# FINITE ELEMENT MODELING OF STREAMFLOW ROUTING 

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Dedicated to members of my<br>family, loved ones, and to all those who appreciate the progress of mankind!

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Thesis Approved:


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## CHAPTER I

INTRODUCTION

## Background

The process of determining the water depths, velocities, and discharges in the channels, rivers, or reservoirs under unsteady conditions arising from flood motions is commonly referred to as flood routing. Interest in flood routing and in part in unsteady flows in water resources stems from the need to plan, design, regulate, and manage our flood prone areas and many other water resource systems.

In a surface water system, runoff, floods, droughts, and stream water quality interact very closely. When excess rainfall occurs over an area, runoff contributes to flooding along rivers whereas drought occurrence due to lack of rainfall results in minimal streamflow, reduced water supply, and less naviagation. In addition, the water quality of the stream becomes poor because of the low flows of the streams. Thus, the occurrence of lack of streamflow or drought affects the management of surface waters in the stream.

The State of Oklahoma experiences runoff ranging from 0.2 inches in the Panhandle to 20 inches in the southeast corner, which reflects the dramatic contrast in precipitation. In the northwestern region an average runoff amounts to about 820,000 acre-feet per year compared to $6,000,000$ acre-feet per year in the southeastern region. Annual average
runoff for the entire state is approximately $22,000,000$ acre-feet (Oklahoma Water Resources Board, 1980).

Flooding has been experienced over the years in Oklahoma. The Water Resources Council estimates that without increased flood management programs average annual flood damages will increase from $\$ 2.3$ billion in 1975 to $\$ 3.6$ billion in the year 2000. These damages occur over a flood plain of some 140 to 180 million acres. The Arkansas River Basin and the Red River Basin experiencedan estimated $\$ 167,000,000$ in flood damages in the state between 1955 and 1975, with the majority of that attributed to the Arkansas River (Oklahoma Water Resources Board, 1980).

Some floods occur gradually, as when prolonged steady rainfall saturates a river basin until most of it runs off, creating a greater volume of water than the natural channels and drainage structures can carry. Others are the result of sudden heavy rains occurring in a short time from thunderstorms. The latter is usually experienced in Oklahoma. In either case, floods are considered a problem only when the result is widely spread damage to agriculture and structures or when the normal activities of man are seriously interrupted.

Like other Great Plains States, Oklahoma has scores of extended droughts on an approximately 20-year cycle (Oklahoma Water Resources Board, 1980). An analysis of drought conditions from 1931 to 1971 indicates that drought occurred somewhere in the state about $51 \%$ of the time, more frequently in the panhandle and less frequently in northeastern and southcentral areas.

Water quality of Oklahoma's streams is adversely affected by natural and man-made pollution. In the west, natural salt springs and salt flats emit into local streams large quantities of chlorides that are subse-
quently carried downstream, polluting other major streams as they pass. In central and eastern Oklahoma, municipal and industrial effluents degrade many streams, restricting their beneficial use.

Thus, the interrelationships among runoff from rainfall, floods, droughts, and stream quality are predicted only when mathematical models to simulate the depth of flow and discharge in a stream resulting from rainfall are available.

## Study Objectives

The purpose of this study is to evaluate the discharge in the streams under the varying conditions of rainfall. The results of the mathematical models developed using the finite element methods will predict the depth of flow, velocity of flow, and the discharge in the streams.

Three mathematical flow models have been developed in this study. The first two are approximate models while the third is a complete model. They are presented below in the order of increasing complexity:

1. The kinematic flow model, KFM, solved explicitly and implicitly by Galerkin's weighted residual finite element method. The implicit version is implemented using a time weighting factor, and the resulting nonlinear system of a tridiagonal matrix equation is solved iteratively by the generalized Newton-Raphson method.
2. The diffusion flow model, DFM, implemented similarly to the implicit kinematic flow model, except that the resulting non-linear system is a bi-tridiagonal matrix equation. Solution is obtained by the Newton-Raphson
technique.
3. The complete flow model, CFM, produces a matrix equation similar to the implicit diffusion flow model and is solved using the same technique.

