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## TIME DOMAIN IDENTIFICATION OF THE <br> ZEROS OF LINEAR SYSTEMS

A DISSERTATION<br>SUBMITTED TO THE GRADUATE FACULTY<br>in partial fulfillment of the requirements for the degree of<br>DOCTOR OF PHILOSOPHY

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TIME DOMAIN IDENTIFICATION OF THE ZEROS OF LINEAR SYSTEMS

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## ABSTRACT

In recent research in the Process Control Laboratory of the University of Oklahoma a time domain technique for the identification of linear systems was formulated, at which time the identification of the system's poles was investigated. The major goal of the present work was to complete the development and verification of the technique with particular reference to determination of the system zeros.

The procedure for determining the system's zeros is based on the analysis of the input and response functions. For an $n^{\text {th }}$ order system with $n-1$ zeros, $n$, linearly-independent input-output records are required. A prior knowledge of the number of zeros, perhaps from theoretical considerations, is nelpful but not necessary. The correct number of zeros can be determined from interpretation of the identification results.

The investigation of the technique was conducted utilizing analog and digital computer simulations of second and third order systems. Studies were made to determine the sensitivity of the identification program to many factors which might be encountered in chemical processes. In adaition a laboratory heat exchange process was also used to confirm the experimental applicability of the technique.

The results of the identification of the laboratory process were compared to models obtained from frequency response testing and pulse testing. An analog simulation of the identified model was also used to compare the response of the model to the response of the actual process.

This technique has two major advantages over the common frequency domain methods. These advantages are: the system parameters are determined explicitly, and an error propagation analysis enables one to determine the uncertainty of the identified model.

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## TABLE OF CONTENTS

Page
LIST OF TABLES ..... viii
LIST OF ILLUSTRATIONS ..... ix
Chapter
I. INTRODUCTION ..... 1
II. THEORY ..... 6
III. ERROR PROPAGATION ANALYSIS ..... 13
Expected Error in the Coefficients of Expected Error in the System Zeros
IV. COMPUTER STUDIES ..... 23
Effect of the Forcing Function
Effect of the Relative Locationof the poles and Zeros
Effect of the order of the system
Effect of Error in the IdentifiedValues of the Poles
Effect of Steady State Errors
Effect of Transport Delay
Effect of the Number of SignificantFigures in the Data
Effect of Noise
Effect of System Non-linearities
Conclusion About the Identificationof Non-linear Systems
V. EXPERIMENTAL STUDIES ..... 66
Experimental EquipmentThe ReactorConstant Temperature BathsFlow Measurement and ControlSupport Equipment
Theoretical Description of the Experimental process
Chapter Page
Experimental Determination of the System by Standard Methods Transient Response Tests Frequency Response Tests Pulse Tests
Time Domain Identification
Pole Identification Zero Identification
VI. CONCLUSIONS AND RECOMMENDATIONS ..... 116
ConclusionsRecommendations for Future Work
REFERENCES ..... 119
APPENDICES
A. NOMENCLATURE ..... 122
B. NUMERICAL EXAMPLE OF ZERO CALCULATION ..... 124
C. LISTINGS OF IDENTIFICATION PROGRAMS ..... 129

## LIST OF TABLES

Table Page
4-1. Effect of the Relative Location of the Poles and Zeros ..... 27
4-2. Effect of the Relative Location of the Poles for and Assumed System Order Less than the Actual Order ..... 35
4-3. Comparison of the Magnitude of the Non- linearities in the Analog Model and the Experimental Process ..... 53
5-1. List of System Constants ..... 82
5-2. DC Gain of Experimental Equipment as Determined by Transient Response Tests ..... 84
5-3. DC Gain of Experimental Equipment as Determined by Frequency Response Tests ..... 87
5-4. DC Gain of Experimental Equipment as Determined by Pulse Tests ..... 88
5-5. Forcings Used for Pole Identification ..... 91
5-6. Results of Pole Identification ..... 93
5-7. Forcings Used for Zero Identification ..... 98
5-8. Test Combinations Used for Zero Identification ..... 101
5-9. Results of Zero Identification of Experimental Equipment ..... 102
5-10. Final Mean Results of Zero Identification of Experimental Equipment ..... 108
5-ll. Summary of DC Gains ..... 115
Figure Page
2-1. General Chemical Process ..... 6
4-1. Effect upon the Identification of the Frequency of the Forcing Function ..... 26
4-2. Approximate Identification of a Three Pole System Using Lesser Order Models ..... 30
4-3. Approximate Identification of a One Zero-Three Pole System Using Lesser Order Models ..... 31
4-4. Approximate Identification of a
Two Zero-Three Pole System
Using Lesser Order Models ..... 32
4-5. Effect of Magnitude of Neglected Pole on an Identification of Incorrect Order ..... 34
4-6. Identification of a Third Order System with a Five Percent Error in All of the Poles ..... 37
4-7. Effect of Input Steady State Error of One Percent ..... 39
4-8. Effect of Input Steady State Error of Ten Percent ..... 40
4-9. Effect of Steady State Error of Twenty Percent in the Response ..... 41
4-10. Effect of One Second Transport Delay ..... 43
4-11. Effect of Ten Second Transport Delay ..... 44
4-12. Effect upon the Identification of the Number of Significant Figures in the Data ..... 46
4-13. Schematic Diagram of Experimental Equipment . ..... 47
Figure Page
4-14. Effect of Five Percent Noise in the Response Function ..... 49
4-15. Analog Computer Circuit for Simulation of the Experimental Equipment ..... 52
4-16. Identification of Poles of Analog Simulation of Theoretical Non-linear Experimental Equipment Model ..... 55
4-17. Identification of the Oil-forced Non-linear Model of the Experimental Equipment, Positive Ramp and Step Forcing ..... 57
4-18. Identification of the Oil-forced Non-linear Model of the Experimental Equipment, Negative Ramp, Positive Step Forcing ..... 58
4-19. Identification of the Oil-forced Non-linear Model of the Experimental Equipment, Sinusoidal Forcing ..... 60
4-20. Predicted Error for the Identification Shown in Figure 4-19 ..... 61
4-21. Identification of the Coolant-forced, Non-linear Model with Forcing Function of a Step and a Ramp ..... 63
4-22. Identification of the Coolant-forced, Non-linear Model with One of the Forcing Functions a Sine Wave ..... 64
5-1. Detail of Reactor ..... 68
5-2. Oil Flow Control System ..... 73
5-3. Coolant Flow Control System ..... 74
5-4. Wall Temperature Measurement System ..... 75
5-5. Schematic Flow Sheet of Experimental System ..... 76
5-6. Frequency Response Tests of Experimental System ..... 86
5-7. Results of Pulse Testing of Experimental System ..... 89
Figure Page
5-8. Identification of the Poles of the Experimental System, Tests 1,3,5 ..... 95
5-9. Identification of the Poles of the Experimental System, Tests 2,4,6 ..... 96
5-10. Identification of the Coolant-forced, Experimental System with Step and Ramp Forcings ..... 103
5-1l. Identification of the Coolant-forced, Experimental System with Step and Sinusoidal Forcings ..... 105
5-12. Identification of the Oil-forced, Experimental System with Step and Ramp Forcings ..... 106
5-13. Identification of the Oil-forced, Experimental System with Step and Sinusoidal Forcings ..... 107
5-14. Bode Magnitude Plot for Comparison of Results of the Identification of the Coolant-forced System ..... 110
5-15. Phase Diagram of Results of Identification of the Coolant-forced System ..... 111
5-16. Bode Magnitude Plot for Comparison of
Results of the Identification of the Oil-forced System ..... 112
5-17. Phase Diagram of Results of Identification of the Oil-forced System ..... 113
5-18. Comparison of the Identified Model to Experimental Data for the Oil-forced System ..... 114
C-1. Block Diagram of Identification Technique ..... 130

# TIME DOMAIN IDENTIFICATION OF THE ZEROS OF LINEAR SYSTEMS 

CHAPTER I

## INTRODUCTION

In recent years great advances have been made in the techniques for the design and analysis of process controls. Basic to these techniques is a complete description of the dynamic behavior of the process. Theoretical considerations are useful in determining the general topology of the equations describing the process; however, it is frequently impossible to evaluate the actual parameters numerically with any degree of confidence. Since the theoretical approach may not yield a satisfactory dynamic model of the process, it is frequently necessary to rely upon experimental methods for determining é model.

Several analysis and testing techniques for system identification have been developed in the past. Although there is diversity in the mathematical techniques employed for data manipulation and in the types of test signals applied, all of the testing methods have been based on the analysis of input-output relationships. The basic methods may be classified as: l. sinusoidal (frequency response) testing, 2. pulse testing and Fourier transformation, and 3. random
or statistical testing. It is possible to consider all of these identification methods as frequency-domain techniques because at some stage of the analytic process, the data is transformed into the frequency domain. Each of these methods depend on frequency domain representations, usually magnitude and phase lag versus frequency diagrams (Bode plots), for the final plant evaluation.

Direct frequency response testing was adapted from electrical engineering. Although it has been widely used in the past, Hougen (H5) points out that it is generally unsatisfactory in the process industries. The principal virtue of this method is the simplicity with which the data can be interpreted. General drawbacks of the method are: (1) the lengthy tests may cause extreme deviations of the system from its normal mode of operation, (2) many tests over a wide range of frequencies are required regardless of the complexity of the system, and (3) even slight nonlinearities will distort the expected sinusoidal response (S3).

The necessity for evaluating the system model from the Bode plots or other frequency domain representations is a limitation which will be discussed later. Although work has been done on the utilization of direct frequency response testing, usually it has been used for model verification, rather than identification.

Pulse testing and Fourier transformation was first introduced by Hougen (H5). The method was developed to overcome some of the major problems encountered with direct
sinusoidal testing. The forcing variable is excited by a single pulse and the input and response are recorded. They are then transformed numerically into the Fourier domain. The frequency response function is obtained as the ratio of the transform of the response to the transform of the input. This method has a notable advantage over direct sinusoidal testing. The power spectral density of an impulse covers the whole frequency range; therefore a single test is theoretically adequate to evaluate the complete frequency response function, and it can replace a whole series of sinusoidal tests.

In practice the above statement must be qualified. It is physically impossible to generate an impulse, necessitating the use of a pulse with a power spectral density less than one. Dreifke (D2) has demonstrated the sensitivity of the identification to pulse shape and duration. Therefore some experimentation is necessary to determine the proper pulse for any given system. Dreifke, Hougen, and Mesmer (D3) have studied the problems of truncation error due to the use of limited (finite) record segments, i.e., use of response records which end before returning to steady-state. The effects of this truncation error and of improper input pulses appear in the resulting frequency response function as a "cutoff frequency". Cycling occurs at frequencies above the cut-off frequency and the scatter of the results becomes excessive, indicating that the response function is unreliable. To eliminate the problem of truncation one might record the
response function until steady state is "practically" reached; however, in chemical processes where the time constants are long, the possible drifts in steady state operating conditions may offset the decrease in truncation error.

Noise will effectively increase the truncation error. In the later stages of response curves, where the amplitudes are low, the effect of noise may completely overshadow the signai, and the frequency response function is unreliable.

The limitations due to noise have been incorporated into another identification technique, random testing followed by statistical analysis of the response curves. Since this present work does not involve this technique the reader is referred to Gallier (Gl) for a further discussion of random testing.

The methods discussed have an inherent disadvantage in the representation of the results. This representation is usually in the form of a magnitude versus frequency plot and phase lag versus frequency plot, comonly called the Bode plot. From the Bode plot it may be very difficult, if not impossible, to evaluate numerically the parameters of complex systems (system order greater than one).

Generally, the useful range of the magnitude ratio plot does not exceed one to two decades on the logarithmic frequency scale. Thus the identifiable poles and zeros are limited to this range. The identifiable order of the system is also limited to this range. Since the phase lag versus
frequency is usually less reliable than the magnitude plot, it offers little additional information.

Heymann ( $\mathrm{H} 2, \mathrm{H} 3$ ) has proposed a technique for linear system identification based exclusively on time domain analysis. This technique circumvents the problem of obtaining the system parameters from frequency response representations. The system parameters are explicitly determined; in addition, error bounds associated with each of the system parameters are ascertained. However, Heymann has formulated and verified only a portion of the technique, the pole identification. It is the major goal of this work to complete the formulation of the technique and to demonstrate its usefulness for system identification.

The completion of the time domain identification technique is accomplished by implementation of the theory describing the determination of the system zeros. An error propagation analysis is also developed to predict the reliability of the identified results.

The test for the applicability of the technique is conducted in two parts. Computer studies are used to determine the sensitivity of the method to the many factors which are encountered in actual chemical processes. The technique is then used to identify a laboratory process, and the results of the identification are compared to the results obtained from direct frequency response testing and pulse testing.

## CHAPTER II

THEORY

A synoptic review of the background theory necessary Eor the zero identification problem is presented here. For a complete analysis of the theory of the time domain identification technique the reader is referred to Heymann (H2,H3).

Any one of the outputs $\left(i^{\text {th }}\right.$ ) of a general linear system of the form


FIGURE 2-1 General Chemical Process
can be described in terms of an $n^{\text {th }}$ order linear ordinary differential equation of the form

$$
\begin{align*}
& \frac{d^{n} y_{i}(t)}{d t^{n}}+b_{1}(t) \frac{d^{n-1} y_{i}(t)}{d t^{n-1}}+\cdots+b_{n}(t) y_{i}(t)= \\
& \quad \sum_{j=1}^{m} g_{j 1}(t) \frac{d^{n-1} x_{j}(t)}{d t^{n-1}}+\cdots+g_{j n}(t) x_{j}(t) . \tag{2-1}
\end{align*}
$$

The system can also be expressed in operator notation as

$$
\begin{equation*}
L^{n}\left[y_{i}(t)\right]=\sum_{j=1}^{m} M_{j}\left(k_{j}\right)\left[x_{j}(t)\right] \tag{2-2}
\end{equation*}
$$

where $L^{n}[]$ and $M_{j}\left(k_{j}\right)[$ ] are linear differential operators, with time dependent parameters, of order $n$ and $k_{j}$ respectively, where $k_{j} \leq n-1$. If the right hand side of Equation (2-2) is lumped into a general time function $f(t)$, Equation (2-1) transforms into

$$
\begin{equation*}
L^{\mathrm{n}}\left[y_{i}(t)\right]=f(t) \tag{2-3}
\end{equation*}
$$

The response of the system given by Equation (2-3) can also be expressed in terms of the system weighting function, $W(t, \lambda)$, i.e. the impulse response, by the following equation:

$$
\begin{equation*}
y_{i}(t)=\int_{-\infty}^{t} w(t, \lambda) f(\lambda) d \lambda \tag{2-4}
\end{equation*}
$$

If $y_{i}(t)=0$ prior to $t=0$, i.e. the system is operating at steady state, then

$$
\begin{equation*}
y_{i}(t)=\int_{0}^{t} W(t, \lambda) f(\lambda) d \lambda \tag{2-5}
\end{equation*}
$$

Considering Equations (2-2) and (2-3), Equation (2-5) can be rewritten as

$$
\begin{align*}
y_{i}(t) & =\int_{0}^{t} W(t, \lambda) \sum_{j=1}^{m} M_{j}\left(k_{j}\right)\left[x_{j}(\lambda)\right] d \lambda \\
& =\sum_{j=1}^{m} \int_{0}^{t} w(t, \lambda) M_{j}\left(k_{j}\right)\left[x_{j}(\lambda)\right] d \lambda \tag{2-6}
\end{align*}
$$

Equation (2-6) can be expressed in an equivalent, but more convenient, form as

$$
\begin{equation*}
y_{i}(t)=\sum_{j=1}^{m} \int_{0}^{t} \dot{w}_{j}(t, \lambda) x_{j}(\lambda) d \lambda \tag{2-7}
\end{equation*}
$$

where the weighting function* $W_{j}(t, \lambda)$ is related to the weighting function $W(t, \lambda)$ in Equation $(2-6)$ by the expression

$$
\begin{gather*}
w_{j}(t, \lambda)=(-1) k_{j} \frac{d^{k}{ }_{j}}{d t^{k}}\left[g_{j 1}(t) w(t, \lambda)\right]+\ldots+  \tag{2-8}\\
g_{j k_{j}}(t) w(t, \lambda) .
\end{gather*}
$$

*The weighting function $W_{j}(t, \lambda)$, which relates a particular input variable $x_{j}(t)$ to the output $y_{i}(t)$, will be called the particular Weighting Function.

