# THE GENERAL STRING POLYGON 

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## THE GENERAL STRING POLYGON

 METHOD

## PREFACE

The material presented in this thesis is the outgrowth of the seminar lectures presented by Professor Jan J. Tuma in the Spring of 1960. The literature survey and the general theory recorded in the introduction were prepared by Professor Tuma.

The application of string polygon method to the analysis of single span rigid frames, with members of variable cross-section, was reported by John T. Oden.

The general theory of the string polygon, in terms of the energy due to bending moments, shearing forces, and normal forces, is presented in this thesis.

The writer wishes to express his indebtedness and gratitude to Professor Jan J. Tuma for his invaluable aid and guidance in preparing this thesis. The writer also expresses his appreciation to Professor Roger L. Flanders for his acting as the author's major adviser, and for his advice and thorough reading of the manuscript.

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Paul C. M. Wu

## TABLE OF CONTENTS

Chapter Page
I. INTRODUCTION ..... 1
II. THEORY OF GENERAL STRING POLYGON ..... 3

1. Basic Derivation ..... 3
2a. Angular Load Function due toBending Moment.15
2b. Angular Load Function due to Shear ..... 18
2c. Linear Load Function due to
Normal Force ..... 21
2d. Change in Slope of the StringPolygon . . . . . . . . . . . . . 232e. Linear Displacements of StringPolygon Influenced by NormalForces . . . . . . . . . . . . . 24
III. EIASTIC WEIGHTS ..... 26
IV. NUMERICAL EXAMPIES ..... 30
2. General Notes ..... 30
3. Example No. I ..... 31
4. Example No. 2 ..... 36
V. CONCLUSION ..... 42
A SELECTED BIBLIOGRAPHY ..... 44

## LIST OF TABIES

TablePage2-1. Interpretation of Angular Constants ..... 11
3-1. Relation of the end Condition BetweenReal Structure and the Corresponding
Conjugate Structures ..... 29
4-1. a. End Moments, end Shears and Normal
Forces ..... 32
b. Elastic Constants ..... 32
4-2. a. Elastic Constants ..... 37
b. Moments, Shearing Forces and Normal Forces ..... 38
c. Linear Deformations due to Normal Force ..... 39
Figure Page
2-1. Bent Member ..... 4
2-2. Moment, Shear, Normal Force Diagram due to Actual Loading ..... 5
2-3. Virtual Loads ..... 6
2-4. Shear and Bending Diagrams due to VirtualLoads8
2-5. Interpretation of $\tau^{(M)}$ jix ..... 15
2-6. Interpretation of $\tau^{(M)}$ jiy ..... 17
2-7. Interpretation of $\tau^{(V)}{ }_{j i x}$ ..... 18
2-8. Interpretation of $\tau^{(V)}$ jiy ..... 202-9. Interpretation of $\Omega_{1 j x}^{(N)}$. . . . . . . . . . 212-10. Interpretation of $\Omega^{(N)}$ ijy22
2-1l. Linear Displacement Under General Loading ..... 24
3-1. The Relation Between a Real Structure and the Corresponding Conjugate Structure ..... 27
4-1. Three Member Cantilever Beam ..... 31
4-2. Conjugate Structure ..... 34
4-3. Symmetrical Gabled Frame ..... 36

NOMENCLATURE
b Width of Beam
$d_{j} \quad$ Length of Bent Member ij
$\mathrm{d}_{\mathrm{jx}} \quad$ Horizontal Projection of $\mathrm{d}_{j}$
$\mathrm{d}_{j y} \quad$ Vertical Projection of $\mathrm{d}_{j}$
$v^{u}, u^{\prime}, \mathrm{x}, \mathrm{y}$, Coordinates of Cross-Section
BM Bending Moment due to Load
BV Shearing forces due to Loads
$F_{i j}, F_{j i}$ Angular Flexibilities
$G_{i j}, G_{j i}$ Carry-Over Value
$M_{i}, M_{j}, M_{k}$ End Moments
$N_{i j}, N_{j i}$,
${ }^{N}{ }_{j k}, \mathbb{N}_{k j} \quad$ End Thrusts
$V_{i j}, V_{j i}$,
$\nabla_{j k}, V_{k j}$ End Shears
$\overline{\mathrm{M}}$ Moment of the Elastic Weights
$\bar{P} \quad$ Elastic Weight
BN Normal Force due to Loads
$\tau$ Angular Load Function
$\pi \quad$ Angle That a Bent Member Makes With the Horizontal
$P_{i} \quad \underset{j k}{\text { Angle Between Extensions of Bent Members ij and }}$ $\phi_{j} \quad$ Change in Change of Slope of Bent Line ijk

| $\Delta_{i j}$ | Linear Deformation of Member ij |
| :--- | :--- |
| $\mathscr{H}$ | Shape Factor |



## CHAPTER I

## INTRODUCTION

The idea of elastic weight and the application of the elastic weights was introduced in the middle of the last century by Otto Mohr (1). The extension of the application of elastic weights and a methodical classification of elastic weights was performed by Mỉler Breslau (2), (3). The study of deformation of beams by means of elastic weights was extensively presented by Wanke (4) and Chmelka (5). The development of the joint elastic weights, in terms of end moments for strips of small length, may be found in work of Kaufmann (6).

In this country, the application of finite elastic weights was shown by Hardy Cross as his Column Analogy (7) and by Michalos as the Column, Shear and Torsion Analogy (8).

The generalization of the joint elastic weight expression and the application of these joint elastic weights, in connection with the string polygon, was developed by Tuma (9) and extended by his students, Chu (10), Oden (11), and Boecker (12), to the solution of many special problems.

The application of the string polygon method requires calculation of angular constants, which are now available
in various publications. (13) (14).
In this thesis, the effort has been made to derive the general expressions for the elastic weights in terms of the bending moments, shearing forces, and normal forces. This leads to the representation of the elastic weight as a vector force and vector moment. This elastic weight is then applied to the conjugate structure. The shear of the conjugate structure is equal to the slope of the real structure and the bending moment of the conjugate structure is the deflection of the real structure along the line of the vector bending moment. The application is illustrated by two examples.