Model performance is evaluated using two forms of channel geometries. The first is comprised of an artificial stream channel of constant geometry with a hypothetical flood hydrograph imposed at the upstream end of the reach. Simulated flow is compared with the Viessman's solution using the explicit finite difference scheme. The second model test involves flow in the Illinois River, a natural river in Oklahoma. The Illinois River, a tributary of the Arkansas River, originates in northwestern Arkansas as Osage Creek and flows westward into Oklahoma. The flood recorded on April 10, 1979, for Watts and Tahlequah guaging stations, 50.4 miles apart, and that for Flint Creek, a tributary approximately 13.2 miles downstream of the Watts Station are utilized.

The choice of the natural channel is limited due to lack of adequate hydraulic data. Thought the Illinois River seems to exemplify varying channel geometric and hydraulic properties inherent in many other natural rivers in Oklahoma, the availability of data and the excellent flood hydrographs of 1979 record make it the best choice.

The objective of the first test with an idealized river channel is to explore the basic principles and to make some appraisal of the sensitivity of the controlling flow parameters in the mathematical models. The model application to a natural channel checks on the capability of simulating natural floods of long durations for use in the design of hydraulic structures as well as for the flood plain zoning.

## CHAPTER II

## LITERATURE REVIEW

Hydraulic And Hydrologic Routing Methods

Significant studies of unsteady flow in an open channel date back to the early works of the French mathematicians, Laplace (1775-76) and Lagrange (1783). The Lagrange celerity formula for small waves in shallow water provided the first impetus for subsequent studies. Later the British School of Mathematical Physicists gave some attention to fluid flow problems with contributions being made by Stokes, Kelvin, Rayleigh and Lamb (Water Waves, 1965).

The more advanced mathematical treatment of unsteady flow in an open channe1 is credited to Barré de Saint Venant (1871), a French mathematician who developed the complete one-dimensional equations of unsteady flow. These are two nonlinear hyperbolic partial differential equations of motion (conservation of mass and conservation of momentum) that very accuately describe the gradually varied flows in open channels. The original form of these equations is:

$$
\begin{align*}
& B \frac{\partial y}{\partial t}+A \frac{\partial v}{\partial x}+v \frac{\partial A}{\partial x}=0  \tag{2.1}\\
& \frac{\partial v}{\partial t}+v \frac{\partial v}{\partial x}+g\left(\frac{\partial y}{\partial x}+S_{f}-S_{0}\right)=0 \tag{2.2}
\end{align*}
$$

Where:

$$
A=\text { channel cross-sectional area, } \mathrm{ft}^{2} ;
$$

$B=$ width of channel water surface, ft;
$y=$ depth of flow, ft;
$x=$ distance along the channel, ft;
$\mathrm{t}=\mathrm{time}, \mathrm{sec} ;$
$S_{f}=$ friction slope, $\mathrm{ft} / \mathrm{ft}$;
$v=$ mean velocity across the section, $\mathrm{ft} / \mathrm{sec}$;
$S_{0}=$ longitudinal bottom channel slope, ft/ft

Two basic techniques for unsteady flow simulation are (1) methods which approximate a solution to the basic equations of unsteady flow (Eq. 2.1 and 2.2), and (2) methods which solve the basic equations. The first methods are sometimes referred to as "hydrologic" routing methods and the second kind, "hydraulic" routing methods (Thomas, 1975).

The importance of the hydraulic routing method has become increasing evident in the light of the modern high-speed digital computer for solutions of unsteady partial differential equations that have no closed form or analytical solution. Numerous unsteady flow phenomena such as surges, effects of tidal fluctuations, backwater resulting from channel junctions or reservoirs, and normal flood waves from excessive rainfall can be analyzed using the hydraulic routing and numerical methods such as finiteelement or finite difference methods.

On the other hand, the early development of the hydrologic routing in the form of a continuity equation is credited to Rippl (1883). In working on reservoir capacity problems, he utilized the concept of successive approximations to routing streams where the data are average daily flows, rather than slope, stage, and velocity measurements. Hydrologic routing is handy when data for hydraulic routing are not available.