In operator form, Equation (2-8) can be written as

$$
\begin{equation*}
w_{j}(t, \lambda)=M_{j}\left(k_{j}\right) *[w(t, \lambda)] \tag{2-9}
\end{equation*}
$$

where $M_{j}\left(k_{j}\right){ }^{*}[]$ is the adjoint operator of $M_{j}\left(k_{j}\right)[]$.
Equation (2-7) relates the $i^{\text {th }}$ response to all of the input variables. If all of the input variables except $x_{e}(t)$ are kept zero, then Equation (2-7) can be written as:

$$
\begin{equation*}
y_{i}(t)=\int_{0}^{t} w_{e}(t, \lambda) x_{e}(\lambda) d \lambda \tag{2-10}
\end{equation*}
$$

where

$$
\begin{equation*}
w_{e}(t, \lambda)=(-I) k^{k} e \frac{d^{k} e}{d t^{k} e}\left[g_{e l}(t) w(t, \lambda)\right]+\ldots+ \tag{2-11}
\end{equation*}
$$

$$
g_{e k_{e}}(t) w(t, \lambda)
$$

$W(t, \lambda)$ is determined by the methods described by Heymann [H2, H3]; thus, in order to determine $W_{e}(t, \lambda)$ the operator $M_{e}{ }^{\left(k_{e}\right)}$ [ ] remains to be evaluated.

If the process is stationary (constant parameter), the evaluation of the operator $M_{e}\left(k_{e}\right)$ [ ] is simplified. Equation (2-11) can be written as:
$W_{e}(t, \lambda)=(-1) k_{e} \frac{d^{k} e}{d t^{k} e}\left[g_{e 1} W(t-\lambda)\right]+\cdots+g_{e k} W(t-\lambda)$
The coefficients $g_{e j}$ can be determined numerically by the
following procedure: the system is forced $k_{e}$ times with linearly independent forcing functions, which are recorded together with their corresponding responses. By numerical solution of $k_{e}$ relations of the type,

$$
\begin{gathered}
y_{i}(t)=(-1)^{k} e g_{e l} \int_{0}^{t} \frac{d^{k} e}{d t^{k} e} w(t-\lambda) x_{e}(\lambda) d \lambda+\cdots+ \\
g_{e k} \int_{0}^{t} w(t-\lambda) x_{e}(\lambda) d \lambda
\end{gathered}
$$

the coefficients $g_{e j}$ are evaluated.
For simplicity in the following development, only one of the input-output relationships will be considered. Equation (2-13) then reduces to
$y(t)=(-1)^{k} g_{I} \int_{0}^{t} \frac{d^{k}}{d t^{k}} w(t-\lambda) x(\lambda) d \lambda+\ldots+g_{k} \int_{0}^{t} w(t-\lambda) x(\lambda) d \lambda$

The computation of these integrals is accomplished by first calculating the coefficients of the weighting function and k-1 of its derivatives.

The homogeneous transfer function as determined by the pole identification is of the form

$$
\begin{equation*}
H(s)=\frac{1}{\left(s+\rho_{1}\right) \cdots\left(s+o_{n}\right)} \tag{2-15}
\end{equation*}
$$

## 11

The weighting function is then

$$
\begin{aligned}
W(t-\lambda) & =a_{1} e^{-\rho_{i}(t-\lambda)}+\cdots+a_{n} e^{-\rho_{n}(t-\lambda)} \\
& =\sum_{i=1}^{n} a_{i} e^{-\rho_{i}(t-\lambda)}
\end{aligned}
$$

where
$a_{i}=\frac{1}{\left(\rho_{1}-\rho_{i}\right) \cdots\left(\rho_{i-1}-\rho_{i}\right)\left(\rho_{i+1}-o_{i}\right) \cdots\left(\rho_{n}-\rho_{i}\right)}$.
The $j^{\text {th }}$ derivation of the weighting function is then given by

$$
\left.\left.\begin{array}{rl}
\frac{d^{j} W(t-\lambda)}{d t^{j}} & =(-1)^{j}\left(a_{1} \rho_{1}^{j} e^{-\rho_{1}}(t-\lambda)\right.
\end{array}\right) \ldots+a_{n} \rho_{n}^{j} e^{-\rho_{n}(t-\lambda)}\right)
$$

Equation (2-14) can now be written as a general term

$$
\begin{align*}
y(t) & =\sum_{j=0}^{k-1}(-1){ }^{j} g_{k-j} \int_{0}^{t} \sum_{i=1}^{n}(-1)^{j} a_{i} \rho_{i}^{j} e^{-\rho_{i}}(t-\lambda) x(\lambda) d \lambda  \tag{2-19}\\
& =\sum_{j=0}^{k-1} g_{k-j} \int_{0}^{t} \sum_{i=1}^{n} a_{i} \rho_{i}^{j} e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda
\end{align*}
$$

Consider Equation (2-19) as composed of two parts, the integral and a coefficient. The integral for any i is the same for all of the coefficients. Therefore, only $k$ integrations are required where $k^{2}$ integrations would be required by a straightforward solution of Equation (2-14).

To determine the unknowns $g_{k-j}, k$ simultaneous equations of the type of Equation (2-19) are necessary. These equations are generated by varying $x(\lambda)$ (the forcing function). The only constraint upon these $x(\lambda)$ 's is that they must all be linearly independent.

Theoretically the solution of the $k$ simultaneous equations at any time $T$ is sufficient to determine the particular weighting function. In practice, however, the uncertainty of the data portends the danger of relying upon the one solution. It is therefore desirable to obtain a solution at many points of time and to have some measure of the reliability of the results at these various points.

The error propagation analysis to be discussed in the next chapter provides this measure of the reliability.

## CHAPTER III

ERROR PROPAGATION ANALYSIS

Two types of error exist in a numerical computation. One is the error associated with the uncertainty of the data, which is primarily responsible for the uncertainty of the results. It is the purpose of the error propagation analysis to predict when the interaction of the uncertainties of the data have a minimum effect upon the final identification results. The other source of error is computer round-off or truncation error. This error effects the numerical value of the computation, but it has no appreciable effect upon the uncertainty of the results. This type of error is discussed more fully in the computer studies section on the effect of the number of significant figures in the data.

If a true value were known for the results of a computation, then a true error could be calculated. Whenever the true result is not known it is necessary to predict an error limit for each stage of the computation. This error limit may be either an extreme or an expected error. The present work is concerned with the latter approach.

Before proceeding with the details of the error propagations analysis, it may be helpful to review the sources of error and the computational steps involved.

The sources of error in the data are:

1. The accuracy of the homogeneous weighting function, i.e. the poles of the system.
2. The reliability and accuracy of the input and response functions.

Several factors enter into the reliability of the input and response functions. They are: values of the steady-state, measurement error, and noise. In this work it is assumed that any noise present is stationary and ergodic random noise. The identification values will therefore have an equal chance of being distributed above or below their mean value and essentially no drift will occur because of the integration of the input function. It is further assumed that the steady-state and measurement error can be lumped together into a percentage error term which is relatively small.

The computational steps involved in the identification are:

1. Calculation of the integral coefficients of $g_{e j}$ in the operator, $M_{e}^{(k)}[]$.
2. Solution of the simultaneous algebraic equations thus obtained.

For the purposes of the error analysis it is assumed that the calculation of the coefficients is done in one step. The solution of the simultaneous equations is performed in two steps, inversion of the coefficient matrix and premultiplication of the response vector by the coefficient inverse matrix.

## Expected Error in the Coefficients of $g_{e i}$

If $F$ is a function of $x, y$, and $z$, the linear term in the Taylor series can be used to express the effect of a small error in $x, y$, and $z$. Thus, if $\Delta x$ is the error in $x$ and $\Delta y$ is the error in $y$ and $\Delta z$ is the error in $z$, the expected error in $F$ is $\Delta F$. These errors may be related by

$$
\begin{equation*}
\Delta F=\frac{\partial F}{\partial x} \Delta x+\frac{\partial F}{\partial y} \Delta y+\frac{\partial F}{\partial z} \Delta z \tag{3-1}
\end{equation*}
$$

Applying this relationship then to the integral part of Equation (2-19)

$$
\begin{equation*}
G_{k}=\sum_{i=1}^{n} \int_{0}^{t} a_{i} e^{-\theta_{i}(t-\lambda)} x(\lambda) d \lambda \tag{3-2}
\end{equation*}
$$

yields

$$
\begin{equation*}
\Delta G_{k}=\frac{\partial G_{k}}{\partial \rho_{1}} \Delta \rho_{1}+\cdots+\frac{\partial G_{k}}{\partial \rho_{n}} \Delta \rho_{n}+\frac{\partial G_{k}}{\partial x} \Delta x \tag{3-3}
\end{equation*}
$$

where $\Delta \rho_{1}, \ldots, \Delta \rho_{n}$ are the expected errors in the poles of the homogeneous weighting function and $\Delta x$ is the expected error in the input. The coefficient is given by Equation (2-17) as

$$
\begin{equation*}
a_{i}=\frac{1}{\left(\rho_{1}-\rho_{i}\right) \ldots\left(\rho_{i-1}-\rho_{i}\right)\left(\rho_{i+1}+\rho_{i}\right) \ldots\left(\rho_{n}-\rho_{i}\right)} \tag{2-17}
\end{equation*}
$$

The various partial derivatives are required to evaluate the expected error $\Delta G_{k}$. Differentation of the general term of Equation (3-2)

$$
\begin{equation*}
G_{k}(i)=\int_{0}^{t} a_{i} e^{-p_{i}(t-\lambda)} x(\lambda) d \lambda \tag{3-4}
\end{equation*}
$$

with respect to $\rho_{j}$ yields two results:
when $\quad j \neq i$

$$
\begin{equation*}
\frac{\partial G_{k}(i)}{\partial \rho_{j}}=\frac{\partial a_{i}}{\partial \rho_{j}} \int_{0}^{t} e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda=\frac{\partial a_{i}}{\partial \rho_{j}} \frac{G_{k}(i)}{a_{i}} \tag{3-5}
\end{equation*}
$$

and when $j=i$
$\frac{\partial G_{k}(i)}{\partial \rho_{i}}=\frac{\partial a_{i}}{\partial \rho_{i}} \int_{0}^{t} e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda+a_{i} \int_{0}^{t}(t-\lambda) e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda$

$$
\begin{equation*}
=\frac{\partial a_{i}}{\partial \rho_{i}} \frac{G_{k}(i)}{a_{i}}+t G_{k}(i)-a_{i} \int_{0}^{t} \lambda e^{-\theta_{i}}(t-\lambda) x(\lambda) d \lambda . \tag{3-6}
\end{equation*}
$$

From the definition of $a_{i}$, Equation (2-17), the various partials of $a_{i}$ can be obtained. When $i \neq j$ only one term of $a_{i}$ contains $\rho_{j}$ so the derivative is straightforward and after rearrangement is

$$
\begin{equation*}
\frac{\partial a_{i}}{\partial \rho_{j}}=\frac{a_{i}}{\rho_{j}-\rho_{i}} \tag{3-7}
\end{equation*}
$$

When $i=j, \rho_{i}$ is contained in every term. It is necessary to differentiate the function as a group of successive products. After differentiation and rearrangement the result is

$$
\begin{align*}
\frac{\partial a_{i}}{\partial \rho_{i}}= & -a_{i}\left(\frac{1}{\rho_{1}-\rho_{i}}+\cdots+\frac{1}{\rho_{i-1}-\rho_{i}}\right.  \tag{3-8}\\
& \left.+\frac{1}{\rho_{i+1}-\rho_{i}}+\cdots+\frac{1}{\rho_{n}-\rho_{i}}\right)
\end{align*}
$$

For simplicity let

$$
\begin{equation*}
d a_{i}=\frac{1}{\rho_{1}-\rho_{i}}+\cdots+\frac{1}{\rho_{n}-\rho_{i}} \tag{3-9}
\end{equation*}
$$

and Equation (3-8) becomes

$$
\begin{equation*}
\frac{\partial a_{i}}{\partial \sigma_{i}}=-a_{i} d a_{i} \tag{3-10}
\end{equation*}
$$

The last term of Equation (3-6)

$$
\begin{equation*}
H_{K}(i)=a_{i} \int_{0}^{t} \lambda e^{-p_{i}(t-\lambda)} x(\lambda) d \lambda \tag{3-11}
\end{equation*}
$$

can be expressed in terms of $G_{k}(i)$. When the right hand side of Equation (3-11) is integrated by parts, the resultant equation, with the use of Equation (3-4), becomes

$$
\begin{equation*}
H_{k}(i)=t G_{k}(i)-\int_{0}^{t} G_{k}(i) d \lambda . \tag{3-12}
\end{equation*}
$$

After substitution of Equations (2-7), (3-10), and (3-12) into Equation (3-5) and (3-6), the resultant partial derivatives are

$$
\begin{align*}
& \frac{\partial G_{k}(i)}{\partial \rho_{j}}=\frac{G_{k}(i)}{\rho_{j}-\rho_{i}}  \tag{a}\\
& \frac{\partial G_{k}(i)}{\partial \rho_{i v}}=-G_{k}(i) d a_{i}+\int_{0}^{t} G_{k}(i) d \lambda . \tag{3-13}
\end{align*}
$$

The next term of the operator $M^{(k)}[$ ]is

$$
\begin{equation*}
G_{k-1}=\int_{0}^{t} \frac{d w(t-\lambda)}{d t} x(\lambda) d \lambda \tag{3-14}
\end{equation*}
$$

The form of the general term of this expression can be expressed by

$$
\begin{align*}
G_{k-1} & =\int_{0}^{t} a_{i} \frac{d e^{-\rho_{i}}}{d t}(t-\lambda) \\
& =-a_{i} \rho_{i} \int_{0}^{t} e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda  \tag{3-15}\\
& =-\rho_{i} G_{k}(i)
\end{align*}
$$

This procedure can be carried out for all of the terms of the operator $M^{(k)}[]$, and general expressions can be developed to describe the coefficients and their partial derivatives. These general expressions are

$$
\begin{align*}
G_{n}(i) & =(-1)^{k-n} \rho_{i}^{k-n} G_{k}(i)  \tag{a}\\
\frac{\partial G_{n}(i)}{\partial \rho_{j}} \Delta \rho_{j} & =(-1)^{k-n} \rho_{i}^{k-n} \frac{\partial G_{k}(i)}{\partial \rho_{j}} \Delta \rho_{j} \tag{b}
\end{align*}
$$

$\frac{\partial G_{n}(i)}{\partial \rho_{i}} \Delta \rho_{i}=(-1) \rho_{i}^{k-n-1}-\left[\rho_{i} \frac{\partial G_{k}(i)}{\partial \rho_{i}}+(k-n) G_{k}(i)\right] \Delta \rho_{i} .(c)$

All that remains to describe completely the error propagation in the computation of the integral coefficients is to define the last term in Equation (3-3), ( $\partial \mathrm{G} / \Delta \mathrm{x}) \Delta \mathrm{x}$. From the general term of the describing equation for $G_{k}$

$$
\begin{equation*}
G_{k}(i)=\int_{0}^{t} a_{i} e^{-\rho_{i}(t-\lambda)} x(\lambda) d \lambda \tag{3-4}
\end{equation*}
$$

the error caused by a small change in $x(\lambda)$ can be determined. This error is

$$
\begin{equation*}
\frac{\partial G_{k}(i)}{\partial x} \Delta x=\int_{0}^{t} a_{i} e^{-\theta_{i}(t-\lambda)} \Delta x(\lambda) d \lambda \tag{3-18}
\end{equation*}
$$

The errors caused by the input for the other terms of the operator $M^{(k)}[]$ can be found in a similar manner. The general
expression is

$$
\begin{equation*}
\frac{\partial G_{n}(i)}{\partial x} \Delta x=a_{i}(-1)^{k-n} \rho_{i}^{k-n} \int_{0}^{t} e^{-\rho_{i}(t-\lambda)} \Delta x(\lambda) d \lambda . \tag{3-19}
\end{equation*}
$$

$G_{n}(i)$ is calculated from the general term of the weighting function and is obtained by summing all of the terms of the weighting function:

$$
\begin{equation*}
G_{n}=\sum_{i=1}^{k} G_{n}(i) \tag{a}
\end{equation*}
$$

$$
\begin{equation*}
\frac{\partial G_{n}}{\partial \rho_{j}} \Delta \rho_{j}=\sum_{i=1}^{k} \frac{\partial G_{n}(i)}{\partial \rho_{j}} \Delta \rho_{j} . \tag{3-20}
\end{equation*}
$$

Substituting Equations (3-19) and (3-20) into Equation (3-3), the total error in $G_{n}$ is obtained

$$
\begin{equation*}
\Delta G_{n}=\sum_{j=1}^{k} \sum_{i=1}^{k} \frac{\partial G_{n}(i)}{\partial \rho_{j}} \Delta \rho_{j}+\sum_{i=1}^{k} \frac{\partial G_{n}(i)}{\partial x} \Delta x . \tag{3-21}
\end{equation*}
$$

As $n$ goes from $l$ to $k, k$-element row vector is generated which corresponds to the row vector for the integral coefficients $G_{n}$. When $k$ inputs have been computed in this manner, a $k \times k$ matrix of error values is obtained corresponding to the coefficient matrix.