The nomenclature is assembled in the front part of this thesis.

The sign convention of statics is used in formation of equilibrium conditions and elasto-static equations.

The sign convention of deformation is used for the calculation of cross-section elements. The signs of vectors are governed by the right hand rule.

## CHAPTER II

## THEORY OF GENERAL STRING POLYGON

The general string polygon theory for bent members is developed in this chapter. All the influences of the bending moments, shearing forces and normal forces are considered.

## 2-1 Basic Derivation

A bent member, ijk, loaded by a general system of loads is considered (Figure 2-1). Whe cross-section of the member ij(jk) is given by ordinates $u, u^{\prime}\left(v, v^{\prime}\right)$ measured from the respective ends. The cross-sectional elements at a given section are:
bending moments $M_{u}\left(M_{v}\right)$,
shearing forces $\mathrm{V}_{\mathrm{u}}\left(\mathrm{V}_{\mathrm{v}}\right)$, and
normal forces $N_{u}\left(N_{v}\right)$.
The geometry of each member is given by the slope $\pi$ and the length $d$. The horizontal projection of each $d$ is $d_{x}$ and the vertical projection of the same length is $d_{y}$. Due to the action of loads, the bent member ijk displaces to the position $i^{\prime} j^{\prime} k$ ' and the change in change of the slope at $j$ is denoted by $\varnothing_{j}$ as shown in Figure 2-1. The


Figure 2-1. Bent Member
calculation of the $\emptyset_{j}$ is accomplished by means of the virtual work.

The moments, shearing forces and normal forces at the section, due to the loads, are shown in Figure 2-2.

The end moments, end shears and end thrusts are designated as:

$$
M_{i}, M_{j}, M_{k} \quad=\text { end moments }
$$

$V_{i j}, V_{j i}, V_{j k}, V_{k j}=$ end shears
$N_{i j}, N_{j i}, N_{j k}, N_{k j}=$ end thrusts.
The bending moment at the section $u$ of the bent member ij is:

$$
M_{u}=M_{i} \frac{w}{\alpha_{j}}+M_{j} \frac{u}{d_{j}}+B M_{u}
$$

and at section $v$ of the bent member $j k$ is:

$$
\begin{equation*}
M_{v}=M_{j} \frac{v^{\prime}}{d_{k}}+M_{k} \frac{v}{d_{k}}+B M_{v} \tag{2-1b}
\end{equation*}
$$



Figure 2-2. Moment, Shear, Normal Diagram due to Actual Loading
where $\mathrm{BM}_{u}$ and $\mathrm{BM}_{v}$ are the bending moments due to loads at the section $u$ and $v$ respectively.

The shearing force at the section $u$ is;

$$
\begin{equation*}
v_{u}=V_{i}+B V_{u} \tag{2-2a}
\end{equation*}
$$

and, at section v is:

$$
\begin{equation*}
V_{v}=V_{k}+B V_{v}, \tag{2-2b}
\end{equation*}
$$

where $\mathrm{BV}_{u}{ }^{B V} V_{v}$, are the shearing forces due to loads on the segment $u$ and $v^{\prime}$ respectively.

The normal force at the section $u$ 1s:

$$
\begin{equation*}
N_{u}=N_{i}+B N_{u} \tag{2-3a}
\end{equation*}
$$

and at section $v$ is:

$$
\begin{equation*}
N_{v}=N_{k}+B N_{v}, \tag{2-3b}
\end{equation*}
$$

where $B N_{u} u^{B N}{ }_{v}$, are the normal forces due to loads on the segment $u$ and $v$ ' respectively.

For the purpose of determining $\varnothing_{f}$, the virtual loads

$$
\frac{1}{d_{j}} \text { and } \frac{1}{d_{k}}
$$

are applied on the member ijk as shown at the Figure 2-3.


The normal force at $u$ and $v$ due to these virtual loads is equal to zero, which indicates that the normal force has no direct influence on the formation of $\varnothing_{j}$. On the other hand, the shearing force and bending moment do influence the formation of $\varnothing_{j}$ and their diagrams are shown in Figure 2-4。

The bending moment at section $u$ due to the virtual loads is:

$$
\begin{equation*}
\left(M_{u}\right)=\frac{u}{d_{j}} \tag{2-4a}
\end{equation*}
$$

and at section $v$ is:

$$
\begin{equation*}
\left(M_{v}\right)=\frac{v^{\prime}}{a_{k}} \tag{2-4b}
\end{equation*}
$$

And, the shearing force at section $u$ due to the virtual loads is:

$$
\left(v_{u}\right)=\frac{1}{d_{j}}
$$

and, at section $v$ is:

$$
\begin{equation*}
\left(v_{v}\right)=-\frac{1}{d_{k}} \tag{2-5b}
\end{equation*}
$$

From the theory of virtual work, the change in slope due to bending moments and shearing forces is:


Shear Diagram Figure 2-4. Shearing Bending Diagrams due to Virtual Loads

$$
\begin{align*}
\varnothing_{j}= & \int_{i}^{j} \frac{M_{u}\left(M_{u}\right) d u}{E I}+\chi_{u} \int_{i}^{j} \frac{v_{u}\left(v_{u}\right) d u}{A_{u}^{G}}+\int_{j}^{k} \frac{M_{v}\left(M_{v}\right) d v}{E I_{v}} \\
& +\int_{j}^{k} \frac{v_{v}\left(v_{v}\right) d v}{A_{v} G} \tag{2-6}
\end{align*}
$$