The associated continuity equation is:

$$
\begin{equation*}
I-0=\frac{d S}{d t} \tag{2.3a}
\end{equation*}
$$

or

$$
\begin{equation*}
\frac{1}{2}\left(I_{1}+I_{2}\right) \Delta t-\frac{1}{2}\left(O_{1}+O_{2}\right) \Delta t=S_{2}-S_{1} \tag{2.3b}
\end{equation*}
$$

where, $I, 0, t,\left(S_{2}-S_{1}\right)$ are the inflow into a given reach, outflow from that reach, time period for the flow to travel through that reach, and the change in storage during that time period in the reach, respectively. The subscripts 1 and 2 represent conditions at the beginning and end of the routing periods.

The hydrologic routing is sometimes referred to as hydrograph routing because of the graphical relationship established between storage and outflow yields a feasible solution to Equation (2.3b) having two unknowns, $0_{2}$ and $\mathrm{S}_{2}$. Puls (1928) established a curve of relation between inflow and outflow versus storage for a variety of flood on the Tennessee River.

With some modifications of the continuity Equation (2.3b) to include the local inflows along the channel, Wisler and Brater (1931) presented a revised graphical scheme of Puls. This method was the first to use computed inflow hydrographs from tributaries and unmeasured areas for which no flow records are available. A number of hydrologic routing methods have emerged over the years including various coefficient routing procedures such as the Muskingum technique (McCarthy, 1938). Interested readers are referred to basic texts on hydrology (Chow, 1964; Viessman, 1972).

## Numerical Methods in Hydraulic Routing

## Finite Difference Methods

A large number of schemes have evolved from the finite difference methods over the years and have been applied with success to equations of unsteady flows and other engineering problems. For instance, the explicit, characteristic, and implicit schemes are the major categories. However, varieties of each group exist, for example the leap frog, diffusion and staggered explicit schemes, method of characteristics with fixed or characteristic grids, and the implicit scheme with weighted four-point or six point (Gunaratnan, 1970; Thomas, 1975; Fread, 1976).

A survey of previous literature indicates that many investigators to date have employed the finite difference schemes in flood routing problems. Isaacson et al. $(1954,1956)$ investigated flood routing in their peoneering work in the Ohio River. Amein (1966) used the method of characteristics to solve the streamflow problem in an attempt to study the effects of friction on peak flows. Amein and Fang (1969) also used an implicit scheme in solving the streamflow routing problem in natural channels in North Carolina. Pinder and Sauer (1971) employed the explicit method in simulating the flood wave modification due to bank storage effects. Fread (1971, 1973, 1974, 1976, 1978) investigated the routing problems using the implicit four-point and wieghted four-point finite difference schemes. Chaudhry and Contractor (1973), Liggett and Wollhiser (1967), Viessman et al. (1972), and many others have in turn used finite difference methods to solve approximate and complete routing equations.

It is interesting to remark that some of the finite difference schemes have some limitations often associated with convergence and stability problems. The explicit method is subject to a stringent stability condition imposing a limiting value for the time step in relation to distance step (Amein and Fang, 1969). The maximum time step that can be used in the explicit scheme to insure numerical stability when frictional effects are relatively small is computed using the Courant condition (Fread, 1973) as:

$$
\begin{equation*}
\Delta t_{c} \leq \frac{\Delta x_{i}}{\left(\left|v_{i}\right|+\left(g^{A} i_{i}\right)^{\frac{1}{2}}\right.} \tag{2.4}
\end{equation*}
$$

where:
$A_{i}$ and $B_{i}=$ area of flow and width of the water surface in the $i^{\text {th }}$ cross-section, respectively;
$\Delta X_{i}=$ the $i^{\text {th }}$ distance step;
$\Delta t_{C}=$ the computational time step;
$(A / B)_{i}=$ the hydraulic depth;
$v_{i}=$ velocity of flow in the $i^{\text {th }}$ cross-section.

Although the explicit scheme would not pose much difficulty for investigation of short time flows, it becomes cumbersome and inefficient for large flood flows in large rivers.

The method of characteristics is highly suitable for rapidly varied flows (Amein, 1966). It can be used for flood studies. However, the scheme is inconvenient in that the results are not obtained at fixed times and locations. A modification of the scheme employing a fixed mesh has been applied by Baltzer and Lai (1968) to tidal flows, but it
has no significant advantage over the explicit method for large river flows.

