\mathbf{F}\mathbf{M}\mathbf{P}$	ROOOl				
			FDV	TWELV				
			FDV	R0002				
			FDV	R0002				
1			FAD	P0005				
			STU	P0006				DL TBA
			LDD STD	P0005 P0007				UL TAB
			LDD	P0007				UL IND
			STD	P0008				DL TAB
			LDD	ZERO				SET FOR
			STD	P0009		IDEN		NO DL OP
1		IDEN	RAU	ROOOL				STORE
		102314	FSB	ONE				FOR
			BMI	SFT1		SFT2		IDENTIF
-		SFT1	RAL	ROOOL				
			SRT	0003				
			SLT	0002				
			STL	C0035		IDB		OMEGA
1		SFT2	RAL	R0001				
			SRT	0002	-			0.7
			SLT	.0002		TDD		OR
		מתד	STL	C0035		IDB		OMEGA
		IDB	RAU FSB	ROOO2 ONE				
			BMI	SFT3		SFT4		
		SFT3	RAL	R0002		₩± ± 4		
			SRT	0008				
			SLT	0004				
			STL	C0036		IDN		BETA
		SFT4	RAL	R0002				
			SRT	8000				
			SLT	0005				OR
-			STL	C 0036		IDN		BETA
1		IDN'	RAL	R0009				COMPOSE
		11/12		10009				

					T		Т	
Ţ	S N	Location	Oper. Code	Data Address	A	Instr. Address	A G	Remarks
			SRT ALO ALO ALO STL WR1 NZB	0006 8005 C0035 C0036 P0001 P0001 STEPN		SETB		AND STORE ID NUMBER BRN LL OP
1		SETB	RAB NZA	0001 N05		NOL		INIT LL OP BRN BM TP
1		NOL	LDD STD STD	ZERO P0008 R0009		STEPN		TYPE O LL N TO ZERO
1		NO5 GAO	LDD STD STD AXB LDD STD RAU FSB STU NZU STU	ZER0 P0002 P0003 P0008 0002 HALF R0009 ONE R0002 R0002 C0020 CT0 C0002		GAO XRS		TYPE 1 LL INITIAL VALUES STORAGE AND N IS HALF COMPUTE GAMMA IF ZERO LIM IS ZER
		CTO SR2 CTAL	RAU NZA FSB FDV STU	R0009 SR2 R0002 C0020 C0020		CTAL CTAL XRS		COMPUTE AND STORE UPPER LIM FOR CONST I
1		STEPN	RAU FAD STU FSB BMI	ROOO9 DELTN ROOO9 ONE THA		TAL		STEP N AND TEST IF LESS THAN ONE
1		THA	RAU FSB FSB BMI	ONE ROOO2 ROOO9 CTH		NEB		AND GRTR THAN ONE MIN BETA

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<u> </u>				t	<u> </u>	<u> </u>		·
Ţ P	<u>s</u> N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
-		NEB	RAU NZU	COO 20 CTO		СТН		
1		CTH	RAU FSB FSB FDV STU	R0002 ONE R0009 R0002 C0002		BCO		COMPUTE AND STORE UPPER LIM FOR VAR I
1		ICO	AXB RAU	0001 C0018				COMPUTE AND
			FSB FMP FAD	COO15 ROOO9 COO16				STORE
			STU RAU FAD FSB FMP	P0000 C0018 C0019 C0014 R0009	В			LL TBA
1			FAD FSB AXB STU	C0015 P0000 0001 P0000	B B	TESTB		AND LL TAB
1		TESTB	RAU SUP BMI	8006 SVEN STEPN		ADI		IF LAST VALUE IN POOO7
1		AD1	RAB	0001		IDEN		INIT B PCH
L		IFO	AXB RAU FSB FMP FAD STU RAU	0001 C0018 C0012 R0009 C0013 P0000 C0018	В			COMPUTE FOR VAR I AND STORE LL TBA
			FAD FSB FMP FAD FSB	C0019 C0011 R0009 C0012 P0000	В			
			AXB STU	0001 P0000	В	TESTB		AND LL TAB
1		TAL	RAU	R0002				TEST IF

TI P	S N	Location	Oper. Code	Data Address	T A G	Instr. Address	T A G	Remarks
			FAD STU RAU FSB BMI	R0004 R0002 R0006 R0002 TOM		ST1		BETA IS MAX IF NO STEP IF YES
1		том	RAU FAD STU RAU FSB BMI	R0001 R0003 R0001 R0005 R0001 READ		SIA		TEST IF OMEGA IS MAX AND READ OR
1		SIA	RAU STU	R0004 R0002		START		STEP BETA
1		ZERO ONE HALF NINES FOUR THREE TWO OCT DEC TWELV	CON 00 10 50 99 40 30 20 80 75 12	STANT 0000 0000 9999 0000 0000 0000 0000 00	S	0000 0051 0050 9999 0051 0051 0051 0051		
		DELTN SVEN FRAC	10 00 66	0000 0000 6666		0050 0007 6750		

CHAPTER VI

SUMMARY AND CONCLUSIONS

The development of a high speed computer program to evaluate beam constants for use in the carry-over moment procedure has been the purpose of this study.