In terms of equations $2-4$ and $2-5$, the equation $2-6$ will become:

$$
\begin{align*}
& \phi_{j}=M_{i} \int_{i}^{j} \frac{u^{\prime} u d u}{d_{j}{ }^{2} E I_{u}}+M_{j}\left(\int_{i}^{j} \frac{u^{2} d u}{d_{j}{ }^{2} E I_{u}}+\int_{j}^{k_{k}} \frac{v^{\prime 2} \frac{d u}{d_{k}{ }^{2} E I}}{v}\right) \\
& +M_{k} \int_{j}^{e_{k}} \frac{v^{\prime} v d v}{d_{k}{ }^{2} E I}+\int_{i}^{j} \frac{B M_{u} u d u}{d_{j}{ }^{E I} u}+\int_{j}^{e^{k}} \frac{B v^{v} v^{\prime} d v}{d_{k} E I v} \\
& +v_{i} \notint_{i}^{j} \frac{d u}{d_{j}{ }^{A} u^{G}}-v_{k} \chi_{6}^{\int_{j}^{k}} \frac{d v}{d_{k^{A} v^{G}}}+\mu_{i}^{\rho} \frac{B v_{u} d u}{d_{j}{ }^{A} u^{G}} \\
& -\gamma \int_{j}^{k} \frac{B v_{v}{ }^{\prime} v}{d_{k} V^{G}}  \tag{2-7a}\\
& \text { or, } \varnothing_{j}=M_{i} G_{i j}{ }^{(M)}+M_{j} \Sigma F_{j}{ }^{(M)}+M_{k} G_{k j}{ }^{(M)}+\Sigma \tau_{j}^{(M)}+V_{i} G_{i j}{ }^{(v)} \\
& +V_{k} G_{k j}{ }^{(v)}+\Sigma \tau_{j}{ }^{(v)} . \tag{2-7b}
\end{align*}
$$

The angular constants in equation (2-7) can be interpreted in terms of a simple beam as shown in Table 2-1. The normal force will cause a linear elongation or contraction of each member. These deformations are:

$$
\begin{align*}
& \Delta_{i j}(\mathbb{N})=\int_{i}^{j} \frac{N_{u} d u}{E A_{u}} \\
&=\int_{i}^{j} \frac{\mathbb{N}_{i j} d u}{E A_{u}}+\int_{i}^{\frac{B N}{u} d u}  \tag{2-8a}\\
& E A_{u}
\end{align*}
$$

and:

$$
\begin{align*}
\Delta_{j k}(N) & =\int_{j}^{k} \frac{N_{v} d v}{E A_{v}} \\
& =\int_{j}^{k} \frac{N_{j k} d v}{E A_{v}}+\int_{j}^{\frac{k N_{v} d v}{E A_{v}}} \tag{2-8b}
\end{align*}
$$

where $\mathbb{N}_{i j}, \mathbb{N}_{j k}$ are end thrust at $i$ and $j$ respectively,
$B N_{u}$ is the normal components of loads on the segment $u$, $B N_{v}$ is the normal components of loads on the segment $v$.

## TABLE 2-1

INTERPRETATION OF ANGULAR CONSTANTS

| Term | Name | Value | Physical Meaning | Illustration |
| :---: | :---: | :---: | :---: | :---: |
| $F_{j i}(M)$ | The angular flexibility due to moment. | $\int_{i}^{j} \frac{u^{2} d u}{d_{j}{ }^{2} E I_{u}}$ | The end slope of a simple beam ij at $j$ due to a unit moment applied at that end. |  |
| $F_{j k}^{(M)}$ | The angular flexibility due to moment. | $\int_{j}^{k} \frac{v^{2} d v}{d^{2} k^{E I} v}$ | The end slope of a simple beam jk at $j$ due to a unit moment applied at that end. |  |
| $G_{i j}(M)$ | The angular carry over value due to moment. | $\int_{i}^{j} \frac{u u^{\prime} d u}{d^{2} j^{E I} u}$ | The end slope of a simple beam ij at $j$ due もo a unit moment applied at the far end $i$. |  |

TABLE 2-1 (Continued)

| Term | Name | Value | Physical Meaning | Illustration |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{G}_{\mathrm{kj}}$ (M) | The angular carry over value due to moment. | $\int_{j}^{k} \frac{v v^{\prime} d v}{d^{2} k^{E I} v}$ | The end slope of a simple beam jk at $j$ due to a unit moment applied at the far end $k$. |  |
| $\tau_{j k}(M)$ | The angular load function due to moment. | $\int_{i}^{j} \frac{B M}{{ }^{j} u d u}{ }_{\alpha_{j} E I} u$ | The end slope of a simple beam ij at j due to a moment influence of the loads. |  |
| $\tau_{j k}(M)$ | The angular load function due to moment. | $\int_{j}^{k} \frac{B M v^{\prime} v^{\prime} d v}{d_{k} E I v}$ | The end slope of a simple beam jk at $j$ due to a moment influence of the loads. |  |

TABLE 2-1 (Continued)

| Term | Name | Value | Physical Meaning | Illustration |
| :---: | :---: | :---: | :---: | :---: |
| $G_{i j}(\mathrm{~V})$ | The angular carry over value due to shear. | $\not \partial \int_{i}^{j} \frac{d u}{d_{j}{ }^{A} u^{G}}$ | The angular slope of a simple beam ij at $j$ due to a unit shearing force applied at far endi. |  |
| $\mathrm{G}_{\mathrm{kj}}$ (V) | The angular carry over value due to shear. | $-26 \int_{j}^{k} \frac{d v}{d_{k} A_{i} G}$ | The angular slope a simple beam jk at $j$ due to a unit shearing force applied at far end k . |  |
| $\tau_{j i}(\mathrm{~V})$ | The angular load function due to shear. | $\int_{i}^{j} \frac{B V_{u} d u}{d_{j}^{A} u^{G}}$ | The end slope of a simple beam ij at $j$ due to a shearing influence of the loads. |  |

TABIE 2-1 (Continued)

| Term | Name | Value | Physical Meaning | Illustration |
| :---: | :---: | :---: | :---: | :---: |
| $\tau_{j k}(V)$ | The angular load function due to shear. | $-\notint_{j}^{k} \frac{B v_{v} d v}{d_{k^{A} v^{G}}}$ | The end slope of a simple beam jk at $j$ due to a shearing influence of the loads. |  |

The physical interpretation of equations (2-8a) and (2-8b) is self-evident and does not need to be explained. 2-2a Angular Load Function due to Bending Moment $\tau(M)$

Consider the segment jj of the member ijk loaded only by a system of vertical loads (Figure 2-5); it is desirable to evaluate $\boldsymbol{T}_{j i}{ }^{(M)}$ in terms of horizontal or vertical coordinates, since loads are usually applied in these directions. It is necessary to imagine the horizontal projection of member id as i'j' as shown in the Figure 2-5. (M) may be defined as the slope of the simple beam i'j' at $j '$ due to loads, $i ' j '$ being the horizontal projection of $i j$.