## Expected Error in the System Zeros

The next stage of the computation is the inversion of the coefficient matrix and premultiplication of the response vector by the inverse. In matrix notation

$$
\begin{equation*}
g=G^{-1} Y \tag{3-22}
\end{equation*}
$$

where $G$ is the coefficient matrix, $g$ is the solution vector and $Y$ is the response vector. The desired error is the error in the $g$ vector. Again using the total differential

$$
\begin{equation*}
\Delta g=\frac{\partial g}{\partial G^{-1}} \Delta G^{-1}+\frac{\partial g}{\partial Y} \Delta Y \tag{3-23}
\end{equation*}
$$

a relationship for the error is obtained. The partial derivative of $g$ with respect to $G^{-1}$ and $Y$ are easy to find as they are independent functions. They are

$$
\begin{align*}
& \frac{\partial g}{\partial G^{-I}}=Y  \tag{a}\\
& \frac{\partial g}{\partial Y}=G^{-1} . \tag{3-24}
\end{align*}
$$

$\Delta G^{-1}$ is harder to determine as it is derived from the error matrix obtained at stage one of the calculation. It is not simply the inverse of this matrix. The easiest way to determine $\Delta G^{-1}$ is a brute force method

$$
\begin{equation*}
\Delta G^{-1}=(G+\Delta G)^{-1}-G^{-1} \tag{3-25}
\end{equation*}
$$

but it does give a good estimate of the error limits in the inverse. $\Delta Y$ is estimated from the response curves.

## COMPUTER STUDIES

Many tests of the identification technique were conducted using digital and analog computer generated data. This technique was used to provide complete control of all conditions which might affect the identification. In this manner it was possible to study the specific effect caused by a potential identification problem. Some of the conditions of interest were: errors in the poles and steady state values, transportation delays, non-linearities, and noise. Other tests were run to determine the best forcing functions and also to determine how the relative location of the poles and zeros affected the identification. The digitally-generated data were obtained as the solution of the following general transfer function

$$
\begin{equation*}
Y(s)=\frac{G(A s+I)(B s+I)}{(C s+1)(D s+l)(E s+1)} \tag{4-1}
\end{equation*}
$$

with three different forcing functions

$$
\begin{equation*}
X(s)=\frac{K}{s}, \frac{I}{s^{2}}, \frac{M}{s^{2}+\omega^{2}} \tag{4-2}
\end{equation*}
$$

In the time domain these are respectively: a step of amplitude, $K$; a ramp of slope, L; and a sinusoid of amplitude, $M$, and angular frequency, $\boldsymbol{\omega}$.

The time domain expressions including transport delay and steady state were programmed as:

$$
\begin{align*}
Y_{\text {step }}(t+\tau)+Y_{S S}=X_{S S} & +K G\left(1-a_{1} e^{-t / C}-a_{2} e^{-t / D}-a_{3} e^{-t / E}\right) \\
Y_{\text {ramp }}(t+\tau)+Y_{S S}=X_{S S} & +L G\left(A+B-C-D-E+t+C a_{1} e^{-t / C}\right. \\
& +D a_{2} e^{-t / D_{+E a_{3}} e^{-t / E)}}\left(\begin{array}{rl}
(b)
\end{array}\right.  \tag{4-3}\\
Y_{\text {sine }}(t+\tau)+Y_{s S}=X_{S S} & +M G\left[-\left(b_{1}+b_{2}+b_{3}\right) \cos \omega t\right. \\
& +\left(1-\omega q_{1}-\omega b_{2}-\omega b_{3}\right) \sin \omega t \\
& +b_{1} e^{-t / C_{+b_{2}} e^{-t / D}} \\
& +b_{3} e^{-t / E]} \tag{c}
\end{align*}
$$

where

$$
\begin{aligned}
& a_{1}=(A-C) \quad(B-C) /(D-C) \quad(E-C) \\
& a_{2}=(A-D) \quad(B-D) /(C-D) \quad(E-D) \\
& a_{3}=(A-E) \quad(B-E) /(C-E) \quad(D-E) \\
& b_{1}=C \omega a_{1} /\left(1+C^{2} \omega^{2}\right) \\
& b_{2}=D \omega a_{2} /\left(1+D^{2} \omega^{2}\right) \\
& b_{3}=E \omega a_{2} /\left(1+E^{2} \omega^{2}\right)
\end{aligned}
$$

## Effect of Forcing Functions

The first series of tests were conducted to determine the best range of frequencies for a forcing sine wave. Figure 4-1 shows some of the results of the gain identification obtained when the system

$$
\begin{equation*}
\frac{Y}{X}(s)=\frac{80(25 s+1)(5 s+1)}{(10 s+1)(20 s+1)(50 s+1)} \tag{4-4}
\end{equation*}
$$

was forced with a step, a ramp, and a series of sine waves, ranging from 0.0001 radians/second to 1 radian/second.

As can be seen in Figure $4-1$, the results were not very good at the two extremes ( 0.0001 radians/second and 1 radian/second). The results did, however, tend to converge to the correct answer. The best results were obtained when the forcing frequencies were between 0.01 and 0.2 radians/ second. As this condition was also the case when two or three sine functions of different frequencies were used for the input functions, a general rule of thumb is obtained. The angular frequency of the forcing sine functions should be between one half the lowest natural frequency and twice the highest natural frequency.

The identification of the zeros is affected in the same manner as the gain. Part $B$ of Figure 4-l shows the identification results of the zeros.


Figure 4-1. Effect of Frequency of Forcing Sine Function upon the Identification.

## Effect of the Relative Location <br> of Poles and Zeros

The relative natural frequencies of the poles and zeros had very little effect upon the identification. Tests were made on third order systems where factors of as much as 250 existed between the largest and smallest poles, with the zero ranging from less than to greater than the poles. The rate of convergence to the correct answer was the only effect upon the identification. The convergence was the slowest when the zeros were less than the poles (Table 4-1).

TABLE 4-1
Effect of the relative location of the poles and zeros

| Pole 1 Pole 2 Pole 3 Zero 1 | Zero 2 | Time to Converge <br> to Less Than <br> l-Percent Error <br> (Sec) |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :---: |
| 0.1 | 0.05 | 0.02 | 0.04 | 0.01 | 18 |
| 0.1 | 0.05 | 0.005 | 0.2 | 0.01 | 12 |
| 0.5 | 0.05 | 0.02 | 0.2 | 0.04 | 12 |
| 0.1 | 0.05 | 0.02 | 0.1 | 0.04 | 8 |
| 0.1 | 0.05 | 0.02 | 0.2 | 0.05 | 8 |
| 0.1 | 0.05 | 0.02 | 0.2 | 0.02 | 10 |
| 0.1 | 0.02083 | 0.02 | 0.2 | 0.02 | 10 |
| 0.1 | 0.02083 | 0.02 | 0.02083 | 0.02 | 16 |
| 0.1 | 0.05 | 0.02 | 0.2 | 0.04 | 8 |
| 0.5 | 0.02 | 0.002 | 0.2 | 0.04 | 24 |

## Effect of the Order of the System

To determine the effect of an incorrect estimate of the number of poles of a system, i.e. the order of the system, upon the zero identification, a series of identification tests was conducted with poles deleted. In addition a series of tests with poles added was carried out. When the assumed order of the system was greater than the actual order, no problems arose since zeros were determined which cancelled the extra poles.

When the assumed system order was less than the actual order, the quality of the identification depended upon the relative values of the poles and zeros. In all cases the gain and zeros approach their values asymptotically; the rate of convergence and initial displacement again depend upon the relative values of the poles and zeros.

Figure 4-2 shows the results of the gain identification when the system is correctly identified as a third order system,

$$
\begin{equation*}
Y(s)=\frac{\sigma X(s)}{(s+0.5)(s+0.02)(s+0.002)} \tag{4-5}
\end{equation*}
$$

as a second order system

$$
\begin{equation*}
Y(s)=\frac{G X(s)}{(s+0.02)(s+0.002)} \tag{4-6}
\end{equation*}
$$

and as a first order system

$$
\begin{equation*}
Y(s)=\frac{g X(s)}{(s+0.002)} \tag{4-7}
\end{equation*}
$$

When the correct poles are used the result of the identification is within 0.2 percent of the correct value of 1.0 by the time of two seconds. However, when identified as a second order system the value of the gain increases with time, eventually becoming asymptotic to 2.0 , which is the correct gain divided by the magnitude of the neglected pole (1/0.5). The DC gain of the system of Equation (4-5) is determined by replacing $s$ by $j \omega$ and setting the frequency $\omega$ equal to zero. $D C$ gain $=1 /(0.5 \times 0.02 \times 0.002)=50000$. To match this $D C$ gain, the gain, $g$, of the system of Equation (4-6) must be equal to the neglected term (1/0.5). When identified as a first order system, the initial error is very large, but the gain may be approaching the theoretical value of the gain of the third order system divided by the neglected poles, $1 /(0.5)$ $(0.02)=100$.

Figure 4-3 gives the results of the same type of tests when the actual system contains one zero. When identified as a second order system the zero approaches the correct value of 0.01 while the gain approaches the value of the correct gain divided by the omitted pole, 2.0. When identified as a first order system $\mathrm{g} / \mathrm{s}+0.002$, the gain levels off at about I. 6 and does not approach the theoretical value of 1.0 . No explanation has been found for this behavicr.

Figure 4-4 shows the results of these tests on a system which contains two zeros. When one poie is neglected, a zero is also omitted from necessity. The omitted zero is the highest frequercy zero $(s+0.2)$. The zero at 0.04 is


Figure 4-2. Approximate Identification of a Three pole System Using Lesser Order Models.


Figure 4-3. Approximate Identification of a One Zero-Three Pole System Using Lesser Order Models.


Figure 4-4. Approximate Identification of a Two Zero-Three Pole System Using Lesser Order Models.
approached asymptotically. The value of the gain approaches the correct gain times the omitted zero divided by the neglected pole $1.0 \times 0.2 / 0.5=0.4$. When the technique is applied assuming a first order system $\mathrm{g} / \mathrm{s}+0.002$, the gain appears to level out at a value of 0.57. The theoretical result is $0.2 \times 0.04 / 0.5 \times 0.02=0.8$.

To determine the effect of the neglected pole upon the identified gain in relation to the magnitude of the neglected pole, a series of three tests of third order systems was run using two poles ( $\rho_{1}=-0.02, \rho_{2}=-0.01$ ) in the identification. The third pole, which was omitted in the identification, was varied from 0.1 to 0.50 . The results of these tests are given in Figure 4-5. To show the relative rates of convergence, the scales have been normalized. If the result of a correct third order identification were also shown on Figure 4-5, it would be a straight line at 1.0 . When the system was identified correctly, the gain was in error only 0.2 percent at time $t=2$ in the worst case.

The preceding tests were conducted upon a system which had only a gain and no zeros. When a zero is added, convergence is also dependent upon the natural frequency of the zero relative to the omitted pole. The results of the gain identification of a system when the neglected pole had a value four times the magnitude of the zero and one where the neglected pole was only twice the magnitude of the zero are also shown in Figure 4-5. As can be seen, the identification


Figure 4-5. Effect of Magnitude of Neglected Pole on an Identification of Incorrect Order.

TABLE 4-2.
EFFECT ON THE IDENTIFICATION OF THE RELATIVE LOCATION OF THE POLES FOR AN ASSUMED SYSTEM ORDER LESS THAN THE ACTUAL ORDER.

| Pole 1 | Pole 2 | Pole 3 | Zero 1 | Zero 2 | $\frac{\text { Neqlect }}{\text { Pole } 1}$ | $\frac{e d \text { Poles }}{\text { Pole } 2}$ | $\frac{\text { Omitte }}{\text { Zero } 1}$ | $\frac{\text { Zeros }}{\text { Zero } 2}$ | ```Percent Error in Identified Gain at t = 50 seconds``` |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.5 | 0.02 | 0.002 |  |  | 0.5 |  |  |  | 6.5 |
| 0.5 | 0.02 | 0.002 |  |  | 0.5 | 0.02 |  |  | 65.15 |
| 0.5 | 0.02 | 0.002 | 0.01 |  | 0.5 |  |  |  | 15.0 |
| 0.5 | 0.02 | 0.002 | 0.01 |  | 0.5 | 0.02 | 0.01 |  | 16.0 |
| 0.5 | 0.02 | 0.002 | 0.2 | 0.04 | 0.5 |  | 0.2 |  | 30.5 |
| 0.5 | 0.02 | 0.002 | 0.2 | 0.04 | 0.5 | 0.02 | 0.2 | 0.04 | 68.3 |
| 0.5 | 0.02 | 0.002 | 0.04 | 0.01 | 0.5 |  | 0.04 |  | 171 |
| 0.5 | 0.02 | 0.002 | 0.04 | 0.01 | 0.5 | 0.02 | 0.04 | 0.01 | 122 |
| 0.5 | 0.02 | 0.01 | 0.05 | 0.025 | 0.5 |  | 0.05 |  | 154 |
| 0.2 | 0.02 | 0.01 | 0.05 | 0.025 | 0.2 |  | 0.05 |  | 116 |
| 0.1 | 0.02 | 0.01 | 0.05 | 0.025 | 0.1 |  | 0.05 |  | 66.8 |
| 0.1 | 0.02 | 0.01 | 0.05 |  | 0.1 |  |  |  | 77.9 |
| 0.2 | 0.02 | 0.01 | 0.05 |  | 0.2 |  |  |  | 49.5 |
| 0.1 | 0.02 | 0.01 |  |  | 0.1 |  |  |  | 30.0 |
| 0.2 | 0.02 | 0.01 |  |  | 0.2 |  |  |  | 16.4 |
| 0.5 | 0.02 | 0.01 |  |  | 0.5 |  |  |  | 6.5 |

is much worse than when there is a gain only. From other tests which were conducted, where the neglected pole was less than or equal to the zero, it can be concluded that the closer the neglected pole approaches to the magnitude of the largest zero, the worse the approximate identification. When the pole and zero are equal, no approximation is possible.

