

I. A VIBRATIONAL STUDY OF THE CYANOETHYLENES
II. AN ALGORITHM FOR THE DETERMINATION
OF THE SHAPE PARAMETERS OF
INFRARED ABSORPTION BANDS

By

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TABLE OF CONTENTS

Chapter	Page
PART I. A VIBRATIONAL STUDY OF THE CYANOETHYLENES	
I. INTRODUCTION	2
Vibrational Spectra	2
The Potential Function	5
The Central Force Field	6
The Simple and Modified Valence Force Fields	7
The Urey-Bradley Force Field	7
The Scope of the Present Investigation	10
II. SPECTROSCOPIC STUDIES OF THE CYANOETHYLENES	11
III. EXPERIMENTAL	14
IV. INTERPRETATION OF THE SPECTRUM OF 1,1-DICYANOETHYLENE	17
V. FORCE CONSTANT CALCULATIONS	27
Calculations Excluding Tetracyanoethylene	31
Calculations Including Tetracyanoethylene	40
VI. DISCUSSION	51
PART II. AN ALGORITHM FOR THE DETERMINATION OF THE SHAPE PARAMETERS OF INFRARED ABSORPTION BANDS	
I. INTRODUCTION	56
The Problem	56
The Shape of Infrared Absorption Bands	57
The Effect of Finite Slit Widths	62
Studies of Infrared Band Profiles	65
The Scope of the Present Investigation	65

Chapter	Page
II. DEVELOPMENT OF THE ALGORITHM	67
The Calculated Spectrum	67
The Least-Squares Procedure	70
III. EXPERIMENTAL	73
General	73
Data Acquisition	73
Evaluation of the Spectral Slit Width	74
IV. RESULTS AND DISCUSSION	76
Synthetic Spectra	76
Carbonyl Stretching Bands of Aliphatic Esters	88
Carbon-Hydrogen Bending Bands of Simple Aromatic Hydrocarbons	99
Summary and Conclusions	106
SELECTED BIBLIOGRAPHY	109
APPENDIX	112

LIST OF TABLES

Table	Page
PART I	
I. The Cyanoethylenes	12
II. Infrared Spectrum of 1,1-Dicyanoethylene	19
III. Results of Zero-Order Calculation for 1,1-Dicyanoethylene	25
IV. Urey-Bradley Force Constants for the Cyanoethylenes . .	28
V. Force Constants and Dispersions Obtained from Overlay Calculations Excluding Tetracyanoethylene	32
VI. Observed and Calculated Planar Fundamental Frequencies for Acrylonitrile	34
VII. Observed and Calculated Planar Fundamental Frequencies for Fumaronitrile	35
VIII. Observed and Calculated Planar Fundamental Frequencies for Maleonitrile	36
IX. Observed and Calculated Planar Fundamental Frequencies for 1,1-Dicyanoethylene	37
X. Observed and Calculated Planar Fundamental Frequencies for Tricyanoethylene	38
XI. Distribution of Potential Energy among Planar Internal Coordinates of Fumaronitrile	41
XII. Distribution of Potential Energy among Planar Internal Coordinates of Maleonitrile	42
XIII. Distribution of Potential Energy among Planar Internal Coordinates of 1,1-Dicyanoethylene	43
XIV. Observed and Calculated Planar Fundamental Frequencies for Tetracyanoethylene	44

Table	Page
XV. Force Constants and Dispersions Obtained from Calculations Including Tetracyanoethylene	46
XVI. Relative Chemical Shifts for a Series of Vinyl Compounds	53
PART II	
XVII. Results from Analyses of Synthetic Spectra Containing One Band	77
XVIII. Results of Zero-Slit Calculations	80
XIX. Results from Analyses of Synthetic Spectra Containing Two Bands	82
XX. Band Parameters for Synthetic Spectra Containing Three Bands	85
XXI. Refined Parameters for Band 3 of the Three-Band Synthetic Spectra	86
XXII. Band-Shape Parameters for the Carbonyl Stretching Band of Methyl Oleate	91
XXIII. Band-Shape Parameters for the Carbonyl Stretching Band of Ethyl Acetate	95
XXIV. Molar Absorptivity for the Carbonyl Stretching Band of Ethyl Acetate	96
XXV. Refined Band Shape Parameters for the Carbonyl Stretching Region of Di-n-Butyl Oxalate Spectra . . .	97
XXVI. Molar Absorptivities for the Carbonyl Stretching Region of Di-n-Butyl Oxalate Spectra	98
XXVII. Refined Band Shape Parameters for the C-H Bending Bands of Aromatic Hydrocarbons	102
XXVIII. Refined Carbon-Hydrogen Bending Band Parameters for Mixtures of Aromatic Hydrocarbons	103
XXIX. Input Requirements for Program Bandfit	114

LIST OF FIGURES

Figure	Page
PART I	
1. Infrared Spectrum of 1,1-Dicyanoethylene	18
2. Gas Phase Contours for a Planar Asymmetric Top with $I_A/I_B = 0.4$	23
3. Urey-Bradley Force Constants for the Cyanoethylenes	30
4. Some Canonical Forms for Tetracyanoethylene	49
PART II	
5. The Triangular Spectral Slit Function	65
6. Single-Band Synthetic Spectrum with 4 cm^{-1} Spectral Slit	78
7. Two-Band Synthetic Spectrum with 7 cm^{-1} Spectral Slit and $S/N = 100$	84
8. Three-Band Synthetic Spectrum with 4 cm^{-1} Spectral Slit and $S/N = 200$	87
9. Result of Fitting a Two-Band Function to a Three-Band Synthetic Spectrum	89
10. Observed and Calculated Carbonyl Stretching Bands for Methyl Oleate with 1.5 cm^{-1} Spectral Slit	93
11. Observed and Calculated Carbonyl Stretching Bands for Di- <u>n</u> -Butyl Oxalate with 3 cm^{-1} Spectral Slit	100
12. Observed and Calculated Carbon-Hydrogen Bending Bands for Toluene and <u>m</u> -Xylene with 2.2 cm^{-1} Spectral Slit	104
13. Observed and Calculated Carbon-Hydrogen Bending Bands for Toluene and <u>m</u> -Xylene with 1 cm^{-1} Spectral Slit	105
14. Observed and Calculated Carbon-Hydrogen Bending Bands for <u>o</u> -Xylene, <u>m</u> -Xylene and <u>t</u> -Butylbenzene with 2.1 cm^{-1} Spectral Slit	107

PART I

A VIBRATIONAL STUDY OF
THE CYANOETHYLENES

CHAPTER I

INTRODUCTION

Vibrational Spectra

Infrared and Raman spectroscopy furnish one of the most powerful methods for investigating the internal forces of molecules. The vibrational frequencies, derived from spectroscopic studies, are functions of the kinetic and potential energies of the system. The kinetic energy is a function of molecular geometry and atomic masses and the potential energy is determined by the forces within the system. Therefore, if the geometry and mass distribution of a molecule are known, the vibrational frequencies of the molecule can be used to obtain information concerning the forces acting during the vibrational motions.

For a molecule containing N atoms, $3N$ coordinates are necessary to define the system at any point in time. Of these, 3 coordinates are necessary to define the center of mass of the molecule with respect to an external reference and, for a nonlinear molecule, 3 define rotational motions about the major axes of the molecule. The remaining $3N-6$ coordinates are available to describe the motions of the component atoms with respect to the rotating coordinate system; that is, the molecular vibrations.

Each of the above-mentioned $3N-6$ degrees of freedom corresponds to a normal mode of vibration of the molecule. In a normal mode, all

nuclei move in phase and at the same frequency. Amplitudes of displacement from the equilibrium position may be different for each nucleus, but each must reach its maximum displacement at the same time and each atom must pass through its equilibrium position at the same time. The frequency associated with a normal mode is known as a normal, or fundamental, frequency.

In discussing kinetic and potential energy of a vibrating system it is convenient to introduce the concept of mass-weighted cartesian coordinates q_i . These are defined by the equations

$$q_i = \sqrt{M_j} \Delta X_{ij}, \quad i = 1, 2, \dots, 3N \quad (1)$$

$$j = 1, 2, \dots, N$$

where M_j is the mass of the j th atom and ΔX_{ij} is one of the three cartesian displacement coordinates for the j th atom. In terms of these coordinates, the kinetic energy is

$$2T = \sum_{i=1}^{3N} \dot{q}_i^2 \quad (2)$$

where the dot implies differentiation with respect to time. For small displacements, the potential energy can be expressed as

$$2V = \sum_{i=1}^{3N} \sum_{j=1}^{3N} f_{ij} q_i q_j \quad (3)$$

where the f_{ij} are force constants given by

$$f_{ij} = \left(\frac{\partial^2 V}{\partial q_i \partial q_j} \right)_0 \quad (4)$$

The zero subscript implies the equilibrium position, that is,

$$q_i = q_j = 0.$$

Although it would be possible to describe the normal modes of vibration in terms of the $3N$ mass weighted cartesian coordinates, it is more advantageous to introduce a system of normal coordinates.

These are defined as

$$Q_k = \sum_{i=1}^{3N} l_{ki} q_i, \quad k = 1, 2, \dots, 3N \quad (5)$$

in which the coefficients l_{ki} are chosen such that, in terms of the normal coordinates, the kinetic and potential energies have the forms

$$2T = \sum_{k=1}^{3N} \dot{Q}_k^2 \quad 2V = \sum_{k=1}^{3N} \lambda_k Q_k^2 \quad (6)$$

Thus the potential energy involves no cross products while the kinetic term retains the same form as in Equation 2. In normal coordinate space, then, each coordinate corresponds to one of the normal modes of the molecule. The λ_k of Equation 6 are related to the fundamental frequencies ν_k by

$$\lambda_k = 4\pi^2 \nu_k^2 \quad (7)$$

There is no problem in considering $3N$, rather than $3N-6$, normal coordinates since translation and rotation may be considered as normal modes with zero frequency.

The Potential Function

Virtually all studies of molecular force fields are based on the assumption that the vibrational motions are harmonic. This assumption would be expected to be true if the motions involved only infinitesimal displacements of the atoms from their equilibrium positions. Nevertheless, the success of quadratic potential energy functions in describing the force fields of many molecules confirms the validity of the harmonic approximation.

For the purpose of developing the potential function, the vibrational potential energy may be expressed in terms of $3N-6$ internal displacement coordinates. These are most conveniently chosen as the instantaneous changes in the bond lengths and bond angles in the molecule. Under the assumption of harmonic motion, the potential energy in terms of the internal displacement coordinates is

$$2V = \sum_{i=1}^{3N-6} \sum_{j=1}^{3N-6} f_{ij} \Delta R_i \Delta R_j \quad (8)$$

where ΔR is an internal displacement coordinate and f_{ij} is the force constant. This equation is entirely analogous to equation 3, the only difference being that a different coordinate set is used.

Equation 8 defines the general valence force field. This field includes not only the $3N-6$ principal force constants ($i = j$), but also all interaction force constants ($i \neq j$). In the absence of symmetry considerations there are then $(3N-5)(3N-6)/2$ force constants to be evaluated to define this field. Since there are only $3N-6$ vibrational degrees of freedom, the system is clearly underdetermined. In general,

however, considerations of molecular symmetry reduce the number of unique force constants to be determined. Consider, for example, the planar vibrations of tetracyanoethylene. There are 17 degrees of freedom associated with the planar modes so that, if symmetry is neglected, 153 force constants would have to be evaluated to define the general valence force field. When molecular symmetry is taken into account, the number of unique force constants to be determined is reduced to 30. Even so, the problem is still underdetermined in that only 17 fundamental frequencies are available.

In addition to the above mentioned difficulties, the general valence force field does not present a clear picture of the nature of the interactions between nonbonded atoms. In order to decrease the number of force constants and to permit a clearer interpretation of the forces acting within a molecule, various assumptions concerning the nature of the force field have been made. These assumptions have led to the postulation of a number of special force fields.¹ Some of these are described below.

The Central Force Field

The central force field is based on the assumption that interactions exist in a molecule along the lines connecting pairs of atoms and that every atom exerts influence on every other atom, independent of whether the atoms are bonded to each other or not. The number of independent force constants required for this field is $N(N-1)/2$, a number quite small compared to the $(3N-5)(3N-6)/2$ required for the general valence force field. This model, however, has not given satisfactory results² and is now seldom used.

The Simple and Modified Valence Force Fields

The simple valence force field is similar to the central force field in its treatment of forces acting between bonded atoms, but does not include interactions between nonbonded atoms. Instead, it is assumed that a force constant is associated with the angle between bonds which opposes a change in the bond angle. This field can usually account for all molecular bending vibrations and in this respect is superior to the central force field. However, the lack of quantitative agreement between frequencies calculated on the basis of the simple valence force field and those actually observed indicate that this field is best regarded as only a rough approximation.

The modified valence force field consists of the simple valence force field with the addition of one or more judiciously chosen interaction terms. This generally leads to a set of force constants which can be evaluated in such a manner as to reproduce the observed frequencies well. It is not, however, always possible to predict just which of the interaction constants it is necessary to include. In fact, several different sets of force constants may adequately describe the frequencies of a given molecule. In addition to this lack of uniqueness, it is usually difficult to give a meaningful interpretation for the selected interaction constants.

The Urey-Bradley Force Field

The Urey-Bradley force field (UBFF)³ may be considered as a simple valence force field modified to include central forces between nonbonded atoms. The potential energy can be expressed in terms of the

basic UBFF as

$$2V = 2 \sum_i K'_i r_i \Delta r_i + \sum_i K_i (\Delta r_i)^2 + 2 \sum_j H'_j (\Delta \alpha_j) + \sum_j H_j (\Delta \alpha_j)^2 + 2 \sum_k F'_k q_k \Delta q_k + \sum_k F_k (\Delta q_k)^2 \quad (9)$$

where r_i is an equilibrium bond distance, α_j a bond angle and q_k an equilibrium distance between nonbonded atoms. K and H are bond stretching and bending force constants respectively and F represents a force constant between nonbonded atoms. The linear terms in equation 9 are necessary because the set of internal coordinates used, which includes the q_k , is no longer linearly independent. The q_r , however, are related to the r_i and α_j by a simple geometric relationship. It is therefore possible to evaluate the H'_j and the K'_i in terms of the F'_k .

Based on the assumptions that, at nonbonded distances usually encountered, the principal force acting between nonbonded atoms is repulsive and that this repulsive force is proportional to $1/r^9$, it is common practice to reduce the number of necessary force constants further by setting F' equal to $-0.1F$. It has been pointed out that this practice is unwarranted in that the F_k may represent stabilizing as well as destabilizing forces.^{4,5} The elimination of F' probably does little harm, however, since it is probably small compared to F .

The UBFF was first used extensively by Shimanouchi.⁶ Since then it has been used successfully by many workers for a wide variety of molecular types.

Urey-Bradley force constants (UBFC's) are, to some extent, transferable; that is, force constants which have been evaluated for one molecule can often be used to calculate normal frequencies for similar molecules. This transferability has been studied extensively and Shimanouchi has published a compilation of UBFC's useful for many small molecules.⁷ As a direct result of this transferability of force constants, it is possible to treat several molecular species in one normal coordinate calculation, evaluating force constants for similar groupings only once.⁸ This technique, called an overlay calculation, makes it possible to determine UBFC's for a series of related molecules instead of attempting to determine a complete set for one molecule. The agreement between observed and calculated frequencies generally worsens as the number of molecules included in an overlay calculation is increased. This is because one is attempting to fit more data with the same number of parameters. The uncertainties in the calculated force constants decreases, however, and their possible physical significance increases.

Overend and co-workers have written a set of computer programs to perform an overlay normal coordinate analysis in which force constants are refined by an iterative least-squares method.⁹ The algorithm used in the force constant calculations will not be discussed here as it is thoroughly described in a recent reference.¹⁰ The above-mentioned computer program was modified slightly and used extensively throughout this investigation.

The Scope of the Present Investigation

The primary purpose of this investigation was to evaluate a consistent set of Urey-Bradley force constants for the planar modes of the series of cyanoethylenes: acrylonitrile, fumaronitrile, maleonitrile, 1,1-dicyanoethylene, tricyanoethylene and tetracyanoethylene.¹¹ These molecules were selected because they comprise a set of structurally similar molecules in which the degree of electron delocalization varies widely from one end of the series to the other. This provided an ideal situation to test the transferability of UBFC's under conditions of varying resonance. In addition, it was felt that confirmation of certain assignments in the spectrum of tetracyanoethylene could be accomplished by an overlay calculation involving all the cyanoethylenes.

The experimental phase of this study involved the preparation of 1,1-dicyanoethylene and the determination of its infrared spectrum. The infrared data, together with infrared and Raman data from the literature for the other cyanoethylenes, was used in a variety of overlay Urey-Bradley calculations.

The overlay calculations have resulted in a satisfactory assignment of the planar modes of 1,1-dicyanoethylene. In addition, the assignments for two of the tetracyanoethylene fundamentals have been changed.

CHAPTER II

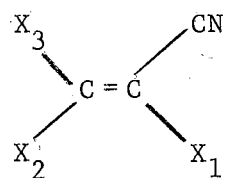
SPECTROSCOPIC STUDIES OF THE CYANOETHYLENES

The cyanoethylenes (Table I) provide an interesting example of a series of structurally similar molecules in which the degree of π -electron delocalization varies smoothly throughout the series. Since these resonance effects should result in a weakening of the $C\equiv N$, $C=C$ and perhaps the $C-H$ bonds with a corresponding strengthening of the $C-C$ bonds, the series provides an interesting ground for the transferability of UBFC's since deviations from transferability should have physical significance in terms of bonding differences. In addition, the possibility of modifying the UBFC of an open-chain molecule to include a resonance parameter is presented. The inclusion of such a parameter has been reported only for cyclic molecules.¹²⁻¹⁶

Halverson and co-workers¹⁷ have reported a complete assignment for the infrared and Raman spectra of acrylonitrile. No vibrational analysis of the planar modes was attempted, but a 4-term modified valence force field was used to calculate approximate frequencies for the four out-of-plane modes.

Spectra and assignments have been reported for fumaronitrile,¹⁸⁻²⁰ maleonitrile,^{19,20} tetracyanoethylene^{19,20} and tricyanoethylene.¹⁸ There are some differences between the spectra reported by Long and George¹⁹ and those reported by Devlin, Overend and Crawford¹⁸ and Miller, et al.²⁰ For those cases where differences occur, the

TABLE I
THE CYANOETHYLENES



Compound	X ₁	X ₂	X ₃
Acrylonitrile	H	H	H
Fumaronitrile	H	CN	H
Maleonitrile	H	H	CN
1,1-Dicyanoethylene	CN	H	H
Tricyanoethylene	CN	CN	H
Tetracyanoethylene	CN	CN	CN

assignments proposed by Miller have been chosen for use in this investigation.

Devlin and co-workers¹⁸ have performed an overlay UBFF vibrational analysis for acrylonitrile, fumaronitrile and tricyanoethylene. As a result of this calculation, a set of 16 force constants were determined from the 35 observed planar fundamental frequencies. As had been previously reported,²¹ they found it necessary to include a quadratic cross term between trans-hydrogen bending coordinates in the potential function. This term was introduced to provide more accurate calculated values for the frequencies of in-plane rocking modes involving hydrogen. Examination of the potential energy distributions for acrylonitrile and fumaronitrile indicate that this force constant is involved mainly in C=C stretching and C-H rocking and deformation vibrations.

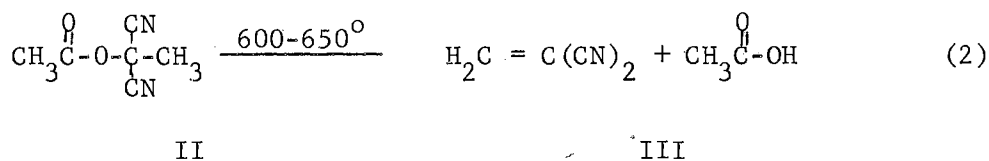
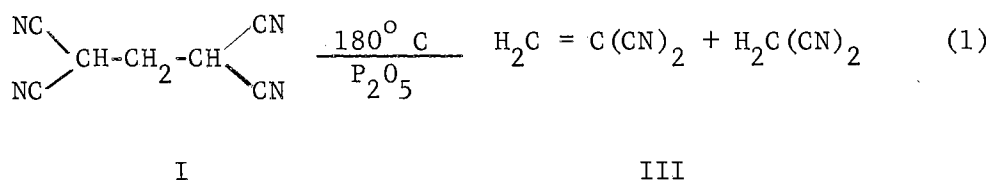
There are thus reasonably firm assignments for the planar modes of all but one of the cyanoethylenes. Vibrational spectra and assignments have more recently been published for 1,1-dicyanoethylene.¹¹ The remaining chapters of this work are concerned with the problem of the vibrational spectrum of this compound and with overlay normal coordinate analyses including all the cyanoethylenes.

CHAPTER III

EXPERIMENTAL

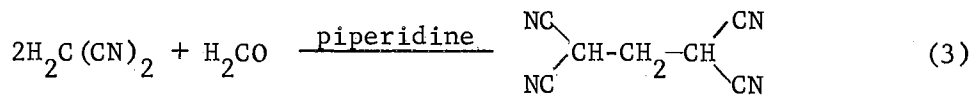
As was noted in the introduction, the experimental phase of this work consisted of the preparation of 1,1-dicyanoethylene and the determination of its infrared spectrum. Some experimental difficulty was encountered in the handling of this compound. It is easily polymerized in the presence of any source of labile protons. Many nucleophiles are also effective in initiating polymerization. As a result of this, extreme care was taken to exclude contaminants, particularly water vapor, during preparation and transfer of 1,1-dicyanoethylene.

The first synthesis of 1,1-dicyanoethylene was reported by Ardis and co-workers.²² Two basic synthetic schemes were presented. These are outlined in equations 1 and 2. In this work, the synthesis involving



the pyrolysis of 1,1,3,3-tetracyanopropane (I) was chosen.

The synthesis of I was accomplished by the method of Diels^{23,24} as outlined in equation 3. Malononitrile (Eastman Practical Grade)



I

was purified by fractional distillation and condensed with formaldehyde (as formalin) in ethanol solution. The resultant solid material was washed with dilute sulfuric acid and with water and recrystallized from a 2:1 mixture of acetonitrile and benzene. The purified I thus prepared melted sharply at 137° C. Yields for this reaction were typically about 40% of theoretical, based on malononitrile as the limiting reagent.

In our early attempts to synthesize III, a closed system involving several ampule-to-ampule distillations was used for the pyrolysis of I. Yields using this apparatus were disappointingly low. We therefore reverted to a method similar to that used by Ardis.²² Dried I and P₂O₅ were placed in a Claisen flask and covered with dry Pyrex wool. The side tube of the flask was connected to a receiver containing P₂O₅ and dry benzene. The receiver was cooled in a Dry Ice-acetone bath and the system was evacuated to about 10 torr. The flask was then heated to 180° C until pyrolysis was completed. The material in the receiver was melted and the benzene solution was fractionally distilled through a 20 cm vacuum-jacketed Vigreux column. The fraction boiling over the range 50.0 - 51.2° C at 10 torr was taken as 1,1-dicyanoethylene. The distilled material was stored in the cold over P₂O₅. No attempt was made to determine the yield in the pyrolysis.

Standard sampling cells, fitted with cesium iodide windows, were used in determining the infrared spectrum of III. All transfers were carried out in a dry box which was well purged with dry nitrogen. The cells were purged separately with dry nitrogen for one hour prior to use. For the determination of the spectrum of gaseous III, the sampling cell was placed in an electrically heated oven and maintained at a temperature in excess of 170°C during the spectral scan. Since it has been reported²⁵ that poly(1,1-dicyanoethylene) undergoes depolymerization at temperatures above 160°C , it was hoped that this would prevent deposition of the polymer in the gas sampling cell. Some polymer, however, was always formed during determination of the spectrum of gaseous III.

Much of the spectrum of 1,1-dicyanoethylene was recorded using a Perkin-Elmer Model 221 prism-grating spectrophotometer. More recent measurements were made with a Perkin-Elmer Model 521 dual-grating instrument. A Beckman IR-7 spectrophotometer with cesium iodide optics was used to determine spectra over the $200 - 250\text{ cm}^{-1}$ region. The parameters of the Perkin-Elmer instruments were set so that the spectral slit width was 1.2 cm^{-1} at 1000 cm^{-1} .

For the region above 500 cm^{-1} , well known absorption bands of polystyrene, atmospheric water vapor, atmospheric CO_2 and indene were used to calibrate the frequency scale of the spectrophotometer.²⁶ No attempt was made to calibrate below this point. It is estimated that spectral features are determined to within 2 cm^{-1} over the range $3500 - 500\text{ cm}^{-1}$ and to within 5 cm^{-1} below 500 cm^{-1} .

CHAPTER IV

INTERPRETATION OF THE SPECTRUM OF 1,1-DICYANOETHYLENE

If it is assumed that 1,1-dicyanoethylene is planar, then it must belong to the point group C_{2v} . Group theoretic considerations would lead one to expect that the 18 normal modes of vibration would be distributed among the symmetry classes as follows:

$$\Gamma_{\text{vib}} = 7A_1 + 2A_2 + 3B_1 + 6B_2$$

Of these, the modes belonging to classes A_1 , B_1 and B_2 will exhibit infrared activity while all of the normal vibrations would be expected to be Raman-active. The planar vibrations belong to symmetry classes A_1 and B_2 and thus will all be infrared-active. As a consequence of the fact that 16 of the 18 normal vibrations are infrared-active, the infrared spectrum (Figure 1 and Table II) is quite rich in intense bands. However, with the aid of gas-phase band contours and a set of frequencies calculated for the planar modes from the overlay force constants used for acrylonitrile, fumaronitrile and tricyanoethylene,¹⁸ it was possible to make a satisfactory assignment of the planar modes. As is the case for the other cyanoethylenes, the assignment of the out-of-plane modes is still unsatisfactory.

1,1-Dicyanoethylene is an asymmetric top, the moments of inertia about the three principal axes being

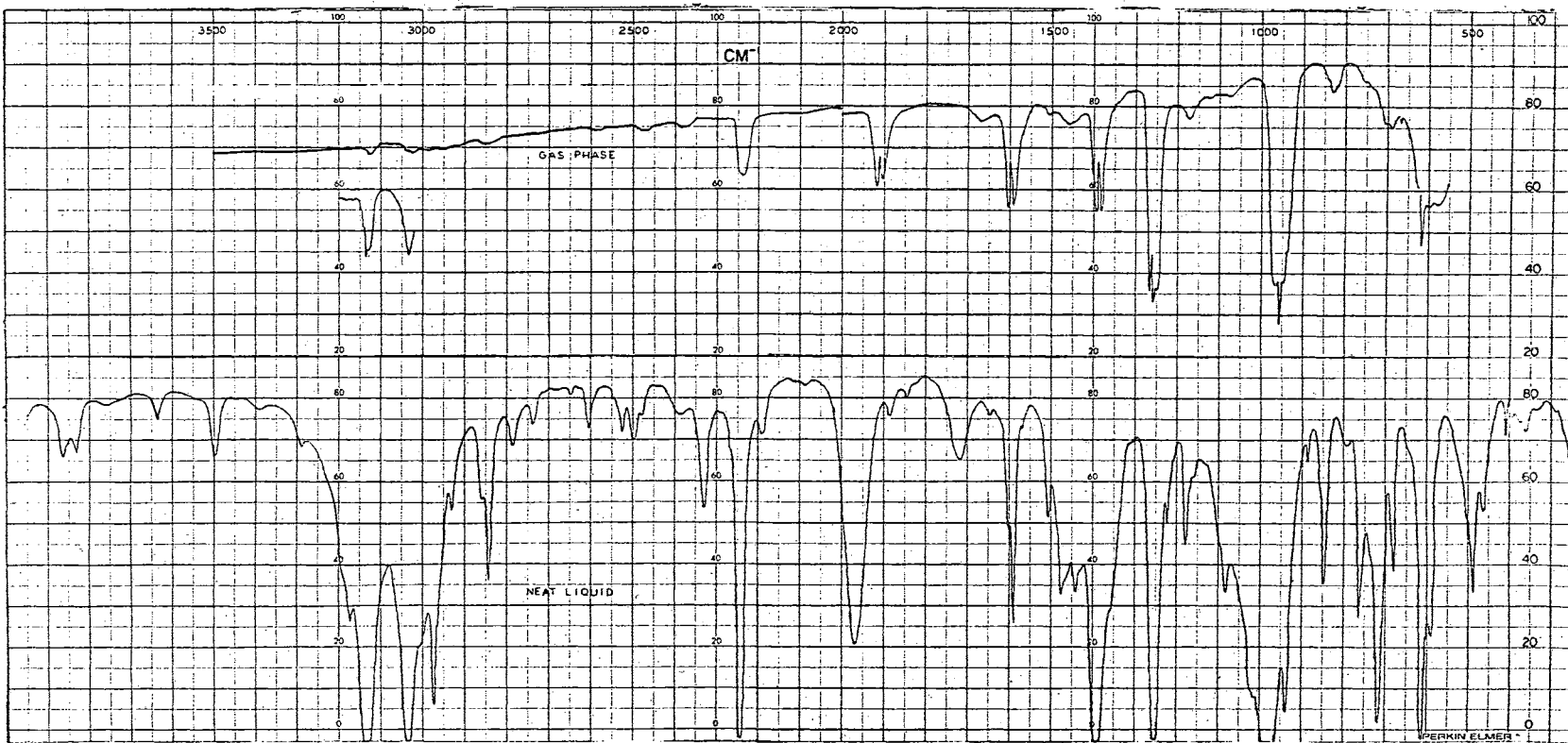


Figure 1. Infrared Spectrum of 1,1-Dicyanoethylene

TABLE II
 INFRARED SPECTRUM OF 1,1-DICYANOETHYLENE

Liquid	Gas	Assignment
3964		$3037 + 939 = 3976$
3932		
3880		
3735		$596 + 3140 = 3736$
3498		$2248 + 1262 = 3510$
3390		
3288		
3175		
3135	3140	ν_{13}
3091		
3036	3037	ν_1
2966		$720 + 2248 = 2968$
2933		
2864		$2255 + 620 = 2875$
2846		$1604 + 1262 = 2866$
2791		$2 \times 1394 = 2788$
2741		$2248 + 494 = 2742$
2653		$1394 + 1262 = 2656$
2610		
2531		$1604 + 939 = 2543$
2503		$2 \times 1394 = 2788$
2482		$2248 + 240 = 2488$

TABLE II (Continued)

Liquid	Gas	Assignment
2334		$1394 + 939 = 2333$
2248	2255	ν_2
2245	2248	ν_{13}
2195		$1262 + 939 = 2201$
1974	1917	$2 \times 983 = 1966$
1890		$1394 + 494 = 1888$
1722		
1605		$983 + 620 = 1603$
1594	1604	ν_3
1511		
1478		
1445		Polymer?
1395	1394	ν_4
1362		
1256	1262	ν_{15}
1227		$2 \times 620 = 1240$
1182		Polymer
1088		
983	959	ν_{10}
939		ν_{16}
890		$720 + 170 = 890$
850		
762	758	$596 + 170 = 766$
720	709	ν_5 (?)

TABLE II (Continued)

Liquid	Gas	Assignment
685		Polymer?
620	620	ν_{11}
606		
596	587	ν_6
494		ν_{17}
469		Polymer?
415		$170 + 240 = 410$
383		
367		
339		$2 \times 170 = 340$
240		$\nu_{18} (?)$
170		ν_7

$$I_A \simeq 77.2 \times 10^{-40} \text{ g cm}^2$$

$$I_B \simeq 176.3$$

$$I_C \simeq 253.5$$

For this molecule, the A-axis lies in the molecular plane and corresponds to a line drawn through the center of mass perpendicular to the line of the C=C bond. The C-axis is perpendicular to the plane of the molecule.