Figure 2-5. Interpretation of $\tau_{\text {jik }}$

If the unloaded member id is naturally inclined at an angle $\pi_{j}$ as shown in Figure 2-5, it follows that:

$$
\begin{equation*}
d u=\frac{d x}{\cos \pi_{j}} \tag{2-9}
\end{equation*}
$$

$\ldots$... and defining $\tau_{j i x}{ }^{(M)}$ by

$$
\begin{equation*}
\tau_{j i x}^{(M)}=\int_{i}^{j} \frac{B M_{x} x d x}{d_{j x} I_{u}} \tag{2-10}
\end{equation*}
$$

it is seen that

$$
\tau_{j i}^{(M)}=\int_{i}^{j} \frac{B M_{u} u d u}{d_{j}{ }^{E I} u}=\frac{I}{\operatorname{CoS} \pi_{j}} \int_{i}^{j} \frac{B M_{x} x d x}{d_{j x}^{E I} u} \quad(2-1 l a)
$$

or

$$
\begin{equation*}
\tau_{j i}(M)=\frac{1}{\operatorname{COS} \pi_{j}} \tau_{j i x} . \tag{2-11b}
\end{equation*}
$$

In a similar manner, the angular load-functions for if due to the action of horizontal loads only may be evaluated. $\tau_{i j y}{ }^{(M)}$ is defined as the slope of the simple beam $i " j "$ at $j "$ due to loads where $i " j "$ is the vertical projection of id as shown in Figure 2-6.


Figure 2-6. Interpretation of $\tau_{j i y}$ (M)
It is seen that

$$
\begin{equation*}
\tau_{j i y}(M)=\int_{i}^{j} \frac{B M y d y}{d_{j y} y^{E I}} \tag{2-12}
\end{equation*}
$$

Since $d u=\frac{d y}{\operatorname{Sin} \pi_{j}}$ it follows that

$$
\begin{equation*}
\tau_{j i}^{(M)}=\frac{1}{\operatorname{Sin} \pi_{j}} \int_{i}^{\frac{B M}{y^{\prime} y d y}} \frac{y^{E I}}{{ }^{E I}} \tag{2-13a}
\end{equation*}
$$

or

$$
\begin{equation*}
\tau_{j i}(M)=\frac{1}{\sin \pi_{j}} \tau_{j i y} \tag{2-13b}
\end{equation*}
$$

## 2-2b Angular Load Function due to Shear $\tau^{(y)}$

By the similar manner, consider the segment ij of the member ijk loaded only by a system of vertical loads as shown in Figure 2-7.

where

$$
\begin{aligned}
& B W_{x}=\text { vertical loads on the segment } x \\
& B V_{u}=\text { perpendicular component of the loads on the }
\end{aligned}
$$

$$
\begin{aligned}
& \text { and } \quad B V_{u}=B W_{x} \cos \pi_{j} \\
& d u=\frac{d x}{\cos \pi_{j}}
\end{aligned}
$$

$$
\tau_{j i x}^{(V)}=\psi_{i}^{j} \frac{B W_{x} d x}{d_{j x} A u G}
$$

It is seen that

$$
\tau_{j i}^{(v)}=\nVdash \int_{i}^{j} \frac{B V_{u} d u}{\partial_{j} A u G}=\psi_{6} \cos \pi_{j} \int_{i}^{j} \frac{B W_{x} d x}{\partial j x A u G}(2-14 a)
$$

or

$$
\begin{equation*}
\tau_{j i}^{(v)}=\tau_{j i x}{ }^{(v)} \cos \pi_{j} \tag{2-14b}
\end{equation*}
$$

also $\quad \tau_{j k}(v)=\tau_{j k x}(v)_{\cos \pi_{k}}$.
$\tau_{i j y}{ }^{(v)}$ is defined as the slope of simple beam $i^{\prime \prime}{ }^{\prime \prime}$ at $j$ " due to horizontal loads where $i$ "j" is the vertical projection of jj.


Figure 2-8. Interpretation of $\tau^{(V)}{ }_{\text {jiy }}$
where $\quad \mathrm{BW}_{\mathrm{y}}=$ Horizontal loads on the segment J

$$
\begin{aligned}
B V_{u}= & \text { Perpendicular loads on the segment } u \text { of } \\
& \text { the member } i j
\end{aligned}
$$

and

$$
B V_{u}=B W_{y} \operatorname{Sin} \pi_{j}
$$

$$
d u=\frac{d y}{\operatorname{Sin} \pi_{j}}
$$

Therefore, $\tau_{j i}(v)=\nLeftarrow \int_{i}^{j} \frac{B V_{u} d u}{d_{j} A u G}=\gamma \sin \pi_{j} \int_{i}^{j_{i}} \frac{B W_{y} d y}{d_{j y} A u G}(2-16 a)$
or

$$
\begin{align*}
& \tau_{j i}(v)=\tau_{j i y}{ }^{(v)_{\operatorname{Sin} \pi_{j}}}  \tag{2-16b}\\
& \tau_{j k}(v)=\tau_{j k y}{ }^{(v)_{\operatorname{Sin} \pi_{k}}} \tag{2-17}
\end{align*}
$$

## 2-2c Linear Load Function due to Normal Force $\Omega_{i j}$

Consider the bent member ij loaded only by a system of vertical load as shown in Figure 2-9.