## Effect of Errors in Identified Values of the Poles

The preceding sections have given an idea of the consequences of an incorrect identification of the order of the system, with the resultant conclusion that it is better to include a pole even though its validity may be in question. Assuming the order of the system is correctly determined, there are still many things which could affect the identification. One of these problems, which is very probable, is an error in the identification of the poles. To determine this effect, a series of tests was run in which an actual $\pm$ 5-percent error was introduced into the poles of the system

$$
\begin{equation*}
Y(s)=\frac{(s+0.04)(s+0.2) X(s)}{(s+0.1)(s+0.05)}(s+0.02) \tag{4-8}
\end{equation*}
$$

The worst effect occurs when the errors in all of the poles are in the same direction, either positive or negative. The other combinations of positive and negative errors are confined within these limits.

The results of these two limiting cases are plotted in Figure 4-6. It is interesting to note that even though



Figure 4-6. Identification of a Third Order System with a Five Percent Error in All of the Poles.
all three poles were in error by 5 percent, the effect on the gain was at most 0.59-percent, and in the largest zero, only 3.5-percent after fifty seconds. The error in the smallest zero was nearly twelve percent at the same time. However, thiserror is decreasing with time while the other errors are increasing.

## Effect of Steady State Errors

Another cause of a poor identification could be that the steady state value of both the forcing and response function has been incorrectly determined. It was expected that a steady state error in the input would have a serious effect on the quality of the identification because of the integration of the homogeneous weighting function times the input, which was required in the identification procedure. However, from the tests which were conducted, it was shown that correct identification is still possible as long as this steady state error is not too great. Even when the steady state error is as much as ten percent of the amplitude of the inputs, the identification appears to converge eventually to the correct value, as shown in Figures 4-7 and 4-8.

When there is an error in the response, its effect upon the identification is not nearly as great as an error in the forcing function. As can be seen from Figure 4-9 the identification did tend to converge to the correct values, with the errors in the identified zeros and gain being less than the error in the response by the time of fifty seconds.


4-7. Effect of Input Steady State Error of Onepercent.


4-8. Effect of Input Steady State Error of Tenpercent.


Figure 4-9. Effect of Steady State Error of Twenty Percent in the Response.

## Effect of Transport Delay

Any transport delay contained in the process or measurement and transducing equipment will normally cause a poor identification. To study the exact effect that transport delay has upon the identification, the response from a three-pole, twozero system was actually shifted to simulate transport delay of one, five, and ten seconds. The identification was then carried out with these shifted data. The results when the shift was one second are plotted in Figure 4-10. The gain and zeros tend to converge to the correct values. The lowest frequency zero is least affected by this time delay.

Figure 4-ll shows the results when the response is shifted by ten seconds. Here again the lowest frequency zero is less affected and appears to converge to the correct value. The other zero and gain, however, converge to a value which is the negative of the correct value.

## Effect of Number of Significant Figures in the Data

To explain some effects which were noted in the results of both the analog computer studies and the identification of the experimental process, a series of studies were conducted to determine the effect of data precision. This series of studies consisted of digitally computing the response of the system

$$
\begin{equation*}
T_{W}=\frac{0.000156(s+0.033) W}{(s+0.0117)(s+0.045)(s+0.0583)} \tag{4-10}
\end{equation*}
$$



4-10. Effect of One Second Transport Delay.


Figure 4-11. Effect of Ten Second Transport Delay.
for inputs of a step of amplitude, 13.0169; a ramp of slope, 0.09918; and a sine wave of amplitude, 30.6253 , and angular frequency, 0.0628.

The data from this system were recorded to seven significant figures and were used as input to the identification program. The data were then truncated to six significant figures and the identification computed. This process of deleting one significant figure and computing the identification was repeated through two significant figures. Partial results of this series of studies for step and sinusoidal forcing are shown in Figure 4-12. The identification has less than 0.1 -percent error when seven or six significant figures are used. With five figures, a deviation of 3.6-percent occurs at 5 seconds and 55 seconds. With four significant figures these deviations are greater (34-percent) with a new deviation occurring at $t=175$ seconds. With two or three significant figures these deviations are very pronounced and a new deviation occurs at $t=130$ seconds. The initial convergence is also seen to be slow.

These deviations, with the exception of the one at $t=130$, are caused by the approximate singularity of the coefficient matrix generated by $k$ solutions of Equation (2-19). In this specific case, the determinant of the coefficient matrix changes sign at $t=55,130$, and 175 seconds. At these points two ten digit numbers are being subtracted to obtain a six digit number, i.e., a loss of four significant figures.


Figure 4-12. Effect upon the Identification of the Number of Significant Figures in the Data.

At $t=130$, the response is also crossing zero which has a greater effect than the truncation of the determinant. The determinant is still an eight digit number.

## Effect of Noise

To investigate the effect of noise on the identification, and in preparation for a study of the effect of a non-linearity in the system, the analog computer was used to simulate the theoretical model of the experimental equipment described in the next chapter. For this discussion and the discussion on non-linearities the experimental equipment can be represented by the schematic diagram of Figure 4-13.


Figure 4-13. Schematic Diagram of Experimental Equipment.

The linear model was used and is described by the differential equations

$$
\begin{align*}
& \dot{T}_{f}=-0.0569 T_{f}+0.006 T_{W}+0.00293 W  \tag{a}\\
& \dot{T}_{W}=0.00532 T_{f}-0.0252 T_{W}+0.00994 T_{c O}  \tag{b}\\
& T_{C O}=0.0273 T_{W}-0.329 T_{C O}-0.208 W_{c} \tag{c}
\end{align*}
$$

The noise was assumed to be caused by either of two mechanisms. The first is the introduction of noise into the system from noise in the forcing input function, with no other noise sources present. The second mechanism of noise generation assumed no noise in the forcing input function. It was further assumed that all of the noise was either generated within the system, or introduced from an input other than the one being forced, or was generated in the transducing and measurement equipment. To study these two mechanisms, a random, square-wave, noise generator designed especially for the Process Control Laboratory was used as a noise source.

For the first case the noise was superimposed on the forcing function to provide the system input. As the noise was a part of the input, and therefore measured as a part of the total input variable for the zero identification, there was no appreciable effect upon the quality of the identification.

For the second mechanism, which is by far the most likely to occur, the result was not so simple. To study this case, the signal from the noise generator was superimposed upon the response signal before recording. The noise



4-14. Effect of Five-percent Noise in the Response Function.
signal was of constant amplitude, independent of the amplitude of the response. The amplitude of the noise signal was approximately $\pm$ five percent of the maximum amplitude of the response signal. Figure 4-14A shows the results of the gain identification for the system of Equations (4-9) when the forcing functions for the hot oil flow (W) input were a step of 100 volts and a ramp of slope l.0. The average frequency of the noise used was 0.77 zero crossing per second. Figure 4-14B shows the zero identification for this same system. As can be seen from these graphs, the identifications tend to oscillate about the correct values with an error of approximately $\pm$ five percent in the gain and $\pm$ one hundred percent in the zero. However, the average value of the gain past the time of twenty seconds is 0.998 , or only 0.2 -percent in error, and the average value for the zero is 0.039 , about eighteen percent in error. From this result it appears that to best determine the gain and zero, it is better to calculate these quantities at many points of time and then average the results, than to depend upon an identification at any one point in time.

## Effect of System Non-linearities

The system used in this study was very similar to the non-linear model of the experimental system. The system is given by

$$
\begin{align*}
& \dot{T}_{f}=0.0569 T_{f}+0.006 T_{W}+0.0293 \mathrm{~W}-0.00307 \mathrm{WT}_{f}  \tag{a}\\
& \dot{T}_{W}=0.00532 \mathrm{~T}_{f}-0.0252 T_{W}+0.00994 T_{c o} \\
& \dot{T}_{C O}=0.0273 T_{W}-0.032 T_{C O}-0.208 W_{c}-0.00844 W_{C} T_{C O}
\end{align*}
$$

The linearized version of this system is:

$$
\begin{align*}
& \dot{T}_{f}=-0.0569 T_{f}+0.006 T_{w}+0.0293 \mathrm{~W}  \tag{a}\\
& \dot{T}_{w}=0.00532 T_{f}-0.0252 T_{w}+0.00994 T_{c o}  \tag{b}\\
& \dot{T}_{w}=0.0273 T_{w}-0.032 T_{c o}-0.208 W_{c} \tag{c}
\end{align*}
$$

When the wall temperature is the response, and the two flow rates are the forcing variables, the transfer function for this system is
$T_{W}=\frac{1.56 \times 10^{-3}(s+0.0329) \mathrm{W}-2.07 \times 10^{-3}(s+0.0569) \mathrm{W}_{\mathrm{C}}}{(\mathrm{s}+0.0117)(\mathrm{s}+0.045)(\mathrm{s}+0.05833)}$

This system was programmed on the analog computer in such a manner as to allow both the linear and non-linear versions to be investigated. The analog program is shown in Figure 4-15.

The same forcing was used for both the linear and non-linear models so that these identifications could be compared. The greatest difference between the forcing of the models and the actual experimental equipment was the magnitude of the forcing function and consequently, the magnitude of the deviation from the steady state and the relative magnitudes of the non-linear term. Table 4-3 gives a quick comparison of these systems for step forcings.


Figure 4-15. Analog Computer Circuit for Simulation of the Experimental Equipment.

TABLE 4-3.
COMPARISON OF THE MAGNITUDE OF THE NON-LINEARITIES IN THE ANALOG SYSTEM AND THE EXPERIMENTAL PROCESS

| Analog System |  |  |  | Experimental Process |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Forcing | Magnitude | $\mathrm{T}_{\mathrm{w}}$ | Ratio | Magnitude | $\mathrm{T}_{\mathrm{w}}$ | Ratio |
| oil | $\begin{aligned} & +13 \mathrm{v} \\ & -13 \mathrm{v} \end{aligned}$ | $\begin{aligned} & +18.82 \mathrm{v} \\ & -21.81 \mathrm{v} \end{aligned}$ | 1.115 | $\begin{aligned} & +0.0125 \mathrm{lb} / \mathrm{sec} \\ & -0.01251 \mathrm{~b} / \mathrm{sec} \end{aligned}$ | $\begin{aligned} & 5.15^{\circ} \mathrm{F} \\ & 5.58^{\circ} \mathrm{F} \end{aligned}$ | 1.082 |
| coolant | $\begin{aligned} & +10 \mathrm{v} \\ & -10 \mathrm{v} \end{aligned}$ | $\begin{aligned} & -26.31 \mathrm{v} \\ & +49.64 \mathrm{v} \end{aligned}$ | 1.89 | $\begin{aligned} & +0.002 \mathrm{lb} / \mathrm{sec} \\ & -0.002 \mathrm{lb} / \mathrm{sec} \end{aligned}$ | $\begin{aligned} & 6.54^{\circ} \mathrm{F} \\ & 8.95^{\circ} \mathrm{F} \end{aligned}$ | 1.37 |

$u$
$\omega$

As can be seen from this table, the analog system, especially for coolant forcing, is quite a bit more non-linear than the actual experimental process.

If in Equation (4-1la) the non-linear term is lumped in with the fluid temperature, and in Equation (4-11c), it is lumped in with the coolant term, a good idea can be obtained as to the extent that the non-linear term affects the system.

$$
\begin{align*}
& T_{f}=-(0.0569+0.00307 W) T_{f}+0.006 T_{C O}+0.0293 W \\
& T_{C O}=0.0273 T_{W}-\left(0.0329+0.00844 W_{c}\right) T_{c O}-0.208 W_{c} \tag{4-14}
\end{align*}
$$

If the magnitude of the step is then substituted into these lumped terms,
$(0.0569 \pm 0.00307 \times 13) \mathrm{T}_{f}=0.0969 \mathrm{~T}_{\mathrm{f}}$ or $0.0169 \mathrm{~T}_{f}$
$(0.0329 \pm 0.00844 \times 10) \mathrm{T}_{\mathrm{CO}}=0.1173 \mathrm{~T}_{\mathrm{CO}}$ or $-0.0515 \mathrm{~T}_{\mathrm{CO}}$
It can easily be seen that the non-linear term has a drastic effect upon these coefficients. In the case of the oil, the coefficient is changed by $\pm 70$-percent and for the coolant by $\pm 250$-percent. The coolant coefficient even changes signs.

The non-linearity, however, has no effect upon the identification of the poles of the model as can be seen in Figure 4-16. This result is to be expected since the tests used in the pole identification are relaxed responses, where both of the forcing variables are equal to zero, and the nonlinear terms drop out of the describing equations. The identified poles are therefore the same for the linear and

non-linear system, even though in Equation (4-15) it is shown that the poles are drastically changed by the non-linearity. This result would lead one to the conclusion that it would be impossible to determine a gain and zero for the non-linear system.

In the case of the oil forcing, which has been shown to be the weaker of the non-linearities, a good linear approximation is obtained. This model is, under certain conditions, very similar to the model for the linear system. This approximation, however, cannot be considered a general model because it applies only under the conditions of the tests which produced it.

Figure 4-17 shows the results of an identification of the non-linear system in which the forcings were such that a very good linear model is obtained. This linear approximation varies little from the results of the identification of the linear system, which also are plotted in Figure 4-17 for comparison purposes.

As can be seen, the system gains are the same until $t=175$ seconds, at which time the gain of the non-linear system begins to decrease. This is probably due to the increasing magnitude of the non-linearity caused by the ramp forcing. The results of another identification of the non-linear system are plotted in Figure 4-18. Again for comparison purposes, the results for the linear system are also shown. Unlike Figure 4-17 there is quite a difference between the models. This case, however, was the worst one encountered in all of the test combinations.



Figure 4-17. Identification of the Oil-forced Nonlinear Model of the Experimental Equipment, Positive Ramp and Step Forcing.


Figure 4-18. Identification of the Oil-forced Nonlinear Model of the Experimental Equipment, Negative Ramp and Positive Step Forcing.

The forcing inputs used in the construction of Figure 4-18 were a positive step and a negative ramp. If a negative step and a positive ramp are used, the deviation in the identification is similar but opposite in direction. The combination of inputs used in Figure $4-17$ was a positive step and a positive ramp. If a negative step and a negative ramp are used, the identification is again similar but with opposite deviations. This trait is very well pointed out in Figure 4-19 where the forcings are a positive step and a sine wave. The identified value alternates above and below the identification results for the linear system at a frequency equal to the forcing frequency. The discontinuities which occur at $t=55$, 130 , and 180 seconds are explained in the section on the effect of the number of significant figures in the data. Figure 4-20 shows how the predicted error varies with time. For both the gain and the zero, the error minimizes within the periods when the gain and the zero are closest to the values for the linear model. The error in the gain minimizes rapidly following a discontinuity, while the error for the zero minimizes just prior to the next discontinuity.

When an identification of the coolant forced system was tried, the results were not nearly as good as for the oil forced system. This behavior was to be expected because of the larger magnitude of the non-linearity. The best identification using step and ramp forcings occurred when these two forcings were in opposite directions. The results of these


Figure 4-19. Identification of Oil-forced Non-linear Model of the Experimental Equipment, Sinusoidal Forcing.



Figure 4-20. Predicted Error for the Identification Shown in Figure 4-19.
identifications are shown in Figure 4-21, along with the identified results for the corresponding linear system. Here again the results of the identified linear model are bounded by the results of the two oppositely forced non-linear model identifications.

It is interesting to note the times at which the error analysis predicted the minimum error would occur. This time was $t=235$ seconds for the identified gain of (A). It is at this point that the identified gain of (A) is equal to the gain of the linear system. For the gain of (B) no minimum predicted error occurred. However, at $t=240$ the predicted error is still decreasing and the identified gain is still approaching the gain of the linear model. The times at which the minimum predicted errors for the identified zeros occurred are ( $A$ ) $t=185$ seconds and ( $B$ ) $t=160$ seconds.

When the forcings were in the same direction, an identification could not be made because of the drift in the results. The error propagation analysis did, however, predict that the error was a minimum at the time the identified gain of the non-linear system was the same as the gain of the linear system.