The band contours to be expected for gas-phase spectra of planar asymmetric top molecules are well known.²⁷ For a molecule such as 1,1-dicyanoethylene, where $I_A/I_B \simeq 0.44$, we would expect that modes for which the change in dipole moment is along one of the major axes would exhibit A-, B- or C-type contours (corresponding to dipole change along the A-, B- or C-axis, respectively) as shown in Figure 2. Since 1,1-dicyanoethylene is of C_{2v} symmetry, it is to be expected that each normal mode will involve a net dipole change along one or the other of the major axes; i.e., there will be no mixing of contour types.

Of the more intense bands only the two at 959 and 620 cm^{-1} showed C-type structure. These were immediately assigned to the H-C-H and C-C-C out-of-plane wagging modes (ν_{10} and ν_{11} , respectively). The frequency of the third infrared active out-of-plane mode was apparently beyond the range of our spectrophotometer. No evidence of the infrared inactive out-of-plane modes has been found. By comparison with other cyanoethylenes the A-type bands at 3135 and 1262 cm^{-1} were assigned to the B_2 species stretch (ν_{13}) and CH_2 rock (ν_{15}). The B-type bands at 1604, 1395 and 709 cm^{-1} were similarly assigned to the C=C stretch (ν_3) and CH_2 deformation (ν_4) and the C-C stretch (ν_5) of the A_1 species.

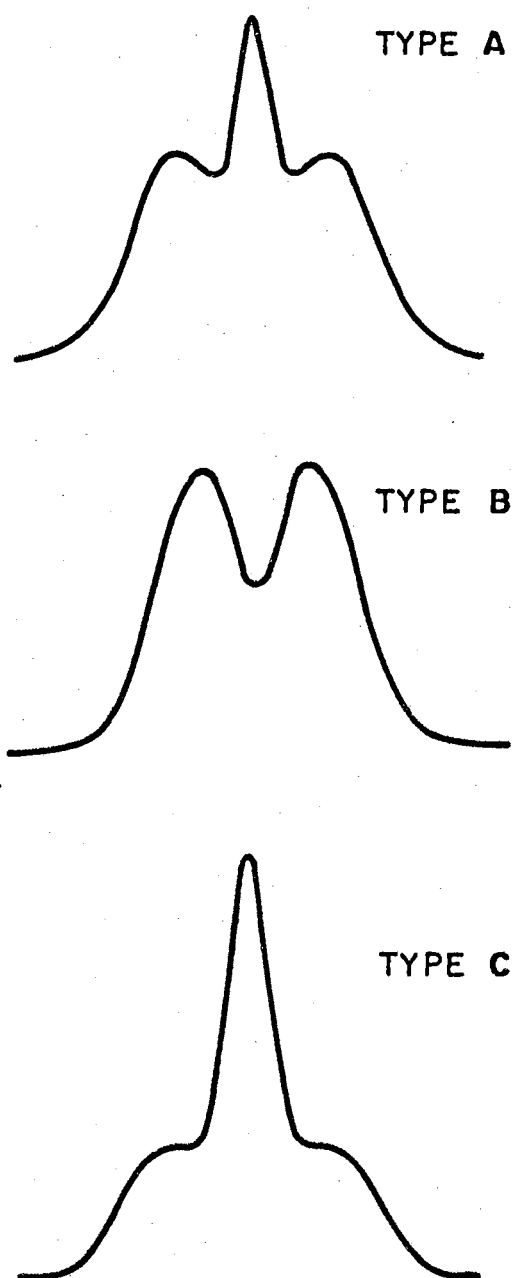


Figure 2. Gas-Phase Contours for a Planar Asymmetric Top with $I_A/I_B = 0.4$

Because of either weak gas-phase band intensities, lack of band structure or overlapping bands, the remaining fundamentals have been assigned primarily on the basis of frequencies calculated using the set of overlay Urey-Bradley force constants (UBFC) previously mentioned.¹⁸ The observed and calculated frequencies, together with mode assignments, are shown in Table III. The calculated frequencies agree very closely with those assigned on the basis of band shapes and, except for three cases, permit immediate assignment of relatively intense bands to the remaining planar modes.

There was considerable hesitation in assigning the medium intensity band at 720 cm^{-1} as ν_5 since the calculated value of 762 cm^{-1} is assignable as $(\nu_6 + \nu_7)$, whereas no combination bands or overtones were found to explain the 720 cm^{-1} band. This, plus the greater intensity of the 720 cm^{-1} band, led to its assignment at ν_5 . Unfortunately ν_7 , used in the assignment of the 762 cm^{-1} band as $(\nu_6 + \nu_7)$, was not itself observed. However, the value 155 cm^{-1} calculated from the transferred UBFC was in fair agreement with the frequency of 170 cm^{-1} required to explain (by combination bands and overtones) the bands at 890, 762 and 339 cm^{-1} . Thus the assignment of ν_7 as 170 cm^{-1} is given. The third difficulty arose in the assignment of ν_{18} and has not been completely resolved. The original calculated value of 278 cm^{-1} for ν_{18} seemed to suggest that a band observed at 240 cm^{-1} should be assigned as ν_{18} . However, subsequent attempts to adjust the force constants so as to give an improved set for the entire molecular series invariably yielded a calculated frequency of about 275 cm^{-1} for ν_{18} , as can be seen in the next chapter. Nevertheless, the 240 cm^{-1} band is currently assigned as ν_{18} and only a normal coordinate treatment of the

TABLE III
 RESULTS OF ZERO-ORDER CALCULATION
 FOR 1,1-DICYANO ETHYLENE

Symmetry Species	No.	$\nu_{\text{obs.}}$	$\nu_{\text{calc.}}$	Assignment ^(a)
A_1	1	3036	2999	C-H Stretch
	2	2248	2238	C \equiv N Stretch
	3	1594	1597	C=C Stretch
	4	1395	1412	HCH Bend
	5	720	765	C-C Stretch
	6	596	566	C(CN) ₂ Scissor
	7	170	157	C-C \equiv N Bend
B_2	13	3135	3103	C-H Stretch
	14	2245	2237	C \equiv N Stretch
	15	1256	1304	HCH Bend
	16	939	956	C-C Stretch
	17	494	521	C(CN) ₂ Rock
	18	240	278	C-C \equiv N Bend

(a) from the calculated potential energy distribution

out-of-plane modes can help in confirming or disproving this assignment.

The assignments discussed herein appear in Table II. In addition, features have been assigned to combination, overtone or impurity bands where such an assignment was logically feasible.

CHAPTER V

FORCE CONSTANT CALCULATIONS

There are several prerequisites necessary before an overlay force constant analysis based on spectroscopic data can be attempted. First, a choice of a molecular force field must be made. As has been mentioned earlier, the Urey-Bradley force field has been selected for this investigation. Within the constraints of the force field a set of meaningful force constants must then be defined and initial values either estimated for them or transferred from similar molecules. The molecular geometry must be known, and finally, firm frequency assignments for at least most of the fundamental modes must be available.

For the planar vibrations of the cyanoethylenes a potential function involving 18 force constants was defined. A list of these constants is shown in Table IV. As before, K implies a bond stretching force constant, H a bending constant and F the interaction between nonbonded atoms. Here F is restricted to the interaction between atoms which are 1,3 with respect to each other, i.e., to the interaction between nonbonded atoms which are bonded to a common atom. The force constant C (for cis) is associated with nonbonded interactions between atoms which are 1,4 to each other and which have a cis relationship with respect to the 2,3 bond. There are, of course, three possible types of cis nonbonded interactions for the cyanoethylenes; H-H, H-C, and C-C. Throughout this investigation it was assumed that only the

TABLE IV
UREY-BRADLEY FORCE CONSTANTS FOR THE CYANOETHYLENES

Force Constant	Initial Value (a)
$K_{C=C}$	6.56
K_{CH}	} 4.83
K^*_{CH}	
K_{CN}	16.73
K_{CC}	5.17
$H_{C=CH}$	0.37
$H_{C=CC}$	0.52
H_{HCH}	0.40
H_{CCH}	0.32
H_{CCN}	0.30
H_{CCC}	0.237
$F_{C=CH}$	0.42
$F_{C=CC}$	0.76
F_{HCH}	0.00
F_{CCH}	0.18
F_{CCC}	0.64
$C_{CC=CC}$	-0.04
Tr	0.06

(a) The units of K, F and C are in mdyne/A and H is in 10^{-11} erg radian⁻².

last of these could be large enough to influence the calculated results. A pictorial representation of these selected force constants is shown in Figure 3 using acrylonitrile and tricyanoethylene as examples. The force constants are shown along the internal coordinates with which they are nominally associated. It should be noted that two CH stretching constants K_{CH} and K^*_{CH} are shown. During the course of this investigation it was found necessary to define one constant for the case where there are two hydrogens on a carbon (K_{CH}) and another for the case where there is only one (K^*_{CH}). The last entry in Table IV represents the interaction between trans-hydrogen bending coordinates previously mentioned.

Refined force constant values for acrylonitrile, fumaronitrile and tricyanoethylene by Devlin, Overend and Crawford¹⁸ were used as a starting point in this work. The initial values used are listed in Table IV. Their value of 4.83 mdyne/A was used for both K_{CH} and K^*_{CH} since they used only one CH stretching constant for the 3 molecules.

Wilcox and co-workers, using microwave spectroscopy, have determined bond lengths and angles for acrylonitrile.^{28,29} They concluded that the molecule is planar and thus belongs to symmetry groups C_s .

The equilibrium bond distances and angles proposed are:

C-H bond	1.07A
C=C bond	1.34A
C-C bond	1.445A
C-N bond	1.159A
C=C-C angle	121.5°
C=C-H angle	120°

Other workers using electron diffraction^{30,31} and microwave spectroscopy^{31,32} have determined slightly different values for these parameters. Bond lengths for tetracyanoethylene have also been reported³³

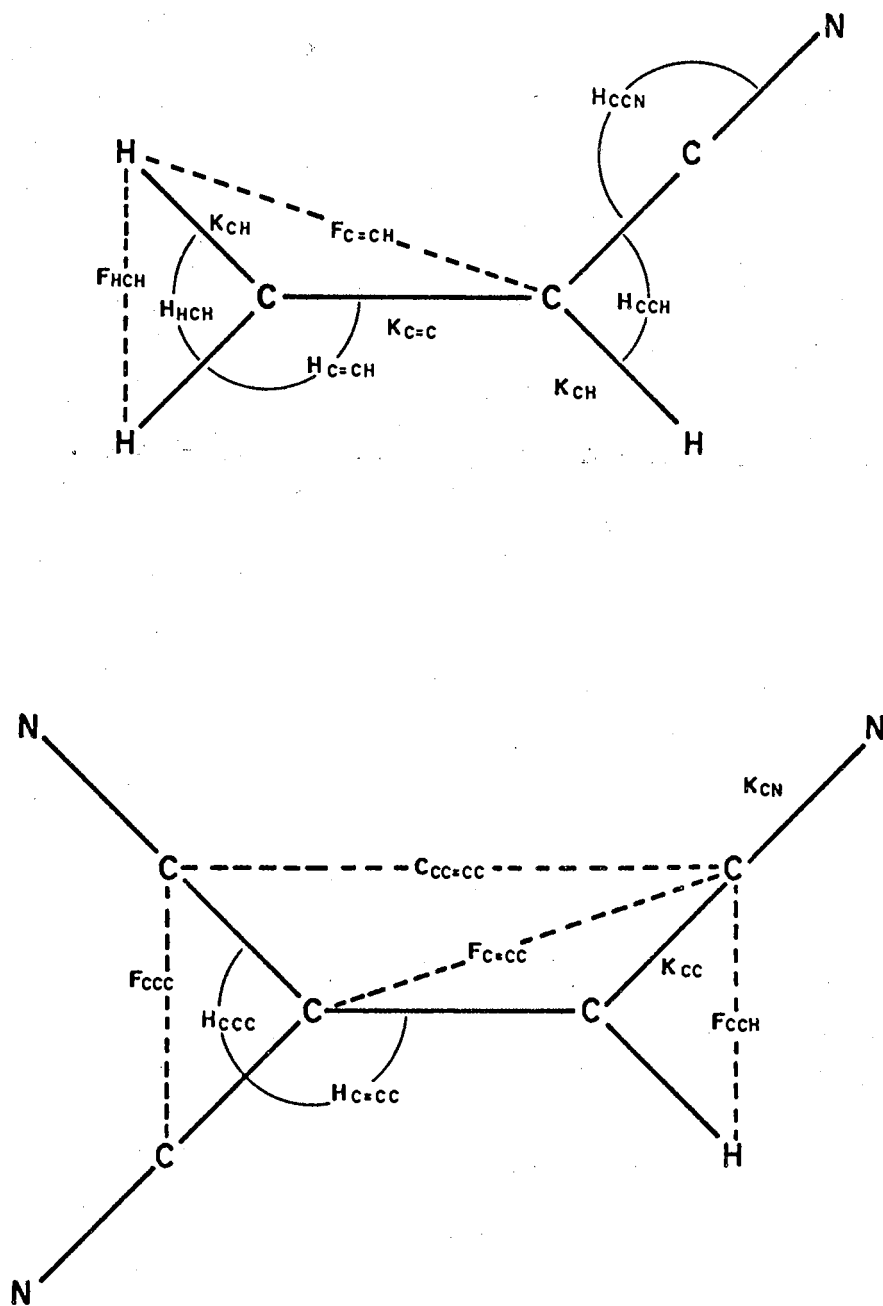


Figure 3. Urey-Bradley Force Constants for the Cyanoethylenes

which agree closely with those given by Wilcox. For the purpose of this investigation the bond lengths and bond angles given by Wilcox were employed. It was not felt that the small differences reported would affect our calculations noticeably. Moreover, the use of these values would make the present calculations more consistent with the earlier cyanoethylene work.

The calculations to be presented in the remainder of this chapter have been divided into two basic groups, calculations excluding tetracyanoethylene and calculations in which tetracyanoethylene is included. This partition is based on early evidence²⁰ that the bonding within tetracyanoethylene is not representative of the molecular series. Thus, force constants for the other five compounds were evaluated prior to consideration of tetracyanoethylene.

Calculations Excluding Tetracyanoethylene

Starting with the force constant values given in Table IV, overlay calculations involving the following sets of molecules have been performed: 1) fumaronitrile and maleonitrile, 2) fumaronitrile, maleonitrile and tricyanoethylene, 3) acrylonitrile, fumaronitrile, maleonitrile and tricyanoethylene and 4) all cyanoethylenes save tetracyanoethylene. The refined force constants resulting from these calculations are listed in Table V as $\Phi(1)$ through $\Phi(4)$. Calculated planar frequencies for each molecule appear in Tables VI-X. Although the percent errors in the frequencies and the stability of the force constants in sets (1), (2) and (4) are quite satisfying, these results were achieved only after extensive preliminary calculations. Set (3)

TABLE V

FORCE CONSTANTS AND DISPERSIONS OBTAINED FROM OVERLAY CALCULATIONS EXCLUDING TETRACYANOETHYLENE^(a)

Force Constant	$\Phi(1)$	$\sigma_{\Phi(1)}$	$\Phi(2)$	$\sigma_{\Phi(2)}$	$\Phi(3)$	$\sigma_{\Phi(3)}$	$\Phi(4)$	$\sigma_{\Phi(4)}$
$K_{C=C}$	6.88	0.13	6.82	0.10	6.72	0.15	6.82	0.09
K_{CH}	—(d)	—	—	—	4.74 ^(b)	0.06	4.94	0.03
K^*_{CH}	4.55	0.06	4.54	0.06			4.56	0.04
K_{CN}	16.99	0.11	16.94	0.10	16.85	0.15	16.94	0.08
K_{CC}	4.52	0.22	4.56	0.16	4.96	0.22	4.63	0.11
$H_{C=CH}$	0.29	0.09	0.29	0.08	0.29	0.06	0.35	0.03
$H_{C=CC}$	0.44	0.07	0.46	0.06	0.55	0.10	0.43	0.05
H_{HCH}	—	—	—	—	0.42	0.03	0.39	0.01
H_{CCH}	0.21	0.10	0.21	0.07	0.33	0.05	0.17	0.03
H_{CCN}	0.30	0.02	0.30	0.02	0.29	0.02	0.30	0.01
H_{CCC}	—	—	0.24	(c)	0.24	(c)	0.24	(c)
$F_{C=CH}$	0.51	0.11	0.49	0.10	0.50	0.09	0.43	0.04
$F_{C=CC}$	1.05	0.11	0.94	0.09	0.80	0.13	0.88	0.07
F_{HCH}	—	—	—	—	0.	(c)	0.	(c)

TABLE V (Continued)

Force Constant	$\phi(1)$	$\sigma_{\phi}(1)$	$\phi(2)$	$\sigma_{\phi}(2)$	$\phi(3)$	$\sigma_{\phi}(3)$	$\phi(4)$	$\sigma_{\phi}(4)$
F_{CCH}	0.52	0.14	0.53	0.10	0.26	0.07	0.57	0.06
F_{CCC}	—	—	0.80	0.10	0.75	0.16	0.86	0.07
$C_{CC=CC}$	0.08	0.05	0.04	0.05	0.05	0.07	0.04	(c)
Tr	0.06	(c)	0.06	(c)	0.06	(c)	0.06	(c)
Av. % Error in ν 's	0.87		1.11		1.31		1.09	

(a) Force constants are in mdyne/A except for the bending constants, which are in 10^{-11} erg/rad².

(b) Only one CH stretching constant was considered in this calculation.

(c) The value of the corresponding force constant was constrained.

(d) A line indicates that the calculation is independent of the indicated force constant.

TABLE VI
OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR ACRYLONITRILE

ν_{obs}	$\nu(3)^{(a)}$	$\nu(4)$
3125	3090	3137
3078	3076	3064
3042	2987	3031
2239	2233	2236
1615	1604	1604
1416	1423	1407
1282	1280	1280
1096	1059	1090
869	862	868
570	574	569
242	250	251
Av. % Error	1.18	0.66

(a) The number in parentheses refers to the corresponding ϕ in Table V.

TABLE VII
OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR FUMARONITRILE

Symmetry	$\nu_{\text{obs.}}$	$\nu(1)^{(a)}$	$\nu(2)$	$\nu(3)$	$\nu(4)$
A_g	3049	3056	3052	3069	3057
	2232	2234	2231	2231	2234
	1611	1616	1610	1615	1613
	1297	1297	1294	1298	1292
	1002	1006	1008	994	1005
	545	530	520	518	514
	261	268	266	265	266
B_u	3067	3073	3068	3087	3069
	2240	2242	2237	2236	2239
	1260	1258	1259	1264	1260
	1001	992	990	982	996
	538	523	518	518	516
	155	154	153	153	152
Ave. % Error		0.86	1.11	1.32	1.24

(a) The number in parentheses refers to the corresponding Φ in Table V.

TABLE VIII
OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR MALEONITRILE

Symmetry	$\nu_{\text{obs.}}$	$\nu(1)^{(a)}$	$\nu(2)$	$\nu(3)$	$\nu(4)$
A ₁	3072	3063	3059	3077	3062
	2230	2238	2235	2234	2237
	1597	1592	1585	1588	1582
	1200	1202	1203	1183	1205
	877	884	880	886	883
	480	480	466	465	463
	126	122	119	119	117
B ₂	3072	3067	3063	3080	3065
	2250	2237	2234	2233	2236
	1371	1371	1372	1393	1381
	1012	1010	1011	996	1009
	714	729	715	715	705
	269	279	278	277	279
Av. % error		0.88	1.22	1.49	1.57

(a) The number in parentheses refers to the corresponding Φ in Table V.

TABLE IX
OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR 1,1-DICYANOETHYLENE

Symmetry	$\nu_{\text{obs.}}$	$\nu(4)^{(c)}$	$\nu(7)$
A_1	3036	3033	3036
	2248	2245	2246
	1594	1606	1598
	1395	1405	1391
	720	763 ^(a)	752
	596	589	592
	170	160 ^(a)	160 ^(a)
B_2	3135	3137	3136
	2245	2233	2235
	1256	1259	1262
	939	934	936
	494	519	481
	240	276 ^(a)	266
Av. % Error		0.91 ^(b)	1.63 ^(b)

(a) Not weighted in calculation.

(b) Calculated on basis of weighted frequencies only.

(c) The number in parentheses refers to the corresponding Φ in Table V or Table XV.

TABLE X
OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR TRICYANOETHYLENE

$\nu_{\text{obs.}}$	$\nu(2)^{(a)}$	$\nu(3)$	$\nu(4)$	$\nu(6)$	$\nu(7)$
3043	3061	3079	3064	3043	3043
2249	2243	2244	2245	2246	2246
2234	2235	2238	2237	2235	2236
2210	2230	2233	2232	2229	2232
1594	1607	1606	1606	1581	1584
1327	1330	1353	1333	1329	1332
1171	1169	1189	1163	1181	1176
1000	998	990	1002	995	995
780	794	797	790	778	770
591	589	588	596	600	601
506	521	518	516	498	494
447	461	461	458	442	443
275	278	276	278	273	273
155	158	157	159	158	158
	119	119	117	111	111
Av. % Error	1.02	1.24	0.99	0.72	0.87

(a) The number in parentheses refers to the corresponding Φ in Table V or Table XV.

is from such a calculation and is included only to demonstrate the guidance obtained from preliminary results.

These calculations were originated with the constraints (a) force constants for similar internal coordinates are equal for all members of the series, (b) $C_{CC=CC} = -0.04$, (c) $Tr = 0.06$, (d) $H_{CCC} = 0.237$ and (e) $F_{HCH} = 0.0$. A comparison of the results of calculations (2) and (3) resulted in a modification of constraint (a). The molecules fumaronitrile, maleonitrile and tricyanoethylene considered in (2) contain only a single type of hydrogen, i.e., one alpha to a cyano group. An average error of 9 cm^{-1} was obtained for the C-H stretching frequencies. However, in calculation (3) acrylonitrile, which contains a (CH_2) group, was also considered and the average error in C-H stretching frequencies jumped to 23 cm^{-1} and K_{CH} increased from 4.54 to 4.74 mdyne/A. At this point it became clear that two C-H stretching force constants should be allowed, one for each type of CH bond. This was done in calculation (4). The average error in the ν_{CH} fell back to 9 cm^{-1} , even though 1,1-dicyanoethylene is also considered in this calculation. Further, two well determined K_{CH} force constants emerged: $K_{CH} = 4.56$ for hydrogen alpha to cyano and $K_{CH}^* = 4.94$ for hydrogen in a (CH_2) grouping.

Multiple values for the force constants associated with other internal coordinates would perhaps be warranted. There are, for example, large errors in the calculated frequencies for ν_6 and ν_{17} in fumaronitrile, for ν_6 in maleonitrile and for ν_{17} in 1,1-dicyanoethylene. The potential energy distributions, based on calculation (4), for these molecules are shown in Tables XI-XIII. It is clear that these modes are highly dependent on $F_{C=CC}$ and that there should perhaps be more

than one force constant describing this interaction.

No attempt has been made to redetermine Tr in this investigation since no additional molecules containing trans hydrogen were considered. However, on the basis of potential energy distributions, constraint (b) was eliminated. In calculation (2), $C_{\text{CC}=\text{CC}}$ adopted a value of 0.04 mdyne/A during the least squares analysis and this value was used in subsequent calculations.

Calculations Including Tetracyanoethylene

A zero-order calculation for tetracyanoethylene using the UBFC's given in Table IV has already been reported.²⁰ Since the calculations described above seemed to yield a set of force constants more representative of the series, a new zero-order calculation was performed using $\Phi(4)$. The results of this calculation are given in Table XIV. Since the average error in frequency for the original zero-order calculation is 4.05%, there is only slight improvement on using the $\Phi(4)$. Nevertheless, on the basis of this calculation, ν_{17} and ν_{24} , originally assigned as 426 and 443 cm^{-1} respectively, were switched. This brings ν_{17} more in line with the calculated value of 451 cm^{-1} .

With this additional evidence that the bonding in tetracyanoethylene is not typical of the series, the approach to the vibrational analysis was altered. Rather than attempt to include this molecule in overlay calculations with other members of the series, a minimum number of force constants were transferred from overlay calculation (2) and the remaining constants were then adjusted in a least squares analysis involving only tetracyanoethylene. Thus $C_{\text{CC}=\text{CC}}$ was assigned the value of 0.04 mdyne/A and H_{CCC} , as for this entire study, was fixed

TABLE XI

DISTRIBUTION OF POTENTIAL ENERGY AMONG PLANAR INTERNAL
COORDINATES OF FUMARONITRILE^(a)

Force Constant	ν_1	ν_2	ν_3	ν_4	ν_5	ν_6	ν_7	ν_{13}	ν_{14}	ν_{15}	ν_{16}	ν_{17}	ν_{18}
$K_{C=C}$	0.4	0.2	59.2	19.7	0.9	2.3	0.0	0.0	0.0	0.0	0.0	0.0	0.0
K_{CH}	89.0	0.1	0.0	0.1	0.4	0.0	0.0	88.3	0.1	0.0	0.2	0.2	0.3
K_{CN}	0.0	86.5	1.9	0.0	7.1	3.9	0.5	0.0	85.9	0.4	13.7	0.0	0.0
K_{CC}	0.1	11.4	10.9	0.7	37.9	21.2	2.8	0.1	11.5	4.6	65.8	1.1	0.8
$H_{C=CH}$	0.1	0.1	10.3	27.3	0.0	0.9	0.4	0.0	0.0	33.6	0.7	6.7	4.1
$H_{C=CC}$	0.2	0.0	0.0	3.6	14.7	11.4	6.9	0.0	0.0	0.3	0.0	18.2	18.3
H_{HCC}	0.0	0.0	5.4	6.0	6.4	2.1	1.4	0.0	0.0	19.2	0.4	0.8	1.6
H_{CCN}	0.0	0.0	0.4	0.4	2.1	22.3	74.8	0.0	0.0	0.4	0.0	49.3	50.3
$F_{C=CH}$	3.9	0.0	2.0	26.9	0.0	0.2	0.3	5.5	0.0	19.2	0.5	3.4	2.0
$F_{C=CC}$	0.1	0.8	3.3	2.1	7.1	34.4	11.4	0.0	1.7	1.8	9.3	19.7	20.2
F_{CCH}	6.3	0.8	4.8	8.6	23.4	1.1	1.3	6.1	0.9	26.4	9.6	1.8	3.1
Tr	0.0	0.0	1.8	4.7	0.0	0.2	0.1	0.0	0.0	-5.8	-0.1	-1.1	-0.7

(a) Based on the results of calculation (4).

TABLE XII

DISTRIBUTION OF POTENTIAL ENERGY AMONG PLANAR INTERNAL
COORDINATES OF MALEONITRILE ^(a)

Force Constant	ν_1	ν_2	ν_3	ν_4	ν_5	ν_6	ν_7	ν_{13}	ν_{14}	ν_{15}	ν_{16}	ν_{17}	ν_{18}
$K_{C=C}$	0.4	0.2	70.1	6.0	5.2	0.7	0.1	0.0	0.0	0.0	0.0	0.0	0.0
K_{CH}	88.7	0.1	0.0	0.1	0.2	0.1	0.4	88.5	0.1	0.0	0.4	0.0	0.0
K_{CN}	0.0	86.2	1.4	1.0	11.1	0.2	0.0	0.0	86.3	1.8	7.6	4.1	0.2
K_{CC}	0.1	11.4	8.8	1.7	62.7	0.2	0.2	0.1	11.4	11.9	40.6	18.9	0.9
$H_{C=CH}$	0.0	0.0	4.9	28.8	1.7	2.8	3.2	0.1	0.1	37.0	0.2	4.0	0.4
$H_{C=CC}$	0.0	0.0	0.7	0.0	1.1	12.7	19.2	0.2	0.0	1.3	12.4	17.7	4.6
H_{HCC}	0.0	0.0	4.3	13.7	0.1	1.1	2.3	0.0	0.0	12.4	6.5	1.5	0.8
H_{CCN}	0.0	0.0	0.1	0.0	1.0	56.3	42.1	0.0	0.0	0.0	1.6	13.4	85.0
$F_{C=CH}$	4.6	0.0	2.5	22.4	2.4	1.0	1.7	4.8	0.0	21.1	0.2	2.4	0.2
$F_{C=CC}$	0.1	1.1	2.8	0.2	10.1	19.2	21.5	0.1	1.2	0.7	6.0	37.1	7.1
F_{CCH}	6.1	0.9	4.0	26.1	4.4	1.5	3.7	6.2	0.8	13.7	24.4	0.8	0.8
$C_{CC=CC}$	0.0	0.0	0.0	0.0	0.2	4.1	5.5	0.0	0.0	0.0	0.2	0.0	0.0

(a) Based on the results of calculation (4).

TABLE XIII

DISTRIBUTION OF POTENTIAL ENERGY AMONG PLANAR INTERNAL
COORDINATES OF 1,1-DICYANOETHYLENE ^(a)

Force Constant	ν_1	ν_2	ν_3	ν_4	ν_5	ν_6	ν_7	ν_{13}	ν_{14}	ν_{15}	ν_{16}	ν_{17}	ν_{18}
$K_{C=C}$	0.0	0.2	58.1	17.0	6.0	0.1	0.1	0.0	0.0	0.0	0.0	0.0	0.0
K_{CH}	93.8	0.0	0.5	1.1	0.0	0.0	0.0	95.1	0.0	0.0	0.2	0.2	0.1
K_{CN}	0.0	86.2	0.9	1.1	9.1	2.7	0.1	0.0	85.9	7.6	4.8	1.4	0.2
K_{CC}	0.0	10.1	4.6	4.8	44.3	9.6	0.0	0.0	12.8	42.3	32.8	4.4	0.0
$H_{C=CH}$	0.1	0.0	8.7	17.6	0.3	0.0	0.0	0.1	0.0	22.8	38.2	2.8	0.7
$H_{C=CC}$	0.0	0.0	0.6	0.7	2.2	7.9	5.3	0.1	0.1	9.9	0.4	25.0	11.3
H_{HCH}	0.2	0.0	19.3	39.2	0.7	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
H_{CCN}	0.0	0.0	0.5	0.6	3.7	27.7	67.5	0.0	0.0	0.8	0.2	26.1	72.8
H_{CCC}	0.0	0.0	0.7	0.7	2.5	8.6	5.8	0.0	0.0	0.0	0.0	0.0	0.0
$F_{C=CH}$	5.2	0.0	2.3	16.5	1.3	0.0	0.0	4.6	0.0	12.5	20.8	1.4	0.3
$F_{C=CC}$	0.1	1.1	3.1	0.2	25.5	5.6	6.1	0.1	1.2	4.5	2.7	38.7	14.5
F_{HCH}	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
F_{CCC}	0.0	2.3	0.6	0.7	4.4	37.7	15.3	0.0	-0.1	-0.4	-0.3	0.0	0.0

(a) Based on the results of calculation (4).

TABLE XIV

OBSERVED AND CALCULATED PLANAR FUNDAMENTAL
FREQUENCIES FOR TETRACYANOETHYLENE

Symmetry	$\nu_{\text{obs.}}$	$\nu(4)$	$\nu(5)$	$\nu(6)$	$\nu(7)$	$\nu(8)$
A_g	2235	2243	2248	2243	2246	2230
	1569	1622	1570	1582	1576	1570
	677	617	637	633	632	639
	535	529	515	518	514	516
	130	130	130	130	130	130
B_{1u}	2263	2248	2251	2247	2248	2268
	958	982	982	985	975	979
	579	594	604	599	601	603
	165	170	170	170	170	170
B_{2u}	2230	2236	2240	2237	2237	2241
	1155	1109	1150	1148	1142	1151
	443	451	435	437	437	436
	119	112	104	104	104	103
B_{3g}	2247	2231	2236	2232	2233	2236
	1282	1299	1282	1278	1283	1282
	490	509	481	485	482	481
	254	276	269	270	270	269
Av. % Error		2.88	2.74	2.69	2.76	2.66

at 0.237×10^{-11} erg/rad². This left 7 UBFC's to be determined by 17 frequencies.