The linear load function

$$
\begin{align*}
\Omega_{i j}(N) & =\int_{i}^{j} \frac{B N_{u} d u}{E A_{u}}=\int_{i}^{j} \frac{B W_{x} \sin \pi_{j} \frac{d x}{\cos \pi_{j}}}{E A_{u}} \\
& =\tan \pi_{j} \int_{i}^{\frac{B W_{x} d x}{E A_{u}}} \tag{2-18}
\end{align*}
$$

where $\quad B N{ }_{x}=$ Vertical loads on the segment $x$

$$
B N_{u}=\mathbb{N o r m a l ~ l o a d s ~ o n ~ t h e ~ s e g m e n t ~} u \text { of the member }
$$

ij
and

$$
\begin{aligned}
B N_{u} & =B W_{x} \operatorname{Sin} \pi_{j} \\
d u & =\frac{d x}{\cos \pi_{j}}
\end{aligned}
$$

In a similar manner, consider the bent member ij loaded only by a system of horizontal loads as shown in Figure 2-10.


$$
\begin{align*}
& \text { The linear load function } \\
& \begin{aligned}
\Omega_{i j}(N) & =\int_{i}^{j} \cdot \frac{B N_{u} d u}{E A} u
\end{aligned} \int_{i}^{j} \frac{B W_{y} \cos \pi_{j} \frac{d y}{\operatorname{Sin} \pi_{j}}}{E A_{u}} \\
&  \tag{2-19}\\
& =\cot \int_{y}^{j} \frac{\operatorname{BW}_{i} d y}{E A} u
\end{align*}
$$

where $\quad B W_{\mathbf{y}}=$ Horizontal loads on the segment $y$

$$
\begin{aligned}
& B N_{u}=\text { Normal loads on the segment } u \text { of the member } \\
& i j \\
& B N_{u}=B W_{y} \cos \pi_{j} \\
& d u=\frac{d y}{\operatorname{Sin} \pi_{j}}
\end{aligned}
$$

$\frac{\text { 2-2d Change in Slope of the String Polygon } \emptyset_{j} \text { for Vertical }}{\text { and Horizontal Loads }}$

The change in slope of the string polygon due to the vertical loads becomes:

$$
\begin{align*}
\emptyset_{j} & =M_{i} G_{i j}(M)+M_{j} \Sigma F_{j}(M)+M_{k} G_{k j}(M)+\frac{\tau_{j i x}(M)}{\cos \pi_{j}} \\
& +\frac{\tau_{j k x}(M)}{\cos \pi_{k}}+V_{i} G_{i j}(V)+V_{k x} G_{k j}(V)+\tau_{j i x} \cos \pi_{j} \\
& +\tau_{j k x} \cos \pi_{k} \tag{2-20}
\end{align*}
$$

And the change in slope of the string polygon due to the horizontal loads becomes:

$$
\begin{align*}
\emptyset_{j} & =M_{i} G_{i j}(M)+M_{j} \Sigma F \\
j & (M)+M_{k} G_{k j}(M)+\frac{\tau_{j i y}}{\operatorname{Sin} \pi_{j}}+\frac{\tau_{j k y}}{\operatorname{Sin} \pi_{k}}  \tag{2-21}\\
& +V_{i} G_{i j}(V)+V_{k} G_{k j}(V)+\tau_{j i y} \operatorname{Sin} \pi_{j}+\tau_{j k y} \operatorname{Sin} \pi_{k}
\end{align*}
$$

2-2e Linear Displacements of String Polygon Influenced by Normal Forces


Figure 2-11. Linear Displacement Under General Loading

The defined linear displacement due to the influence of the normal forces, which the general loading condition of vertical and horizontal is considered in Figure 2-11.

$$
\Delta_{i j x}{ }^{(\mathbb{N})}=\Delta_{i j}{ }^{(\mathbb{N})} \cos \pi_{j} \quad \mid \quad \Delta_{i j y}{ }^{(\mathbb{N})}=\Delta_{i j}{ }^{(\mathbb{N})} \operatorname{Sin}_{j}
$$

or

$$
\Delta_{i j x}=N_{i j} \int_{i}^{j} \frac{d x}{A_{u} E}+\sin \pi_{j} \int_{i}^{j} \frac{B W_{x} d x}{A_{u} E}+\cos \pi_{j} \cot \pi_{j}
$$

$$
\begin{equation*}
\int_{i}^{j} \frac{B W_{y} d y}{A_{u} E} \tag{2-23}
\end{equation*}
$$

and

$$
\begin{align*}
\Delta_{i j y} & =N_{i j} \int_{i}^{j} \frac{d y}{A_{u} E}+\sin _{j}^{\pi} \tan ^{\pi} \int_{i}^{j} \frac{B W_{x} d x}{A_{u}^{E}}+\cos \pi_{j} \\
& \int_{i}^{\frac{B W_{y} d y}{A_{u} E}} \tag{2-24}
\end{align*}
$$

## CHAPTER III

## ELASTIC WEIGHPS

The change in change of slope $\varnothing_{j}$ of bent line ijk of Figure l-l at $j$ is an angular deformation at $j$ and it can be treated as a vector force normal to the plane of the bent member:

$$
\begin{equation*}
\phi_{j}=\bar{P}_{j} \tag{3-1}
\end{equation*}
$$

The linear deformations $\Delta_{i j}$ and $\Delta_{j k}$ can be represented as vector moments acting at $i$ and $j, ~ r e s p e c t i v e l y$.

$$
\begin{equation*}
\Delta_{i j}=\bar{M}_{i j} \tag{3-2}
\end{equation*}
$$

$$
\Delta_{j k}=\bar{M}_{j k}
$$

The angular and linear deformations can be treated as elastic weights applied on the conjugate structure and used in elasto-static equations for the calculation of deformations.

The relationship between a real structure and a conjugate structure is shown in Figure 3-1.


Figure 3-1. The Relation Between a Real Structure and the Corresponding Conjugate Structure

From the relationship of the real structure to the conjugate structure, the following analogies may be stated:

1. The shear of the conjugate structure at any point represents the slope of the real structure at the same point.
2. The bending moment of the conjugate structure at the certain point about a given axis is the displacement of the real structure along that axis.

This two laws hold for any polygon of any set of members without regard to its end conditions and type of loading. The end conditions of the real structure and the end conditions of the corresponding conjugate structure are related to each other as shown in Table 3-1.

The application of the string polygon equations (2-7) and (2-8) as elastic weights defined by equations (3-1) and (3-2) to a solution of a specific problem is shown in the last part of this thesis.

## TABLE 3-1 <br> RELATION OF THE END CONDITION BETWEEN REAL STRUCTURE AND THE CORRESPONDING CONJUGATE STRUCTURE



## CHAPTER IV

## NUMERICAL EXAMPLES

4-1. General Notes

The following illustrative examples comprise this chapter of the thesis and describe the numerical procedure of analysis by the general string polygon method. Dimensions are chosen to simplify calculation and units for various values are in terms of Kips, feet and Kip-feet.