Figure 4-22 shows a representative identification of the coolant forced non-linear system when one of the forcing functions was a sine wave.


Figure 4-21. Identification of the Coolant-forced, Non-linear Model with Forcing Functions of a Step and a Ramp.



TIME, SECONDS
Figure 4-22. Identification of the Coolant-forced, Nonlinear Model with One of the Forcing Functions a Sine Wave.

Conclusions About the Identification of Non-linear Systems
Because of the results which have been discussed in this section, and the fact that the system is so strongly non-linear, the author feels that it is possible to determine a linear model of a non-linear system using this technique. This model could be either a general model or a specific model. For example, a specific model could be used when there is to be a programmed change in plant operating conditions, such as a new steady state flow. It would only be necessary to test the plant between these steady state conditions. A general model would be useful for control purposes where the flow changes are small. The best model in this case would be the linear system. This model could be obtained by averaging the results of the various combinations of opposing identifications. It could also be easily obtained when one of the forcing functions was a sinusoid because of the fact that the identification for sinusoidal forcing gives the two distinct levels which bound the linear model.

## CHAPTER V

## EXPERIMENTAL STUDIES

In the Heymann (H2) dissertation and other chapters of this work the identification technique has been thoroughly computer tested on analog and digitally simulated processes. This technique was used in the identification of a laboratory process to evaluate its performance under conditions similar to those which would be encountered in a chemical processing plant.

Although the laboratory process was a simple stirred tank heat exchanger, it served as a good test of the identification procedure because of the unknown quantities present. Some of these unknowns include system non-linearities, steady state drifts, heat losses, measurement errors, and noise from several sources: the process, instrumentation, and transmission lines. Although every possible means was employed to eliminate or minimize these conditions some of these effects were still noticed.

## Experimental Equipment

The experimental work described in the following sections was performed on equipment located in the process Control Laboratory at the University of Oklahoma. This
equipment is the evolutionary result of several previous identification studies (B1, F2, G1, S3) and has been used by two researchers ( $\mathrm{H} 1, \mathrm{~L} 2$ ) in the investigation of control systems. Because many changes have been made in the equipment since the most recent identification study, especially in the heat transfer-reactor simulator, a complete description of the equipment will be given.

The heart of the experimental system is a simulated continuous stirred tank reactor. Here, by simulated, is meant that heat transfer is the only rate operation occurring. This simulation is done as a matter of convenience in that it is easily accomplished and that no Arrenhius type non-linearity is needed in the model.

## The Reactor

The reactor (Figure 5-1) is similar to that used in the identification study by Bishop (BI) and the invariance investigation of Haskins (Hl). It is of a thick-walled design so that the heat capacitance of the wall must be considered creating a third order system. Many improvements have been made in the design of the reactor since these studies.

The new reactor wall has been precisely machined fol.lowing the rough casting. The wall construction material is Type metal and contains Lead $75 \%$, Antimony $15 \%$, Tin $8 \%$, Copper $2 \%$. Because of this careful construction the volumes and heat transfer areas are well defined. Eight Iron-Constantan thermocouples were inserted in approximately the middle of


5-1. Detail of Reactor.
the wall thickness at $45^{\circ}$ intervals around the wall and equally spaced along the height of the wall. In the experiments which were conducted, all of these thermocouples were connected in parallel to obtain an average wall temperature. The outer wall of pyrex provides an annulus for the flow of the coolant solution. Pyrex was chosen for its low thermal capacitance and because of its relatively low thermal conductivity ( $1: 40$ ) compared to the wall, as well as for the visibility provided.

The two walls are terminated on the ends by lucite plates, chosen for the same thermal properties as the pyrex outer wall. Also, the machinability of the lucite was desirable for several reasons: both the hot oil and the coolant enter through these ends, as do the thermocouples and the impeller shaft. These end plates include "distributors" for the coolant solution. These distributors are annular rings within the lucite with holes to distribute the coolant solution flowing concurrently with the oil.

The hot oil enters the bottom of the reactor and leaves through the top of the reactor. The oil in the reactor is stirred by a four-bladed, paddle-type impeller, 2.25 inches in diameter, located at the vertical midpoint of the reactor. The stirrer was driven by a $1 / 10 \mathrm{Hp} .1800$ rpm motor.

## Constant Temperature Baths

The light turbine oil is maintained at the desired temperature in two thirty gallon insulated stainless steel
tanks, each containing heating and cooling coils. The temperature in, the tanks are independently controlled. The cooling water flow is changed manually to adjust the control system characteristics. With steam used as a heating medium, the flow rate to each tank is controlled by Research Control valves type 75 G with D trim. The pneumatic control signals for the valves are generated by Minneapolis-Honeywell model 152 Pl 4 P recording controlling pyrometers. The measurement signals are obtained from copper constantan thermocouples located near the outlets of the tanks.

The two tanks are used in series because of the nature of the control loops. If used separately a drift of about $\pm 1^{\circ} \mathrm{F}$ in steady state temperature occurs. However, when in series, the temperature of the first tank is allowed to oscillate over this range at a fairly high frequency (approximately $l$ cycle per minute). The second tank is then used to smooth this variation and to correct for any long term drifts. In this manner the oil temperature to the reactor never drifted more than $\pm .25^{\circ}$ F. Both tanks were stirred with $1 / 8$ Hp Lightning model NC2 mixers. The oil was circulated by a Gould $1 / 2$ inch helical gear pump.

The coolant, a 33-percent mixture of ethylene glycol in water, was maintained at its desired temperature in a 25 gallon bath. The coolant was stirred with a Lightning type RR $1 / 4 \mathrm{HP}$ l00-1800 rpm mixer. The glycol mixture was cooled by a Copeland 1.5 ton, Freon-12 compressor with evaporator coils located in the glycol bath. A certain amount of manual
temperature control was provided by means of the expansion valve, a Hoke type 4RB281, 20 turn, l/l6 inch orifice valve with a micrometer adjust handle. This valve could be used to change the evaporator pressure, and thus, the evaporator temperature.

Automatic control of the temperature was provided by another set of heat exchange coils in the bath. Another Minneapolis Honeywell pyrometer, Model 152Pl4, was used as the recording controller on this bath, with the temperature being sensed by a copper constantan type thermocouple. The final control element was a Research Controls type, 175 D, G trim, l/4 inch air-to-close control valve which controlled hot water flow through the heat exchanger. In operation, the temperature of this bath never drifted over $\pm .25^{\circ} \mathrm{F}$. The glycol solution was circulated by a $1 / 2$ inch, gear pump.

## Flow Measurement and Control

To take advantage of the operating characteristics of the Research Controls control valves used to regulate both the oil and coolant flow rates, it was necessary to utilize a flow splitting arrangement on both of these flow systems.

The trim used in the control valves was linear, i.e., for a constant differential pressure across the valve, the flow is proportional to the stem position which in turn is proportional to the pressure on the top. To maintain the constant differential pressure across the valves the flow through the pumps had to remain constant; therefore, a bypass arrangement was used. Flow to the reactor was varied
by an air-to-open valve. The remaining portion of the flow was returned to the feed tanks through an air-to-close-valve. The air-to-open and air-to-close valves were positioned by the same pneumatic signal.

All of the valves were Research Controls, type 75B, l/4 inch valves, with 3 to 15 psi range springs, I trim for coolant, G trim for oil. The pneumatic signals for these valves were generated by Taylor Transet electro-pneumatic transducers, type 701TFlllSISl24, which had a pneumatic output range of 3 to 15 psi with a 9 psi center. The electrical signal required was $\pm 2.5$ milliamps. These transducers are specially designed for good frequency response.

The desired electrical signals were generated on the analog computer. Use of this computer made it possible to develop controllers to overcome the sluggishness and hysteresis of the control valves.

The flow rate of the oil was measured by a Waugh, type 6FLS, turbine flow meter located in the constant hot temperature stream and a Waugh, type Flll, pulse rate converter with an output of $0-250$ millivolts. This output was then amplified on the analog computer and was used as a measure of the flow rate and as an input to the valve position controller (Figure 5-2).

It was necessary to have the same sort of controller for the coolant flow system (Figure 5-3). However, because of equipment availability, the flow rate measurement was quite different. The measurement here was the pressure drop




Figure 5-4. Wall Temperature Measurement System.


Figure 5-5. Schematic Flow Sheet of Experimental System.
across a length of capillary tubing, which was used instead of an orifice because of the linear relation between flow rate and differential pressure. The capillary consisted of the annulus formed by inserting a 10 gauge, vinyl-insulated wire (O.D. $=0.165$ ) through a 10 foot length of Imperial polyflo tubing (I.D. $=0.190$ ). This capillary was located in the coolant feed stream where the temperature was constant to eliminate errors caused by viscosity or density changes. This location required the measurement of both the upstream and downstream pressures. These measurements were accomplished by use of a strain-gage type, differential pressure transducer obtained from a Beckmann, model ll2, Data Logger. The supply voltage of 7.000 volts was also obtained from this data-logger. The output of the transducers was -3 to 12 millivolts for 0 to 15 psig with this supply voltage. These millivolt signals were then amplified differentially using a Sanborn model 350-1500 preamplifier. This amplified signal representing the flow rate was transmitted to the analog computer for use in a feedback controller similar to the one used for the oil flow rate (Figure 5-3).

The controllers for both the oil and coolant flow rates consisted of a proportional controller to overcome the friction and a small amount of integral control to correct for variations from the set point (hysteresis). The time constants of these integral controllers were about one twentieth the size of the smallest time constant of the system.

Because of the high frequency noise developed in the transducers, the thermocouples and their transmission lines, it was necessary to provide filtering for these signals. These filters were programmed on the analog computer and were simply exnonential dampers. The time constants of these exponential filters were several orders of magnitude less than the time constants of the system so that they would not interfere, or even be noticed in the identification.

## Support Equipment

The basic support equipment consisted of the analog computer, its associated panel board and transmission lines, and the analog to digital conversion and recording equipment located in one room. The analog computer, panel board, transmission lines, etc. have been described in detail by Bishop and Sims (B3). The analog computer is a Donner, Model 3400, thirty-amplifier, $\pm 100 \mathrm{~V}$ computer with six electronic multiplıers. The computer has been modified (B2) to accommodate thirty additional amplifiers, four, variablebase diode function generators, two transport delay generators, and a quarter square multiplier.

The panel board, located in an adjacent room containing the process equipment, includes the various temperature recorder controllers, transducers, and process equipment power control switches. Also on the panel board are the terminations of coaxial cables which are used to transmit low level signals between the process equipment room and the analog computer room. The other ends of these
lines terminate within the analog computer programming area, making it easy to patch into the computer, signals from the remote equipment.

The analog to digital converter made it possible to process the large volumes of data that were taken. The converter consisted of a Dymec, Model 42900B, Input Scanner; a Dymec, Model 3440A, Digital Voltmeter with a Model 3443A, Range Unit; a Hewlet Packard, Model V562A, line printer; a Dymec, Model 2540, coupler; and a Friden, Model SP2, paper tape punch. With this equipment it was possible to scan the input and response signals and to record them at the rate of one point per second on the paper tape. It was possible to analyze these data directly with the Osage High Speed computer using a special read-paper-tape program.

## Theoretical Description of the Experiment Process

A mathematical model of the heat transfer process is obtained from energy balances on the oil, the coolant, and the wall. The following assumptions are used to simplify the balances.

1. The fluid within the reactor vessel is all at the same temperature, i.e., perfect mixing.
2. The oil inlet temperature is constant.
3. The coolant inlet temperature is constant.
4. Densities, heat capacities, volumes, heat transfer coefficients and areas are constant.
5. The heat transfer to the coolant occurs at the mean temperature of the coolant; $T_{c m}=\left(T_{c o}+T_{c i}\right) / 2$.

With these assumptions the following equations are obtained from the energy balances:

$$
\begin{align*}
& \left(\rho V C_{p}\right) \dot{f}^{T_{f}^{*}}=h_{i} A_{i} T_{W}^{*}-h_{i} A_{i} T_{f}^{*}+C_{p f} T_{i n} W^{*} \\
& -C_{p f} W^{*} T_{f}^{*}  \tag{a}\\
& \left(o V C_{p}\right) \dot{W}_{W}^{*}=h_{i} A_{i} T_{f}^{*}+\frac{h_{O} A_{O}}{2} T_{C O}^{*} \\
& -\left(h_{i} A_{i}+h_{o} A_{o}\right) T_{W}^{*}  \tag{5-1}\\
& \left(\rho V C_{p}\right) C^{T^{*}}{ }^{*}=2 h_{o} A_{o} T_{w}^{*}-h_{o} A_{o} T_{w}^{*}+2 C_{p c} T_{c i} W_{c}^{*} \\
& -2 C_{p c}{ }^{W} c^{T} T_{c o}{ }^{*}-h_{o} A_{o}{ }^{T} C_{i} \tag{c}
\end{align*}
$$

$$
\begin{aligned}
& \text { where } \\
& \rho=\text { density } \\
& \mathrm{V}=\text { volume } \quad \mathrm{w}=\text { wall } \\
& C_{p}=\text { specific heat } \quad c=\text { coolant } \\
& T^{*}=\text { temperature } \\
& i=\text { inside } \\
& h \quad=\text { heat transfer coefficient } \\
& 0 \text { = outside } \\
& \text { A = heat transfer area } \\
& \text { in }=\text { oil in } \\
& \mathrm{W}=\text { flow rate } \quad \mathrm{ci}=\text { coolant in } \\
& Q \text { = heat loss } \quad \text { ss = steady state } \\
& \text { superscript * }=\text { total variable }
\end{aligned}
$$

Equations (5-1) are the non-linear, total variable model of the system. For convenience these equations can be changed to the perturbation model by assuming that each of the variables is composed of a transient and steady-state portion. Thus

$$
\begin{array}{lll}
T_{f}^{*}=T_{f}+T_{f S S}^{*} & \text { (a) } & W^{*}=W+W_{s s}^{*} \\
T_{W}^{*}=T_{W}+T_{W S S}^{*} & \text { (b) } & W_{C}^{*}=W_{C}+W_{c s s}^{*}  \tag{5-2}\\
T_{C O}^{*}=T_{C O}+T_{c o s s}^{*} & \text { (c) } &
\end{array}
$$

When Equations (5-2) are substituted into Equations (5-1) and the steady state equations corresponding to Equations (5-1) (the time derivatives equal to zero) are subtracted from these substituted equations, the pertubation model results:

$$
\begin{align*}
\left(\rho V C_{p}\right)_{f}=- & \left(h_{i} A_{i}+C_{p f} W_{S S}^{*}\right) T_{f}+h_{i} A_{i} T_{W}+\left(C_{p_{f}} T_{i n}-C_{p f} T_{f S S}^{*}\right) \\
& W-C_{p f} W T_{f}  \tag{a}\\
\left(\rho V C_{p}\right)_{w} \dot{T}_{w}= & h_{i} A_{i} T_{f}-\left(h_{i} A_{i}+h_{o} A_{o}\right) T_{W}+\frac{h_{o} A_{O}}{2} T_{c o} \\
\left(\rho V C_{p}\right)_{c} \dot{T}_{w}= & 2 h_{o} A_{o} T_{W}-\left(h_{o} A_{o}+2 C_{p c} W_{c S S}^{*}\right) T_{c o}-2 C_{p c}\left(T_{c o s s}^{*}-T_{c i}\right) \\
& -2 C_{p c} W_{c} T_{w} .
\end{align*}
$$

$$
\text { (b) } \quad(5-3)
$$

These pertubation equations contain product-type nonlinearities involving the flow rates and their respective temperatures. A linear model is obtained by neglecting the product of these terms, which was shown by Stewart (S3) as being equivalent to using a Taylor Series expansion and retaining the constant and first order terms.