The resulting force constants ($\Phi(5)$ in Table XV) do not differ greatly from the $\Phi(4)$. The differences which are present are easily explained. For example, the C=C and C≡N bonds are apparently significantly weaker and the C-C bond much stronger than in the other cyanoethylenes, as would be expected on the basis of increased electron delocalization. A comparison of the $\Phi(1)$ and $\Phi(2)$ indicates that this trend is detectable in tricyanoethylene before becoming critical in tetracyanoethylene. The agreement of the calculated with the observed frequencies for this calculation (Table XIV) is still not striking.

Force constants for overlay calculations involving tetracyanoethylene and tricyanoethylene [$\Phi(6)$] and tetracyanoethylene, tricyanoethylene and 1,1-dicyanoethylene [$\Phi(7)$] also appear in Table XV. Calculated frequencies corresponding to these Φ 's are shown in Tables IX, X and XIV. These calculations were of no significant value in determining either the force constants or the vibrational assignment for tetracyanoethylene. Rather (7) is included because it is the only calculation in which F_{HCH} was allowed to vary. The refined value of 0.02 mdyne/A supports the choice of $F_{\text{HCH}} = 0.0$ as used in the other calculations. Calculation (6) is of some value in confirming that bonding in tricyanoethylene is intermediate to that of tetracyanoethylene and the other four molecules of the series. Furthermore, this calculation hints that the effect of an alpha cyano group on the C-H bond in tricyanoethylene is somehow attenuated, since the C-H stretching constant adopts a value midway between those found for K_{CH} and K_{CH}^* in calculation (4).

TABLE XV

FORCE CONSTANTS AND DISPERSIONS OBTAINED FROM CALCULATIONS
INCLUDING TETRACYANOETHYLENE ^(a)

Force Constant	$\phi(5)$	$\sigma_{\phi(5)}$	$\phi(6)$	$\sigma_{\phi(6)}$	$\phi(7)$	$\sigma_{\phi(7)}$	$\phi(8)$	$\sigma_{\phi(8)}$
$K_{C=C}$	6.26	0.26	6.37	0.16	6.36	0.17	6.34	0.26
K_{CH}	— (c)	—	—	—	5.01	0.07	—	—
K^*_{CH}	—	—	4.62	0.10	4.71	0.11	—	—
K_{CN}	16.76	0.22	16.70	0.14	16.77	0.12	16.88	0.23
K_{CC}	5.22	0.26	5.17	0.19	5.12	0.17	5.19	0.25
$H_{C=CH}$	—	—	0.25	0.09	0.40	0.04	—	—
$H_{C=CC}$	0.31	0.08	0.30	0.06	0.34	0.06	0.31	0.08
H_{HCH}	—	—	—	—	0.42	0.04	—	—
H_{CCH}	—	—	0.40	0.17	0.34	0.09	—	—
H_{CCN}	0.31	0.04	0.31	0.02	0.31	0.02	0.31	0.04
H_{CCC}	0.24	(b)	0.24	(b)	0.24	(b)	0.24	(b)
$F_{C=CH}$	—	—	0.49	(b)	0.29	0.08	—	—
$F_{C=CC}$	0.67	0.19	0.73	0.11	0.65	0.09	0.66	0.18

TABLE XV (Continued)

Force Constant	$\Phi(5)$	$\sigma_{\Phi(5)}$	$\Phi(6)$	$\sigma_{\Phi(6)}$	$\Phi(7)$	$\sigma_{\Phi(7)}$	$\Phi(8)$	$\sigma_{\Phi(8)}$
F_{HCH}	—	—	—	—	0.02	0.06	—	—
F_{CCH}	—	—	0.27	0.16	0.32	0.16	—	—
F_{CCC}	0.97	0.13	0.93	0.08	0.94	0.07	0.96	0.13
$C_{\text{CC=CC}}$	0.04	(b)	0.04	(b)	0.04	(b)	0.04	(b)
Tr	—	—	—	—	—	—	—	—
ρ	—	—	—	—	—	—	0.06	0.05

(a) Force constants are in mdyne/A except for the bending constants, which are in 10^{-11} erg/rad².

(b) The value of the corresponding force constant was constrained.

(c) A line indicates that the calculation is independent of the indicated force constant.

Seemingly the UBFF cannot adequately encompass the effects which the resonance stabilization in tetracyanoethylene has on the vibrational energy. Scherer and Overend¹³ have developed a method whereby the UBFF is modified to correct for resonance effects in aromatic rings. This method should be useful in considering electron delocalization effects in tetracyanoethylene, as will be developed below.

Several canonical forms for tetracyanoethylene are shown in Figure 4. Many more forms could be written on the basis of molecular symmetry. Considering Figure 4, stretching of the bond labeled (1) would tend to stabilize structures (b) and (c) and thus increase the energy necessary to stretch bonds (2), (3), (4) and (5). This can be expressed by the inclusion of cross terms $\rho\Delta R_1\Delta R_2$, $\rho\Delta R_1\Delta R_3$, $\rho\Delta R_1\Delta R_4$, and $\rho\Delta R_1\Delta R_5$ in the potential energy function. Here ρ is the Kekule' constant to be adjusted in the least squares calculation and the ΔR are bond length internal displacement coordinates corresponding to the numbered bonds in Figure 4. Conversely, compression of bond (1) would stabilize structure (a), giving rise to cross terms $\rho(-\Delta R_1)\Delta R_6$, $\rho(-\Delta R_1)\Delta R_8$ and $\rho(-\Delta R_1)\Delta R_9$. These, together with similar terms which arise from consideration of the remaining bond stretching coordinates, comprise a set of cross terms which is added to the basic UBFF. The results of such a calculation appear as $\Phi(8)$ in Table XV and the calculated frequencies are shown in Table XIV. This calculation was performed exactly as was calculation (5) except that the Kekule' constant ρ was added to the set of force constants to be refined. Although there is some improvement in the agreement between calculated and observed frequencies for calculation (8) over calculation (5), the agreement is still not satisfactory. Better results might be had if more

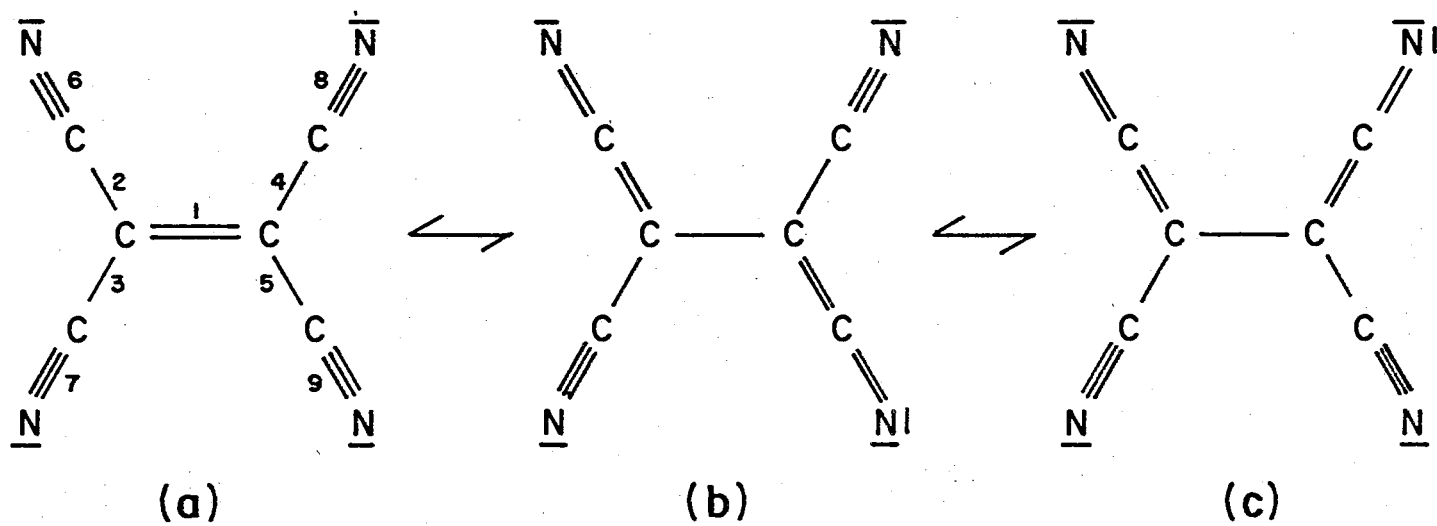


Figure 4. Some Canonical Forms for Tetracyanoethylene

than one Kekule' constant were used. This would remove the inherent assumption that all canonical forms contribute to the bond-bond interactions to the same extent. One calculation which included 3 Kekule' constants would not converge, possibly because, with 9 stretching modes primarily determined by 6 stretching force constants, the problem was too nearly exactly determined.

Nevertheless, the results of calculation (8) indicate that the modification is in the right direction. Compared to the $\Phi(5)$, $K_{C=C}$ and K_{CN} increase slightly while K_{CC} decreases slightly. As would be expected, the bending and nonbonded interaction constants show little or no change.

CHAPTER VI

DISCUSSION

As a result of the overlay calculations reported in the previous chapter, a set of force constants $\Phi(4)$ has been determined which adequately determines the planar fundamental modes of all cyanoethylenes with the exception of tetracyanoethylene. Satisfactory results for tetracyanoethylene were not derived, even when effects of electron delocalization were considered.

Considering the results of calculations (1), (2) and (4), it is surprising that the error in the calculated frequencies is nearly constant regardless of the number of molecules included in the overlay. The average percent errors in calculated frequencies for these three calculations are 0.87, 1.11 and 1.09 respectively. The only difference between calculations (1) and (2) lies in the fact that tricyanoethylene is included in (2) but not in (1). It thus appears again that the inclusion of tricyanoethylene does introduce some bonding features not common to the other four molecules. This effect is undoubtedly a result of increased conjugation.

The average percent error (1.09%) quoted for calculation (4) is somewhat misleading. During that calculation neither ν_5 , ν_7 nor ν_{13} of 1,1-dicyanoethylene were weighted so that error in these modes was not included in the average. If the presently assigned values of 720, 170 and 240 cm^{-1} are used, the average percent error is increased to 1.49

for the entire calculation. Also, it would seem that the results of calculation (4) confirms the assignment of ν_7 as 170 cm^{-1} while favoring the alternate assignment of ν_5 as 762 cm^{-1} and suggesting that ν_{13} , tentatively assigned as 240 cm^{-1} , has not as yet been observed.

Examination of stretching force constants calculated for other small molecules indicates that resonance effects are important even in lower cyanoethylenes. In calculation (4) a value of 6.82 mdyne/A was adopted by $K_{\text{C=C}}$. A similar calculation for ethylene yields a value of 7.46 mdyne/A .²¹ Likewise, for a series of saturated nitriles $K_{\text{C=N}} = 18.2-18.5$ and $K_{\text{C=C}} = 3.1-3.2 \text{ mdyne/A}$ as opposed to 16.94 and 4.63 mdyne/A in this investigation.⁷ It is clear that, for the cyanoethylenes, the C=N and C=C bonds are weaker and the C-CN bond is stronger than would be expected in the absence of delocalization.

It is quite apparent that the C-H bond is significantly weakened by an alpha-cyano group. Consideration of the force constants for cis-, trans- and gem-difluoroethylenes shows the same phenomenon for an alpha-fluorine,³⁴ whereas an alpha-bromine has very little effect.²¹ In fact the weakening of the C-H bond is definitely correlated with the magnitude of the F_{HX} interaction which has the values 1.10 , 0.57 and 0.33 mdyne/A when X is F, CN and Br respectively. It is also known that the C-H bond is lengthened by an adjacent fluorine.³⁵ Further, the relative chemical shifts in the nuclear magnetic resonance spectra of a large number of vinyl compounds show that alpha electronegative groups have the effect on the electron density about hydrogen which is typical of hydrogen bonding.^{36,37} These data, summarized in Table XVI, show that a strong correlation also exists between these electron densities and the magnitude of F_{HX} . These facts indicate the possible

existence of an intramolecular hydrogen bond. If this is so, then the interaction force constant F_{HX} is a measure of the strength of this hydrogen bond, not a measure of the H-X repulsive force.

TABLE XVI
RELATIVE CHEMICAL SHIFTS FOR A SERIES OF VINYL COMPOUNDS

X	$\delta_B - \delta_C$ (a)	$\delta_A - \delta_C$	F_{HX} (b)
F	77	89	1.10
Cl	32	37	0.57
CH ₃	34	37	0.55
Br	25	19	0.33
OCH ₃	91	99	—

(a) Units for chemical shifts are hertz.

(b) Units are mdyne/A.

The positive value adopted by $C_{CC=CC}$ may be indicative of a stabilizing interaction between cis carbon atoms. Thus one might expect ΔH for the cis-trans conversion to be greater than zero, as is the case for the 1,2-difluoroethylenes.³⁸ However, fumaronitrile is

predominant at equilibrium, suggesting that ΔH is less than zero. This indicates that, for the cyanoethylenes, resonance stabilization, which favors the trans configuration, dominates over any stabilizing cis nonbonded interaction.

An alternative explanation for the magnitude of the nonbonded interaction constants is suggested by the work of Bruns et al.³⁹ The results of a semi-empirical molecular orbital calculation for OF_2 raise the possibility that many nonbonded interaction force constants are greatly inflated by electron delocalization. Thus F_{HX} may not be an accurate measure of a nonbonded interaction but may arise mainly from a (CH,CX) bond-bond interaction term. If this is true, then no conclusions concerning the nature of the nonbonded interaction can be drawn from considerations of the sign or magnitude of the force constant since its effect would be effectively masked by that of the bond-bond interaction.

PART II

AN ALGORITHM FOR THE DETERMINATION OF
THE SHAPE PARAMETERS OF INFRARED
ABSORPTION BANDS

CHAPTER I

INTRODUCTION

The Problem

There are basically two types of information which may be derived from an infrared spectrum, the positions of absorption bands and their intensities and shape factors. Band positions are indicative of the energy associated with the vibrational or vibrational-rotational transitions causing the absorption and, because of this, have been used for decades in studies of inter- and intramolecular forces. The work in Part I is but one example of the use of spectroscopic data for this purpose.

Intensity and shape information have not been used extensively as a tool for the interpretation of infrared spectra. Whereas primary frequency standards make it possible to exchange band position data between laboratories, no intensity standards have been developed for use with ordinary laboratory spectrophotometers. Also, observed intensity values are highly dependent on instrumental parameters. As a result, it is almost impossible to compare uncorrected data from different instruments. Finally, any given infrared absorption band, being of finite line width, is often overlapped with an adjacent band. This overlap adds to the difficulty in evaluating the intensity and shape factors for either band.

The purpose of this investigation was to develop a mathematical technique whereby accurate intensity and shape information could be derived from ordinary infrared spectra.

The Shape of Infrared Absorption Bands

Seshadri and Jones⁴⁰ have published a definitive review article concerning the shapes of infrared absorption bands. The discussion in this section is largely based on this review, with additional references as cited.

Consideration of absorption and emission spectra of materials in gas phase indicates that there are three factors which contribute to the line width of spectroscopic transitions. These are radiation damping, the Doppler effect and collision broadening.

According to classical theory,⁴¹ radiation damping arises because a vibrating electric charge radiates, and thus constantly loses, energy. As a consequence the amplitude of vibration is diminished while the natural frequency ν_0 of the oscillator is maintained. A damped vibration of this type is not monochromatic; its frequency distribution may be determined by a Fourier analysis of the electric moment as a function of time. This yields

$$I_{\nu} = \frac{\Delta\nu/2\pi}{(\nu-\nu_0)^2 + (\Delta\nu/2)^2} \quad (1)$$

where I_{ν} is the intensity at frequency ν and $\Delta\nu$ is the line width at half-height in cm^{-1} . The line width is given by

$$\Delta\nu = \frac{4\pi e^2}{3mc^2} \nu_0^2 \quad (2)$$

where m and e are the mass and charge, respectively, of the oscillator and c is the speed of light. In terms of wavelength the line width is

$$\Delta\lambda = \frac{4\pi e^2}{3mc^2} = 1.17 \times 10^{-4} \text{ \AA} \quad (3)$$

According to quantum mechanics, the natural line width (equivalent to radiation damping in the classical treatment) arises because each of the two energy levels E_1 and E_2 , between which a transition occurs, is not sharp but has a finite width ΔE_1 and ΔE_2 respectively. The expression for the half width of the transition is then made up of the term widths of the two states:

$$\Delta\nu = \Delta E_1/h + \Delta E_2/h \quad (4)$$

where h is the Planck constant. Applying the Heisenberg uncertainty principle:

$$\Delta E_1 \cdot \Delta t_1 \sim h/2\pi \quad (5a)$$

$$\Delta E_2 \cdot \Delta t_2 \sim h/2\pi \quad (5b)$$

where Δt_1 and Δt_2 are the average times that the molecule remains in energy states E_1 and E_2 . Under the assumption that only infrared fundamental bands are to be considered; i.e., that all transitions are between the ground and first excited states, the band width can be expressed as

$$\Delta\nu = \frac{4\pi e^2}{mc^2} (\nu_{2,1}^2 f_{2,1}) \quad (6)$$

where $f_{2,1}$ is the oscillator strength of the transition whose frequency is $\nu_{2,1}$.

A comparison of equations 2 and 6 shows that the quantum mechanical treatment differs from the classical approach by a factor of $3f_{1,2}$. The oscillator strengths of infrared transitions are less than unity⁴⁰ so that the value derived from the classical model may be regarded as an upper limit to the true natural line width. Line widths calculated by this method are of the order of 10^{-6} cm^{-1} at 1000 cm^{-1} . Since observed band widths are of the order of $1\text{-}20 \text{ cm}^{-1}$, it may be concluded that radiation damping is relatively unimportant in the evaluation of infrared absorption band shapes.

Doppler broadening results from thermal motion of the molecules. Molecules having a velocity component, V , toward or away from the radiation source will absorb or emit at the frequency ν according to

$$\nu = \nu_0(1 \pm V/c) \quad (7)$$

where ν_0 is the frequency which would be observed for the molecule at rest. The fraction of molecules within a velocity range ∂V is given by

$$\frac{dn}{n} = \left[\frac{M}{2\pi RT} \right]^{\frac{1}{2}} \exp \left[-(M/2RT)v^2 \right] dv \quad (8)$$

where M is the molecular weight of the gas, R is the gas constant and T is the absolute temperature. The total intensity at frequency ν is then

$$I_\nu = \left[\frac{Mc^2}{2\pi RT\nu_0^2} \right]^{\frac{1}{2}} \exp \left[-\left(\frac{Mc^2}{2RT\nu_0^2} \right) (\nu - \nu_0)^2 \right] \quad (9)$$

The Doppler effect thus predicts a Gaussian band shape with a band width given by

$$\Delta\nu = 2(\ln 2)^{\frac{1}{2}} \left[\frac{2RT}{Mc^2} \right]^{\frac{1}{2}} \nu_0 \quad (10)$$

Using equation 10 with $\nu_0 = 1000 \text{ cm}^{-1}$ and a molecular weight of 100 at 20°C , $\Delta\nu$ is calculated to be $1.2 \times 10^{-3} \text{ cm}^{-1}$. This is less than the limits of precision in infrared spectroscopy and it may then be concluded that Doppler broadening is too small to affect band shape measurements.

Collision broadening arises from a perturbation of the energy levels of an absorbing or emitting molecule caused by close approach of other molecules. The effect of such perturbations on a band profile was first derived by Lorentz⁴² and later refined by Van Vleck and Weisskopf.⁴³ Classically, the equation of motion for an oscillator of natural frequency ω_0 radians/sec., mass m and charge e with an electrical field $E \cos \omega t$ is

$$m \frac{d^2x}{dt^2} + \omega_0^2 x = eE \cos \omega t \quad (11)$$

This differential equation has a solution of the form:

$$x = C_1 \exp(i\omega t) + C_2 \exp(i\omega_0 t) + C_3 \exp(-i\omega_0 t) \quad (12)$$

where

$$C_1 = \frac{eE}{m(\omega_0^2 - \omega^2)} \quad (13)$$

and C_2 and C_3 are determined by the nature of the collision. As the result of a collision, x and dx/dt have taken on values which determine the behavior of the oscillator until the occurrence of the next collision. In order to evaluate an average value of x over the

varying time intervals from the last collision, equation 12 is multiplied by the probability that the molecule will collide again after a lapsed time θ . The probability is given by

$$n(t) = \frac{1}{V} \exp\left(-\frac{t}{V}\right) dt \quad (14)$$

where V is the mean time between collisions. The resulting equation for \bar{X} (the average value) is a function containing real and imaginary parts and, for the imaginary part, it can be shown that the absorption coefficient α at frequency ω is

$$\alpha_{\omega} = \frac{2\pi n e^2}{mc} \cdot \frac{\omega}{\omega_0} \left[\frac{1/V}{(\omega - \omega_0)^2 + (1/V)^2} - \frac{1/V}{(\omega + \omega_0)^2 + (1/V)^2} \right] \quad (15)$$

where n is the number of molecules in the system. This equation defines a profile which will be referred to in succeeding portions of this work as a Lorentz profile. On modifying the expression so that it is more useful for the analysis of infrared absorption spectra, equation 15 becomes⁴⁰

$$\alpha_{\nu} = \frac{K}{\pi} \frac{\Delta\nu/2}{(\nu - \nu_0)^2 + (\Delta\nu/2)^2} \quad (16)$$

where K is a function of the particular band. At $\nu = \nu_0$, equation 16 yields

$$\alpha_{\nu_0} = \frac{2K}{\pi\Delta\nu} \quad (17)$$

or

$$\frac{K}{\pi} = \frac{\alpha_{\nu_0} \cdot \Delta\nu}{2} \quad (18)$$

Making the appropriate substitution and designating α_{ν_0} as A_0 , the peak absorbance of the absorbing species, and α_{ν} as A_{ν} , the absorbance at frequency ν , equation 16 may be written

$$A_{\nu} = \frac{A_0 (\Delta\nu/2)^2}{(\nu-\nu_0)^2 + (\Delta\nu/2)^2} \quad (19)$$

Finally, multiplying numerator and denominator by $(2/\Delta\nu)^2$ yields

$$A_{\nu} = \frac{A_0}{1 + (2/\Delta\nu)^2 (\nu-\nu_0)^2} \quad (20)$$

Although the above derivations arise from consideration of the behavior of gases, many workers have used the Lorentz function or modifications of it in studies involving condensed materials.

The Effect of Finite Slit Widths

When a spectrometer is set to transmit radiation of a given frequency a frequency band of finite width is actually transmitted. The energy distribution of this band is generally symmetrical about the nominal frequency and its width is a function of the mechanical slit settings, resolving power and optical aberrations of the spectrometer. The transmittance observed for a sample at any frequency ν' is thus not only dependent on the absorption coefficient at ν' but is also a function of the absorption coefficients across the detected frequency band. If the energy distribution over the detected frequencies is designated as $g(\nu, \nu')$, the apparent transmittance value, $T_{\nu'}^a$, is given by⁴⁴

$$T_{\nu'}^a = \frac{\int T_{\nu} g(\nu, \nu') d\nu}{\int g(\nu, \nu') d\nu} \quad (21)$$

where the convolution is performed over the range of detected frequencies. It is apparent from this equation that the magnitude of error introduced by a symmetric slit function is dependent on $\partial^2 T_\nu / \partial \nu^2$.

Where the slope of the observed spectrum is relatively constant, T_ν^a will approach T_ν . However, where the slope is changing rapidly, as is the case near the center of an absorption band, considerable error results. Moreover, the relative error varies for bands of different intensity^{45,46} giving rise to the curvature often observed in Beer's law plots derived from data taken at large slit widths.

The function $g(\nu, \nu')$ may be visualized by considering the effect of scanning a spectrometer with a finite spectral slit width across a region of the spectrum containing an infinitely narrow absorption band. The signal observed under these conditions could be regarded as the locus of $g(\nu, \nu')$, where ν' is the position of the absorption. For very narrow slits $g(\nu, \nu')$ may be regarded as a diffraction function with a large maximum at ν' . If the slits are wide compared to the wavelength of the transmitted radiation, as is usually the case in infrared spectrometry, the secondary maxima of the diffraction function may be disregarded, and the distribution regarded as approximately triangular,⁴⁴ as shown in Figure 5. The triangular distribution actually applies only to the situation where the entrance and exit slits are of equal size, but this configuration is used in all commercial infrared spectrometers. The function depicted in Figure 5 may be described by the equation

$$g(\nu, \nu') = \begin{cases} S - |\nu' - \nu|, & |\nu' - \nu| \leq S \\ 0 & , |\nu' - \nu| \geq S \end{cases} \quad (22)$$

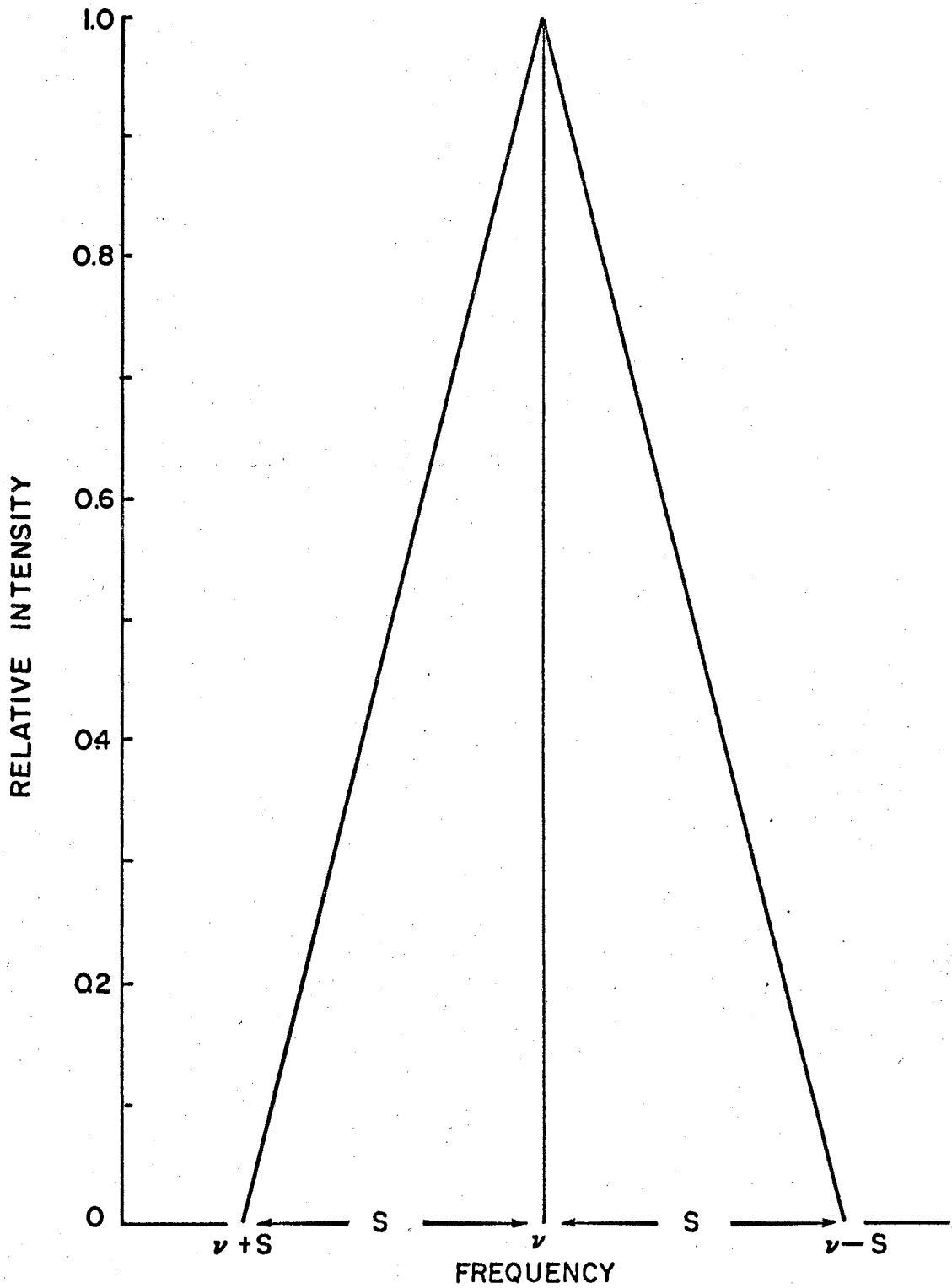


Figure 5. The Triangular Spectral Slit Function.

where S , the spectral slit width, is defined as half the frequency interval subtended by the spectrometer when set to ν' .

Studies of Infrared Band Profiles

Most of the significant work in the study of the shapes of infrared absorption bands has been done by Jones and co-workers.^{40,47-49} Their study has included the Lorentz and Gauss functions and the sum and product functions of these two. At present, Jones is using the Lorentz-Gauss sum function almost exclusively.⁵⁰ It should be pointed out that the primary object of this group is to reduce infrared data to a minimal number of numerical indices, not to extract true intensity data.

Abramowitz and Bauman⁵¹ and Fraser and Suzuki⁵² have also used Lorentz-Gauss sum and product functions, but the former have concluded that, at narrow slit widths, most absorption bands are almost purely Lorentzian. Early work in this laboratory indicates that this conclusion is correct. In a study of the hydrogen bonding of esters and alcohols, overlapped carbonyl stretching absorption bands were successfully resolved on the basis of a Lorentz band shape.^{53,54} The same algorithm was used to resolve carbon-hydrogen stretching bands in the spectra of aliphatic molecules.⁵⁵

The Scope of the Present Investigation

For the purposes of this investigation infrared absorption bands are regarded as Lorentz functions convolved by a triangular slit function. This mathematical model was chosen because (1) unlike the earlier models used in this laboratory and elsewhere, the effect of

finite slits is included as a separable factor and (2) it is felt, in light of the above discussion of the factors influencing band shape, that the parameters derived from this model would have more physical significance than those calculated using more complex models such as the Lorentz-Gauss sum function.

A computer program was designed and written for the purpose of performing the calculations necessary for the analysis. Several spectra, containing varying amounts of noise, were synthesized and used in testing the convergence characteristics of the program.

The experimental phase of this study consisted of the acquisition of data from infrared spectra of several esters and simple aromatic compounds. The former were chosen because their carbonyl stretching bands are well isolated from other bands while the latter provided a system of relatively narrow out-of-plane hydrogen bending bands which are highly susceptible to slit broadening. In addition, the various aromatic compounds provided the opportunity to determine the effectiveness of the model in the separation of overlapping bands.

The infrared data was then fitted to the model described above. The results clearly show that the slit-convolved Lorentz function is quite adequate to fit the observed data. In addition, analysis of data taken at several different slit widths indicate that the mathematical analysis is reasonably effective in removing distortion from this cause.

CHAPTER II

DEVELOPMENT OF THE ALGORITHM

The Calculated Spectrum

An infrared absorption spectrum containing one or more bands may be regarded, in the absence of slit perturbations, as a summation of functions of the type given in equation 20. Under this assumption, the absorbance calculated for any frequency ν is then

$$A_{\nu}^0 = \sum_{i=1}^N \frac{A_{0i}}{1 + (2/\Delta\nu_i)^2 (\nu - \nu_{0i})^2} + \beta \quad (23)$$

where A_{0i} , $\Delta\nu_i$, and ν_{0i} are the peak absorbance, line width and band position, respectively, for the i -th band and N is the number of bands contributing to the spectrum. The zero superscript on the left hand side denotes that the absorbance has been calculated under the assumption of zero slit width. The constant β is included to absorb small baseline errors.