## 4-2. Example No. 1

A three-member cantilever beam of constant cross section is considered. (See Figure 4-1.) The angular and linear deflections at the free end 1 are required in terms of modulus of elasticities, $E, G$, moment of inertia $I$, and area of the constant cross section $A$. The shape factor $\mathcal{H}$ is assumed to be 1.


Figure 4-1. Three-Member Cantilever Beam
(a) End moments, end shears and end normal forces:

| Evaluating joint | End moment$(k-f t)$ |  |  | End shearing force(kip) |  | End normal force(kip) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (j) | $M_{i}$ | $M_{j}$ | $\mathrm{M}_{\mathrm{k}}$ | $\mathrm{V}_{i j}$ | $\mathrm{V}_{\mathrm{kj}}$ | $\mathrm{N}_{\text {ij }}$ | $\mathrm{N}_{\mathrm{kj}}$ |
| 1 | 0 | 0 | - 18 | 0 | $-\frac{18}{5}$ | 0 | $+\frac{24}{5}$ |
| 2 | 0 | - 18 | - 98 | 0 | $-\frac{56}{5}$ | 0 | $+\frac{42}{5}$ |
| 3 | -18 | - 98 | -288 | - $\frac{24}{5}$ | - 24 | $+\frac{18}{5}$ | 0 |
| 4 | -98 | -288 | 0 | - 14 | 0 | 0 | 0 |

(b) Elastic constants

| Evaluating | $G^{M}(E I)$ |  | $\mathrm{F}^{\mathrm{M}}$ (EI) |  | $\tau^{M}$ (EI) |  | $G^{V}(A G)$ |  | $\tau^{\mathrm{V}}$ (AG) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (j) | ij | kj | ji | jk | ji | jk | ij | kj | ji | jk |
| 1 | 0 | $\frac{5}{3}$ | 0 | $\frac{10}{3}$ | 0 | $\frac{45}{3}$ | 0 | -1 | 0 | $-\frac{9}{5}$ |
| 2 | $\frac{5}{3}$ | $\frac{5}{3}$ | $\frac{10}{3}$ | $\frac{10}{3}$ | $\frac{45}{3}$ | $\frac{80}{3}$ | 1 | -1 | - $\frac{9}{5}$ | $-\frac{16}{5}$ |
| 3 | $\frac{5}{3}$ | $\frac{5}{3}$ | $\frac{10}{3}$ | $\frac{10}{3}$ | $\frac{80}{3}$ | $\frac{125}{3}$ | 1 | -1 | $-\frac{16}{5}$ | $-\frac{25}{5}$ |
| 4 | $\frac{5}{3}$ | 0 | $\frac{10}{3}$ | 0 | $\frac{125}{3}$ | 0 | 1 | 0 | - $\frac{25}{5}$ | 0 |

(c) Elastic weights

$$
\begin{aligned}
& \overline{\mathrm{P}}_{1}=\phi_{1}=(-18)\left(\frac{5}{3 \mathrm{EI}}\right)+\frac{45}{3 \mathrm{EI}}+\left(-\frac{18}{5}\right)\left(-\frac{1}{\mathrm{AG}}\right)-\frac{9}{5 \mathrm{AG}}=-\frac{45}{3 \mathrm{EI}}+\frac{9}{5 \mathrm{AG}} \\
& \stackrel{\rightharpoonup}{P}_{2}=\phi_{2}=(-18)\left(\frac{20}{3 \mathrm{EI}}\right)+(-98)\left(\frac{5}{3 \mathrm{EI}}\right)+\frac{125}{3 \mathrm{EI}}+\left(\frac{-56}{5}\right)\left(-\frac{1}{\mathrm{AG}}\right)-\frac{25}{5 \mathrm{AG}} \\
& =\frac{-725}{3 E I}+\frac{31}{5 A G} \\
& \overline{\mathrm{P}}_{3}=\phi_{3}=(-18)\left(\frac{5}{3 \mathrm{EI}}\right)+(-98)\left(\frac{20}{3 \mathrm{EI}}\right)+(-288)\left(\frac{5}{3 \mathrm{EI}}\right)+\frac{205}{3 \mathrm{EI}} \\
& +\left(-\frac{24}{5}\right)\left(\frac{1}{A G}\right)+(-24)\left(-\frac{1}{A G}\right)-\frac{41}{5 A G}=-\frac{3285}{3 E I}+\frac{55}{5 A G} \\
& \bar{P}_{4}=\phi_{4}=(-98)\left(\frac{5}{3 E I}\right)+(-288)\left(\frac{10}{3 \mathrm{EI}}\right)+\frac{125}{3 \mathrm{EI}}+(-14)\left(\frac{1}{\mathrm{AG}}\right)-\frac{25}{5 \mathrm{AG}} \\
& =-\frac{3245}{3 E I}-\frac{95}{5 A G} \\
& \bar{M}^{N}{ }_{12}=\Delta_{12}=\tan \pi \int_{0}^{6} \frac{B W_{x} d x}{E A}=\frac{8}{6} \int_{0}^{6} \frac{x d x}{E A}=\frac{8}{6 E A}\left|\frac{x}{2}\right|_{0}^{6}=\frac{8 x 18}{6 E A}=\frac{24}{E A} \\
& \Delta_{12 x}=\frac{24}{E A} \times \frac{3}{5}=\frac{72}{5 \mathrm{EA}} \\
& \Delta_{12 \pm}=\frac{24}{E A} \times \frac{4}{5}=\frac{96}{5 E A}
\end{aligned}
$$

$$
\begin{aligned}
& \overline{\mathrm{M}}_{23}^{\mathbb{N}}=\Delta_{23}=\left(\frac{18}{5}\right) \int_{0}^{10} \frac{d u}{E A}+\tan \pi_{2} \int_{0}^{8} \frac{x d x}{E A} \\
&= \frac{180}{5 E A}+\frac{6 \times 32}{8 E A}=\frac{36}{E A}+\frac{24}{E A}=\frac{60}{E A} \\
& \Delta_{23 x}=\frac{60}{E A} \times \frac{4}{5}=\frac{240}{5 E A} \\
& \Delta_{23 y}=\frac{60}{E A} \times \frac{3}{5}=\frac{180}{5 E A}
\end{aligned}
$$