When the values from Table 5-1 are substituted into Equation (5-3) the resultant non-linear model for these specific operating conditions are

TABLE 5-1
LIST OF SYSTEM CONSTANTS

| Symbol | Nomenclature | Value | Units | Source |
| :---: | :---: | :---: | :---: | :---: |
| $c_{p c}$ | coolant heat capacity | 0.823 | $B T U / 1 b^{\circ} \mathrm{F}$, | 1 |
| $C_{\text {pw }}$ | wall heat capacity | 0.037 | BTU/ $1 b^{\circ} \mathrm{E}$ | 1 |
| $c_{p f}$ | oil heat capacity | 0.405 | $\mathrm{BTU} / 1 \mathrm{~b}^{\circ} \mathrm{F}$ | 2 |
| $\mathrm{h}_{\mathrm{i}}$ | oil side heat transfer coefficient | 0.0078 | $\mathrm{BTU} / \mathrm{sec}^{\circ} \mathrm{F}$ | 3 |
| $h_{0}$ | coolant side heat transfer coefficient | 0.0166 | $\mathrm{BTU} / \mathrm{sec}^{\circ} \mathrm{F}$ | 3 |
| $A_{i}$ | inside heat transfer area | 0.322 | $f t^{2}$ | 2 |
| $\mathrm{A}_{0}$ | outside heat transfer area | 0.444 | $f t^{2}$ | 2 |
| $\mathrm{T}_{\text {ci }}$ | coolant inlet temperature | 46.8 | ${ }^{\circ} \mathrm{F}$ | 3 |
| $\mathrm{T}_{\text {coss }}$ | coolant steady state outlet temperature | 71.5 | ${ }^{\circ} \mathrm{F}$ | 3 |
| $\mathrm{T}_{\text {fss }}$ | oil steady state temperature | 146.8 | ${ }^{\circ} \mathrm{F}$ | 3 |
| $\mathrm{T}_{\text {in }}$ | oil inlet temperature | 156.2 | ${ }^{\circ} \mathrm{F}$ | 3 |
| $\mathrm{T}_{\text {wss }}$ | steady state wall temperature | 73.5 | ${ }^{\circ} \mathrm{F}$ | 3 |
| $\mathrm{V}_{\mathrm{c}}$ | coolant volume | 0.0102 |  | 2 |
| $\mathrm{V}_{f}$ | oil volume | 0.0170 |  | 2 |
| $\mathrm{V}_{\mathrm{w}}$ | wall volume | 0.0153 | $f t^{3}$ | 2 |
| $\mathrm{W}_{\text {ss }}$ | steady state oil flow | 0.0467 | 1b/sec | 3 |
| $\mathrm{W}_{\text {css }}$ | steady state coolant flow | 0.00667 | lb/sec | 3 |
| $\rho_{c}$ | coolant density | 65.41 | $\underline{b / f t}{ }^{3}$ | 2 |
| $\rho_{f}$ | oil density | 53.0 1 | $b / f t^{3}$ | 2 |
| $\rho_{\text {w }}$ | wall density | 640.5 | $1 b / f t^{3}$ | 2 |
| $Q_{L}$ | heat loss | 0.035 | BTU/sec | 3 |

Sources: 1. Handbook; 2. Laboratory measurements;
3. Steady state data.
$\dot{T}_{C}=-0.0583 \mathrm{~T}_{\mathrm{f}}+0.00688 \mathrm{~T}_{\mathrm{W}}+10.54 \mathrm{~W}-1.11 \mathrm{WT}_{\mathrm{f}}$
$\dot{T}_{W}=0.00693 T_{c}-0.0273 T_{W}+0.0102 T_{C O}$
$\dot{\mathrm{T}}_{\mathrm{CO}}=0.0268 \mathrm{~T}_{\mathrm{W}}-0.0334 \mathrm{~T}_{\mathrm{CO}}-74.05 \mathrm{~W}_{\mathrm{C}}$
$-3.00 \mathrm{~W}_{\mathrm{c}} \mathrm{T}_{\mathrm{co}}$.
The linearized process model is then

$$
\begin{align*}
& \dot{T}_{f}=-0.0583 T_{f}+0.00688 T_{W}+10.34 \mathrm{~W}  \tag{a}\\
& \dot{T}_{W}=0.00693 T_{f}-0.0273 T_{W}+0.0102 T_{c o}  \tag{b}\\
& \dot{T}_{c O}=0.0268 T_{w}-0.0334 T_{c o}-74.05 W_{c} . . \tag{c}
\end{align*}
$$

When the wall temperature is the response variable, the transfer function corresponding to the linear model is:

$$
\begin{equation*}
T_{W}=\frac{0.073(s+0.0334) W-0.752(s+0.0587) W_{C}}{(s+0.0129)(s+0.0458)(s+0.0607)} \tag{5-6}
\end{equation*}
$$

This equation represents the transfer function for which the time domain identification technique will be testing.

## Experimental Determination of the System Characteristics by Standard Methods

To evaluate the performance of the identification technique, which is the basis of this work, two other identification procedures were used also: frequency response testing and pulse testing. Transient response tests were also conducted for an independent determination of the system gains. The individual techniques and their ramifications and results will be discussed first, and a general comparison will be made along with pertinent conclusions.

## Transient Response Tests

This technique is called transient response testing because the experimental system is forced away from the steady state and is allowed to come to a new steady state. However, no data were taken on the actual transient portion of the tests, which served only as a convenient means for determining DC gains of the system for comparison purposes. The results of these tests are listed in Table 5-2.

## TABLE 5-2

DC GAIN OF EXPERIMENTAL EQUIPMENT AS DETERMINED BY TRANSIENT TESTS

| Forcing | Direction | DC Gain <br> $\circ \mathrm{F} / \mathrm{lb} / \mathrm{min}$. |
| :--- | :---: | :---: |
| coolant | + | 19.7 |
| coolant | - | 42.9 |
| oil | + | 1.35 |
| oil | - | 2.97 |

The data of these tests were recorded in two ways: the voltage output of the transducers and amplifier systems versus potentiometric readings of the temperatures and versus scale and stop watch measurements on the flow systems. These recordings provided calibration of the measurement systems.

## Frequency Response Tests

Frequency response tests were conducted following the standard procedure of forcing the system with sinusoidal variations of many different frequencies. The frequency of the forcings ranged from 0.0118 cycles/minute to 17.65 cycles/ minute. The electrical signals for these forcing functions were generated on the analog computer by solving the equations

$$
\begin{align*}
& x=-\omega \sin \omega t \\
& y=w \cos \omega t  \tag{5-7}\\
& y(0)=A(\text { volts })
\end{align*}
$$

where $\ddot{x}$ was the desired sine function and $A$ was the desired amplitude.

These electrical signals from the analog computer were then converted to pneumatic signals by the Taylor transducers, and the flow rate was varied by the control valves in response to the pressure signal. The signals from the flow rate transducers and amplifiers were then recorded on the Sanborn 6-channel recorder. The amplitude and phase lag of the resonse function were then determined from these response records as a function of the forcing frequency The results of the tests are given in the form of a standard Bode Plot in Figure 5-6 and Table 5-3.

The difficulties with these tests are many and varied, three of which will be listed here.

1. The length of time required to conduct the test is very great. The lowest frequency used in the tests required


5-6. Frequency Response Tests of Experimental System.
about ninety minutes for just one cycle, and extrapolation still was necessary to determine the DC gain.
2. The disturbances to the system were great. The variations do tend to average out, however, due to the plus-minus nature of the forcing. However, there is still a transient response associated with this forcing.
3. Results obtained in the form of the Bode Plot are limited in usefulness for further theoretical considerations.

TABLE 5-3

DC GAIN OF EXPERIMENTAL EQUIPMENT AS DETERMINED BY FREQUENCY RESPONSE TESTS

| Forcing | DC gain <br> ${ }^{\circ} \mathrm{F} / \mathrm{lb} / \mathrm{min}$. |
| :--- | :---: |
| coolant | 33.2 |
| oil | 2.36 |

Pulse Tests

Due to the product non-linearities associated with flow forcing of the reactor, pulse testing provides a means of estimating the extent to which these non-linearities affect the system description. This estimate is given as a result of forcing the flow in both directions from the steady state. The pulses used in the tests were sawtooth-shaped with an approximate duration of seventy five seconds. This magnitude represents a long pulse, far from an impulse. Because of the
physical limitations imposed by the flow systems, the magnitude of the forcing pulse could not be made great enough to provide sufficient energy to the system in a shorter period of time. However, this long pulse did not disturb the system greatly. The wall temperatures changed about six degrees, or about eight percent of the steady state temperature.

As can be seen from Figure 5-7, which is a standard Bode Plot, i.e., magnitude ratio and phase lag versus frequency, the non-linearity did not have a great effect on the system dynamics. However, there was a significant effect upon the DC gains of the system (Table 5-4) when both coolant and oil were forced. In both cases there was a factor of approximately 2 between the respective DC gains.

TABLE 5-4
DC GAIN OF EXPERIMENTAL EQUIPMENT AS DETERMINED BY PULSE TESTING

| Forcing | Direction | DC Gain <br> ${ }^{\circ} \mathrm{F} / \mathrm{lb} / \mathrm{min}$. |
| :--- | :---: | :---: |
| coolant | + | 22.8 |
| coolant | - | 39.8 |
| oil | + | 1.66 |
| oil | - | 2.98 |

The pulse testing technique seems to be very sensitive to pulse duration, pulse shape, noise, and amount of energy developed in the system. In the pulse tests conducted


5-7. Results of Pulse Testing of Experimental System.
on this system, it was impossible to obtain good information over more than one decade of magnitude in the Bode plot. This limitation is not intended to imply that pulse testing technique is not a valuable tool in general. Because the application of this technique is not the prime purpose of this work, essential refinements in the technique, such as smoothing the data and using a wide variety of forcing pulses, were not employed.

The computations involved in pulse testing are much more difficult than those of frequency response testing, but once a computer program is developed, the results are quickly obtained. Theoretically the identification can be made with only one short test. The computational procedures used in this work were presented by Dreifke and Hougen (D2) and were programmed by the author for the Osage computer.

## Time-Domain Identification

The determination of a model for the experimental process was conducted in two stages. First the pole identification was carried out and then the major concern of this present work, the zero identification, was executed.

Pole Identification
The theoretical model of the process describes it as having three poles. In order to identify the system as third order, it is necessary to have three linearly independent responses of the unforced process (H2, p. 43). In an effort to obtain at least three independent responses, six
tests were conducted using various combinations of the two flow streams to force the process from steady state. These combinations are listed in Table 5-5.

TABLE 5-5
FORCINGS USED FOR POLE IDENTIFICATION

| Test <br> Number | Forcing (Flow Rate) |  |
| :--- | :--- | :--- |
|  | Oil | Coolant |
| 2 | decreased |  |
| 3 |  | increased |
| 4 | increased | decreased |
| 5 | decreased | increased |
| 6 |  | decreased |

With each of these forcings, the system was driven a small amount from steady state (a wall temperature change of about six degrees F$)$. The forcing was then removed and the wall temperature recorded as it again approached steady state.

All twenty possible combinations of these six tests (taken three at a time) were used in the digital computations for the poles. Many of these combinations yielded no significant results. The precise reasons are now known, but there are several possibilities.

1. The tests were not linearly independent: $A$ probable cause, but not conclusive, since the tests appear
to have different initial values, initial slopes, and initial rates of change of slope which would indicate they were linearly independent.
2. The noise level was too great: An improbable cause as an inspection of the recorder traces shows them to have a very low noise level (especially in the early portions of the tests where the energy levels were high) of no more than one or two percent. Heymann (H2) was able to obtain identifications with as much as eight percent noise.
3. The steady state conditions changed while the tests were being conducted: A most probable cause of trouble as there were basically four factors controlling the steady state. They include the oil flow rate, the coolant flow rate, the oil temperature and the coolant temperature. Every possible effort was made to control these factors as closely as possible, but there were limitations on the equipment available. In controlling the oil and coolant temperatures a deadband existed in the controllers, making it impossible to hold these temperatures any closer to the desired values than $=0.25$ degrees. Because of the measurement technique required by available equipment for the coolant flow, the pressure drop and the recorded voltage could be maintained constant, but the flow rate changed because of air bubbles adhering to the walls of the capillary.

All of these factors add up to what the author believes to be a rather unreliable steady state, and therefore
the most probable cause of trouble. Heymann (H2, p.199) states, "Low frequency disturbances such as changes or drifts of steady state operation level cannot be tolerated for this identification." (He makes no other reference to this problem except for one other short qualitative statement as to the effect of steady state changes.) There exists no quantitative measure of this effect, so that it is impossible to tell the extent of the corruption of the identifications by the oil and coolant temperature changes, etc. It was possible, however, to obtain from many of the tests good identification results. The results of a mean and standard deviation analysis of the identifications are listed in Table 5-6.

TABLE 5-6
RESULTS OF POLE IDENTIFICATION

| Pole | Mean Value | Standard <br> Deviation |
| :---: | :---: | :---: |
| 1 | -0.0498 | 0.0098 |
| 2 | -0.0182 | 0.0055 |
| 3 | -0.00987 | 0.0024 |

This averaging was necessitated by the general instability of the identifications, and it is believed that they provide a reliable identification of the process dynamics in this range of operation.

The instability of the identifications can be seen in the plots of Figures (5-8) and (5-9). In both cases the high frequency pole is the most unstable, but this result does not seem unreasonable when compared with Figure 34 of Heymann (H2, p. 184) which shows that even under ideal computer conditions this pole is not identified as well as the others.

During the early parts of the identifications, the two lowest frequency poles are shown as being equal because they were identified as being complex. The complex part was very unstable, and consequently it was assumed that these poles were real and close to being equal. Heymann (H2, p. 167) found that when two poles were very nearly equal the identifications tended to oscillate between real and complex poles with the complex part being unstable.

In Figure 5-9 there is an abrupt change in the identification at delta $=61$, which corresponds to time $=$ 183 seconds. This identification is with tests 2,4 , and 6. A close look at the trace of response 4 showed a sudden shift in the response at about 180 seconds, which probably causes this instability in the identification.

To summarize the results of the pole identification, this author believes that a reliable identification has been made through the use of the mean of all twenty combinations of tests even though some instability exists within each test.


Figure 5-8. Identification of the Poles of the Rxperimental System, Tests $1,3,5$.


Figure 5-9. Identification of the Poles of the Experimental System, Tests 2, 4, 6.

## Zero Identification

In all of the computer studies the zero identification technique worked very well and an identification was obtained, even when there was a quite large error in some of the factors. However, in all of these tests there was only one error, which had a known magnitude. In order to test the technique under more realistic circumstances, the experimental process was used in place of the computer. In the process there were a multiplicity of errors, the magnitude of which were unknown. These include all of the errors assumed in the computer studies: misidentification of poles, incorrect steady state, noise, transportation delay, non-linearities, etc., and probably some items which were not considered in the computer studies.

Here again, as with the pole tests, in order to assure a good selection of independent tests, seven tests were conducted with both the oil flow rate forcing and coolant flow rate forcing. Table 5-7 is a summary of these forcings.

Unlike the pole tests where the required information is the relaxed response for the zero test, the forced response is the desired information along with the forcing. The system was therefore forced from steady state by each of the forcings of Table 5-7, while the response and input were alternately recorded versus time as output of the analog to digital conversion equipment.