The basic physical quantity measured by most commercial spectrometers is transmittance. It is thus more meaningful to perform the band shape analysis in terms of this variable. Also, for reasons which will become clear later, it is more convenient to regard the spectrum not as a continuous function but as a vector composed of discrete samples taken at equal increments of the independent variable (frequency).

The transmittance calculated for the k-th data point of the calculated spectrum is thus

$$T_k^o = \exp \left[- \sum_{i=1}^N \frac{A'_{0i}}{1 + \rho_i (v_k - v_{0i})^2} - \beta' \right] \quad (24)$$

where $A'_{0i} = 2.303 A_{0i}$, $\beta' = 2.303 \beta$, $\rho_i = (2/\Delta v_i)^2$ and v_k is the frequency corresponding to the k-th datum.

In terms of discrete samples, the expression for the slit function (equation 22) becomes

$$g(v, v_k) = \begin{cases} S - |v_k - v_{k+j}| = S - jx; & j = 0, 1, 2, \dots, m \\ 0, & j \geq m \end{cases} \quad (25)$$

where m is the number of samples contained in $S \text{ cm}^{-1}$ and x is the sampling interval. The slit convolution can then be expressed as a summation

$$T_k^{\text{calc}} = a_0 T_k^o + \sum_{j=1}^m a_j (T_{k+j}^o + T_{k-j}^o) \quad (26)$$

where T_k^{calc} is the value calculated for the k-th datum. The coefficients a_j are easily shown from equation 25 to be

$$a_j = \frac{S - jx}{F}; \quad j = 0, 1, 2, \dots, m-1 \quad (27a)$$

$$a_m = \frac{S - mx}{2F} \quad (27b)$$

where F is the normalization factor given by

$$F = 2mS - m^2 x \quad (28)$$

The factor of 2 is included in the numerator of equation 27b to make the numerical integration in equation 26 conform to the trapezoidal rule.

An alternative model for the slit convolution may be developed if it is assumed that any $2 S \text{ cm}^{-1}$ region of the calculated spectrum is well approximated by a simple polynomial in $S - (v_k - v)$, such as

$$T_Z^0 = \sum_{i=0}^m a_i Z^i, \quad -S \leq Z \leq S \quad (29)$$

where $Z = S - (v_k - v)$ and M is the degree of the polynomial. The convolution is then

$$T_k^{\text{calc}} = \frac{\int_0^S T_Z^0 (S-z) \partial Z + \int_{-S}^0 T_Z^0 (S+Z) \partial Z}{\int_0^S (S-Z) \partial Z + \int_{-S}^0 (S+Z) \partial Z} \quad (30)$$

Substituting the expression for T_Z^0 , the integrals are easily evaluated and the expression for T_k^{calc} becomes

$$T_k^{\text{calc}} = 2 \sum_{i=0}^m \delta_i \left(\frac{1}{i+1} - \frac{1}{i+2} \right) a_i S^i \quad (31)$$

where δ_i is unity for even i and zero for odd i . The a_i are readily determined from the T_k^0 by use of the simplified least-squares procedure outlined by Savitsky.⁵⁶ The computer program developed in this investigation allows the use of either slit convolution model.

The Least-Squares Procedure

The object of this algorithm is to choose a set of $3N+1$ band parameters A_{0i} , $\Delta\nu_i$, ν_{0i} , and β for a spectrum containing N bands such that the best fit is obtained between the observed and calculated transmittance values. According to the least-squares criterion, this is achieved when the sum of squares of the deviations is a minimum:

$$\sum_{k=1}^P (T_k^{\text{obs}} - T_k^{\text{calc}})^2 = \text{minimum} \quad (32)$$

where P is the number of elements in the vector of observed transmittance values. Since the T_k^{calc} are not linearly related to the band parameters, the Gauss-Newton nonlinear regression method was chosen for the analysis.⁵⁷

For the sake of convenience, the various band parameters are now regarded as elements of a parameter vector whose elements are Q_i , where $i=1,2,\dots,3N+1$. A set of initial parameter values, Q_i' , are estimated and used to calculate a vector of the T_k^{calc} . If the initial estimate is reasonably close to the true value, the residual function may be linearized by using a first order Taylor's expansion

$$R_k' = (T_k^{\text{obs}} - T_k^{\text{calc}})' = \sum_{i=1}^{3N+1} \left(\frac{\partial T_k}{\partial Q_i} \right)' \Delta Q_i \quad (33)$$

where $\Delta Q_i = Q_i - Q_i'$ and the primes indicate that the indicated terms are to be evaluated for the approximate values of the parameters. This equation is linear and may be solved for the ΔQ_i by ordinary least-

squares methods, viz.:

$$\overline{\Delta Q} = B^{-1} \overline{Z} \quad (34)$$

where $\overline{\Delta Q}$ is a vector of the ΔQ_i , B^{-1} is the inverse of a matrix B whose i,j-th element is

$$B_{ij} = \sum_{k=1}^P \left(\frac{\partial T_k^{\text{calc}}}{\partial Q_i} \right) \left(\frac{\partial T_k^{\text{calc}}}{\partial Q_j} \right) \quad (35)$$

and \overline{Z} is a vector whose i-th component is

$$Z_i = \sum_{k=1}^P \left(R_k \frac{\partial T_k^{\text{calc}}}{\partial Q_i} \right) \quad (36)$$

The Q_i 's are adjusted by the ΔQ_i and the process is repeated until minimum for equation 32 is obtained.

A computer program (BANDFIT) has been written to implement the algorithm outlined in this chapter. The digitized infrared spectrum together with a set of band parameter estimates and control information comprise the input data for this routine. From the parameters and the specified spectral slit function, a calculated spectrum is generated and the error sum of squares is evaluated. The partial derivative sums necessary for the construction of matrix B are calculated at this time. Matrix B is then inverted, the ΔQ 's are evaluated and the original estimate set is updated. The problem is then recycled using the new parameter set. This process is repeated until either a minimum in the error sum of squares has been reached or the fractional improvement in the sum of squares for two successive cycles falls below

the limit set by the user. A listing of program BANDFIT and instructions for its use may be found in the Appendix.

CHAPTER III

EXPERIMENTAL

General

All materials, with the exception of solvents, used in this investigation were purified by preparative gas chromatography before use. Solvents were used without further purification.

Infrared sampling cells were calibrated in the normal manner from their diffraction pattern. Standard deviations in these calibrations were determined to be less than 0.4 percent relative to the path length.

The true zero reading for the slit micrometer of the infrared spectrometer was determined by plotting the square root of single beam energy vs. the micrometer index. The intercept of this plot was taken as the zero offset error and was used to correct subsequent micrometer readings.

Data Acquisition

All spectra were determined using a Perkin-Elmer Model 621 double-beam infrared grating spectrometer. The spectrometer was equipped with an encoding device so that digitization of transmittance values could be commanded at equal frequency intervals. These intervals could be varied from 0.1 to 10.0 cm^{-1} . All data for this investigation were

acquired at intervals of either 0.5 or 0.2 cm^{-1} . The transmittance data were recorded as 3-digit numbers on computer-compatible magnetic tape for processing at a later time.

The scanning rate for all spectra was kept below 1 $\text{cm}^{-1}/\text{sec}$. This was done to minimize servo errors which might become important at higher rates.

Solution and solvent spectra were determined separately. During computer processing each solution spectrum was divided by the appropriate solvent spectrum to yield the equivalent of a differential spectrum determined with perfectly matched cells. This was done to cancel out all effects arising from cell aberrations and differences between the refractive indices of the salt plates and the solution.

Evaluation of the Spectral Slit Width

The equation given by Seshadri and Jones⁴⁰ was used to determine spectral slit width values in this investigation. For a spectrometer set at frequency ν with a mechanical slit width w_s , S is given by

$$S = \frac{\nu^2 d}{NnF} \left[1 - \left(\frac{n}{2\nu d} \right)^2 \right]^{\frac{1}{2}} \left\{ w_s + \left[\left(\frac{F}{B\nu} \right)^2 + w_A^2 \right]^{\frac{1}{2}} \right\} \quad (37)$$

where d is the grating spacing, N is the number of grating passes, n is the order of the spectrum, F is the focal length of the monochromator, B is the limiting aperture of the spectrometer and w_A is the virtual mechanical slit width arising from aberrations. All terms save B and w_A were available from either the instrument settings or the manufacturer's literature. The effective aperture was determined from a measurement of the illuminated portion of the grating and w_A was

assigned the value of 0.004 cm as recommended by Seshadri and Jones.

A constant physical slit width was employed for spectra taken in the region 1950-1550 cm^{-1} since this procedure gives more flexibility in choosing a spectral slit width than does the instrumental slit program and since $\partial S/\partial \nu$ for a constant w_s is small in this region. For spectra in the 800-650 cm^{-1} , however, use of the slit program was necessary because the size of the derivative precluded the use of fixed slits.

CHAPTER IV

RESULTS AND DISCUSSION

Synthetic Spectra

As was mentioned previously, synthetic spectra were used to evaluate the band shape analysis computer program. Before discussing the results of these analyses, it would be appropriate to describe the method used to generate the spectra. For each synthetic spectrum, a vector of transmittance values was generated using equations 24 and 26. A "noise" vector was also generated by indexing a table of random normal deviates with a pseudo-random number generator. This vector was then scaled to the desired signal-to-noise ratio (S/N) and combined with the transmittance vector to yield the synthetic spectrum. Before analysis, each datum was truncated to three decimal places so as to conform to the data format of the real spectrometer.

Three sets of synthetic spectra were analyzed using program BANDFIT. The first set was composed of 24 spectra, each containing a single band with $A_0 = 1.0$, $\Delta\nu = 1.0$, $\Delta\nu = 10.0 \text{ cm}^{-1}$ and $\nu_0 = 1000 \text{ cm}^{-1}$. Transmittance vectors spanned the range $1050\text{-}950 \text{ cm}^{-1}$ with a sampling interval of 0.5 cm^{-1} . Results of the least-squares procedure appear in Table XVII, and a graph of the synthetic spectrum and that calculated on the basis of the refined parameters for the case where $S = 4.0 \text{ cm}^{-1}$ and $S/N = 200$ is shown in Figure 6. Refined values for ν_0 and β do not

TABLE XVII

RESULTS FROM ANALYSES OF SYNTHETIC
SPECTRA CONTAINING ONE BAND

Spectral Slit Width, cm^{-1}	S/N	Refined Parameters		Error Sum of ₂ Squares x 10^2	Least-Squares Cycles
		A_0	$\Delta\nu, \text{cm}^{-1}$		
0.	400	1.0012	9.98	0.105	3
	200	1.0034	9.95	0.410	3
	100	1.0069	9.90	1.615	3
	50	1.0142	9.80	6.397	3
1.	400	1.0018	9.97	0.099	5
	200	1.0035	9.95	0.407	4
	100	1.0080	9.89	1.621	4
	50	1.0160	9.78	6.420	4
2.	400	1.0019	9.97	0.102	7
	200	1.0044	9.94	0.401	6
	100	1.0081	9.89	1.603	6
	50	1.0159	9.78	6.352	5
4.	400	1.0016	9.98	0.103	10
	200	1.0042	9.95	0.394	9
	100	1.0080	9.90	1.563	8
	50	1.0162	9.79	6.231	7
7.	400	1.0012	9.98	0.092	12
	200	1.0018	9.97	0.361	11
	100	1.0043	9.94	1.435	10
	50	1.0082	9.88	5.698	9
10.	400	0.9994	10.00	0.081	15
	200	1.0009	9.98	0.320	14
	100	1.0007	9.97	1.286	13
	50	1.0010	9.93	5.075	12

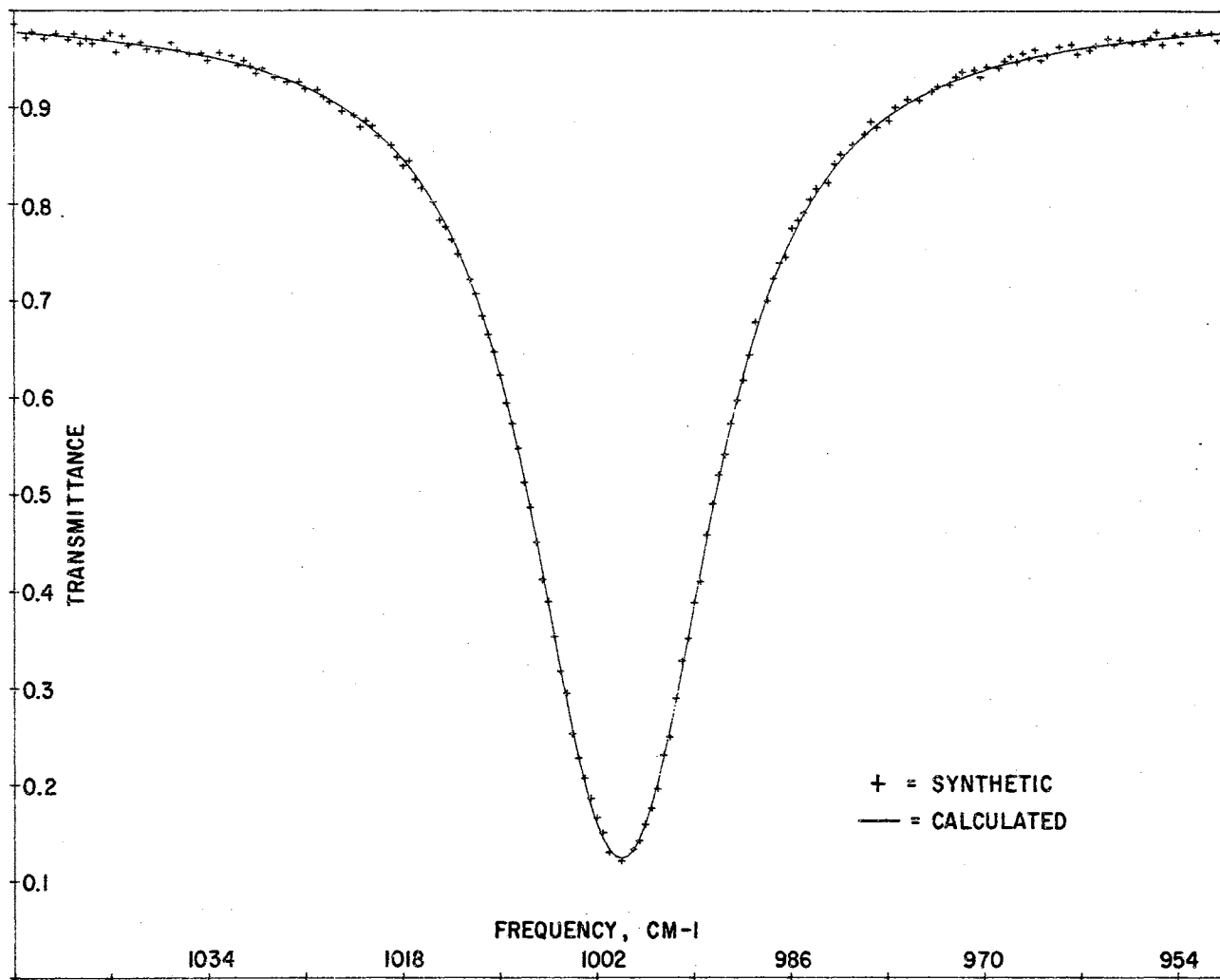


Figure 6. Single-Band Synthetic Spectrum with 4 cm^{-1} Spectral Slit

appear in the table since ν_0 converged to the true value in all cases and β was constrained at zero. For the most part, the other two parameters converged to reasonably accurate values. There does, however, appear to be a relationship between S/N and the calculated values for A_0 and $\Delta\nu$. The reason for this dependence is not completely understood, but it is believed to arise from a local bias in the noise vector. This is plausible since a single vector, scaled to give the various S/N values, was used for all spectra in the set. An examination of the noise pattern indicated some bias in the vicinity of the band center but it was not of sufficient magnitude to be responsible for the observed deviations. In any case, the effect is probably due to an artifact and is unimportant at moderate noise levels.

The relationship between the spectral slit width and the number of cycles taken in the analysis is a direct result of the fact that initial estimates of the parameter values were taken from the simulated spectrum. For wide spectral slits, therefore, the estimates are farther from the true values than for narrow slits and the least-squares procedure consequently requires more iterations to converge. Also, at smaller values of S/N, the error sum of squares levels off sooner because of the uncertainty introduced by the noise. This is the cause of the apparent relationship between S/N and the number of calculation cycles.

Five of the spectra referred to in Table XVII (those with S/N = 200) were re-analyzed under the assumption that $S = 0$; that is, that the spectra could be fitted by an unperturbed Lorentz function. Although the results of these calculations (Table XVIII) indicate that, for spectral slit width values of 4 cm^{-1} and below, the error sums of

TABLE XVIII
RESULTS OF ZERO-SLIT CALCULATIONS^a

Spectral Slit Width, cm^{-1}	Refined Parameters			Error Sum of Squares $\times 10^2$	Least-Squares Cycles
	A_0	$\Delta\nu, \text{cm}^{-1}$	β		
1.	0.9991	9.98	-0.0001	0.407	3
2.	0.9818	10.14	-0.0003	0.402	3
4.	0.9169	10.77	-0.0015	0.458	3
7.	0.7895	12.37	-0.0042	0.862	4
10.	0.6702	14.54	-0.0078	1.578	4

a) S/N = 200 for all spectra.

squares are comparable with those reported in Table XVII, it is apparent that A_0 and $\Delta\nu$ have converged toward the apparent parameter values rather than the true values. Thus, a simple Lorentz function may be used to fit experimental data taken at narrow spectral slit widths but the calculated band parameters derived will not be independent of the instrument slit width.

The second set of 24 synthetic spectra considered each contain two bands with $A_{01} = 1.0$, $A_{02} = 0.5$, $\Delta\nu_1 = \Delta\nu_2 = 10.0 \text{ cm}^{-1}$, $\nu_{01} = 1010 \text{ cm}^{-1}$ and $\nu_{02} = 990 \text{ cm}^{-1}$. For these spectra the range was 1060-940 cm^{-1} with a sampling interval of 0.5 cm^{-1} . Results of the least-squares analyses for these spectra, shown in Table XIX, display the same types of trends and dependencies as were observed for the set of single band spectra. Again, the ν_0 values are not listed since these parameters converged to the theoretical values in all cases. A plot of synthetic and calculated spectra for a representative member of the set appears in Figure 7.

A third set of synthetic spectra, each consisting of three bands, was also subjected to band shape analysis. The parameters used in the synthesis of the spectra are listed in Table XX. The results of the analyses for bands 1 and 2 are similar to those already discussed. The refined parameter values for the third band are noteworthy, however, and appear in Table XXI. Even though this band appears as an unresolved shoulder on the 1000 cm^{-1} band (see Figure 8) the least-squares analysis resulted in reasonably accurate parameter values in most cases. Exceptions are for those spectra generated with large spectral slit widths and/or high noise levels. Since the presence of band 3 was not obvious in plots of the 10 cm^{-1} spectral slit spectra, these

TABLE XIX

RESULTS FROM ANALYSES OF SYNTHETIC SPECTRA CONTAINING TWO BANDS

Spectral Slit Width, cm^{-1}	S/N	Refined Parameters				Error Sum of Squares $\times 10^2$	Least-Squares Cycles
		A_{01}	$\Delta\nu_1, \text{cm}^{-1}$	A_{02}	$\Delta\nu_2, \text{cm}^{-1}$		
0.	400	1.0038	9.95	0.5002	10.04	0.107	4
	200	1.0071	9.90	0.5002	10.07	0.441	4
	100	1.0149	9.81	0.5005	10.14	1.784	4
	50	1.0304	9.61	0.5010	10.27	6.880	4
1.	400	1.0037	9.95	0.5002	10.03	0.114	8
	200	1.0065	9.92	0.5007	10.04	0.447	8
	100	1.0136	9.83	0.5010	10.10	1.761	7
	50	1.0271	9.66	0.5021	10.20	6.844	6
2.	400	1.0022	9.97	0.5005	10.01	0.113	10
	200	1.0055	9.93	0.5007	10.03	0.433	9
	100	1.0106	9.87	0.5015	10.06	1.738	8
	50	1.0214	9.74	0.5031	10.13	6.803	7
4.	400	1.0001	10.00	0.5006	9.99	0.107	11
	200	1.0008	9.99	0.5010	9.99	0.421	10
	100	1.0016	9.98	0.5023	9.97	1.685	9
	50	1.0037	9.96	0.5041	9.94	6.597	9
7.	400	0.9970	10.03	0.4998	9.99	0.097	13
	200	0.9941	10.07	0.4997	9.97	0.383	12
	100	0.9881	10.13	0.4999	9.93	1.537	11
	50	0.9768	10.26	0.5001	9.85	5.962	10

TABLE XIX (Continued)

Spectral Slit Width, cm^{-1}	S/N	Refined Parameters				Error Sum of Squares $\times 10^2$	Least-Squares Cycles
		A_{01}	$\Delta\nu_1, \text{cm}^{-1}$	A_{02}	$\Delta\nu_2, \text{cm}^{-1}$		
10.	400	0.9966	10.02	0.4978	10.05	0.085	17
	200	0.9937	10.04	0.4963	10.09	0.332	15
	100	0.9876	10.07	0.4915	10.21	1.331	14
	50	0.9743	10.17	0.4850	10.37	5.293	13

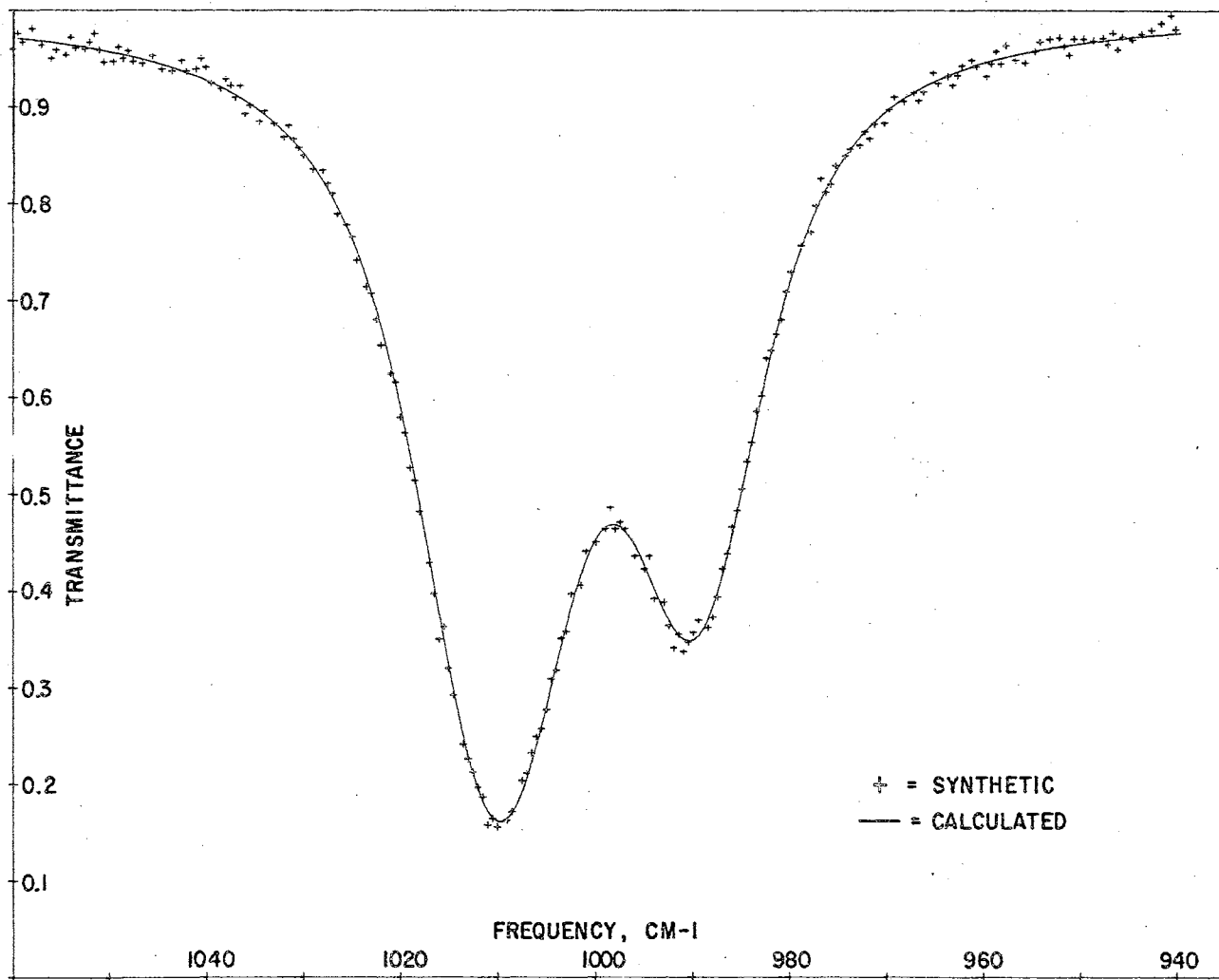


Figure 7. Two-Band Synthetic Spectrum with 7 cm^{-1}
Spectral Slit and $S/N = 100$

TABLE XX
BAND PARAMETERS FOR SYNTHETIC SPECTRA
CONTAINING THREE BANDS^a

<u>Band</u>	<u>A₀</u>	<u>Δν, cm⁻¹</u>	<u>ν₀, cm⁻¹</u>
1	0.4	15.0	1020
2	1.0	10.0	1000
3	0.1	10.0	985

a) Spectra consist of 271 data points over the region 1070-935 cm⁻¹.

TABLE XXI

REFINED PARAMETERS FOR BAND 3 OF THE
THREE-BAND SYNTHETIC SPECTRA

Spectral Slit Width, cm^{-1}	S/N	A_{03}	$\Delta\nu_3$, cm^{-1}	ν_{03} , cm^{-1}
0.	400	0.0991	10.06	985.0
	200	0.0979	10.13	984.9
	100	0.0960	10.19	984.8
	50	0.0919	10.39	984.7
1.	400	0.0993	10.11	985.0
	200	0.0982	10.25	985.0
	100	0.0964	10.53	985.0
	50	0.0929	11.11	985.0
2.	400	0.0993	10.19	985.0
	200	0.0987	10.35	985.1
	100	0.0976	10.71	985.1
	50	0.0951	11.49	985.2
4.	400	0.0999	10.18	985.1
	200	0.0999	10.42	985.1
	100	0.1001	10.82	985.3
	50	0.1008	11.81	985.6
7.	400	0.1007	10.26	985.1
	200	0.1013	10.47	985.2
	100	0.1030	10.91	985.3
	50	0.1072	11.89	985.7
10.	400	0.0972	9.29	984.6
	200	0.0961	9.11	984.5
	100	0.0916	7.67	984.0
	50	0.0922	5.88	983.5

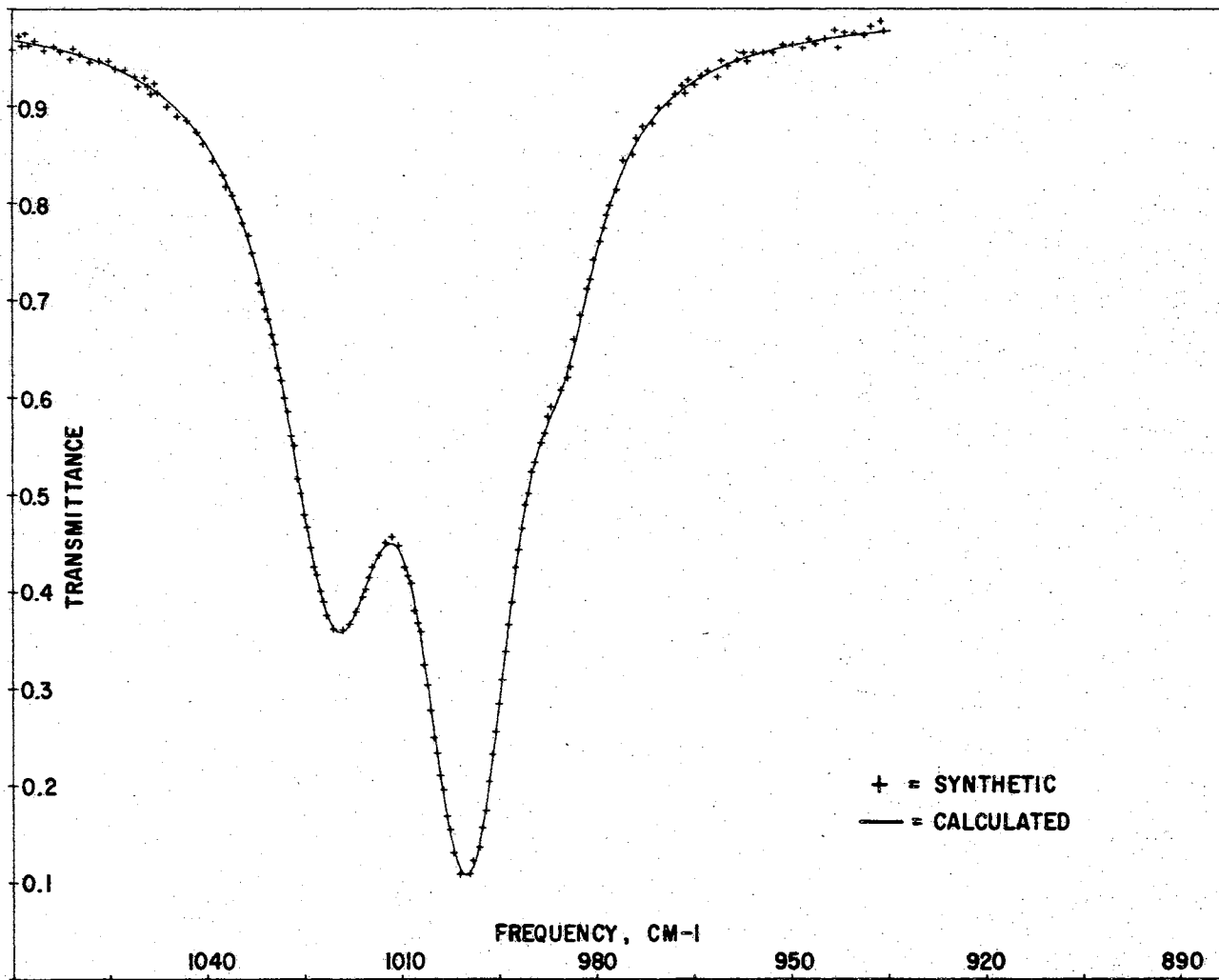


Figure 8. Three-Band Synthetic Spectrum with 4 cm^{-1}
Spectral Slit and S/N = 200

spectra were analyzed again under the assumption that only two bands were present in the spectrum. The result of one such analysis is shown in Figure 9. The difference between the synthetic and calculated curves makes it apparent that an error of this type would not go undetected in an actual analysis.

The algorithm developed in this investigation has thus been proved capable of evaluating the true band parameters for a variety of synthetic spectra. It must be noted, however, that these spectra are ideal cases in that a Lorentzian band shape, triangular slit function and normally distributed noise are presumed. The following two sections of this chapter examine the applicability of the algorithm to the analysis of real spectra.

Carbonyl Stretching Bands of Aliphatic Esters

Four solutions of methyl oleate in *n*-tridecane were prepared and their infrared spectra determined at two different spectral slit widths over the 1800-1700 cm^{-1} region. Data from these spectra were analyzed in three phases. In the first phase, the data taken within 10 cm^{-1} of the carbonyl stretching band were fitted to a single band model. This was done to diminish the effect of any possible unresolved bands on the calculations. The refined parameters from these analyses were then used to calculate a spectrum over the entire 1800-1700 cm^{-1} range. Differences between the observed and calculated curves were examined to determine if further bands should be incorporated in the model. Finally, least-squares calculations based on the revised estimates were performed.