One requirement for the explicit scheme and method of characteristics is the use of equal distance intervals. This appears disadvantageous for rivers with irregular geometry (Fread, 1974). Thus, the development of the implicit schemes arise not only toovercome the equal distance requirement but also as a means of negating the restriction of small time steps imposed on the explicit and characteristic methods for reasons of stavility. The four-point implicit finite difference method appears most advantageous since it can readily be used with unequal distance intervals (Fread, 1973, 1974, 1976).

In the light of the inherent advantages of the implicit four-point scheme, Fread (1973) investigated the influence of the time weighting factor, $\theta$, for spatial variables along with those of the channel parameters, such as the length of the reach, bed slope, roughness coefficient, and surface width, on the numerical distortion (dispersion and attenuation of computed stage hydrographs). The definition of $\theta$ is presented in Figure 1. Among other things, the following observations were made. The lower range of allowable $\theta$ values minimizes the distortion which results from the use of large time steps in the integration of the implicit difference equation. A value of $\theta=0.55$ was chosen to minimize distortion while conservatively insuring theoretical stability criteria. The tendency for the stability of the numerical computations to decrease with increasing value of $\theta$ exists.

On the other hand, numerical distortion increases when the channel length, $L$, or the Manning roughness factor, $n$, increases; and it decreases when the magnitude of the initial depth of flow, $y_{0}$, or the


Figure 1. Definition of the Dimensionless Time
channel bottom slope, $S_{0}$, increases. The channel width, B, was observed to have little or no effect on the magnitude of the numerical distortion.

In general, it is expected that the results obtained by the implicit method would be no different than those obtained by other numerical methods for the solution of the complete equation of unsteady flow in open channels. The main difference is that the implicit method provides the result faster (Amein and Fang, 1969).

## Finite Element Methods

## Evolution and Extension to Fluid Dynamics

The evolution of the present-day finite element methods has followed a long but imprecise history (Zienkiewicz, 1977). To date, the unified efforts of the early mathematicians and those of engineers mostly in the structural discipline have given rise to a complete picture of finite element methods. The contributions of the mathematicians are seen in the area of formula development (the governing differential equations) for the physical problems and solution techniques such as the variational principle, Gurtin principle, and the weighted residual principles.

On the other hand, the engineers tend to approach the problem by establishing a direct analogy between the real discrete element and finite portions of a continuum domain. As Zienkiewicz (1977) puts it, it is from this "direct analogy" view that the term finite element was born. The existence of a unified treatment of the "standard discrete problems" leads to the first definition of the finite element process as the method of approximation to continuum problems such as:

1. The continuun is divided into a finite number of eicments whose behavior is specified by a finite number of parameters, and
2. The solution of the complete system as an assembly of its elements follows precisely the same rules as those applicable to standard discrete problems (p. 3).

For the simple reason that a number of classical mathematical procedures of approximation fall into this category as well as the various direct approximations in engineering, Zienkiewicz (1977) states that the origin of the finite element procedures and the precise moment of its invention are difficult to determine. A supporting point of view is held by Oden (1972) who comments on the piecewise approximations and rudiments of the idea of interpolation supposedly used in ancient Babylonia and Egypt that preceded the calculus over 2000 years ago.

More recently, the practice of representing a structural system by a collection of discrete elements was utilized in the early works of the aircraft structural engineers (Courant, 1943). The formal presentation of the finite element methods together with the direct stiffness method for assembling elements is attributed to Turner, Clough, Martin, and Topp (1956). It was Clough (1960) who first used the term "finite elements" in a later paper devoted to plane elasticity problems.

The application of the finite element method to fluid flows began to assume a degree of importance in the mid-sixties following the early works of Zienkiewicz et al. (1965, 1966), Javandel and Witherspoon (1968), and Tyagi (1971) in porous media flow. For the last decade scores of papers have emerged applying the finite element methods to surface water systems for estuaries, reservoirs and streams, and groundwater systems for flow in saturated and unsaturated zones and groundwater quality (Gallagher, et al., 1974, 1976; Gray, Pinder, and Brebbia, 1976; Ciriani, Maione and Wallis, 1974; Tyagi, 1975a, 1975b, 1975c).