In the analog computer tests of the non-linear system the coolant forced process was much more non-linear than the

TABLE 5-7

## FORCINGS USED FOR ZERO IDENTIFICATION

| Test | 0 Oil flow $\mathrm{lb} / \mathrm{sec}$ | Coolant flow $\mathrm{lb} / \mathrm{sec}$ |
| :--- | :--- | :--- |
| 1 | -0.0125 | +0.001 |
| 2 | +0.0125 | -0.001 |
| 3 | $-0.00025 t$ | -0.00001 t |
| 4 | $+0.00025 t$ | $+0.00001 t$ |
| 5 | $-0.025 \sin (0.091 t)$ | $-0.002 \sin (0.091 t)$ |
| 6 | $-0.025 \sin (0.050 t)$ | $-0.002 \sin (0.050 t)$ |
| 7 | $-0.025 \sin (0.010 t)$ | $-0.002 \sin (0.015 t)$ |

oil forced system. In the actual experimental process the opposite is true. If one assumes that the theoretical model of the experimental system given by Equations (5-4) is correct, then an idea of the degree of non-linearity can be garnered by lumping the non-linear terms with the respective temperatures as

$$
\begin{align*}
& \dot{\mathrm{T}}_{f}=-(0.0583+1.11 \mathrm{~W}) \mathrm{T}_{\mathrm{f}}+0.00688 \mathrm{~T}_{\mathrm{w}}+10.54 \mathrm{~W} \quad(\mathrm{a}) \\
& \dot{\mathrm{T}}_{\mathrm{CO}}=0.0268 \mathrm{~T}_{\mathrm{W}}-\left(0.0334+3.00 \mathrm{~W}_{\mathrm{C}}\right) \mathrm{T}_{\mathrm{CO}}-74.05 \mathrm{~W}_{\mathrm{c}} \tag{5-8}
\end{align*}
$$

(b)

Substituting the values of the step forcing into these lumped terms,

$$
\begin{aligned}
& 0.0583+1.11 W^{2}=0.0583 \pm 0.0139 \\
& 0.0334+3.00 W_{C}=0.0334=0.003
\end{aligned}
$$

(a)
(b) ${ }^{(5-9)}$

It can be seen that this term for oil forcing will vary $\pm 24$ percent, and the coolant forced term will vary $\pm 9$-percent, thus indicating that the oil is almost 3 times more non-linear than the coolant side. On the basis of this information, the coolant side identification should be better than the oil side. The truth of this prediction will be seen as the results are discussed.

Nineteen combinations of these tests were used in the identification of each transfer function. Table 5-8 summarizes these combinations.

Table 5-9 gives a summary of the mean values of the identified gains and zeros for the nineteen combinations of tests with both the oil and coolant forced systems. These values are all within about $\pm 50$-percent of the overall mean value and approximately half of them are within $\pm$ lO-percent of the overall mean value.

A close inspection of some of the extreme cases will show why their individual mean values should be eliminated from the overall, or at least modified. One good example is in the case of sets eleven and fourteen, where the forcings were ramps and a very low frequency sine wave. As was determined in the computer studies of linear systems, this condition did not produce good results (page 25). The early part of the sine forcing appears very similar to a ramp, and therefore the linear independence of these tests is tenuous. In practice these identifications were poor with a large drift and slow convergence. In the case of set fourteen of
the oil system, no identification at all was made for the gain of the system. These tests should therefore be eliminated from the overall identification.

Sets eighteen and nineteen of the coolant tests, which have the extreme values of the gain are valid in that taken as a pair, they represent the worse cases of step and ramp forcing, but their identified results when averaged give the linearlized model.

Sets eleven and fourteen are the only ones for which it seems reasonable to eliminate the mean values. The corrected overall mean values are:

$$
\begin{array}{ll}
\text { oil gain } & =5.47 \times 10^{-3} \\
\text { oil zero } & =3.75 \times 10^{-5} \\
\text { coolant gain } & =4.08 \times 10^{-3} \\
\text { coolant zero } & =5.88 \times 10^{-5}
\end{array}
$$

Figure 5-10 shows the results of a pair of identifications of the coolant system. For a negative step and negative ramp the gain and zero curves have positive slopes, while for positive forcings the slopes are negative. It is possible to see that these curves bound the overall mean value with the gain curves crossing at almost the exact value of the mean. This result would indicate that the mean value is a very good linear identification for this pair of tests.

In the computer analysis of a non-linear system, when step and sine forcings were used, the identification divided into distinct groups bounding the linear system. The same is

## TABLE 5-8

TEST COMBINATIONS USED FOR ZERO IDENTIFICATION


TABLE 5-9
RESULTS OF ZERO IDENTIFICATION OF EXPERIMENT EQUIPMENT

|  | Oil Forced System |  | Coolant Forced System |  |
| :---: | :---: | :---: | :---: | :---: |
|  | Gain $\times 10^{3}$ | Zero $\times 10^{5}$ | Gain $\times 10^{3}$ | Zero $\times 10^{5}$ |
| 1 | 5.55 | 3.66 | 4.64 | 3.86 |
| 2 | 5.62 | 3.80 | 3.41 | 6.69 |
| 3 | 5.53 | 3.20 | 3.60 | 4.25 |
| 4 | 4.54 | 2.61 | 5.79 | 7.10 |
| 5 | 7.05 | 3.77 | 4.14 | 7.19 |
| 6 | 5.84 | 4.05 | 3.97 | 5.07 |
| 7 | 5.24 | 3.32 | 4.02 | 6.45 |
| 8 | 3.41 | 3.98 | 5.34 | 6.17 |
| 9 | 5.93 | 3.88 | 3.00 | 6.98 |
| 10 | 5.71 | 3.75 | 3.92 | 6.08 |
| 11 | 4.92 | 2.44 | 3.67 | 9.67 |
| 12 | 6.14 | 4.67 | 2.52 | 5.21 |
| 13 | 6.16 | 4.46 | 3.81 | 4.96 |
| 14 | * | 5.25 | 4.68 | 9.60 |
| 15 | 5.29 | 5.17 | 3.98 | 6.97 |
| 16 | 5.39 | 4.02 | 4.21 | 7.01 |
| 17 | 5.29 | 3.27 | 4.30 | 7.21 |
| 18 | 6.73 | 4.34 | 6.33 | 1.91 |
| 19 | 4.54 | 3.98 | 2.48 | 7.95 |
| Mean | 5.45 | 3.78 | 4.09 | 6.28 |

[^0]

Figure 5-10. Identification of Coolant-forced, Experimental System with Step and Ramp Forcings.
true with the results from the experimental process, which can be seen in Figure 5-11. The overall mean values fit pretty well within the identification results, indicating that these mean values are also valid for these tests.

In the identification for the oil system it appeared that the ramp was the predominant test when the forcings were a ramp and a step. Therefore, Figure 5-12 shows the results of the pair of ramp forcings with a positive step. Although these tests do not show the clear cut opposition in the gain, as was noticed with the coolant system, the overall mean value does appear to be a good linear identification of the tests.

In Figure 5-13 the identification for positive step and sine forcings is plotted. Here again, as in Figure 5-12, there is not the distinct separation in the identification of the gain as was noticed with the coolant system. However, again the overall mean value does appear to be a good linear identification, in that the mean value of the 45 points which are plotted for the gain is 0.000524 , plus the nine points which are off the graph, is 0.00101 .

With this verification that these overall mean values comprise a good identification, they will be used to determine the actual model in terms of the wall temperature and flow rates. The actual zero is merely the identified zero divided by the gain. The values are given in Table 5-10. To determine the gain in terms of ${ }^{\circ} \mathrm{F} / \mathrm{lb} / \mathrm{sec}$, it is necessary to multiply the identified gain, which is in terms of volts


Figure 5-11. Identification of Coolant-forced, Experimental System with Step and Sinusoidal Forcings.


Figure 5-12. Identification of Oil-forced, Experimental System with Step and Ramp Forcing.


Figure 5-13. Identification of Oil-forced, Experimental System with Step and Sinusoidal Forcings.
temperature $\sec ^{2} /$ volts flow by a factor of ${ }^{\circ} \mathrm{F} /$ volt temp x volt ${ }_{\text {flow }} / \mathrm{lb} / \mathrm{sec}$.

For the oil system this factor is 333 , and for the coolant system it is 2083, as determined by the measurement and amplification equipment. The results are given in Table 5-10.

TABLE 5-10
FINAL MEAN RESULTS OF ZERO IDENTIFICATION OF EXPERIMENTAL EQUIPMENT

| System | Gain | Zero |
| :--- | :--- | :--- |
| oil | 0.181 | -0.00675 |
| coolant | 0.852 | -0.0154 |

The linearized transfer function as obtained by the pole and zero identification is then

$$
\begin{equation*}
T_{W}=\frac{0.181(s+0.00675) W-0.852(s+0.0154) W_{C}}{(s+0.498)(s+0.0183)(s+0.00987)} \tag{5-10}
\end{equation*}
$$

## Comparison of Results

The Bode diagram is used to present the results of the three iaentification methods for comparison. This presentation was chosen because it is the final form of two of the identification techniques, and the time-domain results are easily converted to this form. The reverse is not nearly as easy or as accurate.

Figures 5-14 and 5-15 give the Bode diagram of the coolant forced system. These graphs show that the results of the time domain identification compare very favorably with the results of the other testing methods and the theoretical model. The Bode diagrams of the oil forced system (Figures 5-16 and 5-17) do not show this good comparison. Examination of the curve for the time domain identification reveals a peak in the magnitude ratio at approximately 0.1 cycles per minute. This peak is a result of the low natural frequency of the identified zero.

To obtain some idea of the quality of this identified system, analog computer tests were run to compare the identified system to the theoretical system and to actual data from the plant. The procedure used in these tests was to match first, as closely as possible, the actual plant data with the theoretical non-linear model. The response of the theoretical linear model was then obtained by removing the product term. The last step was to determine the response of the identified model. To compare the dynamics of the models closely, the amplitude of the response of the identified model was adjusted to equal the amplitude of the theoretical linear model over the first quarter cycle.

As can be seen in Figure 5-18, the identified model matches the actual data much better than does the theoretical linear model--almost as well as the non-linear model. The dissimilarity between the model in question and the models obtained by other means can therefore be attributed to an





5-17. Phase Diagram of Results of Identification of the Oil-forced system.

attempt by the time domain identification technique to approximate the dynamics of the non-linear process. For this reason it is assumed that the identified model is as good a linear approximation as can be determined.

The DC gains of the oil forced system as determined by all of the experimental methods are very similar if an average is used for the pulse and transient tests. The average of the $D C$ gains of the frequency response, pulse, and transient tests is 2.28. The DC gain, as determined by the time domain identification, is 2.27, which appears to be a very reasonable value.

Although the DC gains of the coolant forced systems compare less favorably with the average of the results from the other experimental methods than the DC gain of the oil forced system does, the identified gain lies within the range of values experimentally determined for positive and negative forcings. The identified gain of 24.33 is therefore believed to be a valid result.

> TABLE 5-11
> SUMMARY OF DC GAINS ( $\left.{ }^{\circ} \mathrm{F} / \mathrm{lb} / \mathrm{min}.\right)$

| Testing Method | Oil | Coolant |
| :--- | :---: | :---: |
| Theoretical | 1.18 | 20.51 |
| Time domain | 2.27 | 24.33 |
| Frequency response | 2.36 | 33.21 |
| Pulse positive | 1.66 | 22.82 |
| Pulse negative | 2.98 | 39.81 |
| Transient positive | 1.35 | 19.70 |
| Transient negative | 2.97 | 42.90 |

CHAPTER VI

## CONCLUSIONS AND RECOMENDATIONS

## Conclusions

This technique, when used in conjunction with Heymann's pole identification technique, appears to be a useful method for system identification. The model is obtained in its most useful form, the poles and the zeros, which one can readily transform into a differential equation. The error propagation analysis is a good measure of the reliability of the identified zeros, but it should not be used as the ultimate criterion for determining the best value.

For most real processes, the best values for the zeros will be the arithmetic means of the identified values at the points where the predicted error is relatively small. The use of the mean minimizes the effect of noise.

In most cases errors in the data are not magnified by the computation. Most errors do, however, delay the convergence of the results to the correct value. The error propagation analysis does not indicate this slow convergence; therefore caution should be exercised in the interpretation of the results when the time is less than one-half the longest time constant.

In the computer tests of the technique that were conducted on the Osage High Speed Computer the final form of the computation was the factored-zeros polynomial. For real systems it was found that this polynomial should not be factored until the results are interpreted and the mean values determined. The factoring of this polynomial greatly magnifies the errors in the zeros.

No great limitations of the time domain technique were detected in the investigation. In applying the technique, however, great care should be taken in the preparation of the experimental tests, especially in regards to the steady state and transport delay.

## Recommendations for Future Work

The present investigation has focused primarily upon the development and verification of the identification technique. Further work should be done in identification of actual processes. There are several points that should be investigated in this connection. These apply to both the pole and the zero identification.
(1) The linearization of nonlinear systems should be analyzed in greater detail.
(2) A study of time-varying systems should be conconducted.
(3) The use of data-smoothing techniques to obtain more stable identifications should be investigated.

118
To facilitate the interpretation of the results, the zero identification program should be modified to determine the mean values of the identified results at the times when the predicted error is small. These mean values would then constitute a set of smoothed coefficients of the polynomial, the solution of which determines the system zeros.

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APPENDICES

## APPENDIX A

NOMENCLATURE*

```
a}\mp@subsup{i}{i}{=}\mp@subsup{i}{}{\mathrm{ th }}\mathrm{ coefficient of homogenous weighting function (2-16)
    b
        dependent variable in an n}\mp@subsup{}{}{\mathrm{ th }}\mathrm{ order scalar
        differential system (2-1)
da}\mp@subsup{i}{i}{}=\mathrm{ sum of the (n-1) terms of the form l/( }\mp@subsup{\rho}{1}{}-\mp@subsup{\rho}{i}{})(3-9
f(t) = general time function (general forcing
        function of a process) (2-3)
g = k element zero vector (3-22)
G = k x k matrix of G G (i) (3-22)
G
    (k-n)}\mp@subsup{}{}{\mathrm{ th}}\mathrm{ derivative of the weighting function (3-2)
gjk
    independent variable of a scalar differential
    system (2-1)
H(s) = homogeneous transfer function (2-15)
H}\mp@subsup{n}{n}{}(i)=\mathrm{ integral of the ith term of the (k-n)
    derivative of the homogeneous weighting function
    times the input times the variable of integration (3-11)
    *The numbers at the end of the definitions are the
equation numbers in which the symbols are either defined or
first used.
```



## APPENDIX B

## NUMERICAL EXAMPLE OF ZERO CALCULATION

The calculation is performed in five basic steps. These steps are:

1. Calculation of the coefficients of the weighting function and $n-1$ of its derivatives for the equations shown below.
2. Calculation of the integrals.
3. Multiplication of the integrals of step 2 by their coefficients and summation on i, and calculation of the expected error in these results.
4. Inversion of the coefficient matrix and calculation of the expected error in the inverse.
5. Multiplication of the inverse matrix times the response vector to determine the systems zeros and calculation of the expected error in the zeros.