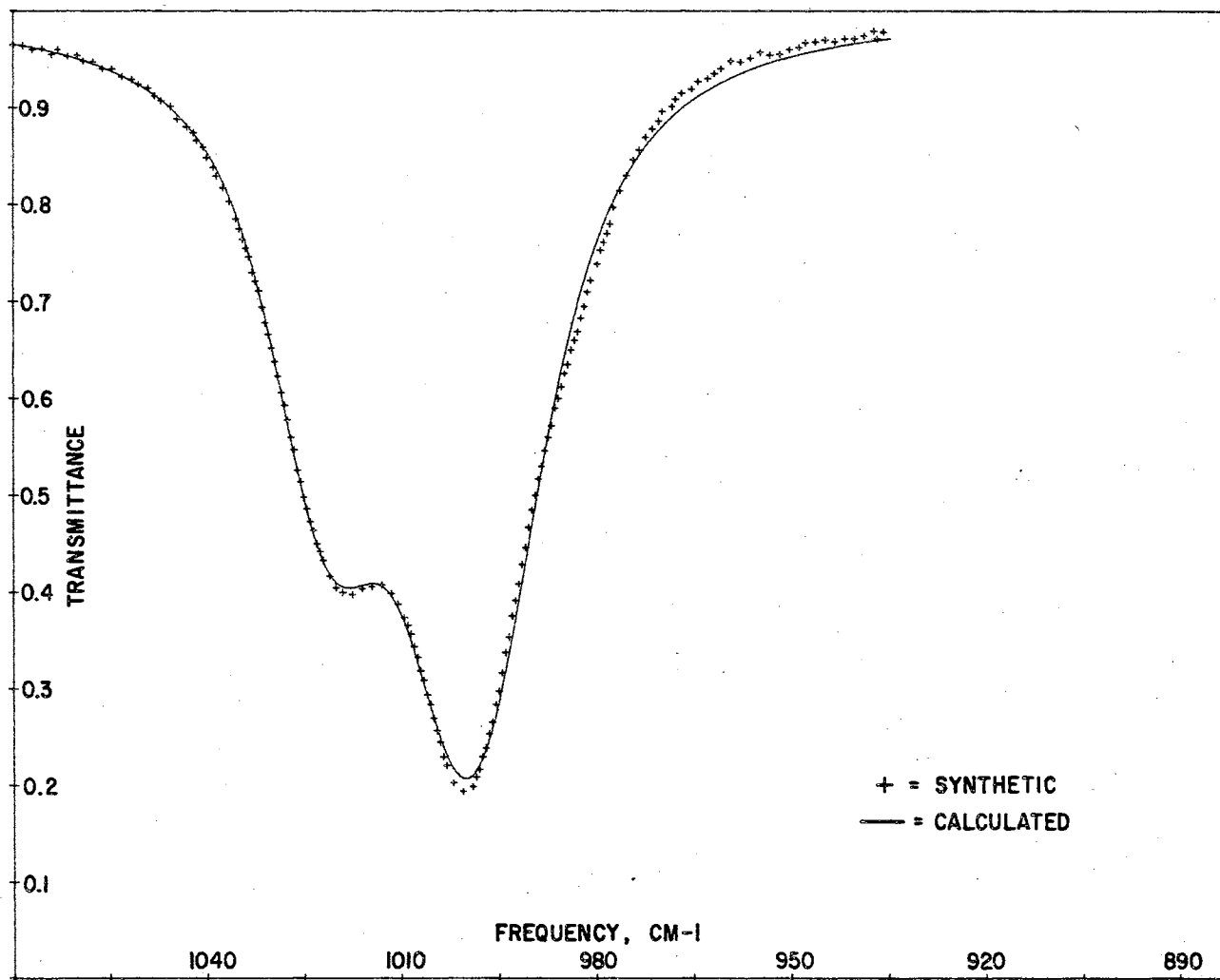


Figure 9. Result of Fitting a Two-Band Function to a Three-Band Synthetic Spectrum

Results for the analyses involving data in the vicinity of the methyl oleate carbonyl stretching band are shown in Table XXII. On the basis of the observed intensities, the molar absorptivity (ϵ) for the data at each slit value was calculated. Thus, with 1.5 cm^{-1} slits $\epsilon = (57.70 \pm 0.36) \text{ M}^{-1} \text{ mm}^{-1}$ while for the data taken at 3.2 cm^{-1} , $\epsilon = (55.99 \pm 0.36) \text{ M}^{-1} \text{ mm}^{-1}$. Using the refined intensities, however, the molar absorptivities are (59.87 ± 0.39) and $(60.84 \pm 0.68) \text{ M}^{-1} \text{ mm}^{-1}$, respectively. Although no "true" value exists for comparison, the indications are that the molar absorptivities based on the refined parameters tend toward a single, and presumably correct, value. Also, there is a distinct improvement in the refined $\Delta\nu$ values over those observed. The mean values for the measured band widths are $(11.66 \pm 0.04) \text{ cm}^{-1}$ and $(12.17 \pm 0.07) \text{ cm}^{-1}$, respectively, for the 1.5 cm^{-1} and 3.2 cm^{-1} data while the corresponding means for the refined parameters are (11.34 ± 0.06) and $(11.28 \pm 0.05) \text{ cm}^{-1}$.

Montigny⁴⁶ has published a table of coefficients whereby, for given values of apparent band intensity, apparent band width and spectral slit width, the ratio of apparent to true band width for a slit-broadened infrared absorption band may be evaluated. Using these coefficients, expected values for the refined $\Delta\nu$'s were calculated and appear as the rightmost column of Table XXII. The mean deviation between the refined and calculated band widths is 1.4 percent relative to the calculated value, certainly within acceptable limits for this parameter. Since Montigny's calculations were based on the assumption that an absorption band may be regarded as the convolution of a Lorentz function with a triangular slit function, the agreement observed here serves as justification for the choice of this model in the present

TABLE XXII
 BAND-SHAPE PARAMETERS FOR THE CARBONYL STRETCHING
 BAND OF METHYL OLEATE^a

Spectral Slit Width, cm ⁻¹	Concentration M.	Observed Parameters			Refined Parameters			Calculated Values for $\Delta\nu$, cm ⁻¹
		A ₀	$\Delta\nu$, cm ⁻¹	ν_0 , cm ⁻¹	A ₀	$\Delta\nu$, cm ⁻¹	ν_0 , cm ⁻¹	
1.5	0.00976	0.2786	11.73	1747.2	0.2846	11.42	1747.5	11.55
	0.01294	0.3642	11.66	1747.3	0.3750	11.39	1747.5	11.48
	0.01667	0.4694	11.68	1747.4	0.4843	11.32	1747.6	11.51
	0.01966	0.5513	11.58	1747.3	0.5675	11.25	1747.5	11.42
3.2	0.00976	0.2700	12.26	1747.0	0.2866	11.34	1747.3	11.51
	0.01294	0.3558	12.22	1747.1	0.3778	11.32	1747.4	11.48
	0.01667	0.4537	12.07	1747.2	0.4906	11.21	1747.4	11.37
	0.01966	0.5362	12.13	1747.0	0.5734	11.26	1747.3	11.45

a) Sample path for all spectra was 0.4788 mm.

investigation.

The differences between the observed and refined band position values are the result of a small amount of skew in the observed absorption bands. Because of this, the center of the band calculated by the least squares does not correspond exactly to the position of zero slope.

Spectra calculated using the refined parameters from Table XXII were compared with the observed spectrum over the 1800-1700 cm^{-1} region. One such comparison is shown in Figure 10. Differences between the curves indicate that two more bands, one at 1726.5 cm^{-1} and one at 1706.5 cm^{-1} , should be included in the model. This was done and the calculations repeated. The results for the expanded model were less than satisfactory in that the parameters calculated for the additional bands varied markedly from spectrum to spectrum and the calculated band widths often assumed unrealistically large (50-100 cm^{-1}) values. As would be expected, however, inclusion of the additional bands did result in a better fit between observed and calculated spectra. The error sums of squares for the three-band model were lower by a factor of about 30 than equivalent calculations based on a single band. This behavior could indicate that another factor, such as spectrometer imbalance in a region of atmospheric absorbance, is instrumental in causing the deviations. The fact that all carbonyl stretching spectra examined in this investigation exhibited some structure in this region lends credence to this explanation.

Data from the carbonyl stretching region of spectra of ethyl acetate in n-tridecane were analyzed in a slightly different fashion. Preliminary calculations indicated the presence of a small band at 1729 cm^{-1} in addition to the carbonyl absorbance at 1748 cm^{-1} . The

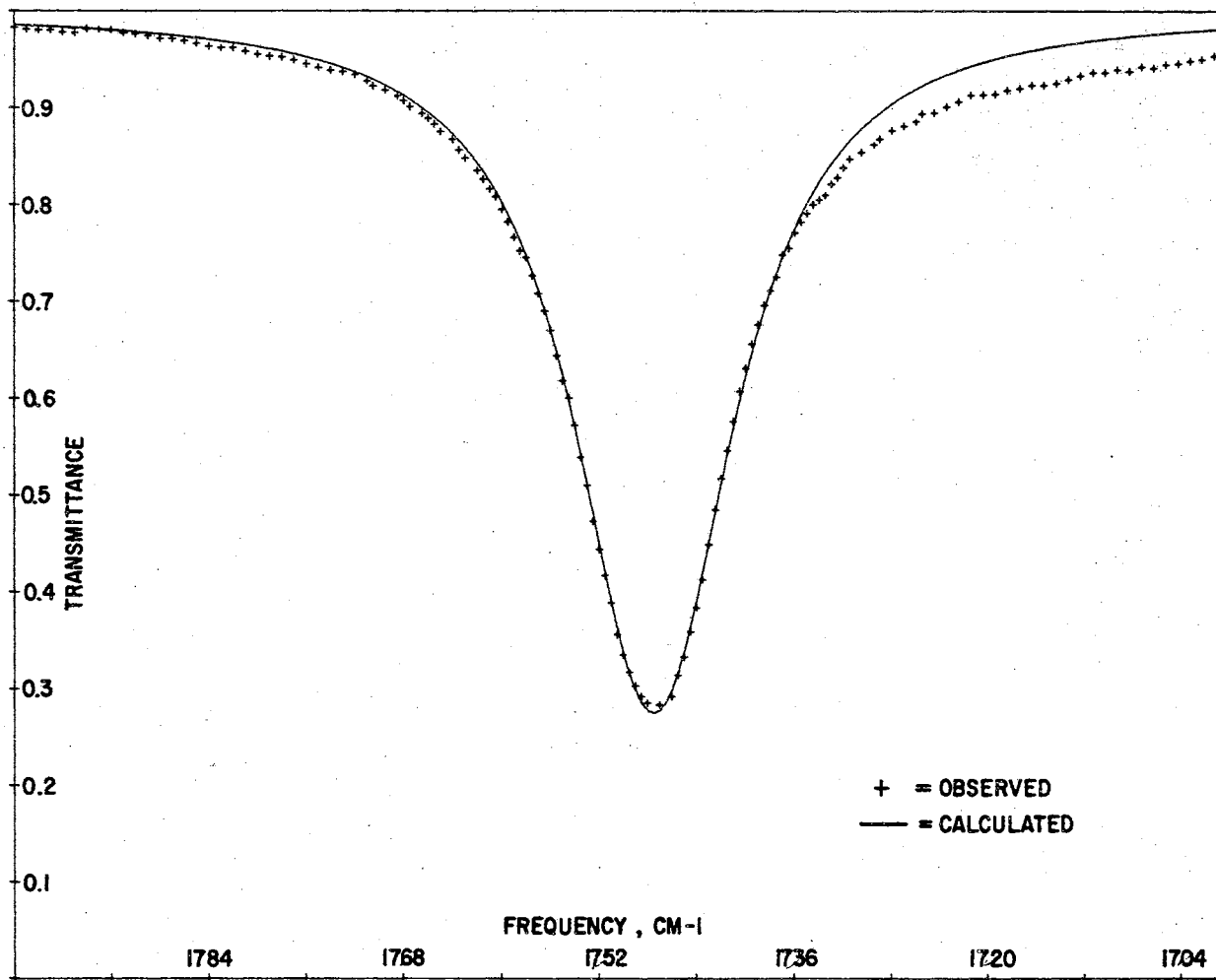


Figure 10. Observed and Calculated Carbonyl Stretching Bands for Methyl Oleate with 1.5 cm^{-1} Spectral Slit

band width for this minor feature was determined graphically as was an approximate value for the molar absorptivity. These values were then included in least-squares analyses of the 1800-1700 cm^{-1} region of the spectrum but were constrained to their original values so that only parameters for the carbonyl stretching band were refined. The results of these calculations appear in Table XXIII. A comparison of the observed and refined values for $\Delta\nu$ clearly demonstrates the power of the algorithm. Whereas the means of the observed values are 9.19, 10.45 and 12.48 cm^{-1} for spectra taken at spectral slit widths of 2.6, 5.0 and 7.5 cm^{-1} respectively, the mean for all 18 refined values is $(8.66 \pm 0.16) \text{cm}^{-1}$. Also, the refined band widths agree reasonably well with the theoretical values shown in the last column. The behavior of the refined intensity parameters is shown in Table XXIV. Although there is a 26 percent difference between the observed molar absorptivities for the 2.6 and 7.5 cm^{-1} data, the refined intensities clearly define a single molar absorptivity.

Carbonyl stretching spectra for n-tridecane solutions of di-n-butyl oxalate were also studied. This compound exhibits strong bands at 1774 and 1748 cm^{-1} arising from out-of-phase and in-phase motions of the two carbonyl groups. In addition, there is a small unresolved band in the vicinity of 1720 cm^{-1} . Transmittance data were gathered over the 1825 to 1700 cm^{-1} region and analyzed on the basis of this three-band model. The refined parameters from these analyses are listed in Table XXV and the molar absorptivity values calculated from the refined intensities appear in Table XXVI. No observed parameters are listed since comparison of these with the refined values is meaningless in the presence of extensive overlap.

TABLE XXIII

BAND-SHAPE PARAMETERS FOR THE CARBONYL
STRETCHING BAND OF ETHYL ACETATE

Spectral Slit Width, cm^{-1}	Concentration M	Observed Parameters			Refined Parameters			Calculated Values for $\Delta\nu$, cm^{-1}
		A_0	$\Delta\nu$, cm^{-1}	ν_0 , cm^{-1}	A_0	$\Delta\nu$, cm^{-1}	ν_0 , cm^{-1}	
2.6 ^a	0.00324	0.1332	9.38	1747.8	0.1421	8.64	1747.9	8.64
	0.00648	0.2629	9.21	1748.2	0.2789	8.41	1748.3	8.53
	0.01027	0.4151	9.12	1747.8	0.4333	8.76	1748.0	8.49
	0.01381	0.5367	9.10	1748.0	0.5626	8.69	1748.1	8.51
	0.01717	0.6786	9.17	1747.8	0.7015	8.93	1748.0	8.61
	0.02054	0.8130	9.18	1747.8	0.8391	8.96	1748.1	8.66
5.0 ^a	0.00324	0.1175	10.52	1747.3	0.1414	8.21	1747.6	8.13
	0.00648	0.2313	10.67	1747.6	0.2742	8.56	1747.7	8.42
	0.01027	0.3624	10.47	1747.6	0.4282	8.64	1747.7	8.29
	0.01381	0.4722	10.36	1747.6	0.5582	8.67	1747.8	8.26
	0.01717	0.6022	10.33	1747.6	0.7060	8.84	1747.8	8.34
	0.02054	0.7082	10.37	1747.6	0.8310	8.95	1747.8	8.57
7.5 ^b	0.00324	0.1010	12.73	1747.2	0.1389	8.32	1747.6	8.14
	0.00648	0.1919	12.86	1747.5	0.2651	8.49	1747.7	8.49
	0.01027	0.3097	12.52	1747.6	0.4255	8.65	1747.8	8.24
	0.01381	0.4090	12.31	1747.6	0.5639	8.67	1747.8	8.10
	0.01717	0.4803	12.27	1747.6	0.6673	8.74	1747.8	8.15
	0.02054	0.5955	12.16	1747.5	0.8269	8.87	1747.7	8.16

a) Sample path was 0.4801 mm.

b) Sample path was 0.4789 mm.

TABLE XXIV

MOLAR ABSORPTIVITY FOR THE CARBONYL
STRETCHING BAND OF ETHYL ACETATE

Spectral Slit Width, cm^{-1}	Molar Absorptivity, $\text{M}^{-1} \text{mm}^{-1}$	
	Observed Intensities	Refined Intensities
2.6	81.31 ± 0.91	83.26 ± 0.73
5.0	71.22 ± 0.83	83.12 ± 0.80
7.5	59.92 ± 0.74^a	83.29 ± 0.69^a

a) A_0 for the 0.01771M solution was not used in determination of this value.

TABLE XXV

REFINED BAND SHAPE PARAMETERS FOR THE CARBONYL STRETCHING
REGION OF DI-N-BUTYL OXALATE SPECTRA^a

Spectral Slit Width, cm^{-1}	Concentration M.	A_{01}	$\Delta\nu_1, \text{cm}^{-1}$	ν_{01}, cm^{-1}	A_{02}	$\Delta\nu_2, \text{cm}^{-1}$	ν_{02}, cm^{-1}	A_{03}	$\Delta\nu_3, \text{cm}^{-1}$	ν_{03}, cm^{-1}
3.0	0.00627	0.1077	13.01	1774.3	0.2230	10.91	1748.0	0.0184	36.11	1732.7
	0.01152	0.1975	13.75	1774.4	0.4230	10.75	1748.2	0.0103	25.24	1718.1
	0.01783	0.3027	13.80	1774.6	0.6451	10.78	1748.3	0.0175	25.36	1714.2
	0.02318	0.3924	13.91	1774.4	0.8279	10.82	1748.1	0.0228	23.67	1713.7
4.4	0.00627	0.1083	13.53	1773.9	0.2270	11.03	1747.8	0.0096	22.91	1720.7
	0.01152	0.1974	13.75	1774.2	0.4195	10.90	1748.1	0.0136	25.04	1717.5
	0.01783	0.3033	13.67	1774.1	0.6449	10.74	1747.9	0.0196	28.64	1714.0
	0.02318	0.3921	13.83	1774.2	0.8278	10.80	1747.9	0.0252	25.93	1713.2
6.0	0.00627	0.1084	13.01	1773.9	0.2293	10.02	1747.8	0.0227	30.89	1734.8
	0.01152	0.1982	13.65	1774.1	0.4268	10.55	1747.8	0.0127	21.80	1717.1
	0.01783	0.3042	13.74	1774.1	0.6606	10.45	1747.8	0.0175	25.66	1714.7
	0.02318	0.3947	13.74	1774.1	0.8506	10.48	1747.8	0.0239	23.53	1713.8
10.0	0.00627	0.1105	12.95	1773.7	0.2410	9.81	1747.7	0.0198	30.71	1731.9
	0.01152	0.2052	13.35	1774.0	0.4486	9.96	1747.8	0.0132	23.04	1718.8
	0.01783	0.3111	13.64	1774.0	0.7093	9.68	1747.8	0.0197	29.19	1714.6
	0.02318	0.4038	13.66	1774.0	0.9184	9.74	1747.7	0.0258	23.77	1713.8

a) Sample path for all spectra was 0.4780 mm.

TABLE XXVI

MOLAR ABSORPTIVITIES FOR THE CARBONYL STRETCHING
REGION OF DI-N-BUTYL OXALATE SPECTRA^a

Spectral Slit Width, cm ⁻¹	ϵ_1	ϵ_2	ϵ_3^b
3.0	35.18 ± 0.13	74.68 ± 1.24	2.25 ± 0.09
4.4	35.11 ± 0.11	74.37 ± 0.77	2.08 ± 0.06
6.0	35.39 ± 0.09	76.94 ± 0.66	2.00 ± 0.26
10.0	36.14 ± 0.35	84.13 ± 0.63	2.26 ± 0.07

a) The units for ϵ are M⁻¹mm⁻¹.

b) Data for the 0.00627M solution were not used in the calculation of these entries.

Examination of the table entries for the two major bands indicates that, for the data taken at spectral slit widths of 3.0 and 4.4 cm^{-1} , all parameters tend to converge to a common value. The uncertainty in the parameters for the small feature at about 1715 cm^{-1} was not unexpected, considering the small magnitude of the molar absorptivity and the fact that it appears on the wing of a much stronger band. They are more stable, however, than the small bands postulated for methyl oleate. It is therefore concluded that the band is real. A plot of observed and calculated spectra for this ester is given in Figure 11.

Intensity and band width parameters for the 1748 cm^{-1} band calculated using the 6.0 cm^{-1} slit data deviate somewhat from those calculated from the data taken with 3.0 and 4.4 cm^{-1} slits. For the spectra determined with $S = 10 \text{ cm}^{-1}$, parameters for all but the 1715 cm^{-1} band have deviated significantly from the narrower slit width values. It thus appears that, for medium to strong absorption bands, shape parameters calculated using data taken at wide slit widths are unreliable. It should be noted, however, that slits of 6.0 and 10.0 cm^{-1} are far wider than those used in normal spectrometric work. One may conclude, therefore, that the least-squares algorithm has successfully described the carbonyl stretching region of the butyl oxalate spectrum.

Carbon-Hydrogen Bending Bands of Simple Aromatic Hydrocarbons

Aromatic compounds bearing ring hydrogen atoms generally exhibit one or more C-H bending bands in the 900-600 cm^{-1} region of the infrared spectrum. The final phase of this investigation involved a study of these bands for carbon disulfide solutions of toluene, o-xylene,

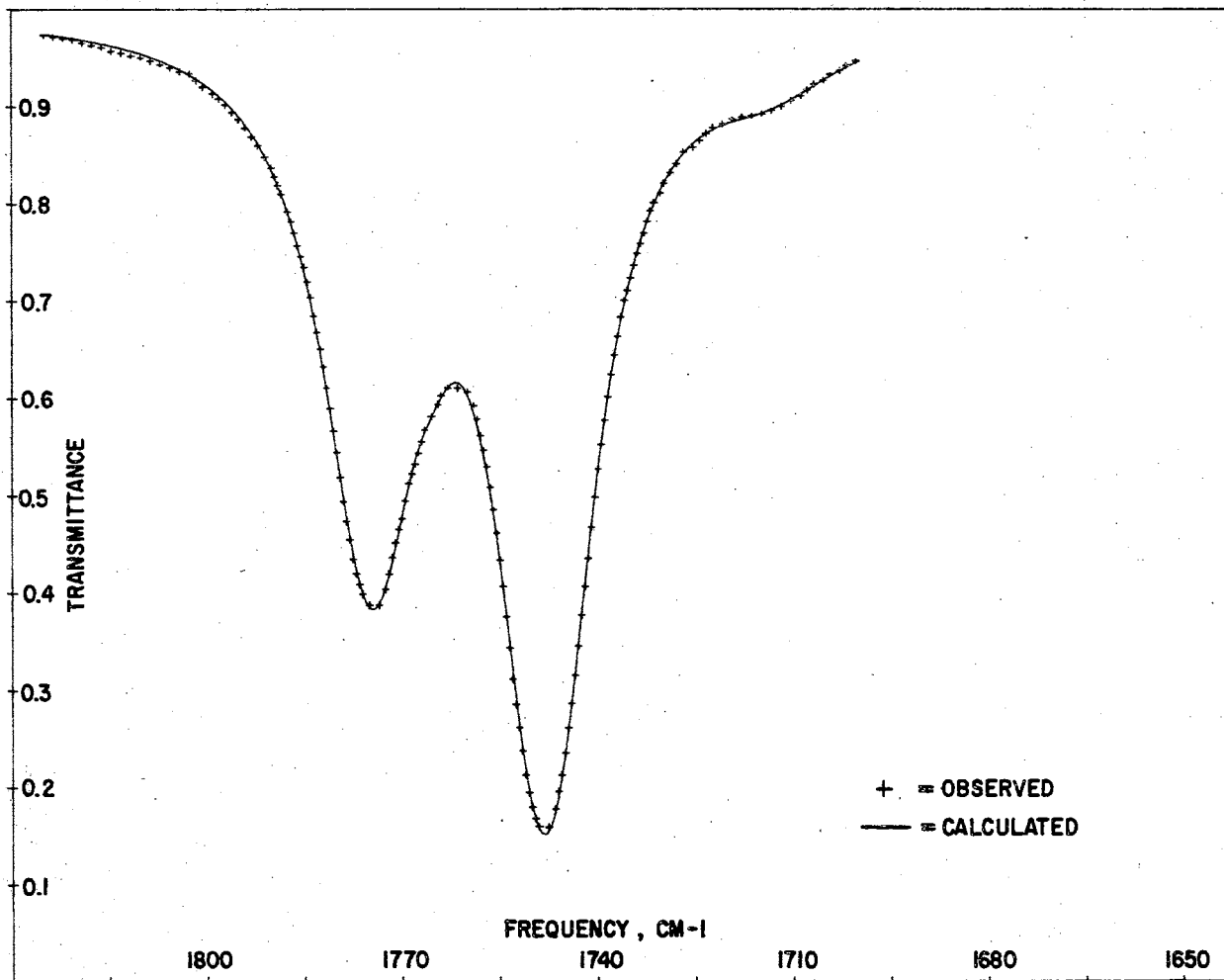


Figure 11. Observed and Calculated Carbonyl Stretching Bands for Di-n-Butyl Oxalate with 3 cm⁻¹ Spectral Slit

m-xylene and t-butylbenzene. Spectra were determined over the region 800-650 cm^{-1} at frequency intervals of 0.2 cm^{-1} . Individual solute concentrations varied from 0.0306 to 0.1044 M and the spectral slit widths employed ranged from 1.0 to 2.2 cm^{-1} .

Several least-squares analyses were performed involving spectra of solutions containing a single solute. The results for 35 such calculations are summarized in Table XXVII. The refined parameters for the individual hydrocarbons were then used as initial estimates for analyses of the spectra of solutions containing two or more solutes.

As can be seen from Table XXVIII, most of the refined parameters for the spectra of mixtures fall within one or two standard deviations of the values derived from the spectra of single component solutions. Notable exceptions are those parameters for the 693.2 cm^{-1} toluene band and the 690.1 m-xylene band which were calculated from data acquired at a spectral slit width of 2.2 cm^{-1} . Although the observed data are apparently well described by the calculated model (see Figure 12), the molar absorptivity and band width values calculated for these absorbances do not agree with those listed in Table XXVII. It therefore seems that the algorithm cannot accurately resolve bands which are so extensively overlapped. Better values for these parameters are calculated from the data taken with 1.0 cm^{-1} slits. Here the apparent overlap is not so severe (see Figure 13) and the algorithm is more effective. It is also possible to calculate accurate parameters from the 2.2 cm^{-1} data if $\Delta\nu$ and ν_0 for the overlapped bands are constrained to the initial values for 5 or 6 cycles. The success of this procedure, however, is somewhat predicated on the accuracy of the initial estimates. For the case of the mixture of o-xylene, m-xylene and t-butylbenzene,

TABLE XXVII

REFINED BAND SHAPE PARAMETERS FOR THE C-H BENDING
BANDS OF AROMATIC HYDROCARBONS

Compound	ν_0, cm^{-1}	$\epsilon, \text{M}^{-1} \text{mm}^{-1}$	$\Delta\nu, \text{cm}^{-1}$
Toluene ^a	727.9 ± 0.1	52.52 ± 0.86	4.15 ± 0.08
	693.2 ± 0.1	17.17 ± 0.43	4.49 ± 0.09
<u>o</u> -Xylene ^b	741.5 ± 0.0	61.55 ± 0.38	3.87 ± 0.04
<u>m</u> -Xylene ^a	767.8 ± 0.0	34.29 ± 0.65	4.53 ± 0.09
	690.1 ± 0.1	17.22 ± 0.17	3.86 ± 0.21
<u>t</u> -Butylbenzene ^b	761.7 ± 0.1	27.34 ± 0.12	3.92 ± 0.14
	697.1 ± 0.1	31.04 ± 0.67	5.65 ± 0.18

a) Spectral slit width varies from 1.0 to 2.2 cm^{-1} .

b) Spectral slit width is 2.1 cm^{-1} .