## (d) Conjugate system



Figure 4-2. Conjugate Structure
(e) Deformations at end 1

$$
\begin{aligned}
\phi_{14} & =\Sigma \bar{P}=\frac{-\frac{7300}{3 E I}-\frac{9}{5 A G}}{\Delta_{14 \mathrm{X}}}
\end{aligned}=\Sigma\left(\begin{array}{c}
\text { Angular deformation at } \\
\text { end 1 } \\
\text { end } 4
\end{array}\right)
$$

$$
\Delta_{l \Delta y}=\Sigma \bar{P}_{x}+\Sigma \bar{M}_{y}
$$

$$
=\left(-\frac{725}{3 E I}+\frac{31}{5 \mathrm{AG}}\right)(6)+\left(-\frac{3285}{3 \mathrm{EI}}+\frac{55}{5 \mathrm{AG}}\right)(14)
$$

$$
+\left(-\frac{3245}{3 \mathrm{EI}}-\frac{95}{5 \mathrm{AG}}\right)(24)-\frac{96}{5 \mathrm{EA}}-\frac{180}{5 \mathrm{EA}}=-\frac{138145}{3 \mathrm{EI}}
$$

$$
-\frac{1324}{5 A G}-\frac{276}{5 E A}
$$

## 4-3. Example 2.

Find the reactions at $A$ and $E$ of the symmetrical gabled frame shown in Figure 4-3.


Figure 4-3. Symmetrical Gabled Frame

## A. Elastic Constants

| Joint | $\mathrm{G}^{(\mathrm{M})}{ }_{i j} \mathrm{E}$ | $\Sigma F^{(M)}{ }_{j}{ }^{\text {E }}$ | $\mathrm{G}^{(M)} \mathrm{kj}^{\mathrm{E}}$ | $\mathrm{G}^{(V)}{ }_{i j}{ }^{\mathrm{E}}$ | $G^{(V)}{ }_{k j}{ }^{\text {E }}$ | $\Sigma \tau^{(M)}{ }_{j} \mathrm{E}$ | $\tau^{(V)}{ }_{j}{ }^{\text {E }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | - | 55.94 | 14.86 | - | -1.79 | - | - |
| B | 14.86 | 20.14 | 4.86 | 1.79 | -1. 22 | 163.50 | $-4.00$ |
| C | 4.86 | 16.66 | 4.86 | 1.22 | -1.22 | 150.20 | +4.00 |
| D | 4.86 | 20.14 | 14.86 | 1.22 | -1.79 | - | - |
| E | 14.86 | 55.94 | - | 1.79 | - | - | - |

## B. Moments, Shearing Forces, and Normal Forces

| Joint | Moment | Shearing force |  | Normal force |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $\mathrm{V}_{\text {ij }}$ | $\mathrm{V}_{\mathrm{kj}}$ | $\mathrm{N}_{\text {ij }}$ | $\mathrm{N}_{\mathrm{kj}}$ |
| A | $\mathrm{M}_{\mathrm{A}}$ | - | $-\mathrm{R}_{\text {AX }}$ | - | $-_{\text {R }}$ AY |
| B | $M_{A}-20 R_{A X}$ | $-\mathrm{R}_{\text {AX }}$ | $+\begin{aligned} & -.4 R_{A X} \\ & +.916 R_{A Y}-6.4 \end{aligned}$ | ${ }^{-R} A Y$ | $-.916 R_{A X} .$ |
| C | $\left\{\begin{array}{l} M_{A}-28 R_{A X} \\ +18.33 R_{A Y}-64 \end{array}\right.$ | $-.4 \mathrm{R}_{\mathrm{AX}}+.916 \mathrm{R}_{\mathrm{AY}}+$ | $\left\|\begin{array}{l} +.4 \mathrm{R}_{\mathrm{AX}}+.916 \mathrm{R}_{\mathrm{AY}} \\ +6.4 \end{array}\right\|$ | $-.916 \mathrm{R}_{\mathrm{AX}}-.4 \mathrm{R}_{\mathrm{AY}}$ | $\begin{aligned} & -.916 R_{A X} \\ & +.4 R_{A Y}-14.6 \end{aligned}$ |
| D | ${ }^{M_{A}-20 R_{A X}} \begin{aligned} & \text { a } \\ & +36.67 R_{A Y}+64\end{aligned}$ | $\left\|\begin{array}{l} +.4 R_{A X}+.916 R_{A Y} \\ +6.4 \end{array}\right\|$ | $\mathrm{R}_{\text {AX }}+16$ | $\begin{aligned} & -.916 R_{A}+.4 R_{A Y} \\ & -16.4 \end{aligned}$ | $+\mathrm{R}_{\text {AY }}$ |
| E | $\begin{aligned} & { }^{M_{A}}+36.67 R_{A Y} \\ & +384 \end{aligned}$ | $R_{A X}+16$ | - | $+\mathrm{R}_{\text {AY }}$ | - |