The system chosen for this example is the same as was used in determining the effect of the number of significant figures in the data. This system is

$$
\begin{equation*}
Y(s)=\frac{0.00156 s+0.000051324}{(s+0.0117)(s+0.045)(s+0.05833)} X(s) \tag{A1-1}
\end{equation*}
$$

The coefficients of the homogeneous weighting function are:

$$
\begin{align*}
& a_{1}=644.006 \\
& a_{2}=-2252.815  \tag{Al-2}\\
& a_{3}=1608.807 .
\end{align*}
$$

The system contains one zero, therefore, $k=1$ and the coefficients of the first derivative are needed. These coefficients are:

$$
\begin{align*}
& a_{1} \rho_{1}=-7.535 \\
& a_{2} \rho_{2}=101.327  \tag{A1-3}\\
& a_{3} \rho_{3}=-93.842 .
\end{align*}
$$

The inputs to be considered are:

$$
\begin{align*}
& x_{1}(t)=13.017 \\
& x_{2}(t)=0.09914 t . \tag{A1-4}
\end{align*}
$$

If the time is assumed to be fifteen seconds, the responses of the system (Al-1) to the inputs of (Al-4) are:

$$
\begin{align*}
& Y_{1}(15)=1.54386 \\
& Y_{2}(15)=0.064790 . \tag{A1-5}
\end{align*}
$$

The second step of the identification is the calculation of the integrals of Equation (2-19). This calculation is done in two parts. First, calculation of the function $e^{\rho_{i}}{ }^{\lambda} x(\lambda)$ at one second intervals of time (equal to the sampling
time of the input and response data) up to the maximum desired time. The integral of this function is then calculated at each time point of interest. The program allows for selection of the initial time point and for the spacing of all subsequent ones. The integration technique used in the identification is a fifth order quadrature which uses six data points in the calculation of the integral at each point. This use of many points provides a ceriain degree of smoothing.

$$
\text { The results of the integrations at } t=15 \text { are: }
$$

| Input | Pole | $e^{15 P_{i}} \times(15)$ | $\int_{0}^{15} e^{\rho_{i} \lambda} x(\lambda) d \lambda$ |  |
| :---: | :---: | :---: | :---: | :---: |
| I | 1 | 15.5142 | 213.436 | (A1-6) |
|  | 2 | 25.5658 | 278.862 |  |
|  | 3 | 31.1146 | 312.147 |  |
| 2 | 1 | 1.77239 | 12.5482 |  |
|  | 2 | 2.92071 | 17.7076 |  |
|  | 3 | 3.56719 | 20.3979 |  |

Muitiplication of the integrals by the proper coefficients and summation 0 (he three products to yield the coefficients of the unknowns is the third step of the computational scheme. From Equation (2-19) an example of one of these calculation is obtained by setting $j=0$.

$$
\begin{equation*}
y_{0}(15)=g_{1} \sum_{i=1}^{3} a_{i} e^{-\rho_{i} t} \int_{0}^{15} e^{\rho_{i} \lambda} x(\lambda) d \lambda \tag{i}
\end{equation*}
$$

The expansion of this summation is

$$
\begin{align*}
y_{0}(15) & =g_{1}\left(a_{1} e^{-\rho_{1} t} \int_{0}^{15} e^{\rho_{1} \lambda} x(\lambda) d \lambda+a_{2} e^{-\rho_{2} t} \int_{0}^{15} e^{\rho_{2} \lambda} x(\lambda) d \lambda\right. \\
& +a_{3} e^{-\rho_{3} t} \int_{0}^{15} e^{\left.\rho_{3} \lambda^{x} x(\lambda) d \lambda\right) .} \tag{Al-8}
\end{align*}
$$

Substitution of the values from (Al-1), (Al-2), and (Al-7) into Equation (Al-8) yields

$$
\begin{equation*}
Y_{0}(15)=4817.00 g_{1} \tag{A1-9}
\end{equation*}
$$

The other coefficients can be determined in a similar manner. In matrix form the results are

$$
G=\left|\begin{array}{cc}
4817.00 & 833.027  \tag{A1-10}\\
149.684 & 36.6890
\end{array}\right|
$$

At the same time the error matrix for an assumed one-percent error in the poles and input as calculated from Equation (3-20) is

$$
\Delta G=\left|\begin{array}{cc}
166.977 & 50.5095  \tag{Al-11}\\
3.02734 & 2.22371
\end{array}\right|
$$

The inverse of the coefficient matrix is

$$
G^{-1}=\left|\begin{array}{ll}
0.000705012 & -0.0160074  \tag{Al-12}\\
-0.00287631 & 0.0925630
\end{array}\right|
$$

and the expected error in the inverse is

$$
\Delta G^{-1}=\left|\begin{array}{ll}
0.00000772113 & -0.000175677  \tag{Al-13}\\
-0.0000123617 & -0.00127542
\end{array}\right|
$$

The final step of the computation is multiplication of the inverse times the response vector to give the zeros of the system.

$$
\begin{gather*}
g=G^{-1} \mathrm{X} \\
S=\left|\begin{array}{l}
0.0000513219 \\
0.00155654
\end{array}\right| \tag{Al-14}
\end{gather*}
$$

The expected errors in these zeros are calculated from Equation

$$
\begin{equation*}
\Delta g=\Delta G^{-1} X+G^{-1} \Delta X \tag{3-23}
\end{equation*}
$$

When the expected error in the response is assumed to be onepercent, this computation yields

$$
\Delta g=\left|\begin{array}{c}
0.00000105144  \tag{A1-15}\\
-0.0000861536
\end{array}\right|
$$

A more readily understandable form of these errors
is an absolute percentage error. These errors are:

$$
\Delta g=\left|\begin{array}{l}
2.04871  \tag{A1-16}\\
5.53496
\end{array}\right| \text { percent }
$$

This step completes the identification calculation. (Al-14) are the final results and (A1-16) are the expected percentage errors in the respective terms of (Al-14).

## APPENDIX C

## LISTINGS OF IDENTIFICATION PROGRAMS

In this appendix the listings of the programs developed for the identification of the system zeros are given. These programs are written in g-level Fortran IV. Comments are inserted throughout the programs to identify the inputs and to explain the computational steps. Figure $\mathrm{B}-1$ is a block diagram of the computational scheme.

The programs used in the pole identification studies were developed by Heymann and are listed in his dissertation. These programs are written in a modified Algol language called Osage Algol (Wl).


C-I. Block Diagram of Identification Technique




```
        CAL& INTEG(NPI,CN,SY;SCOS:I,J,U,CV) ZERO1190
        CALL INTEGINPI:CN,SS,SSIN,T,J,U,CVI ZEROL200
    10 CONILNUE
    CT4R1 = CT4+1 -U/CV
    IF(SLUN(1).EQ.O) GOTO 12
    DO 37 J=1,CN
    17 CONIINUE
    12 CONJINUE
        M=0
        DO 13 CK=U,NP1,CV
        M=M+1
        ERR(N=CI(CK)*CER(L)/100
        DO d8 I=1,CN
        CALL CALC&CN,CK,NPI,N,I,CAI,CBI,CREX,CIMX,SSIN,SCOS,GP,
        1 PERERR,CDP,ERRIN,U,GVJ
            CG(I,L,M)=CP
    CE{I,L,M\=CDP
    18 CONUINUE
    13 CONJINUE
    IF(A.EQ.CT) GOTO 19
    L=L+1
    GUTO }2
    1s P=0
    PAGE=1
    IF(SLON(1).EQ.O) GOTQ 22
    DO 23 J=1,CN
    DO 23 L=1,CT
    WRIGE(3,11118)(I,CG(H,L,I),CE(J,L,I), I= I,M)
    23 CONUINUE
    2 CONIINUE
```



```
C THE CALCULATIONS JO THIS POINT BAVE DETERMINED A SET OF
    ZERO1510
    SIMULTANEDUS EQUAIIONS. THE FOLLOWING SOLVE THESE EQUATIONS ZERO1520
    BY INVERTING THE GOEEFICIENT MAYRIX AND MULTIPLYING VECTOR BY IERO1530
    THE RESPONSE VECTOR BY IT.
*** *

NNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNN
 ododx


\(\mathrm{CK}=\mathrm{CK}+1\)
26
Condinue
朗910
\(\mathrm{CK}=\mathbf{C K}+1\)
ZERO1920
DO \(27 \quad \mathrm{I}=1, \mathrm{CN}\)
DO \(27 \mathrm{~J}=1, \mathrm{CN}\)
ZERO1930
ERO1940
\(Q=C N *(P-1)+1\)
ZERO1950
CREX(I,J) \(=\) CG(J.CALOI,CKd
CIMX(L,,\(J)=C E(J, C A(Q), C K)\)
ZERO1960
ZERO1970
ZERO1980
ZERO1990
ZERO1992
ZERO1995
ZERO2000
ZERO2010
ZERO2020
ZER02025
ZER02030
ZERO2040
ZER02050
ZER02060
ZER02070
ZER02080
ZER02090
ZERO2100
ZERO2110
ZERO2120
ZERO2130
ZERO2140
ZERO2150
ZERO2160
ZERO2170
ZERO2180
ZERO2190
98 CONILNUE
ZER02200
ZER02210
ZERO2220
ZERO2230

\section*{33 CONTINUE}

ZERO2240
CALA MADDICN,CIMX,CREX,-11 ZERO2245
\(Q=U *(C K-1) \neq C V \quad\) ZERO2250
IFISLON (21.EQ.O) GOTD 37
IFACOUNT-LE.551 GO TO 96
ZER02260
WRIUE (3, 11123) SERIES,PAGE
PAGE=PAGE +1
ZERO2270
ZERO2280
PAGE-PAGE+1
ZER02290
CQUAT=0
ZERO2300
96 CONJINUE
COWAT \(=\) COUNT \(+C N+3\)
ZERO2310
WRIUE(3.1117) \(Q\)
Z ER02320
WRLJE(3.11111) ZERO2330
37 CONIINUE
DO \(38 \quad 1=1, C N\)
\(M=C N *(P \rightarrow 1) \& I\)
QQ=O-DELAY(M)
ZERO2340
ZERO2350
ZERO2360
ZERO2370
ZERO2380
ZERO2390
ZERO2400
ZERO2410
ZERO2420
CAITI)=CS(I)*CER\&CA(N)J/100
IF (SLON(2).EQ.O) GOTO 39
WRITE(3.11118) CA(M) iCS (I). CAI(I)
39 CONIINUE
38 CONSINUE
IFIGOUNT.LT.57) GOTO 97
WRLIE(3.11123) SERIES.PAGE
PAGE=PAGE+1
COUNT=0
97 CONILNUE
COUAT = COUNT \(+C N+1\)
ZERO2430
ZERO2440
ZERO2450
ZERO2460
ZER02470
ZERO2480
ZERO2490
CAL.t MUATIP(CN,CV,CREX,CS)
C, is MULTIP\{CN,CESI,CIMX,CS)
: AL. MULTIP(CN, CEAI ©CREX, CAI)
DO an \(3 \quad I=1, C N\)
CS(A)=CEAI(I)+CEBI(Id
433: CCINIINUE
QN1 \(=Q-1\)
ZERO2500
ZERO2510
ZERO2520
ZERO2530
ZERO2540
ZERO2550

WHLEE(3.1117) QML
ZERO2570
ZERO2580
\(\ddagger\) \(\stackrel{\infty}{\infty}\)
fif
\(\stackrel{\star}{\omega}\)


NNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNNN





```

11111 FORMAT1" *"TEST*.12X,"RESPONSE*./)
11113 FORMAI!: "IDENTIEICATION OF ZEROS

```

```

    113.f TESTS *,12,**,I2,*,*I2,6X,*PAGE *,I41
    11115 FORMATI' ', 'COEFFICIENTS OF',IX,I2,IX,
1 'DERIVAIIVE OF WEIGHTING FUNCTION:)
L1116 FORNAT\&* *,6X.*GAIN'.6X,3(E13.6,4X))
11117 FORNAT\: , 4X,13.6X,F9.3,6X,F9.3,6X,F9.3,6X,F9.3)
11118 FORNAT4. ©,4X,13,6X,E13.6,6X,E13.6)
11120 FORNAT(* *,6X:ZERO*: REAL !.2(E13.6.4X))
11121 FORMAT(' *10X," IMAG *,2(E13.6,4X))

```

```

11123 FORMAT('1',15X,*SERIES ",I4.28X,'PAGE !.14)
11124 FORNAT(**:15X. SERLES :I4.28X, PAGE :I4)
SFOQ
END
ZERO3330
ZER03350
ZERO3360
ZER03370
ZERO3380
ZER03390
ZERO3400
ERO3400
ZER03410
ZER03420
ZERO3430
ZER03440
ZER03450
ZERO3460
ZER03470

```





\begin{tabular}{|c|c|}
\hline SIART REDNCTION & INV00320 \\
\hline DO \(18=1, N\) & I NV00340 \\
\hline \(L=R(1)\) & I NVO0350 \\
\hline R(L) \(=\) R(PIVI) & INV00360 \\
\hline R(PINI) = L & INV00370 \\
\hline \(L=C \ 11\) & I NV00380 \\
\hline C(L) = C(PIVJ) & INV00390 \\
\hline C(PEVJ)=L & I NV00400 \\
\hline INTEGER ICNT,ICNTI,ICNTJ & INV00410 \\
\hline \(1 \mathrm{CNIT}=\mathrm{R}(1)\) & INV00420 \\
\hline ICNI \(1=C(1)\) & INV00430 \\
\hline IF ADABS(AlICNT.ICNTA,I)-GE.EPS) GO TO 4 & INV00440 \\
\hline SINGUL=1 & INV00460 \\
\hline REIURN & INV00480 \\
\hline CONUINUE & I NV00490 \\
\hline \(006 J=1, N\) & INV00500 \\
\hline \(J K=A-J+1\) & INV00510 \\
\hline IF(NK.EQ-I) GO JO 5 & I NV00520 \\
\hline LCNJJ=C(JKI & I NV00530 \\
\hline A(ICNT,ICNTJ) \(=\) A(ICNT-ICNTJ)/A(ICNT, ICNTI) & INV00540 \\
\hline CONU LNUE & INV00550 \\
\hline AILONT, ICNTL) = 1./A(IONT, ICNTI) & I NV00560 \\
\hline PIVET = 0 & INV00570 \\
\hline DO \(\triangle K=1, N\) & INVO0580 \\
\hline IF AK-EQ.IA GOTO 7 & I NV00590 \\
\hline DO \(6 \mathrm{~J}=1, N\) & INV00600 \\
\hline \(J K=A-J+1\) & INV00610 \\
\hline IF AJK.EQ. 1 ) GO 106 & I NV00620 \\
\hline ICNUJ=C(.JK) & I NV00630 \\
\hline ICNSH \(=\) R(K) & I NV00640 \\
\hline H=AALCNT, ICNTJ)*A\&ICNTK, LCNTLJ & I NV 00650 \\
\hline A (ICNTK, ICNTJ) = A (ICNTK, ICNTJ)-H & INV00660 \\
\hline .IF AI.GE.K.OR.I.GE.JK.OR.DABS(PIVOT).GE.DABS(A)ICNTK, (CNTJ))) & INV00670 \\
\hline 1 GO JO 6 & I NV00680 \\
\hline PIVI=K & I NV00690 \\
\hline P I VIJ \(=. \mathrm{JK}\) & INV00700 \\
\hline
\end{tabular}
ONJ

CONE INUE


```

    W=A\J,P) PERO0335
    A1J,P)=A(K,P) PER00340
    A(K,P,)=W
    2 CONTINUE
GO \$0 6
5 CONTINUE
DO \ P=1,N
W=A|P,J)
A(P,Ji=A(P,K)
A(P,K)=W
7 CONIINUE
6 CONIINUE
TAG(J)=TAG(K)
TAG:K := T
TAGN=TAG(J)
TAGK=TAG(K)
LOCa(T)=LGC(IAGJ)
LOC(TAGIJ)=J
CONTINUE
3
coniINUE
RETURN
END
PERO0345
PER00350
PERO0355
PEROO360
PERO0363
PER00366
PER00370
PER00372
PER00372
PERO0375
PER00377
PERO0380
PERO0390
PER00400
PER00410
PERO0420
PER00430
PER00440
PER00450
PERO0460
PEROO470

```
```

    SUBRQUTINE MULTIP(CN,CA.CB,CC)
    MULT0100
INTEGER CN,I,J
REAE CA(5),CB(5,5),CC(5)
DO \& I=1,CN
CA(E)=0
DO 2 J=1,CN
CA(I) = CA(I) + CB(I,J)*CC(J)
2 CONJINUE
1 CONTINUE
RE TURN
END
MULTO190
MULT0110
MULTO120
MULTO130
MULTO140
MULTO150
MULTO160
MULTO170
MULT0180

```
SUBROUTINE MADD(N,A,B,M)

MADDO 100
MADDOI 10 MADDO120 MADDO1.30
\(001 \quad[=1, N\)

MADOO 140
DO \(\quad J=I, N\)
\(A(I, J)=A(I, J) * M * B(I, W)\)
CONTINUE
MADDO150
RETURN
MADDO 160
END
MADDO170
SUBROUTINE MADD(N,A,B,M)
DIMENSION A《5,5),B(5,5)
OO \(1 \quad I=1, N\)

SUBROUTINE INTERP\&N,A)
```


[^0]:    *no identification obtained.