TABLE XXVIII

REFINED CARBON-HYDROGEN BENDING BAND PARAMETERS
FOR MIXTURES OF AROMATIC HYDROCARBONS

Components	Spectral Slit Width, cm^{-1}	ν_0 , cm^{-1}	ϵ , $\text{M}^{-1} \text{mm}^{-1}$	$\Delta\nu$, cm^{-1}
Toluene + <u>m</u> -Xylene	2.2	767.8	34.72	4.42
		727.9	51.77	4.08
		694.0	12.23	1.99
		690.8	23.16	6.20
Toluene + <u>m</u> -Xylene	1.0	767.7	36.21	4.58
		727.8	53.21	4.31
		693.2	16.92	4.42
		690.1	17.97	3.59
<u>o</u> -Xylene, <u>m</u> -Xylene, <u>t</u> -Butylbenzene	2.1	767.7	33.79	4.35
		761.9	27.90	4.22
		741.5	61.16	3.88
		697.2	30.87	5.60
		689.9	15.16	3.76

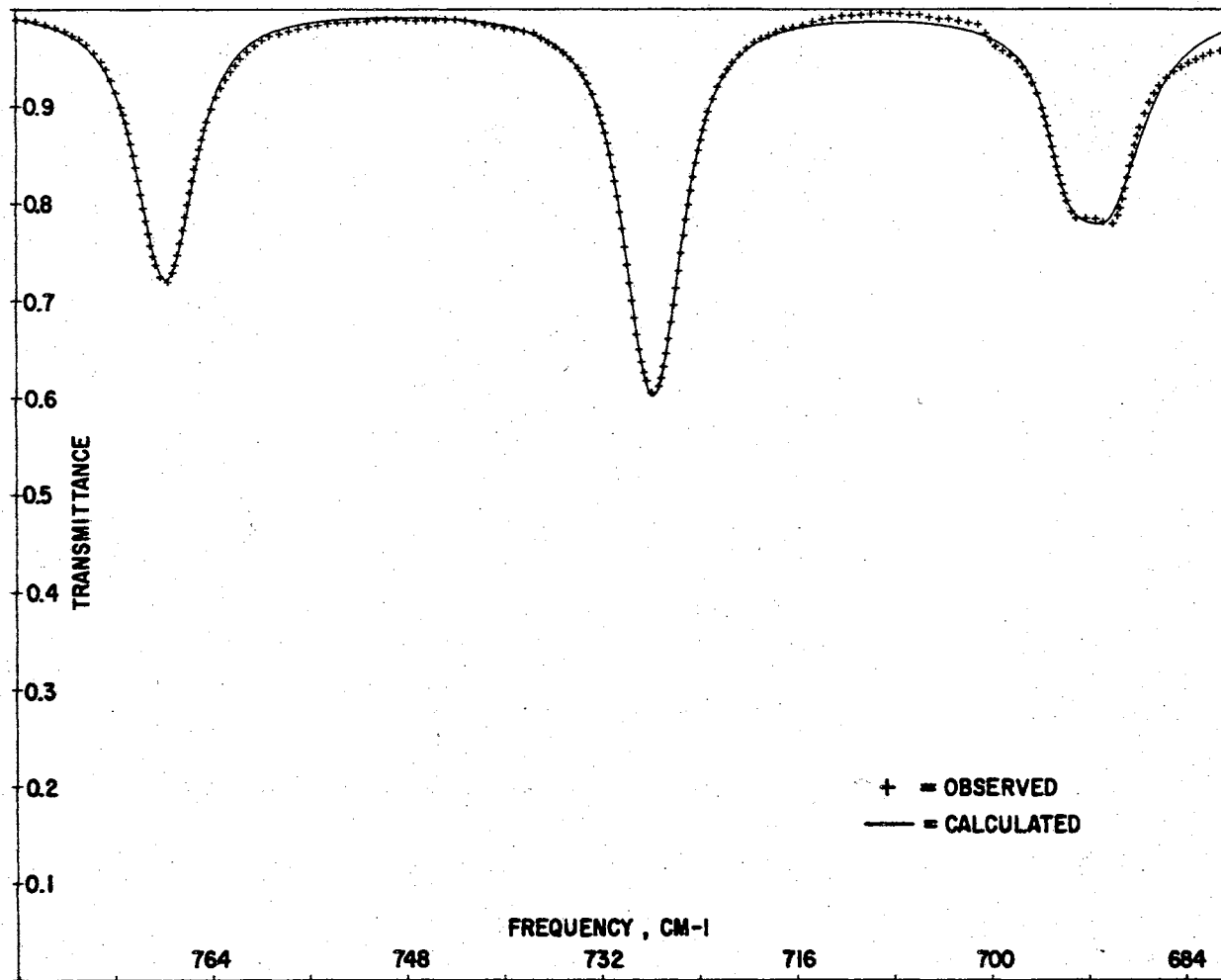


Figure 12. Observed and Calculated Carbon-Hydrogen Bending Bands for Toluene and m-Sylene with 2.2 cm⁻¹ Spectral Slit

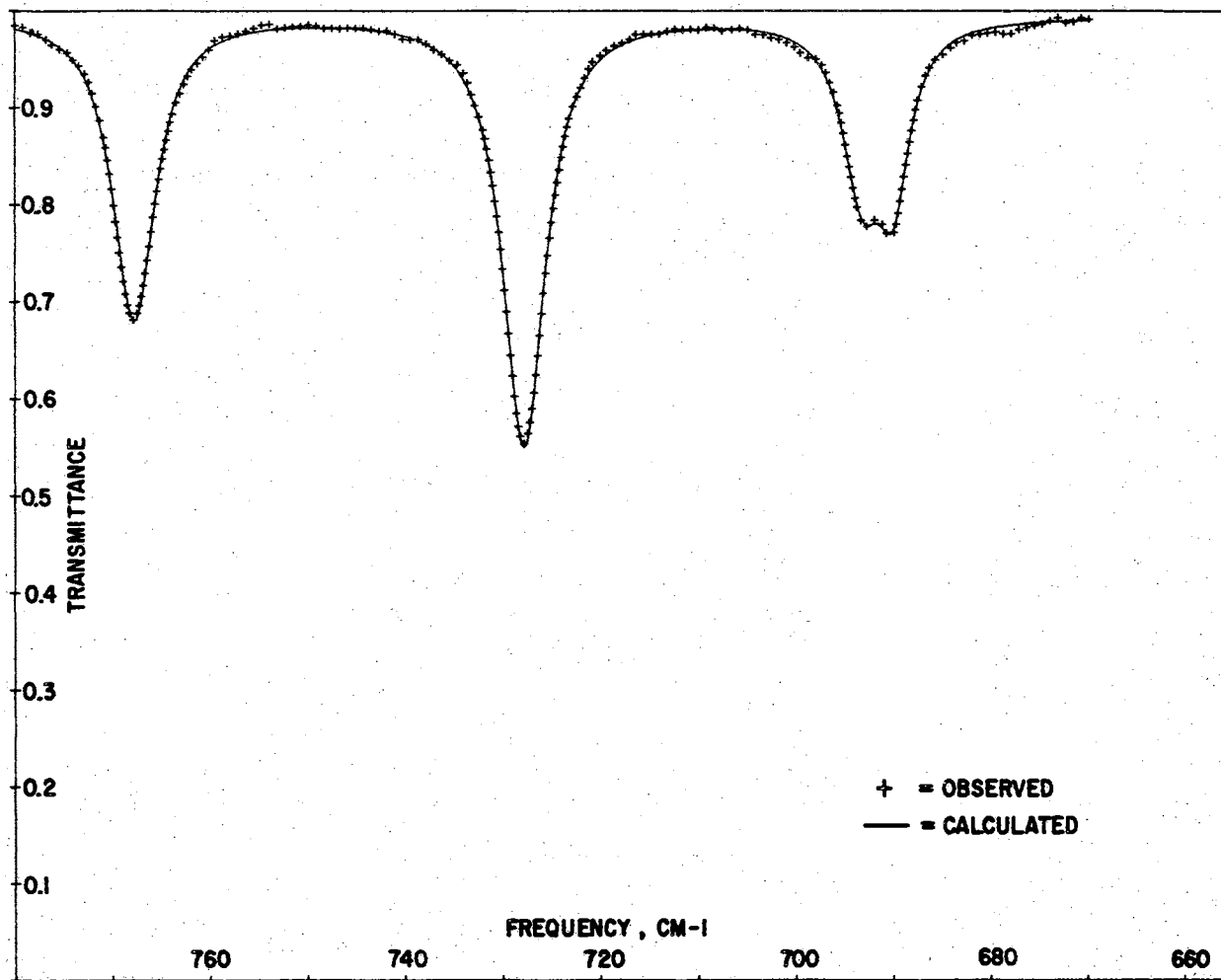


Figure 13. Observed and Calculated Carbon-Hydrogen Bending Bands for Toluene and m-Xylene with 1 cm⁻¹ Spectral Slit

the value calculated for the width of the 690 cm^{-1} m-xylene band is only slightly lower than the single-component value. A plot of observed and calculated spectra for this mixture is shown as Figure 14.

Summary and Conclusions

The purpose of this investigation was to develop an algorithm whereby reliable values for the shape parameters of infrared absorption bands could be calculated. In achieving this end, a mathematical model was devised wherein an infrared spectrum is represented by a set of one or more Lorentz functions convolved with a triangular spectral slit function. This model was implemented in the form of a digital computer program (BANDFIT) which serves to calculate the shape parameters for each component band of a spectrum.

The effectiveness of program BANDFIT was evaluated initially on the basis of calculations involving synthetic spectra. In the second phase of the study real data, taken from the infrared spectra of aliphatic esters and simple aromatic hydrocarbons, were subjected to analysis. The results of these calculations demonstrate that band shape parameters derived from these analyses may be used to generate spectra which closely approximate those determined experimentally. In addition, the calculated parameters are physically significant in that, for a given set of spectral features, they are essentially unaffected by changes in spectral slit width.

The algorithm is not entirely successful in the determination of shape parameters for systems of strongly overlapped bands. It appears that unrealistic band width and intensity values result if the features are approximately one band width apart and if the ratio of the spectral

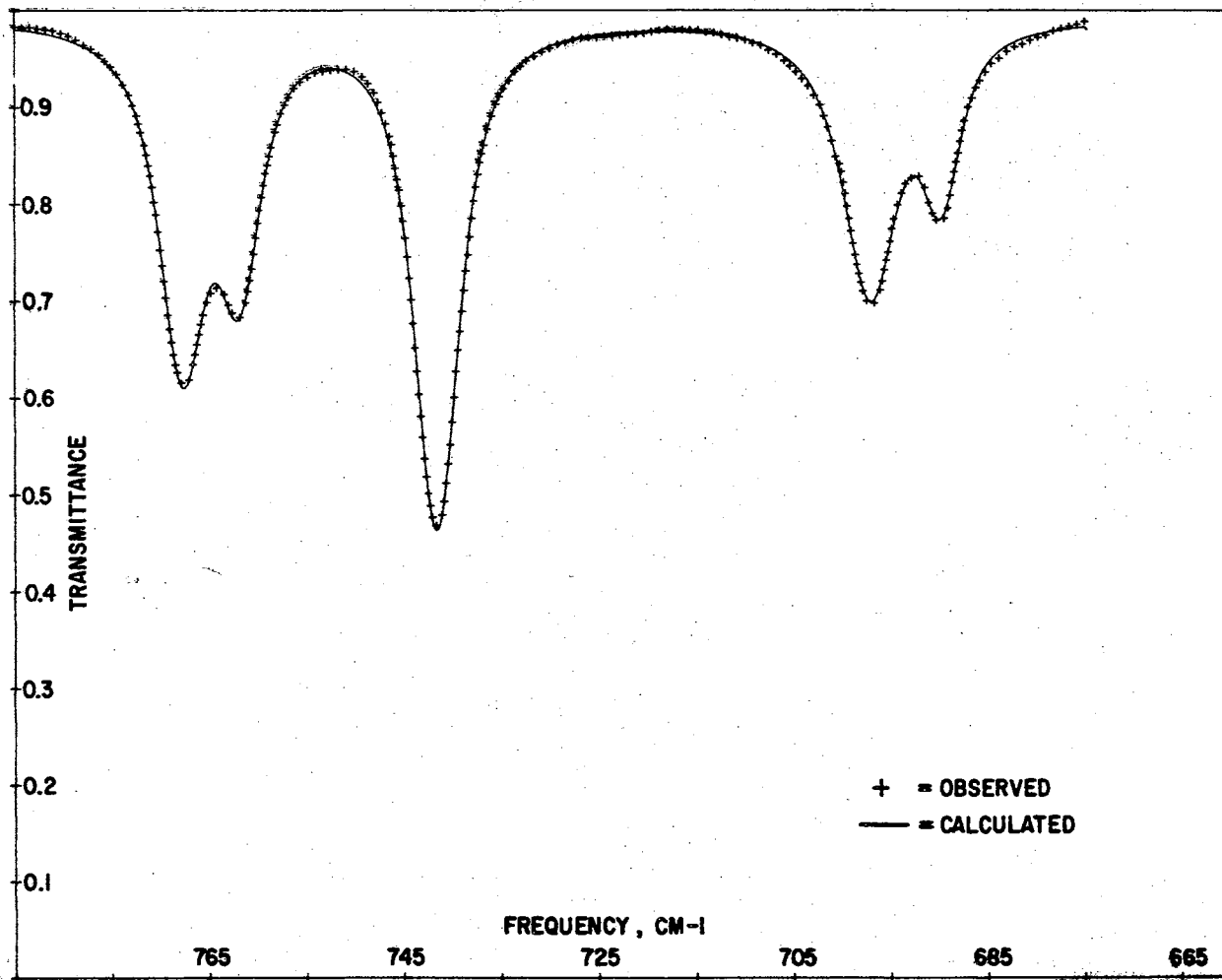


Figure 14. Observed and Calculated Carbon-Hydrogen Bending Bands for o-Xylene, m-Xylene and t-Butylbenzene with 2.1 cm^{-1} Spectral Slit

slit width to the apparent band width is greater than about 0.25. For greater frequency separation, however, accurate determination of shape parameters is easily accomplished, even for bands appearing as unresolved shoulders. This is evidenced by the calculations involving the three band synthetic spectrum and by the set of self-consistent parameters derived for the 1715 cm^{-1} band in di-n-butyl oxalate spectra.

It thus appears that infrared solution spectra are accurately described by the mathematical model devised and implemented during this investigation. Extraction of the shape parameters for a spectrum on the basis of this model is one step toward the removal of distortions introduced by the examining instrument. It is apparent, however, that there are other factors which influence the shapes and positions of infrared absorption bands. Perturbations of the detected signal by the electrical and mechanical systems of the spectrophotometer are probably instrumental in band broadening and in the skewing of peaks. As was previously mentioned, the effects of such distortions can possibly be eliminated by a deconvolution of the observed signal with the inverse of the instrument transfer function. Development of algorithms to evaluate the transfer function and to implement the deconvolution are the next logical steps in this field of research.

SELECTED BIBLIOGRAPHY

1. E. B. Wilson, Jr., J. C. Decius and P. C. Cross, "Molecular Vibrations," McGraw-Hill, New York, 1955.
2. G. Herzberg, "Molecular Spectra and Molecular Structure. II. Infrared and Raman Spectra of Polyatomic Molecules," Van Nostrand, Princeton, N. J., 1945.
3. H. C. Urey and C. A. Bradley, Phys. Rev. 38, 1969 (1931).
4. J. P. Devlin, J. Chem. Phys. 39, 2385 (1963).
5. J. P. Devlin, J. Chem. Phys. 41, 2951 (1964).
6. T. Shimanouchi, J. Chem. Phys. 17, 245, 734 and 848 (1949).
7. T. Shimanouchi, Pure Applied Chem. 7, 131 (1963).
8. J. H. Schachtschneider and R. G. Snyder, Spectrochim. Acta 19, 117 (1963).
9. J. Overend and J. R. Scherer, J. Chem. Phys. 32, 1289, 1296 and 1720 (1960).
10. B. D. Latimer, "A Vibrational Study of Some Trihaloboroxines," (unpublished Ph.D. thesis, Oklahoma State University, 1966).
11. A. Rosenberg and J. P. Devlin, Spectrochim. Acta 21, 1613 (1965).
12. S. Califano and B. Crawford, Jr., Spectrochim. Acta 16, 889 (1960).
13. J. Scherer and J. Overend, Spectrochim. Acta 17, 719 (1961).
14. J. Scherer, J. Chem. Phys. 36, 3308 (1962).
15. G. Zerbi, B. Crawford, Jr. and J. Overend, J. Chem. Phys. 38, 127 (1963).
16. M. Scrous, C. diLauro and S. Califano, Spectrochim. Acta 21, 571 (1965).
17. F. Halverson, R. F. Stamm and J. J. Whalen, J. Chem. Phys. 16, 808 (1948).

18. P. Devlin, J. Overend and B. Crawford, Jr., *Spectrochim. Acta* 20, 23 (1964).
19. D. A. Long and W. O. George, *Spectrochim Acta* 19, 1717 (1963).
20. F. A. Miller et al., *Spectrochim. Acta* 20, 1233 (1964).
21. J. R. Scherer and J. Overend, *J. Chem. Phys.* 33, 1681 (1960).
22. A. E. Ardis et al., *J. Am. Chem. Soc.* 72, 1305 (1950).
23. O. Diels, H. Gartner and R. Kasck, *Ber.* 55, 3439 (1922).
24. O. Diels and B. Conn, *Ber.* 56, 2076 (1923).
25. A. E. Ardis and H. Gilbert, U. S. Patent 2,535,827 (Dec. 26, 1950); *Chem. Abstr.* 45, 3196^b (1951).
26. International Union of Pure and Applied Chemistry, "Tables of Wavenumbers for the Calibration of Infra-red Spectrometers," Butterworths, Washington, D. C., 1961..
27. G. Herzberg, "Molecular Spectra and Molecular Structure. II. Infrared and Raman Spectra of Polyatomic Molecules," Van Nostrand, Princeton, N. J., 1945, pp. 468-484.
28. W. S. Wilcox, J. H. Goldstein and J. W. Simmons, *J. Chem. Phys.* 22, 516 (1954).
29. "Tables of Interatomic Distances and Configurations in Molecules and Ions," Special Publication No. 11, The Chemical Society, London, 1958.
30. M. Aoualt and C. St. Arnaud, *Can. J. Phys.* 35, 504 (1957).
31. "Tables of Interatomic Distances and Configurations in Molecules and Ions, Supplement 1956-1959," Special Publication No. 18, The Chemical Society, London, 1965.
32. C. C. Costain and B. P. Stoicheff, *J. Chem. Phys.* 30, 777 (1959).
33. D. A. Bekoe and K. N. Trueblood, *Z. Krist.* 113, 1 (1960).
34. N. C. Craig et al., Ohio State Symposium on Molecular Structure and Spectroscopy, June 1964.
35. V. W. Laurie and D. T. Pence, *J. Chem. Phys.* 38, 2693 (1963).
36. V. W. Brugel, T. Anhel and F. Frucksberg, *Z. Elektrochem.* 64, 1121 (1960).
37. A. A. Bothner-By, *J. Am. Chem. Soc.* 83, 231 (1961).

38. N. C. Craig and E. A. Entemann, *J. Am. Chem. Soc.* 83, 3047 (1961).
39. R. Bruns, L. Raff, and J. P. Devlin, *Theoret. Chim. Acta (Berl.)* 14, 232 (1969).
40. K. S. Seshadri and R. N. Jones, *Spectrochim. Acta* 19, 1013 (1963).
41. H. Morgenau and W. W. Watson, *Revs. Modern Phys.* 8, 22 (1936).
42. H. A. Lorentz, *Koninkl. Ned. Akad. Wetenschap. Proc.* 8, 591 (1906).
43. J. H. VanVleck and V. F. Weisskopf, *Revs. Modern Phys.* 17, 227 (1945).
44. J. R. Nielsen, V. Thornton and E. B. Dale, *Revs. Modern Phys.* 16, 307 (1944).
45. D. A. Ramsay, *J. Am. Chem. Soc.* 74, 72 (1952).
46. P. Montigny, *Spectrochim. Acta* 20, 1373 (1964).
47. R. N. Jones, K. S. Seshadris, N. B. W. Jonathan and J. W. Hopkins, *Can. J. Chem.* 41, 750 (1963).
48. J. Pitha and R. N. Jones, *Can. J. Chem.* 44, 3031 (1966).
49. J. Pitha and R. N. Jones, *Can. J. Chem.* 45, 2347 (1967).
50. R. N. Jones, personal communication, 1969).
51. S. Abramowitz and R. P. Bauman, *J. Chem. Phys.* 39, 2757 (1963).
52. R. D. B. Fraser and E. Suzuki, *Anal. Chem.* 41, 37 (1969).
53. H. F. Smith and A. S. Rosenberg, *J. Chem. Soc.*, 5391 (1963).
54. A. S. Rosenberg and H. F. Smith, *J. Chem. Soc.*, 5395 (1963).
55. A. S. Rosenberg and H. F. Smith, Pittsburgh Conference on Analytical Chemistry and Applied Spectroscopy, March 1964.
56. A. Savitzky and M. J. E. Golay, *Anal. Chem.* 36, 1627 (1964).
57. H. O. Hartley, *Technometrics* 3, 269 (1961).

APPENDIX

Instructions for Use of the Computer Program

Program BANDFIT was designed to implement the algorithm developed in this investigation. It is written in the Fortran IV programming language for execution on an IBM System 360/65 digital computer with full operating system. In its present configuration the program is limited to the analyses of spectra containing 10,000 or fewer digital values and which are composed of a maximum of 10 absorption bands.

The program consists of a main routine and seven subroutines; SPCTRM, SMTH1, SCAN, CONMAT, LORN3, PLOT1 and SCALAB. Also, a local subroutine (GDATE) is used to determine the date of the analysis. If this routine is not available the call to it and references to the variable DATE should be deleted. The functions of the other subroutines are briefly outlined below.

Subroutine SPCTRM is used to read ordinate values for the observed spectrum either from punched cards or from a user disk file. A call to entry point DISK of this subroutine serves to initialize the disk file. The program is written so that the method of reading the observed spectrum may be changed by simply modifying SPCTRM. Subroutine SMTH1 is used to digitally smooth the observed spectrum, SCAN generates the vector of coefficients employed in the slit convolution and CONMAT is called to calculate the orthogonal polynomial coefficients required by both SMTH1 and SCAN. The calculated spectrum,

error sum of squares and the sums of partial derivatives used in the least-squares calculation are evaluated in LØRN3. Finally, subroutines PLØT1 and SCALAB are used to generate graphs of the observed and calculated spectra and of the error function. In order to use PLØT1, the Calcomp (California Computer Products, Inc.) plotter subroutine package must be available. If it is not, calls to all entry points of subroutine PLØT1 should be eliminated from the main program.

Input requirements for program BANDFIT are detailed in TABLE XXIX. Although two types of spectrum header cards are described in the table, only one is required for a given problem. Also, it should be noted that band intensity estimates are input in terms of molar absorptivity rather than absorbance. This is done so that the same estimates may be used for the data from a series of solutions. The absorbance estimate is calculated in the program from the molar absorptivity, concentration and cell path values.

Printed output for a typical problem consists of the error sum of squares and parameter set for each calculation cycle and a list of the final refined parameter values together with estimates of their uncertainties. Listings of the observed and calculated spectra, punched and plotted output are optional as selected by control variables.

TABLE XXIX

INPUT REQUIREMENTS FOR PROGRAM BANDFIT

Card	Columns	Format	Variable Name and Function
Record	1-2	I2	KARD = 1 for the first and KARD = 2 for the second of the 2 record cards.
	5-80	19A4	ALPHA, any alphanumeric information desired to describe the problem.
Control	1-2	I2	KARD = 3
	3-8	I6	NØPRØB, the problem number.
	9-11	I3	If NPLT = 2, plots will be generated for the observed and calculated spectra and for the error function. If NPLT = 1, only the observed and calculated spectra will be plotted. If NPLT = 0, no plotting is done.
	12-14	I3	LTYPE = 1 if the observed spectrum is in transmittance units, LTYPE = 2 if it is in absorbance units.
	15-17	I3	NPEAK, the number of absorption bands in the spectrum.
	18-20	I3	IHØLD is nonzero if any of the parameter estimates are to be constrained. If IHØLD = -1, all parameters are constrained and a zero-order calculation is performed.
	21-23	I3	If MØDE = 1 the observed spectrum is to be input from cards, if MØDE = 2 it is accessed from the disk file.

TABLE XXIX (Continued)

Card	Columns	Format	Variable Name and Function
Control	24-26	I3	NGATE; if greater than zero the observed spectrum will be smoothed with an NGATE-point polynomial function of degree NPØLY (see below).
	27-29	I3	NPØLY
	30-32	I3	IFPNCH, if nonzero the final calculated spectrum will be punched on cards.
	33-35	I3	IFPRNT controls printing of the observed and calculated spectra. There are options for no listing (IFPRNT = 0), a single-spaced (IFPRNT = 1), or a double-spaced (IFPRNT = 2) listing.
	36-45	F10.5	TSTP; if the fractional improvement in the error sum of squares for two successive cycles falls below this value, the problem is terminated.
	46-55	F10.5	DX, the abscissa interval for the calculated spectrum. If DX is not input, the value for EDX (see below) is used.
	56-60	F5.2	FACTOR, used to scale the vector of corrections before the parameters are updated. If not input, a value of 1.0 is assumed.
Spectrum Header (MØDE = 1) ^a	1-5	I5	IDENT, the spectrum number.
	9-15	F7.0	XSTART, the starting frequency for the spectrum.
	16-20	F5.1	EDX, the abscissa sampling interval for the observed spectrum.

TABLE XXIX (Continued)

Card	Columns	Format	Variable Name and Function
Spectrum Header (MODE = 1) ^a	21-25	I5	NIPTS, the number of data points in the observed spectrum.
	31-40	F10.5	BEGIN and END (see below) define the frequency range to be considered in the analysis. If not input, the range is determined from the spectrum.
	41-50	F10.5	END
Spectrum Header (MODE = 2) ^b	1-5	I5	IDENT, the spectrum number.
	11-20	F10.5	BEGIN, see above.
	21-30	F10.5	END, see above.
Parameter ^c	1-2	I2	KARD = 4
	11-20	F10.5	EO, the initial molar absorptivity value in $M^{-1} \text{mm}^{-1}$.
	21-25	I5	IHEO, the number of cycles for which EO is to be constrained.
	31-40	F10.5	DELNU, the initial value for $\Delta\nu$ in cm^{-1} .
	41-45	I5	IHRHØ, the number of cycles for which DELNU is to be constrained.
	51-60	F10.5	XNU, the initial value for ν_0 in cm^{-1} .
	61-65	I5	IHNU, the number of cycles for which XNU is to be constrained.

TABLE XXIX (Continued)

Card	Columns	Format	Variable Name and Function
Instrument	1-2	I2	KARD = 5
	11-20	F10.5	BETA, the initial value for the constant β .
	21-25	I5	IHBETA, the number of cycles for which BETA is to be constrained.
	26-30	I5	MPØLY, the degree of polynomial to be used for the slit integration. If MPØLY = 0, a trapezoidal rule integration is performed.
	31-40	F10.5	SLIT, the spectral slit width in cm^{-1} .
	41-50	F10.5	XL, the path of the sampling cell in mm.
Concentration	1-2	I2	KARD = 6
	11-80	10F7.4	C(1) through C(10), the concentration to be associated with each absorption band. NPEAK values must be entered.

- a) This card must be followed by cards containing the observed spectrum punched in 2014 format.
- b) For this mode of input, the spectrum and header information are expected to be on a direct access disk file. Refer to the listing of subroutine SPCTRM for the required format.
- c) There must be NPEAK parameter cards.

Program Listings

C....	PROGRAM BANDFIT	BAND0010
C		BAND0020
C....	REQUIRES SUBROUTINES SCAN, LORN3, PLOT1, SMTH1, CONMAT,	BAND0030
C....	SCALAB, SPCTRM	BAND0040
C		BAND0050
	DEFINE FILE 7(2500,1288,L,KEY)	BAND0060
	DIMENSION ALPHA(38),B(31,32),BSAVE(31),C(10),D(51),DATE(2),	BAND0070
1	DELNU(10),EO(10),HOLDEO(10),HOLDRH(10),HOLDXN(10),	BAND0080
2	IHEO(10),IHNU(10),IHRHD(10),NDX(32),NHOLD(31,2),	BAND0090
3	VAR(31),WS(55),PLBL(9)	BAND0100
	DIMENSION IR(5000),PLBL1(6)	BAND0110
	EQUIVALENCE (IR(1),FREQ(1))	BAND0120
	COMMON A(31,32),A0(10),CON(51),DELSQ,DX,ENU(5000),ETR(5000),	BAND0130
1	FREQ(5000),FREQHI,NCOLS,NMAX,NPEAK,NROWS,NTPTS,NXPTS,	BAND0140
2	PRTSUM(31),RHO(10),R(5000),T(5000),XNU(10),ISLIT,KEY,BETA	BAND0150
	REAL*8 A,B,WS	BAND0160
	DATA PLBL/'CALC','D AN','D OB','SVD ','SPEC','TRA ',' P',	BAND0170
1	'ROB ',' /, BLANK,ZERO/' ','0'/'	BAND0180
	DATA PLBL1/'ERRO','R CU','RVE ',' P','ROB ',' ' /	BAND0190
C		BAND0200
C....	FORMAT STATEMENTS.	BAND0210
C		BAND0220
	5000 FORMAT(I2,2X,19A4)	BAND0230
	5001 FORMAT (I2,I6,9I3,2F10.5,F5.2)	BAND0240
	5003 FORMAT (I2,8X,F10.5,2I5,2F10.5)	BAND0250
	5005 FORMAT(I2,8X,10F7.4)	BAND0260
	5006 FORMAT(I2,8X,3(F10.5,I5,5X))	BAND0270
	6000 FORMAT(1H0,/,/,2X,15(1H*),13H CARD NUMBER 15,13H OUT OF ORDER,15(1H	BAND0280
	1*))	BAND0290
	6001 FORMAT(1H1,19A4///1X,19A4//)	BAND0300
	6002 FORMAT (I4)	BAND0310
	6003 FORMAT(1H0,5X,50HSPREAD TOO GREAT BETWEEN FHI AND FLO	BAND0320
1	,5X,9HTHEY ARE 2(2X,F10.5))	BAND0330
	6004 FORMAT (/,/'0',13X,'SLIT WIDTH',7X,'FREQLO',7X,'FREQHI',7X,	BAND0340
1	'CELL PATH'/'0',15X,F5.2,8X,F8.2,5X,F8.2,8X,F7.4///	BAND0350

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2          '0',21X,'NUMBER OF EXPERIMENTAL POINTS =',I5//          BAND0360
3          '0',20X,'PEAK      ABSORBANCE      CONCENTRATION')    BAND0370
6005 FORMAT ('0',21X,12,8X,F6.4,9X,F8.5)                          BAND0380
6006 FORMAT('0',22X,'***** PROBLEM NUMBER' I5,' *****'//20X,'NUMBER OF PBAND0390
1EAKS = ' I2//15X,'THE INITIAL ESTIMATES OF (THE PARAMETERS ARE ....'BAND0400
2//3X,'PEAK',7X,'EO',7X,'HOLD'6X,'DELNU',5X,'HOLD'7X,'NU'7X,'HOLD')BAND0410
6007 FORMAT ('0',15X,'..... IMPROVEMENT IN FIRST PHASE')        BAND0420
6008 FORMAT (/ '0*** ENTERING 10THS CYCLE ON ITERATION',I3,' ***'/          BAND0430
1          '0',5X,'PREVIOUS SUM SQ. = ',E14.7,5X,                BAND0440
2          'CURRENT SUM SQ. = ',E14.7)                            BAND0450
6009 FORMAT(1H0,15X,35H.... IMPROVEMENT FOUND ON PEAK NO. I4)    BAND0460
6010 FORMAT(1H1)                                                  BAND0470
6011 FORMAT (/ '0',5X,'..... SUM OF SQUARES FOR ITERATION',I3,' = ',          BAND0480
1          E13.6,' .....',19X,F11.6/'0',20X,'... PARAMETER SET ...'//          BAND0490
2          '0',5X,'PEAK',8X,'EO',18X,'DELNU',17X,'NU')          BAND0500
6012 FORMAT(1H0,26H***** SINGULAR MATRIX, SFT I3,14H ON ITERATION I3,4H*BAND0510
1***)                                                            BAND0520
6013 FORMAT(1H0,15X,28H** HALTED AFTER 25 CYCLES **)            BAND0530
6014 FORMAT(1H0,///,10X,27HPROBLEM COMPLETED ON CYCLE I3)      BAND0540
6015 FORMAT(1H0,5X,36H... SUM OF SQUARES STATIONARY AFTER I3,2X, 9HCYCLBAND0550
1ES...))                                                         BAND0560
6016 FORMAT(1H1,20X,27H...RESULTS FOR PROBLEM NO. I6,4H ...///25X,17HSUBBAND0570
1M OF SQUARES = E13.6)                                          BAND0580
6017 FORMAT ('0',5X,'MOLAR',8X,'ABSORBANCE',11X,'LINE WIDTH',12X,          BAND0590
1          'FREQUENCY'// ' ABSORPTIVITY. ',3('VALUE   STD DEV',6X))    BAND0600
6018 FORMAT ('0',2X,2F9.4,1PE12.4,0PF9.2,1PE12.4,0PF10.1,1PE12.4)    BAND0610
6019 FORMAT (///' ',23X,'ABSORBANCE',21X,'TRANSMITTANCE'///          BAND0620
1          ' ',5X,'FREQUENCY',5X,'OBSERVED',5X,'CALCULATED',10X,          BAND0630
2          'OBSERVED',5X,'CALCULATED')                                BAND0640
6020 FORMAT (A1,14,2X,F7.1,2(4X,F10.4),9X,F10.4,4X,F10.4,8X,2F10.4)    BAND0650
6021 FORMAT(1H0,//,5X,29HSUM OF SQUARES FOR ABSORBANCE,10X,32HSUM OF SQBAND0660
1UARES FOR TRANSMITTANCE//10X,E20.8,10X,E20.8)                  BAND0670
6022 FORMAT('0',3X,12,4X,3(F10.4,3X,'(',I2,')',3X))            BAND0680
6023 FORMAT (/ '0 LARGEST ABSORBANCE ERROR AT',F7.1,' CM-1'//          BAND0690
1          ' OBSERVED TRANS =',F7.4,5X,'CALCD TRANS =',F7.4)        BAND0700