The number of instructions in the program has been held to a minimum by expressing the beam constants in terms of recurring integrals which could be evaluated by supplying appropriate upper limits in equivalent algebraic expressions and by expressing succeeding beam constants in terms previously defined.

Using the program presented the computer will evaluate constants for beams with either one parabolic haunch or two symmetrical parabolic haunches for which β is expressed as a multiple of one-tenth and ω does not exceed two. If beam constants are desired for ω greater than two the number of terms in the subroutines for functional evaluation should be checked for accuracy.

The output from the computer may be listed on the IBM 402 Tabulator. Plate I shows typical results of such tabulations for a symmetrical beam for which

 $\omega = 1.0,$ $\beta = 0.1 \rightarrow 0.5$

From Chapter 5 and plate I the beam for which

$$\omega = 1.0,$$

β = 0.3

PLATE I

Tabulated Constants $\omega = 1.0, \beta = 0.1+0.5$

	and the second s						
1000100001	2905140250	1639384150	2905140250	4098460949	4135960949	4098460949	4135960949
1000106001	0000000000	0000000000	6178512349	617851 2 049	6315527449	5541497049	
1000109001	585 2 543049	4504481049	4689558049	3167466049	2726573049	1630451049	
1000200001	2526261050	1562787650	2526261050	3906969049	4040302349	3906969049	4040302349
1000206001	00000000000	0000000000	5964048549	5964049049	6067359749	5360737049	
1000209001	5570671049	4357426049	4373982049	3054115049	2422136049	1558226049	•
1000300001	2 188819950	1444752950	2188819950	3611882249	3874382249	3611882249	3874382249
1000306001	0000000000	0000000000	5606609249	5606609049	5663373649	5049844049	
1000309001	5120138049	4093080049	3896282049	2842467049	2102808049	1441250049	
1000400001	1884941050	1293156150	1884941050	3232890349	3632890349	3232890349	3632890349
1000406001	00000000000	0000000000	5106193849	5106194049	5111444549	4600943549	
1000409001	4526512049	3700761049	3371631049	2549814049	1806545049	1290088049	
1000500001	1606747650	1115873850	1606747650	2789684349	3310517649	2789684349	3310517649
1000506001	0000000000	00000000000	4462803549	4462803549	4424874449	4010423649	
1000509001	3863797249	3203876849	2857954049	2200587049	1532499049	1113038049	

is given by lines 7,8, and 9 where the identification numbers are of the form

100030n001,

and the coefficients (Fig. 6-1) are used with Eqs. (4-13) to (4-19) inclusive and related equations in chapter 3 to obtain the desired constants. The angular live load functions for

are obtained from symmetry.

Beam constants have been computed for all combinations of

$$\omega = 0.1 \rightarrow 2.0$$
$$\beta = 0.1 \rightarrow 1.0$$

for unsymmetrical beams and

$$\beta = 0.1 \rightarrow 0.5$$

for symmetrical beams. Comparisons were made with values presented in (2) and the results published in (3).

f _{BA}	g	f _{AB}	$t_{\rm BA}^{(\rm UL)}$	$t_{\rm BA}^{({ m DL})}$	$t_{AB}^{(UL)}$	$t_{AB}^{(DL)}$
.2189	.1445	.2189	.0361	.0387	.0361	.0387

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Ł	01	
L	aı	

n	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
$t_{\rm BA}^{({ m LL})}$.0144	.0284	.0409	₀0505	.0561	.0566	.0512	.0390	.0210
$t_{AB}^{(LL)}$.0210	.0390	.0512	.0566	.0561	₀0505	.0409	.0284	.0144

(Ъ)

Fig. 6-1.—Beam Constants For $\omega = 1.0$, $\beta = 0.3$.

Computations requiring the use of these constants are normally such that four place accuracy is sufficient.

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VITA

Therman Iveal Lassley

Candidate for the Degree of

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