C. Linear Deformations due to Normal Force

| Member | $A B$ | $B C$ | $C D$ | $D E$ |
| :--- | :---: | :---: | :---: | :---: |
| $\int_{0}^{1} \frac{d s}{A}$ | 13.54 | 9.23 | 9.23 | 13.54 |
| $\int_{0}^{1} \frac{\mathrm{wds}}{A}$ | - | 75.9 | - | - |

$$
\bar{M}_{A B}=\Delta{ }_{A B}^{(N)}=-13.54 R_{A Y}
$$

$$
\bar{M}_{B C}=\Delta{ }_{B L}^{(N)}=\left(-0.916 R_{A X}-0.4 R_{A Y}-14.6\right)(9.23)-75.9
$$

$$
\bar{M}_{C D}=\Delta^{(N)}{ }_{C D}=\left(-0.916 R_{A X}+0.4 R_{A Y}-14.6\right)(9.23)
$$

$$
\bar{M}_{D E}=\Delta^{(N)}{ }_{D E}=13.54 \mathrm{R}_{\mathrm{AY}}
$$

| $\overline{\mathrm{M}}_{\mathrm{ABX}}=0$ | $\overline{\bar{M}}_{\mathrm{ABY}}=-13.54 \mathrm{R}_{\mathrm{AY}}$ |
| :--- | :--- |
| $\overline{\mathrm{M}}_{\mathrm{BCX}}=-7.8 \mathrm{R}_{\mathrm{AX}}-3.4 \mathrm{R}_{\mathrm{AY}}-193.5$ | $\overline{\mathrm{M}}_{\mathrm{BCY}}=-3.4 \mathrm{R}_{\mathrm{AX}}-1.5 \mathrm{R}_{\mathrm{AY}}-84.3$ |
| $\overline{\mathrm{M}}_{\mathrm{CDX}}=-7.8 \mathrm{R}_{\mathrm{AX}}+3.4 \mathrm{R}_{\mathrm{AY}}-124$ | $\overline{\mathrm{M}}_{\mathrm{CDY}}=-3.4 \mathrm{R}_{\mathrm{AX}}+1.5 \mathrm{R}_{\mathrm{AY}}-53.9$ |
| $\overline{\mathrm{M}}_{\mathrm{DEX}}=0$ | $\overline{\mathrm{M}}_{\mathrm{DEY}}=13.54 \mathrm{R}_{\mathrm{AY}}$ |
| $\overline{\overline{\mathrm{M}}_{\mathrm{X}}=-15.6 \mathrm{R}_{\mathrm{AX}}-317.5}$ | $\Sigma \bar{M}_{\mathrm{Y}}=-6.8 \mathrm{R}_{\mathrm{AX}}-138.2$ |

## D. Elastic Weights

$$
\begin{aligned}
\bar{P}_{A}=\phi_{A}= & 55.94 M_{A}+14.84\left(M_{A}-20 R_{A X}\right)+1.79 R_{A X} \\
\bar{P}_{B}=\phi_{B}= & 14.86 M_{A}+20.14\left(M_{A}-20 R_{A X}\right)+4.86\left(M_{A}-28 R_{A X}\right. \\
& \left.+18.33 R_{A Y}-64\right)-1.79 R_{A}-1,{ }_{2}\left(-.4 R_{A X}+.916\right. \\
& \left.\times R_{A Y}-6.4\right)+163.5-4 \\
\bar{P}_{C}=\phi_{C}= & \left(M_{A}-20 R_{A X}\right) 4.86+\left(M_{A}-28 R_{A X}+18.33 R_{A Y}\right. \\
& -64)(16.66)+\left(M_{A}-20 R_{A X}+36.37 R_{A Y}+64\right) 4.86 \\
& +\left(-.4 R_{A X}+.916 R_{A Y}\right)(1.22)+\left(.4 R_{A X}+.916 R_{A Y}\right. \\
& +6.4)(-1.22)+154 \\
\bar{P}_{D}=\phi_{D} & =\left(M_{A}-28 R_{A X}+18.33 R_{A Y}-64\right)(4.86)+\left(M_{A}\right. \\
& \left.-20 R_{A X}+36.67 R_{A Y}+64\right)(20.14)+\left(M_{A}+36.67 R_{A Y}\right. \\
& +384)(14.86)+\left(.4 R_{A X}+.916 R_{A Y}+6.4\right)(1.22) \\
& +\left(R_{A X}+16\right)(-1.79)
\end{aligned}
$$

E. Elastic Equations

$$
\begin{align*}
\Sigma \overline{\mathrm{P}}= & 0 \\
& 247.7 \mathrm{M}_{\mathrm{A}}-2332 \mathrm{R}_{\mathrm{AX}}+4546 \mathrm{R}_{\mathrm{AY}}+28375=0  \tag{1}\\
\Sigma \overline{\mathrm{PY}}+ & \Sigma \overline{\mathrm{M}}_{\mathrm{X}}=0 \\
& 2332 \mathrm{M}_{\mathrm{A}}-40116 \mathrm{R}_{\mathrm{AX}}+42862 \mathrm{R}_{\mathrm{AY}}+113578=0 \tag{2}
\end{align*}
$$

$\Sigma \overline{\mathrm{P} X}+\Sigma \overline{\mathrm{M}}_{\mathrm{Y}}=0$

$$
\begin{equation*}
4542 M_{A}-42739 R_{A X}+154480 R_{A Y}+1055327=0 \tag{3}
\end{equation*}
$$

## F. Resultants

Solving equations 1,2 , and 3, gives the following:

$$
\begin{array}{ll}
R_{A X}=8.4684 & \text { Kip } \\
R_{A Y}=7.5225 & \text { Kip } \\
M_{A}=56.1181 & \text { K. ft. }
\end{array}
$$

Applying the equation of static, gives the following:

$$
\begin{array}{ll}
R_{E X}=7.5316 & \text { Kip } \\
R_{E Y}=7.5225 & \text { Kip } \\
M_{E}=51.3749 & \text { K. ft. }
\end{array}
$$

G. Comparison of Resultants

This problem was solved by Oden (11), using the angular elastic weight due to bending moment only. The comparison of this writer's resultants with Oden's resultants is as follows:


## CHAPTER V

## CONCIUSION

The application of the general string polygon theory to the analysis of co-planar bent members is developed in this thesis.

The significant points of this study are summarized as follows:
l. The general expression for the joint elastic weight due to deformation of bending, shear and normal force is represented by a vector force and a vector moment.
2. These joint elastic weights are applied on the conjugate structure and are in equilibrium with the end conditioning elements (reactions of the conjugate structure).
3. The shear of conjugate structure is the slope of real structure; also, the bending moment of the conjugate structure about a given line is the deflection of the real structure along the same line.
4. Because the end conditioning elements are forming the static equilibrium with the elastic weights, three independent static equations are available for each conjugate structure. These equations are
called the elasto-static equations.
5. The inclusion of the influence of the shearing forces and normal forces in the analysis of deformation of bent members loaded in their own plane does not cause any special difficulty by this method.

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