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6024	FORMAT ('O OBSERVED AND CALCD CURVES HAVE BEEN PLOTTED')	BAND0710
6025	FORMAT ('O',5X,'** DATA SMOOTHED WITH ',I3,' POINT POLYNOMIAL OF DEGREE ',I3,' **')	BAND0720
6026	FORMAT (1H0,5X,44H"THEORETICAL CURVE CALCULATED AT INTERVALS OF 5.28 CM-1)	BAND0730
6027	FORMAT (//1H ,32X,2A4)	BAND0740
6028	FORMAT (' ',7X,I2,4X,3(F10.5,11X))	BAND0750
6029	FORMAT ('O',5X,'.... NO IMPROVEMENT IN 10THS CYCLE')	BAND0760
6030	FORMAT ('O ERROR CURVE HAS BEEN PLOTTED')	BAND0770
6031	FORMAT ('O',24X,'BETA =',F9.5,5X,'HOLD =',I3/	BAND0780
1	'O',29X,'DAMPING FACTOR =',F5.2//)	BAND0790
6032	FORMAT ('O',4X,'BETA =',F9.5)	BAND0800
6033	FORMAT (//'O',18X,'BETA =',F8.4,5X,'STD DEV =',1PE11.4)	BAND0810
801	FORMAT (15,3X,F7.0,F5.1,15)	BAND0820
802	FORMAT (2014)	BAND0830
C		BAND0840
C		BAND0850
C....	INITIALIZE.	BAND0860
C		BAND0870
	CALL GDATE (DATE)	BAND0880
	CALL DISK	BAND0890
	LIGHT = 1	BAND0900
	IFLAG = 1	BAND0910
C		BAND0920
C....	READ TWO ALPHANUMERIC INFORMATION CARDS, WRITE PAGE HEADING.	BAND0930
C		BAND0940
1	READ (5,5000,END=43) KARD,(ALPHA(I),I=1,19)	BAND0950
	IF(KARD.EQ.1) GO TO 3	BAND0960
	IF (KARD.EQ.-9) GO TO 43	BAND0970
2	IF(LIGHT.EQ.1) WRITE (6,6000) KARD	BAND0980
	LIGHT = 2	BAND0990
	GO TO 1	BAND1000
3	READ (5,5000) KARD,(ALPHA(I),I = 20,38)	BAND1010
	IF(KARD.NE.2) GO TO 2	BAND1020
	WRITE (6,6001) (ALPHA(I),I = 1,38)	BAND1030
		BAND1040
		BAND1050

```

WRITE (6,6027) DATE
C
C.... READ PROBLEM CONTROL CARD.
C
  READ (5,5001) KARD,NOPROB,NPLT,LPROB,NPEAK,IHOLD,MODE,NGATE,
  1 NPOLY,IFPNCH,IFPRNT,TSTP,DX,FACTOR
  IF (KARD.NE.3) GO TO 2
  IF (FACTOR.LE.0.0) FACTOR = 1.0
C
C.... READ IN SPECTRAL DATA.
C
  CALL SPCTRM (LIGHT,MODE,EDX,&1)
  IF (DX.LE.0.0) DX = EDX
  IF (NGATE.EQ.0) GO TO 7
C
C.... SMOOTH THE DATA.
C
  CALL SMTH1 (NGATE,NPOLY,WS)
  CALL SMTH2 (NXPTS,ETR,ETR)
C
C.... INITIALIZE PLOT ROUTINE.
C
  7 IF (NPLT.EQ.0) GO TO 31
  GO TO (20,31), IFLAG
  20 CALL PLOT1
  IFLAG = 2
C
C.... INITIALIZE MATRIX NHOLD. NHOLD(I,1) IS A FLAG TO INDICATE IF
C.... A PARAMETER IS TO BE HELD CONSTANT. NHOLD(I,2) CONTAINS
C.... THE NUMBER OF CYCLES FOR WHICH THE PARAMETER IS TO BE HELD
C.... CONSTANT.
C
  31 DO 30 I = 1,31
  NHOLD(I,1) = 1
  NHOLD(I,2) = 0

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BAND1060
BAND1070
BAND1080
BAND1090
BAND1100
BAND1110
BAND1120
BAND1130
BAND1140
BAND1150
BAND1160
BAND1170
BAND1180
BAND1190
BAND1200
BAND1210
BAND1220
BAND1230
BAND1240
BAND1250
BAND1260
BAND1270
BAND1280
BAND1290
BAND1300
BAND1310
BAND1320
BAND1330
BAND1340
BAND1350
BAND1360
BAND1370
BAND1380
BAND1390
BAND1400

```

30 CONTINUE	BAND1410
C	BAND1420
C.... READ PEAK PARAMETER CARDS.	BAND1430
C	BAND1440
DO 4 I = 1,NPEAK	BAND1450
READ(5,5006) KARD,E0(I),IHE0(I),DELNU(I),IHRHO(I),XNU(I),IHNU(I)	BAND1460
RHO(I) = 4.0/(DELNU(I) * DELNU(I))	BAND1470
IF ((KARD.NE.4) GO TO 2	BAND1480
4 CONTINUE	BAND1490
C	BAND1500
C.... READ ESTIMATE OF CONSTANT TERM, NUMBER OF ITERATIONS IT IS TO BE	BAND1510
C.... HELD CONSTANT, THE DEGREE OF POLYNOMIAL TO BE USED IN SLIT CON-	BAND1520
C.... VOLUTION, THE SPECTRAL SLIT WIDTH AND THE CELL THICKNESS.	BAND1530
C	BAND1540
READ (5,5003) KARD,BETA,IHBETA,MPOLY,SLIT,XL	BAND1550
IF(KARD.NE.5) GO TO 2	BAND1560
C	BAND1570
C.... READ CONCENTRATION VALUES TO BE ASSOCIATED WITH EACH PEAK.	BAND1580
C	BAND1590
READ(5,5005) KARD,(C(J),J = 1,10)	BAND1600
IF(KARD.NE.6) GO TO 2	BAND1610
C	BAND1620
C.... CALCULATE PEAK ABSORBANCE ESTIMATES.	BAND1630
C	BAND1640
DO 5 J = 1,NPEAK	BAND1650
AO(J) = E0(J) * XL * C(J)	BAND1660
5 CONTINUE	BAND1670
IF(SLIT.LE.0.0) GO TO 51	BAND1680
C	BAND1690
C.... BUILD VECTOR TO BE USED IN SLIT CONVOLUTION.	BAND1700
C	BAND1710
CALL SCAN(SLIT,DX,MPOLY,NMAX,D,WS)	BAND1720
NMP = NMAX + 1	BAND1730
DO 50 NI = 1,NMP	BAND1740
CON(NI) = D(NI)	BAND1750

50 CONTINUE	BAND1760
GO TO 6	BAND1770
51 NMAX = 0	BAND1780
CON(1) = 1.0	BAND1790
C	BAND1800
C.... WRITE PEAK PARAMETERS.	BAND1810
C	BAND1820
6 WRITE(6,6006) NOPROB,NPEAK	BAND1830
DO 60 I = 1,NPEAK	BAND1840
WRITE(6,6022) I,E0(I),IHE0(I),DELNU(I),IHRHO(I),XNU(I),IHNU(I)	BAND1850
60 CONTINUE	BAND1860
WRITE (6,6031) BETA,IHBETA,FACTOR	BAND1870
WRITE (6,6026) DX	BAND1880
IF (NGATE.NE.0) WRITE (6,6025) NGATE,NPOLY	BAND1890
C	BAND1900
C.... BUILD MATRIX NHOLD.	BAND1910
C	BAND1920
NPK2 = 2 * NPEAK	BAND1930
NDEL = 0	BAND1940
IF(IHOLD.LE.0) GO TO 8	BAND1950
DO 602 L = 1,NPEAK	BAND1960
I = L + 1	BAND1970
IF(IHE0(L).LE.0) GO TO 600	BAND1980
NDEL = NDEL + 1	BAND1990
NHOLD(I,1) = 2	BAND2000
NHOLD(I,2) = IHE0(L)	BAND2010
600 IF(IHRHO(L).LE.0) GO TO 601	BAND2020
NDEL = NDEL + 1	BAND2030
IS = I + NPEAK	BAND2040
NHOLD(IS,1) = 2	BAND2050
NHOLD(IS,2) = IHRHO(L)	BAND2060
601 IF(IHNU(L).LE.0) GO TO 602	BAND2070
NDEL = NDEL + 1	BAND2080
IS = I + NPK2	BAND2090
NHOLD(IS,1) = 2	BAND2100

NHOLD(15,2) = IHNU(L)	BAND2110
602 CONTINUE	BAND2120
IF (IHBETA.LE.0) GO TO 8	BAND2130
NDEL = NDEL + 1	BAND2140
NHOLD(1,1) = 2	BAND2150
NHOLD(1,2) = IHBETA	BAND2160
C	BAND2170
C.... CALCULATE THE LIMITS OF THE FREQUENCY VECTOR ASSOCIATED WITH THE	BAND2180
C.... CALCULATED SPECTRUM.	BAND2190
C	BAND2200
8 ISLIT = SLIT/DX + 0.001	BAND2210
XSLIT = DX * ISLIT	BAND2220
FREQLO = ENU(NXPTS) - XSLIT	BAND2230
FREQHI = ENU(1) + XSLIT	BAND2240
NTPTS = (1.0/DX)*(FREQHI - FREQLO) + 1.0	BAND2250
IF(NTPTS .LT.5000) GO TO 12	BAND2260
WRITE (6,6003) FREQLO ,FREQHI	BAND2270
LIGHT = 2	BAND2280
GO TO 1	BAND2290
12 WRITE(6,6004) SLIT,FREQLO,FREQHI,XL,NXPTS	BAND2300
DO 13 J = 1,NPEAK	BAND2310
WRITE(6,6005) J,A0(J),C(J)	BAND2320
C	BAND2330
C.... SCALE CONCENTRATION VALUES AND PEAK ABSORBANCE ESTIMATES.	BAND2340
C	BAND2350
C(J) = C(J) * XL	BAND2360
A0(J) = A0(J) * 2.302585	BAND2370
13 CONTINUE	BAND2380
BETA = BETA * 2.302585	BAND2390
IF(LPROB.EQ.1) GO TO 17	BAND2400
C	BAND2410
C.... IF THE INPUT SPECTRAL DATA IS IN TERMS OF ABSORBANCE, CONVERT	BAND2420
C.... TO TRANSMITTANCE.	BAND2430
C	BAND2440
DO 14 K = 1,NXPTS	BAND2450

EXPN = - 2.302585 * ETR(K)	BAND2460
ETR(K) = EXP(EXPN)	BAND2470
14 CONTINUE	BAND2480
17 IF ((IHOLD.GE.0) WRITE (6,6010)	BAND2490
C	BAND2500
C.... BUILD FREQUENCY VECTOR FOR THE CALCULATED SPECTRUM.	BAND2510
C	BAND2520
DO 19 K=1,NTPTS	BAND2530
FREQ(K) = FREQHI - (K-1)*DX	BAND2540
19 CONTINUE	BAND2550
C	BAND2560
C.... INITIALIZE FOR FIRST CYCLE OF LEAST-SQUARES.	BAND2570
C	BAND2580
KGO = 1	BAND2590
ITER = 1	BAND2600
NROWS = 3 * NPEAK + 1	BAND2610
NCOLS = NROWS + 1	BAND2620
NPK = 0	BAND2630
KBETA = 0	BAND2640
LAST = 1	BAND2650
IF (IHOLD.LT.0) LAST = 2	BAND2660
SUMSQ = 1.0 E 30	BAND2670
DO 200 I = 1,NROWS	BAND2680
BSAVE(I) = 0.0	BAND2690
200 CONTINUE	BAND2700
C	BAND2710
C.... BEGIN ITERATIVE LEAST-SQUARES LOOP.	BAND2720
C.... CALCULATE DIMENSIONS OF COEFFICIENT MATRIX B FOR CURRENT CYCLE.	BAND2730
C	BAND2740
21 NROWB = NROWS - NDEL	BAND2750
NCOLB = NCOLS - NDEL	BAND2760
C	BAND2770
C.... INITIALIZE AND BUILD MATRIX A.	BAND2780
C	BAND2790
DO 210 I = 1,NROWS	BAND2800

DO 210 J = 1,NCOLS	BAND2810
A(I,J) = 0.0	BAND2820
210 CONTINUE	BAND2830
DELSQ = 0.0	BAND2840
CALL LORN3 (LAST)	BAND2850
C	BAND2860
C..... PROBLEM FINISHED?	BAND2870
C	BAND2880
GO TO (212,40),LAST	BAND2890
C	BAND2900
C.....NO, CALCULATE PERCENT IMPROVEMENT.	BAND2910
C	BAND2920
212 PERCNT = 1.0 - DELSQ/SUMSQ	BAND2930
C	BAND2940
C.....TEST FOR DIVERGENCE OR CONVERGENCE.	BAND2950
C	BAND2960
IF (PERCNT.LT.0.0) GO TO (294,290,29), KGO	BAND2970
IF (PERCNT.EQ.0.) GO TO 28	BAND2980
IF (PERCNT.LT.TSTP.AND.KGO.EQ.1) GO TO 38	BAND2990
GO TO (33,321,32), KGO	BAND3000
C	BAND3010
C.....BUILD A VECTOR CONTAINING THE INDECES OF THOSE ROWS OF MATRIX A	BAND3020
C.....WHICH WILL BE LOADED INTO MATRIX B.	BAND3030
C	BAND3040
22 K = 0	BAND3050
DO 221 I = 1,NCOLS	BAND3060
IGO = NHOLD(I,1)	BAND3070
GO TO (220,221),IGO	BAND3080
220 K = K + 1	BAND3090
NDX(K) = I	BAND3100
221 CONTINUE	BAND3110
C	BAND3120
C.....BUILD MATRIX B.	BAND3130
C	BAND3140
DO 223 I = 1,NROWB	BAND3150

IS = NDX(I)	BAND3160
DO 222 J = 1,NCOLB	BAND3170
JS = NDX(J)	BAND3180
B(I,J) = A(IS,JS)	BAND3190
222 CONTINUE	BAND3200
223 CONTINUE	BAND3210
C	BAND3220
C.....INVERT MATRIX B.	BAND3230
C	BAND3240
DO 27 K = 1,NROWB	BAND3250
IF(B(1,1).LT.0.000001) GO TO 36	BAND3260
TEMP = 1.0/B(1,1)	BAND3270
DO 23 J = 1,NROWB	BAND3280
JP = J + 1	BAND3290
WS(J) = TEMP * B(1,JP)	BAND3300
23 CONTINUE	BAND3310
WS(JP) = TEMP	BAND3320
DO 25 J = 2,NROWB	BAND3330
TEMP = B(J,1)	BAND3340
JM = J - 1	BAND3350
DO 24 L = 1,NROWB	BAND3360
LP = L + 1	BAND3370
B(JM,L) = B(J,LP) - TEMP * WS(L)	BAND3380
24 CONTINUE	BAND3390
B(JM,LP) = - TEMP * WS(LP)	BAND3400
25 CONTINUE	BAND3410
DO 26 J = 1,NCOLB	BAND3420
B(NROWB,J) = WS(J)	BAND3430
26 CONTINUE	BAND3440
27 CONTINUE	BAND3450
C	BAND3460
C.... SCALE CORRECTIONS BY DAMPING FACTOR.	BAND3470
C	BAND3480
DO 18 K=1,NROWB	BAND3490
B(K,1) = FACTOR * B(K,1)	BAND3500

18 CONTINUE	BAND3510
NP = NPEAK	BAND3520
JGO = 1	BAND3530
M = 0	BAND3540
IF (NHOLD(1,1).GT.1) GO TO 269	BAND3550
M = 1	BAND3560
BSAVE(1) = B(1,1)	BAND3570
BETA = BETA + BSAVE(1)	BAND3580
C	BAND3590
C.....SAVE CORRECTIONS FOR THOSE PARAMETERS NOT HELD CONSTANT.	BAND3600
C	BAND3610
269 DD 276 L = 2,NRDS	BAND3620
I = L - 1	BAND3630
IGO = NHOLD(L,1)	BAND3640
GO TO (270,275),IGO	BAND3650
C	BAND3660
C.....UPDATE ESTIMATES.	BAND3670
C	BAND3680
270 M = M + 1	BAND3690
BSAVE(L) = B(M,1)	BAND3700
271 IF (I.LE.NP) GO TO (272,273,274),JGO	BAND3710
NP = NP + NPEAK	BAND3720
JGO = JGO + 1	BAND3730
GO TO 271	BAND3740
272 A0(I) = A0(I) + B(M,1)	BAND3750
GO TO 276	BAND3760
273 IM = I - NPEAK	BAND3770
TEMP = RHO(IM)	BAND3780
TEST = B(M,1)	BAND3790
277 IF ((TEMP+TEST).GT.0.0) GO TO 278	BAND3800
TEST = 0.9 * TEST	BAND3810
GO TO 277	BAND3820
278 TEST1 = TEMP/4	BAND3830
279 TEST2 = TEMP + TEST	BAND3840
IF (TEST2.GT.TEST1) GO TO 280	BAND3850

TEST = 0.9 * TEST	BAND3860
GO TO 279	BAND3870
280 B(M,1) = TEST	BAND3880
BSAVE(L) = TEST	BAND3890
RHO(IM) = RHO(IM) + B(M,1)	BAND3900
GO TO 276	BAND3910
274 IM = I - NPK2	BAND3920
XNU(IM) = XNU(IM) + B(M,1)	BAND3930
GO TO 276	BAND3940
C	BAND3950
C.....TEST TO DETERMINE IF A PARAMETER BEING HELD CONSTANT SHOULD BE	BAND3960
C.....ALLOWED TO VARY.	BAND3970
C	BAND3980
275 IF(ITER.LE.NHOLD(L,2)) GO TO 276	BAND3990
NHOLD(L,1) = 1	BAND4000
NHOLD(L,2) = 0	BAND4010
NDEL = NDEL - 1	BAND4020
276 CONTINUE	BAND4030
IF(NHOLD(1,1).LE.1) GO TO 281	BAND4040
IF(ITER.LE.IHBETA) GO TO 281	BAND4050
NHOLD(1,1) = 1	BAND4060
NHOLD(1,2) = 0	BAND4070
NDEL = NDEL - 1	BAND4080
281 IF(NDEL.LT.0) NDEL = 0	BAND4090
C	BAND4100
C.....RECYCLE.	BAND4110
C	BAND4120
GO TO 21	BAND4130
C	BAND4140
C.....IF THE PROBLEM HAS DIVERGED, TRY 0.1 OF THE CALCULATED	BAND4150
C.....CORRECTIONS.	BAND4160
C	BAND4170
294 DO 293 I=1,NPEAK	BAND4180
IP = I + NPEAK + 1	BAND4190
IPP = IP + NPEAK	BAND4200

AO(I) = HOLDEO(I) + 0.1*BSAVE(I+1)	BAND4210
RHO(I) = HOLDRH(I) + 0.1*BSAVE(IP)	BAND4220
XNU(I) = HOLDXN(I) + 0.1*BSAVE(IPP)	BAND4230
293 CONTINUE	BAND4240
BETA = HOLDBE + 0.1*BSAVE(1)	BAND4250
KGD = 2	BAND4260
WRITE (6,6008) ITER,SUMSQ,DELSQ	BAND4270
GO TO 21	BAND4280
C	BAND4290
C.....IF PROBLEM STILL DIVERGES, TRY 0.1 OF THE CALCULATED	BAND4300
C.....CORRECTIONS FOR EACH PEAK IN TURN.	BAND4310
C	BAND4320
290 DO 292 I=1,NPEAK	BAND4330
AO(I) = HOLDEO(I)	BAND4340
RHO(I) = HOLDRH(I)	BAND4350
XNU(I) = HOLDXN(I)	BAND4360
292 CONTINUE	BAND4370
29 IF ((NHOLD(I,1).GT.1).OR.(KBETA.GT.0)) GO TO 291	BAND4380
BETA = HOLDBE + 0.1*BSAVE(1)	BAND4390
KBETA = 1	BAND4400
GO TO 35	BAND4410
291 NPK = NPK + 1	BAND4420
IF(NPK.GT.NPEAK) GO TO 39	BAND4430
NPK1 = NPK + NPEAK	BAND4440
NPK2 = NPK1 + NPEAK	BAND4450
AO(NPK) = HOLDEO(NPK) + 0.1 * BSAVE(NPK + 1)	BAND4460
RHO(NPK) = HOLDRH(NPK) + 0.1 * BSAVE(NPK1 + 1)	BAND4470
XNU(NPK) = HOLDXN(NPK) + 0.1 * BSAVE(NPK2 + 1)	BAND4480
IF (NPK.GT.1) GO TO 351	BAND4490
BETA = HOLDBE	BAND4500
GO TO 35	BAND4510
351 IT = NPK - 1	BAND4520
AO(IT) = HOLDEO(IT)	BAND4530
RHO(IT) = HOLDRH(IT)	BAND4540
XNU(IT) = HOLDXN(IT)	BAND4550

35 KGD = 3	BAND4560
GO TO 21	BAND4570
C	BAND4580
C.....UPDATE ERROR SUM OF SQUARES, PRINT CURRENT ESTIMATES.	BAND4590
C	BAND4600
32 WRITE (6,6009) NPK	BAND4610
GO TO 33	BAND4620
321 WRITE (6,6007)	BAND4630
33 SUMSQ = DELSQ	BAND4640
WRITE (6,6011) ITER,SUMSQ,PERCNT	BAND4650
KGD = 1	BAND4660
ITER = ITER + 1	BAND4670
NPK = 0	BAND4680
KBETA = 0	BAND4690
DO 34 I1 = 1,NPEAK	BAND4700
EO11 = A0(I1) * 0.4342945/C(I1)	BAND4710
DELX = 2.0/SQRT(RHO(I1))	BAND4720
WRITE(6,6028) I1,EO11 ,DELX ,XNU(I1)	BAND4730
C	BAND4740
C.....STORE RESULTS FROM CURRENT CYCLE.	BAND4750
C	BAND4760
HOLDE0(I1) = A0(I1)	BAND4770
HOLDRH(I1) = RHO(I1)	BAND4780
HOLDXN(I1) = XNU(I1)	BAND4790
34 CONTINUE	BAND4800
BB = 0.434295 * BETA	BAND4810
WRITE (6,6032) BB	BAND4820
HOLDBE = BETA	BAND4830
IF(ITER.GT.25) GO TO 37	BAND4840
C	BAND4850
C.....RECYCLE.	BAND4860
C	BAND4870
GO TO 22	BAND4880
C	BAND4890
C.....SINGULAR MATRIX, DO FINAL CALCULATION ON PREVIOUS ESTIMATES.	BAND4900

C		BAND4910
	36 WRITE (6,6012) K,ITER	BAND4920
	GO TO 390	BAND4930
C		BAND4940
C.....	PROBLEM EXCEEDED MAXIMUM NUMBER OF CYCLES.	BAND4950
C		BAND4960
	37 WRITE (6,6013)	BAND4970
	GO TO 40	BAND4980
C		BAND4990
C....	NO CHANGE IN SUM OF SQUARES, TERMINATE.	BAND5000
C		BAND5010
	28 WRITE (6,6015) ITER	BAND5020
	GO TO 40	BAND5030
C		BAND5040
C.....	PROBLEM COMPLETED.	BAND5050
C		BAND5060
	38 WRITE (6,6014) ITER	BAND5070
	GO TO 40	BAND5080
C		BAND5090
C.....	NO IMPROVEMENT ON TENTHS CYCLE, DO FINAL CALCULATION ON PREVIOUS	BAND5100
C.....	ESTIMATES.	BAND5110
C		BAND5120
	39 WRITE (6,6029)	BAND5130
	390 LAST = 2	BAND5140
	DO 391 I1 = 1,NPEAK	BAND5150
	AO(I1) = HOLDEO(I1)	BAND5160
	RHO(I1) =HOLDRH(I1)	BAND5170
	XNU(I1) = HOLDXN(I1)	BAND5180
	391 CONTINUE	BAND5190
	BETA = HOLDBE	BAND5200
	GO TO 21	BAND5210
C		BAND5220
C.....	PRINT FINAL RESULTS.	BAND5230
C		BAND5240
	40 WRITE (6,6016) NOPROB,DELSQ	BAND5250

WRITE (6,6017)	BAND5260
NDF = NXPTS - NROWB	BAND5270
M = 0	BAND5280
DO 45 J=1,NROWS	BAND5290
VAR(J) = 0.0	BAND5300
IF (IHOLD.LT.0) GO TO 45	BAND5310
JGO = NHOLD(J,1)	BAND5320
GO TO (44,45), JGO	BAND5330
44 M = M + 1	BAND5340
VAR(J) = DELSQ*B(M,M+1)/NDF	BAND5350
45 CONTINUE	BAND5360
DO 41 J = 1,NPEAK	BAND5370
JP = J + NPEAK	BAND5380
JPP = JP + NPEAK	BAND5390
AA = AO(J) * 0.4342945	BAND5400
EOJ = AA/C(J)	BAND5410
DEL = 2.0/SQRT(RHO(J))	BAND5420
SIGAO = 0.434295 * SQRT(VAR(J+1))	BAND5430
TEMP = 1./(RHO(J)**3)	BAND5440
SIGDEL = SQRT(TEMP*VAR(JP+1))	BAND5450
SIGXNU = SQRT(VAR(JPP+1))	BAND5460
41 WRITE (6,6018) EOJ,AA,SIGAO,DEL,SIGDEL,XNU(J),SIGXNU	BAND5470
BB = 0.434295 * BETA	BAND5480
SIGBET = 0.434295 * SQRT(VAR(1))	BAND5490
WRITE (6,6033) BB,SIGBET	BAND5500
CCHAR = BLANK	BAND5510
IF (IFPRNT.GT.1) CCHAR = ZERO	BAND5520
IF (IFPRNT.GT.0) WRITE (6,6019)	BAND5530
IBIG = 1	BAND5540
BIG = 0.0	BAND5550
SUMAB = 0.0	BAND5560
SUMTR = 0.0	BAND5570
DO 42 K = 1,NXPTS	BAND5580
ETRK = ETR(K)	BAND5590
RK = R(K)	BAND5600

XX = -ALOG10(ETRK)	BAND5610
XY = - ALOG10(RK)	BAND5620
SUMTR = SUMTR + (ETRK - RK) * (ETRK - RK)	BAND5630
SUMAB = SUMAB + (XX - XY) * (XX - XY)	BAND5640
ERRA = XX - XY	BAND5650
ERRT = ETRK - RK	BAND5660
TEST = ABS(ERRA)	BAND5670
IF (TEST.LE.BIG) GO TO 52	BAND5680
IBIG = K	BAND5690
BIG = TEST	BAND5700
52 CONTINUE	BAND5710
IF (IFPRNT.EQ.0) GO TO 42	BAND5720
WRITE (6,6020) CCHAR,K,ENU(K),XX,XY,ETRK,RK,ERRA,ERRT	BAND5730
42 CONTINUE	BAND5740
WRITE (6,6021) SUMAB,SUMTR	BAND5750
WRITE (6,6023) ENU(IBIG),ETR(IBIG),R(IBIG)	BAND5760
KGO = 1	BAND5770
LIGHT = 1	BAND5780
C	BAND5790
C.... PUNCH CALCULATED SPECTRUM.	BAND5800
C	BAND5810
IF (IFPNCH.EQ.0) GO TO 47	BAND5820
WRITE (8,801) NOPROB,ENU(1),EDX,NXPTS	BAND5830
DO 46 K=1,NXPTS	BAND5840
46 IR(K) = 1000*R(K) + 0.5	BAND5850
WRITE (8,802) (IR(K), K=1,NXPTS)	BAND5860
C	BAND5870
C.... PLOT RESULTS.	BAND5880
C	BAND5890
47 IF (NPLT.EQ.0) GO TO 1	BAND5900
XHI = ENU(1)	BAND5910
CALL STRING (C(1),4)	BAND5920
WRITE (0,6002) NOPROB	BAND5930
PLBL(9) = C(1)	BAND5940
CALL PLOT2 (XHI,EDX,NXPTS,PLBL,36,1)	BAND5950

CALL PLOT3 (R,XHI,EDX,NXPTS)	BAND5960
CALL PLOT4 (ETR,XHI,EDX,NXPTS,3)	BAND5970
WRITE (6,6024)	BAND5980
IF (NPLT.LT.2) GO TO 1	BAND5990
DO 48 K=1,NXPTS	BAND6000
ETRK = ETR(K)	BAND6010
ETRK = ALOG10(R(K)/ETRK)	BAND6020
IF (ETRK.LT.-0.2) ETRK = -0.2	BAND6030
IF (ETRK.GT.0.2) ETRK = 0.2	BAND6040
48 ETR(K) = ETRK	BAND6050
PLBL1(6) = C(1)	BAND6060
CALL PLOT2 (0.,0.,-1,PLBL1,24,2)	BAND6070
CALL PLOT3 (ETR,XHI,EDX,NXPTS)	BAND6080
WRITE (6,6030)	BAND6090
C	BAND6100
C.....GET A NEW PROBLEM.	BAND6110
C	BAND6120
GO TO 1	BAND6130
C	BAND6140
C.....END OF JOB.	BAND6150
C	BAND6160
43 IF (IFLAG.GT.1) CALL PLOT5	BAND6170
STOP	BAND6180
END	BAND6190
SUBROUTINE SCAN (SLIT,DX,NPOLY,MPNT,SCNSET,X)	SCAN0010
C WRITEN BY A S ROSENBERG 5 12 67	SCAN0020
C MODIFIED 5 7 69	SCAN0030
C	SCAN0040
C CALCULATES A SET OF CONVOLVING NUMBERS, SCNSET (OF DIMENSION	SCAN0050
C MPNT+1) TO BE USED IN SCANNING A DIGITIZED SPECTRUM WITH A TRIANG-	SCAN0060
C ULAR SPECTRAL SLIT FUNCTION.	SCAN0070
C	SCAN0080
C SLIT = SPECTRAL SLIT WIDTH IN CM-1.	SCAN0090

C		SCAN0100
C	DX = ABCISSA INCREMENT, IN CM-1, FOR THE DIGITIZED SPECTRUM.	SCAN0110
C		SCAN0120
C	NPOLY = THE DEGREE OF POLYNOMIAL TO BE USED IN FITTING ORDINATE	SCAN0130
C	VALUES OVER A TWO SPECTRAL SLIT WIDTH INTERVAL OF THE SPECTRUM.	SCAN0140
C	IF NPOLY IS LESS THAN 2, THE SET OF CONVOLVING NUMBERS WILL BE	SCAN0150
C	CALCULATED ACCORDING TO THE TRAPEZOID RULE.	SCAN0160
C		SCAN0170
C	X IS AN ARRAY (DOUBLE WORD) TO BE USED AS SCRATCH AREA FOR MATRIX	SCAN0180
C	INVERSION.	SCAN0190
C		SCAN0200
C	REQUIRES SUBROUTINE CONMAT.	SCAN0210
C		SCAN0220
	DIMENSION SCNSET(51)	SCAN0230
	REAL*8 A(4,55),X(55)	SCAN0240
C	INITIALIZATION	SCAN0250
	TERM = SLIT/DX	SCAN0260
	IF (NPOLY.LT.2) GO TO 5	SCAN0270
	MPNT = TERM + 0.5	SCAN0280
	IF (MPNT.GE.1) GO TO 1	SCAN0290
4	MPNT = 0	SCAN0300
	SCNSET(1) = 1.	SCAN0310
	RETURN	SCAN0320
1	IF (NPOLY.GT.6) NPOLY = 6	SCAN0330
	IF (NPOLY.LT.2) NPOLY = 2	SCAN0340
	IT = NPOLY/2	SCAN0350
	IF (IT.GT.MPNT) IT = MPNT	SCAN0360
C	CALCULATE A MATRIX.	SCAN0370
	NPOLY = 2*IT	SCAN0380
	NROW = IT + 1	SCAN0390
	IPAR = 0	SCAN0400
	LPNT = MPNT + 1	SCAN0410
	CALL CONMAT (NROW,LPNT,IPAR,A,X)	SCAN0420
C	SCALE A AND BUILD SCNSET.	SCAN0430
	TERM = TERM * TERM	SCAN0440

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TEMP = 1.
IT = 1
DO 2 I=1,LPNT
  SCNSET(I) = A(1,I)
2 CONTINUE
DO 3 I=2,NROW
  IT = IT + 2
  TEMP = TEMP * TERM
  FACTOR = IT*(IT + 1)
  FACTOR = 2.*TEMP/FACTOR
  DO 3 J=1,LPNT
    SCNSET(J) = SCNSET(J) + FACTOR*A(I,J)
3 CONTINUE
RETURN
5 MPNT = TERM + 0.001
IF (MPNT.LT.1) GO TO 4
LPNT = MPNT + 1
TEMP = SLIT - MPNT*DX
SUM = TEMP/2
SCNSET(1) = SUM
DO 6 I=2,MPNT
  TEMP = TEMP + DX
  SCNSET(I) = TEMP
  SUM = SUM + TEMP
6 CONTINUE
SCNSET(LPNT) = SLIT
SUM = 2*SUM + SLIT
DO 7 I=1,LPNT
  SCNSET(I) = SCNSET(I)/SUM
7 CONTINUE
RETURN
END

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SCAN0450
SCAN0460
SCAN0470
SCAN0480
SCAN0490
SCAN0500
SCAN0510
SCAN0520
SCAN0530
SCAN0540
SCAN0550
SCAN0560
SCAN0570
SCAN0580
SCAN0590
SCAN0600
SCAN0610
SCAN0620
SCAN0630
SCAN0640
SCAN0650
SCAN0660
SCAN0670
SCAN0680
SCAN0690
SCAN0700
SCAN0710
SCAN0720
SCAN0730
SCAN0740
SCAN0750
SCAN0760

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SUBROUTINE LORN3 (LAST)

LORN0010

COMMON	A(31,32),AO(10),CON(51),DELSQ,DX,ENU(5000),ETK(5000),	LORNO020
1	FREQ(5000),FREQHI,NCOLS,NMAX,NPEAK,NROWS,NTPTS,NXPTS,	LORNO030
2	PRTSUM(31),RHO(10),R(5000),T(5000),XNU(10),ISLIT,KEY,BETA	LORNO040
	REAL*8 A	LORNO050
C		LORNO060
C....	BUILD VECTOR OF ABSORBANCE VALUES FOR CALCULATED SPECTRUM.	LORNO070
C		LORNO080
	NROWM = NROWS - 1	LORNO090
	DO 100 J = 1,NTPTS	LORNO100
	T(J) = BETA	LORNO110
100	CONTINUE	LORNO120
	DO 2 J = 1,NPEAK	LORNO130
	RHOJ = RHO(J)	LORNO140
	XNUJ = XNU(J)	LORNO150
	DO 1 K = 1,NTPTS	LORNO160
	FDIFF = XNUJ - FREQ(K)	LORNO170
	DENK = 1.0 + RHOJ * FDIFF * FDIFF	LORNO180
	T(K) = T(K) + AO(J)/DENK	LORNO190
	1 CONTINUE	LORNO200
	2 CONTINUE	LORNO210
C		LORNO220
C....	CONVERT TO TRANSMITTANCE.	LORNO230
C		LORNO240
	DO 4 K = 1,NTPTS	LORNO250
	T(K) = EXP(-T(K))	LORNO260
	4 CONTINUE	LORNO270
C		LORNO280
C....	CONVOLVE CALCULATED SPECTRUM WITH SLIT FUNCTION, BUILD	LORNO290
C....	MATRIX A.	LORNO300
C		LORNO310
	NMP = NMAX + 1	LORNO320
	CONS = CON(NMP)	LORNO330
	K = ISLIT	LORNO340
	DO 14 J = 1,NXPTS	LORNO350
	K = K + 1	LORNO360

TK = T(K)	LORNO370
SUMTN = 0.0	LORNO380
JUMP = 1	LORNO390
DO 5 L = 1, NROWM	LORNO400
PRTSUM(L) = 0.0	LORNO410
5 CONTINUE	LORNO420
DO 11 I1 = 1, NPEAK	LORNO430
I2 = I1 + NPEAK	LORNO440
I3 = I2 + NPEAK	LORNO450
IF(NMAX .NE.0) GO TO 3	LORNO460
SUMTN = 0.0	LORNO470
GO TO 9	LORNO480
3 DO 8 N = 1, NMAX	LORNO490
NJ = NMP - N	LORNO500
CONIN = CON(NJ)	LORNO510
KPN = K + N	LORNO520
KMN = K - N	LORNO530
IF(KMN.EQ.0) KMN = 1	LORNO540
TKPN = T(KPN)	LORNO550
TKMN = T(KMN)	LORNO560
GO TO (50,51), JUMP	LORNO570
50 SUMTN = SUMTN + CONIN * (TKPN + TKMN)	LORNO580
51 GO TO (53,8), LAST	LORNO590
53 V3 = XNU(I1) - FREQ(KPN)	LORNO600
V4 = XNU(I1) - FREQ(KMN)	LORNO610
V1 = V3 * V3	LORNO620
V2 = V4 * V4	LORNO630
DENP = 1.0 + RHO(I1) * V1	LORNO640
DENM = 1.0 + RHO(I1) * V2	LORNO650
T1 = TKPN / (DENP * DENP)	LORNO660
T2 = TKMN / (DENM * DENM)	LORNO670
7 PRTSUM(I1) = PRTSUM(I1) + CONIN * (TKPN/DENP + TKMN/DENM)	LORNO680
PRTSUM(I2) = PRTSUM(I2) + CONIN * (V1 * T1 + V2 * T2)	LORNO690
PRTSUM(I3) = PRTSUM(I3) + CONIN * (V3 * T1 + V4 * T2)	LORNO700
8 CONTINUE	LORNO710

9	GO TO (90,10), JUMP	LORN0720
90	JUMP = 2	LORN0730
	R(J) = CONS * TK + SUMTN	LORN0740
	DELTA = ETR(J) - R(J)	LORN0750
	DELSQ = DELSQ + DELTA * DELTA	LORN0760
	GO TO (10,14),LAST	LORN0770
10	DENKD = XNU(I1) - FREQ(K)	LORN0780
	DENKV = DENKD * DENKD	LORN0790
	DENK = 1.0 + RHO(I1) * DENKV	LORN0800
	DENS = DENK * DENK	LORN0810
	PRTSUM(I1) = -(TK/DENK + PRTSUM(I1))	LORN0820
	PRTSUM(I2) = A0(I1) * (TK * DENKV/DENS + PRTSUM(I2))	LORN0830
	PRTSUM(I3) = 2.0 * RHO(I1) * A0(I1) * (TK * DENKD/DENS + PRTSUM(I3))	LORN0840
	13)	LORN0850
11	CONTINUE	LORN0860
	DD 13 I1 = 1,NROWM	LORN0870
	I1P = I1 + 1	LORN0880
	DD 12 I2 = I1,NROWM	LORN0890
	I2P = I2 + 1	LORN0900
	A(I1P,I2P) = A(I1P,I2P) + PRTSUM(I1)*PRTSUM(I2)	LORN0910
12	CONTINUE	LORN0920
	A(I1P,NCOLS) = A(I1P,NCOLS) + PRTSUM(I1)*DELTA	LORN0930
	A(1,I1P) = A(1,I1P) - PRTSUM(I1)*R(J)	LORN0940
13	CONTINUE	LORN0950
	A(1,1) = A(1,1) + R(J)*R(J)	LORN0960
	A(1,NCOLS) = A(1,NCOLS) - R(J)*DELTA	LORN0970
14	CONTINUE	LORN0980
	GO TO (140,16),LAST	LORN0990
C		LORN1000
C....	IF NOT THE LAST CYCLE, COMPLETE BUILDING OF MATRIX A.	LORN1010
C		LORN1020
140	DD 15 K = 2,NROWS	LORN1030
	KM = K - 1	LORN1040
	DD 15 L = 1,KM	LORN1050
	A(K,L) = A(L,K)	LORN1060

15 CONTINUE
16 RETURN
END

LORNI070
LORNI080
LORNI090

SUBROUTINE PLOT1
DIMENSION BUFF(750),Y(500),PLABEL(20)
C..... INITIALIZE.
CALL PLOTS (BUFF,3000)
CALL PLOT (0.,-31.,3)
CALL PLOT (0.,-29.,-3)
JOB = 0
SIZE = 0.07
TEST = SIZE + 0.01
TEST = TEST * TEST
RETURN
ENTRY PLOT2 (XHI,DX,NPNT,PLABEL,NLABEL,ITYPE)
C..... START A NEW PAGE.
JOB = JOB + 1
IF (JOB.EQ.1) GO TO 3
IT = MOD(JOB,2) + 1
GO TO (1,2),IT
1 CALL PLOT (-1.0,10.5,-3)
GO TO 3
2 CALL PLOT (13.0,-11.5,-3)
3 CALL PLOT (14.0,0.,2)
CALL PLOT (14.0,11.,2)
CALL PLOT (0.,11.,2)
CALL PLOT (0.,0.,2)
XJOB = JOB
CALL SYMBOL (0.5,0.07,0.14,'PLOT ',0.,5)
CALL NUMBER (1.1,0.07,0.14,XJOB,0.,-1)
CALL PLOT (1.0,0.5,-3)
C..... DRAW AXES.
IF (NPNT.LT.0) GO TO 7

PLOT0010
PLOT0020
PLOT0030
PLOT0040
PLOT0050
PLOT0060
PLOT0070
PLOT0080
PLOT0090
PLOT0100
PLOT0110
PLOT0120
PLOT0130
PLOT0140
PLOT0150
PLOT0160
PLOT0170
PLOT0180
PLOT0190
PLOT0200
PLOT0210
PLOT0220
PLOT0230
PLOT0240
PLOT0250
PLOT0260
PLOT0270
PLOT0280
PLOT0290
PLOT0300

XLO = XHI - (NPNT-1)*DX	PLOT0310
CALL SCALAB (XHI,XLO,DX,12.5,XSTART,XINCR)	PLOT0320
7 GO TO (8,9), ITYPE	PLOT0330
8 CALL AXIS (0.,0.,'TRANSMITTANCE',13,10.,90.,0.,0.1,10.)	PLOT0340
GO TO 10	PLOT0350
9 CALL AXIS (0.,0.,'ABSORBANCE',10,10.,90.,-0.2,0.04,10.)	PLOT0360
10 CALL AXIS (0.,0.,'FREQUENCY, CM-1',-15,12.5,0.,XSTART,XINCR,10.)	PLOT0370
CALL PLOT (12.5,0.,3)	PLOT0380
CALL PLOT (12.5,10.,2)	PLOT0390
CALL PLOT (0.,10.,2)	PLOT0400
CALL SYMBOL (0.5,10.18,0.14,PLABEL,0.,NLABEL)	PLOT0410
RETURN	PLOT0420
ENTRY PLOT3 (Y,XHI,DX,NPNT)	PLOT0430
C.....DRAW A LINE PLOT.	PLOT0440
FACTOR = 10.	PLOT0450
IF (ITYPE.EQ.1) GO TO 11	PLOT0460
DO 6 I=1, NPNT	PLOT0470
6 Y(I) = Y(I) + 0.2	PLOT0480
FACTOR = 25.	PLOT0490
11 YY = FACTOR * Y(I)	PLOT0500
XX = (XHI - XSTART)/XINCR	PLOT0510
DIF = DX/XINCR	PLOT0520
CALL PLOT (XX,YY,3)	PLOT0530
DO 4 I=2, NPNT	PLOT0540
XX = XX - DIF	PLOT0550
YY = FACTOR * Y(I)	PLOT0560
CALL PLOT (XX,YY,2)	PLOT0570
4 CONTINUE	PLOT0580
RETURN	PLOT0590
ENTRY PLOT4 (Y,XHI,DX,NPNT,IBCD)	PLOT0600
C.....DRAW A POINT PLOT.	PLOT0610
XOLD = -1.	PLOT0620
YOLD = -1.	PLOT0630
DIF = DX/XINCR	PLOT0640
XX = (XHI - XSTART)/XINCR + DIF	PLOT0650

DO 5 I=1, NPNT	PLOT0660
XX = XX - DIF	PLOT0670
YY = 10 * Y(I)	PLOT0680
DIST2 = (XOLD-XX)*(XOLD-XX) + (YOLD-YY)*(YOLD-YY)	PLOT0690
IF (DIST2.LE.TEST) GO TO 5	PLOT0700
CALL SYMBOL (XX,YY,SIZE,IBCD,0.,-1)	PLOT0710
XOLD = XX	PLOT0720
YOLD = YY	PLOT0730
5 CONTINUE	PLOT0740
RETURN	PLOT0750
ENTRY PLOT5	PLOT0760
C..... TERMINATE.	PLOT0770
CALL PLOT (18.0,0.,-3)	PLOT0780
CALL PLOT (0.,0.,999)	PLOT0790
RETURN	PLOT0800
END	PLOT0810

SUBROUTINE SMTH1 (NPNT, NPOLY, X)	SMTH0010
C WRITTEN BY A S ROSENBERG 5 12 67	SMTH0020
C FIRST SECTION:	SMTH0030
C CALCULATES A MATRIX B TO BE USED IN A LEAST-SQUARES SMOOTHING OF	SMTH0040
C DIGITAL DATA.	SMTH0050
C NPNT = NUMBER OF DATA POINTS IN SMOOTHING GATE.	SMTH0060
C NPOLY = DEGREE OF POLYNOMIAL TO BE USED TO FIT NPNT ORDINATE	SMTH0070
C VALUES.	SMTH0080
C X IS A DOUBLE-WORD ARRAY USED FOR SCRATCH.	SMTH0090
C REQUIRES SUBROUTINE CONMAT.	SMTH0100
DIMENSION B(7,13)	SMTH0110
REAL*8 A(4,17), X(25)	SMTH0120
C INITIALIZE.	SMTH0130
KOUNT = 1	SMTH0140
IF (MOD(NPNT,2).NE.1) NPNT = NPNT + 1	SMTH0150
IF (NPNT.GT.25) NPNT = 25	SMTH0160
IF (NPNT.LT.5) NPNT = 5	SMTH0170

	IF (NPOLY.GT.6)	NPOLY = 6	SMTH0180
	IF (NPOLY.LT.2)	NPOLY = 2	SMTH0190
	IF (NPOLY.GE.NPNT)	NPOLY = NPNT - 2	SMTH0200
C	COMPUTE PARITY AND DIMENSIONS OF INTERMEDIATE MATRIX A.		SMTH0210
	IPAR = MOD(NPOLY,2)		SMTH0220
	NROW = NPOLY/2 + 1		SMTH0230
	LPNT = (NPNT - 1)/2 + 1		SMTH0240
C	CALCULATE EVEN (OR ODD) ROWS OF B.		SMTH0250
	1 CALL CONMAT (NROW,LPNT,IPAR,A,X)		SMTH0260
	I = IPAR + 1		SMTH0270
	DO 3 J=1,NROW		SMTH0280
	DO 2 K=1,LPNT		SMTH0290
	B(I,K) = A(J,K)		SMTH0300
	2 CONTINUE		SMTH0310
	I = I + 2		SMTH0320
	3 CONTINUE		SMTH0330
C	FINISHED WITH FIRST SECTION?		SMTH0340
	IF (KOUNT.GT.1) RETURN		SMTH0350
C	NO.....CALCULATE ODD (OR EVEN) ROWS OF B.		SMTH0360
	KOUNT = 2		SMTH0370
	I = NPOLY - 1		SMTH0380
	NROW = I/2 + 1		SMTH0390
	IPAR = MOD(I,2)		SMTH0400
	GO TO 1		SMTH0410
C	SECOND SECTION:		SMTH0420
	ENTRY SMTH2 (N,YIN,YOUT)		SMTH0430
C	OPERATES ON AN ARRAY OF DIGITIZED DATA YIN WITH MATRIX B. THE		SMTH0440
C	SMOOTHED DATA IS PLACED IN ARRAY YOUT. YIN AND YOUT MAY REFER TO		SMTH0450
C	THE SAME ARRAY IN THE CALLING PROGRAM.		SMTH0460
C	N IS THE DIMENSION OF ARRAYS YIN AND YOUT.		SMTH0470
	DIMENSION YIN(1),YOUT(1)		SMTH0480
	DIMENSION COEF(7)		SMTH0490
C	STORE FIRST NPNT DATA POINTS IN TEMPORARY STORAGE.		SMTH0500
	DO 4 I=1,NPNT		SMTH0510
	X(I) = YIN(I)		SMTH0520

4	CONTINUE	SMTH0530
	MPNT = LPNT - 1	SMTH0540
C	CALCULATE SMOOTHED VALUES FOR FIRST LPNT DATA POINTS.	SMTH0550
	NROW = NPOLY + 1	SMTH0560
	KOUNT = 1	SMTH0570
5	DO 8 K=1,NROW	SMTH0580
	SUM1 = 0.0	SMTH0590
	IPAR = MOD(K,2)	SMTH0600
	DO 7 L=1,MPNT	SMTH0610
	M = NPNT - L + 1	SMTH0620
	IF ((IPAR.EQ.0) GO TO 6	SMTH0630
	SUM1 = SUM1 + (X(L) + X(M))*B(K,L)	SMTH0640
	GO TO 7	SMTH0650
6	SUM1 = SUM1 + (X(L) - X(M))*B(K,L)	SMTH0660
7	CONTINUE	SMTH0670
	COEF(K) = SUM1 + X(LPNT)*B(K,LPNT)	SMTH0680
8	CONTINUE	SMTH0690
	IF (KOUNT.GT.1) GO TO 14	SMTH0700
	DO 10 I=1,LPNT	SMTH0710
	L = I - LPNT	SMTH0720
	J = 1	SMTH0730
	SUM1 = 0.0	SMTH0740
	DO 9 K=1,NROW	SMTH0750
	SUM1 = SUM1 + J*COEF(K)	SMTH0760
	J = J * L	SMTH0770
9	CONTINUE	SMTH0780
	YOUT(I) = SUM1	SMTH0790
10	CONTINUE	SMTH0800
C	SMOOTH MAIN BODY OF DATA.	SMTH0810
	M = N - 2*MPNT	SMTH0820
	DO 13 I=2,M	SMTH0830
	DO 11 J=2,MPNT	SMTH0840
	K = J - 1	SMTH0850
	X(K) = X(J)	SMTH0860
11	CONTINUE	SMTH0870

X(NPNT) = YIN(NPNT + I - 1)	SMTH0880
SUM1 = 0.	SMTH0890
DO 12 J=1,MPNT	SMTH0900
K = NPNT - J + 1	SMTH0910
SUM1 = SUM1 + (X(J) + X(K))*B(1,J)	SMTH0920
12 CONTINUE	SMTH0930
YOUT(MPNT+I) = SUM1 + X(LPNT)*B(1,LPNT)	SMTH0940
13 CONTINUE	SMTH0950
C OPERATE ON LAST MPNT DATA POINTS.	SMTH0960
KOUNT = 2	SMTH0970
GO TO 5	SMTH0980
14 DO 16 I=1,MPNT	SMTH0990
J = 1	SMTH1000
SUM1 = 0.0	SMTH1010
DO 15 K=1,NROW	SMTH1020
SUM1 = SUM1 + COEF(K)*J	SMTH1030
J = I * J	SMTH1040
15 CONTINUE	SMTH1050
L = N - MPNT + I	SMTH1060
YOUT(L) = SUM1	SMTH1070
16 CONTINUE	SMTH1080
RETURN	SMTH1090
END	SMTH1100

SUBROUTINE CONMAT (NROW,LPNT,IPAR,A,X)	CNMT0010
C WRITTEN BY A S ROSENBERG 5 12 67	CNMT0020
C CORRECTED 7 10 67	CNMT0030
C BUILDS AND INVERTS AN AUGMENTED MATRIX A USED IN CALCULATING A	CNMT0040
C LEAST-SQUARES CONVOLVING SET.	CNMT0050
C THE DIMENSIONS OF A ARE NROW BY (NROW + LPNT).	CNMT0060
C IF IPAR=0, AN EVEN SET OF POLYNOMIALS IS INVOLVED. IF IPAR=1, THE	CNMT0070
C SET IS ODD.	CNMT0080
C X IS A VECTOR USED AS WORKING STORAGE IN THE MATRIX INVERSION	CNMT0090
C ROUTINE.	CNMT0100

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REAL*8  A(4,55),X(55),SUM(7),TEMP
C INITIALIZE AND BUILD FIRST ROW OF AUGMENTING MATRIX.
IT = NROW + 1
MPNT = LPNT - 1
NCOL = NROW + LPNT
IF ((IPAR.NE.0)) GO TO 2
DO 1 I=1,LPNT
  J = I + NROW
  A(1,J) = 1.
  X(I) = 1.
1 CONTINUE
GO TO 4
2 DO 3 I=1,LPNT
  J = I - LPNT
  K = NROW + I
  A(1,K) = J
  X(I) = J*J
3 CONTINUE
C BUILD REMAINDER OF AUGMENTING MATRIX.
4 IF (NROW.LE.1) GO TO 13
DO 5 I=2,NROW
DO 5 J=1,LPNT
  K = LPNT - J
  K = K*K
  L = NROW + J
  A(I,L) = K*A(I-1,L)
5 CONTINUE
C CALCULATE SUMS
13 SUM(1) = 0.0
DO 6 I=1,MPNT
  SUM(1) = SUM(1) + 2.*X(I)
6 CONTINUE
SUM(1) = SUM(1) + X(LPNT)
IF (NROW.LE.1) GO TO 14
IT = 2*NROW - 1

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CNMT0110
CNMT0120
CNMT0130
CNMT0140
CNMT0150
CNMT0160
CNMT0170
CNMT0180
CNMT0190
CNMT0200
CNMT0210
CNMT0220
CNMT0230
CNMT0240
CNMT0250
CNMT0260
CNMT0270
CNMT0280
CNMT0290
CNMT0300
CNMT0310
CNMT0320
CNMT0330
CNMT0340
CNMT0350
CNMT0360
CNMT0370
CNMT0380
CNMT0390
CNMT0400
CNMT0410
CNMT0420
CNMT0430
CNMT0440
CNMT0450

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      DO 7 K=2,IT
      SUM(K) = 0.0
      DO 7 I=1,MPNT
      J = I - LRNT
      J = J*J
      X(I) = J*X(I)
      SUM(K) = SUM(K) + 2.*X(I)
7 CONTINUE
C BUILD LEFT-HAND SIDE OF MATRIX.
14 DO 8 I=1,NROW
      DO 8 J=1,NROW
      K = I + J - 1
      A(I,J) = SUM(K)
8 CONTINUE
C INVERT AUGMENTED MATRIX.
      NCOLM = NCOL - 1
      DO 12 I=1,NROW
      TEMP = 1.0/A(I,1)
      DO 9 J=1,NCOLM
      JP = J + 1
      X(J) = TEMP * A(I,JP)
9 CONTINUE
      X(JP) = TEMP
      IF (NROW.LE.1) GO TO 15
      DO 11 J=2,NROW
      TEMP = A(J,1)
      JM = J - 1
      DO 10 L=1,NCOLM
      LP = L + 1
      A(JM,L) = A(J,LP) - TEMP*X(L)
10 CONTINUE
      A(JM,LP) = -TEMP*X(LP)
11 CONTINUE
15 DO 12 J=1,NCOL
      A(NROW,J) = X(J)

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CNMT0460
CNMT0470
CNMT0480
CNMT0490
CNMT0500
CNMT0510
CNMT0520
CNMT0530
CNMT0540
CNMT0550
CNMT0560
CNMT0570
CNMT0580
CNMT0590
CNMT0600
CNMT0610
CNMT0620
CNMT0630
CNMT0640
CNMT0650
CNMT0660
CNMT0670
CNMT0680
CNMT0690
CNMT0700
CNMT0710
CNMT0720
CNMT0730
CNMT0740
CNMT0750
CNMT0760
CNMT0770
CNMT0780
CNMT0790
CNMT0800

```

12 CONTINUE
RETURN
END

CNMT0810
CNMT0820
CNMT0830

```
      SUBROUTINE SCALAB (XMAX,XMIN,XINCR,S,XSTART,DX)          SCAL0010
C.....WRITTEN BY A S ROSENBERG      7 3 67                SCAL0020
C.....ABCISSA SCALING ROUTINE FOR SUBROUTINE PLOT1.        SCAL0030
C.....USED IN PLACE OF SUBROUTINE SCALE IN CALCOMP PLOTTER PACKAGE. SCAL0040
      DX = (XMAX - XMIN)/S                                    SCAL0050
      IF (DX) 1,7,2                                          SCAL0060
7 DX = 1.0                                                  SCAL0070
      XSTART = XMAX + 0.5                                    SCAL0080
      RETURN                                                SCAL0090
1 TEMP = XMAX                                              SCAL0100
  XMAX = XMIN                                              SCAL0110
  XMIN = TEMP                                              SCAL0120
  DX = -DX                                                 SCAL0130
2 IDX = 0.4343 * ALOG(DX)                                   SCAL0140
  IXMX = XMAX * (10.0**(-IDX)) + 0.99                     SCAL0150
  XSTART = IXMX                                           SCAL0160
  XSTART = XSTART * (10.0**IDX)                           SCAL0170
  DX = ALOG10 ((XSTART - XMIN)/S)                          SCAL0180
  IDX = DX                                                 SCAL0190
  TEMP = IDX                                               SCAL0200
  DX = 10.0**(DX-TEMP)                                     SCAL0210
  TEMP = 1.0                                               SCAL0220
3 IF (DX - 1.0) 4,6,5                                       SCAL0230
4 DX = DX * 10.0                                           SCAL0240
  IDX = IDX - 1                                           SCAL0250
  GO TO 3                                                  SCAL0260
5 TEMP = 1.5                                               SCAL0270
  IF (DX.LE.1.5) GO TO 6                                   SCAL0280
  TEMP = 2.0                                               SCAL0290
  IF (DX.LE.2.0) GO TO 6                                   SCAL0300
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TEMP = 4.0	SCAL0310
IF (DX.LE.4.0) GO TO 6	SCAL0320
TEMP = 5.0	SCAL0330
IF (DX.LE.5.0) GO TO 6	SCAL0340
TEMP = 8.0	SCAL0350
IF (DX.LE.8.0) GO TO 6	SCAL0360
TEMP = 10.0	SCAL0370
6 DX = TEMP * (10.0**IDX)	SCAL0380
DX = -DX	SCAL0390
RETURN	SCAL0400
END	SCAL0410

SUBROUTINE SPCTRM (LIGHT,MODE,EDX,*)	SPCT0010
DIMENSION IT(20),IDNO(500),INDEX(500),XDATE(2),Y(10000)	SPCT0020
COMMON A(31,32),A0(10),CON(51),DELSQ,DX,ENU(5000),ETR(5000),	SPCT0030
1 FREQ(5000),FREQHI,NCOLS,NMAX,NPEAK,NROWS,NIPTS,NXPTS,	SPCT0040
2 PRISUM(31),RHO(10),R(5000),T(5000),XNU(10),ISLIT,KEY,BETA	SPCT0050
EQUIVALENCE (R(1),Y(1)),(T(1),Y(5001))	SPCT0060
REAL*8 A	SPCT0070
IF (MODE.EQ.2) GO TO 6	SPCT0080
C.... READ SPECTRUM HEADER CARD.	SPCT0090
READ (5,501) IDENT,XSTART,EDX,NIPTS,BEGIN,END	SPCT0100
501 FORMAT (I5,3X,F7.0,F5.1,I5,5X,2F10.5)	SPCT0110
C.... CALCULATE INDICES OF INITIAL AND FINAL POINTS TO BE INPUT.	SPCT0120
JBRNCH = 1	SPCT0130
10 XEND = XSTART - (NIPTS-1)*EDX	SPCT0140
IF (BEGIN.GT.XSTART.OR.BEGIN.LT.XEND) BEGIN = XSTART	SPCT0150
IF (END.GT.BEGIN.OR.END.LT.XEND) END = XEND	SPCT0160
IBEGIN = (XSTART - BEGIN)/EDX + 1	SPCT0170
IEND = (XSTART - END)/EDX + 1	SPCT0180
NXPTS = IEND - IBEGIN + 1	SPCT0190
IF (NXPTS.LE.5000) GO TO (1,9), JBRNCH	SPCT0200
C.... ERROR RETURN.	SPCT0210
WRITE (6,601)	SPCT0220

601	FORMAT ('0',5X,'TOO MANY DATA POINTS, PROBLEM TERMINATED')	SPCT0230
	LIGHT = 2	SPCT0240
	RETURN 1	SPCT0250
C....	INPUT LOOP.	SPCT0260
1	IBRNCH = 1	SPCT0270
	L = 1	SPCT0280
	DO 5 I=1,NIPTS,20	SPCT0290
	IN = 20	SPCT0300
	IP = I + 19	SPCT0310
	IF (IP.GT.NIPTS) IN = 20 + NIPTS - IP	SPCT0320
C....	READ ONE DATA CARD.	SPCT0330
	READ (5,502) (IT(J), J=1,IN)	SPCT0340
502	FORMAT (2014)	SPCT0350
	GO TO (2,5), IBRNCH	SPCT0360
2	DO 4 J=1,IN	SPCT0370
	K = I + J - 1	SPCT0380
C....	TEST TO SEE IF DATA POINT IS WITHIN PROPER SPECTRAL REGION.	SPCT0390
	IF (K.LT.IBEGIN) GO TO 4	SPCT0400
	IF (K.LE.IEND) GO TO 3	SPCT0410
	IBRNCH = 2	SPCT0420
	GO TO 5	SPCT0430
C....	TRANSFER VALUE TO TRANSMITTANCE VECTOR AND CALCULATE	SPCT0440
C....	CORRESPONDING ELEMENT OF FREQUENCY VECTOR.	SPCT0450
3	ETR(L) = IT(J)/1000.	SPCT0460
	ENU(L) = BEGIN - (L-1)*EDX	SPCT0470
	L = L + 1	SPCT0480
4	CONTINUE	SPCT0490
5	CONTINUE	SPCT0500
	WRITE (6,603) IDENT,BEGIN,END	SPCT0510
603	FORMAT (/'0',9X,'SPECTRUM',I6,' READ FROM CARDS, RANGE =',	SPCT0520
1	F7.1,' TO',F7.1)	SPCT0530
	RETURN	SPCT0540
6	READ (5,503) IDENT,BEGIN,END	SPCT0550
503	FORMAT (I5,5X,2F10.5)	SPCT0560
	DO 7 I=1,NSPEC	SPCT0570

IF (IDENT.NE.IDNO(I)) GO TO 7	SPCT0580
KEY = INDEX(I)	SPCT0590
GO TO 8	SPCT0600
7 CONTINUE	SPCT0610
WRITE (6,602) IDENT	SPCT0620
602 FORMAT ('0',5X,'SPECTRUM',I6,' NOT ON DISK')	SPCT0630
LIGHT = 2	SPCT0640
RETURN 1	SPCT0650
8 READ (7'KEY) IDENT,XSTART,EDX,NIPTS,INST,XDATE,(Y(I), I=1,315)	SPCT0660
JBRNCH = 2	SPCT0670
GO TO 10	SPCT0680
9 IF (IEND.LE.315) GO TO 13	SPCT0690
DO 11 I=316,IEND,322	SPCT0700
IP = I + 321	SPCT0710
READ (7'KEY) (Y(J), J=I,IP)	SPCT0720
11 CONTINUE	SPCT0730
13 DO 12 I=1,NXPTS	SPCT0740
IM = I - 1	SPCT0750
IP = IM + IBEGIN	SPCT0760
ETR(I) = Y(IP)	SPCT0770
ENU(I) = BEGIN - IM*EDX	SPCT0780
12 CONTINUE	SPCT0790
WRITE (6,604) IDENT,BEGIN,END	SPCT0800
604 FORMAT ('/0',9X,'SPECTRUM',I6,' READ FROM DISK, RANGE =',	SPCT0810
1 F7.1,' TO',F7.1)	SPCT0820
RETURN	SPCT0830
ENTRY DISK	SPCT0840
KEY = 1	SPCT0850
READ (7'KEY) NSPEC,NREC,IDNO,INDEX	SPCT0860
RETURN	SPCT0870
END	SPCT0880

VITA

2

Arthur Stanley Rosenberg

Candidate for the Degree of

Doctor of Philosophy

Thesis: I. A VIBRATIONAL STUDY OF THE CYANOETHYLENES
II. AN ALGORITHM FOR THE DETERMINATION OF THE SHAPE
PARAMETERS OF INFRARED ABSORPTION BANDS

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