## Weighted Residuals Methods, WRM

Elaborate discussion on the basic finite element schemes is found in literature (Zienkiewicz, 1977; Oden, 1972; Finlayson, 1972; Segerlind, 1976; Ames, 1977; Chung, 1978). Norrie and DeVries (1975) presented a bibliography covering over 3800 citations during 1956-1974.

As an approximate method of solving differential equations of initial and/or boundary value problems in engineering and mathematical physics, the finite element can be implemented via variational principle or weighted residual principles. The variational principle is based on the works of Rayliegh (1877) and Ritz (1909). Some classes of problems can not easily be put into variational form, particularly when the governing differential equations are not self-adjoint. Thus, this method has limited application.

The weighted residual methods (WRM), which include the orthogonal collocation, Bubnov-Galerkin, Subdomain, and least-squares, are employed to deal directly with the governing equations of the physical problems (Finlayson, 1972; Ames, 1977). The weighted residuals in general utilize the concept of orthogonal projections of a residual of a differential equation onto a subspace spanned by certain weighting function. Stated differently, the unknown solution in all the WRM is approximated by a set of local basis functions containing adjustable constants or functions (Ames, 1977). These constants or functions are chosen by various criteria to give the "best" approximation for the selected family. For instance, the least-squares method requires higher order interpolation functions in general, even if the physical behavior may be adequately described by lower order (linear) functions (Chung, 1978). This restriction limits its
use. The collocation is the simplest (WRM) to apply, but it has a drawback in terms of the number of nodes needed to achieve the same results as with the Galerkin method.

Of special interest in all the WRM is the Bubnov-Galerkin method (1913 and 1915, respectively). The method (often referred to as Galerkin without Bubnov) is the most popular and widely used. Large numbers of non-linear fluid flow systems are easily transformed into "finite element equations" directly. The classical procedures of the Galerkin assume the weighting function and the trial function to be identical (Zienkiewicz, 1977). Like the variational principle, the Galerkin always yields asymmetric matrix equation for linear differential operators.

## Approximation of Time Derivatives.

The concept of extending the finite element to include the time domain is discussed by Oden $(1969,1972)$ and Chung (1978). The approach is to regard the basis function as being dependent on time as well as the spatial domain such that:

$$
\begin{equation*}
\frac{\partial v(x, t)}{\partial t}=\frac{\partial N_{i}(x, t)}{\partial t} v_{i} \tag{2.5}
\end{equation*}
$$

where:

$$
\begin{aligned}
v(x, t)= & \text { dependent variable, } v \text {, expressed as a } \\
& \text { function of space, } x, \text { and time, } t ; \\
N_{i}(x, t)= & \text { basis function at node, } i \text {, as a function } \\
& \text { of space and time. }
\end{aligned}
$$

Other investigators have extended this idea in many studies in
water resources*. Zienkiewicz and Lewis (1973) investigated two linear finite element formulations in the time domain. Grotkop (1973) applied the Galerkin method in the time domain to estuary modeling; Gray and Pinder (1974) conducted a numerical experiment on the use of the Galerkin finite element method to approximate both the time and space derivatives and thereby study the suitability of using higher order basis functions in the time domain for solving the transient groundwater flow equation. Van Genuchten (1977) employed higher basis function (including Hermitean) to a one-dimensional solute transport equation and studied the accuracy of the resulting schemes.

One major disadvantage of finite element approximations in time derivatives is the enormous increase in the computational time and effort. Gray and Pinder (1974) noted the inherent tradeoff between increased accuracy and decreased computational efficiency associated with the finite element time derivatives. The optimum scheme for approximating the time derivative in a groundwater flow problem is dependent on both the behavior of the solution and the method of time step selection.

The second approach to the time derivative approximation is the so called "semidiscrete method" in which the time derivative of a variable at nodes is replaced by a temporal operator (finite difference operator) from the relation, (Chung, 1978):

$$
\begin{equation*}
\frac{\partial v(x, t)}{\partial t}=\dot{v}_{i}(x, t)=N_{i}(x) \dot{v}_{i}(t) \tag{2.6}
\end{equation*}
$$

where:
$\dot{v}_{i}(t)=$ time derivative of $v$ prescribed at node $i$

[^0]Many of the finite element models of transient problems adopt the previously described semidiscrete method. When the time derivatives are approximated with finite differences, either a central in time (CrankNicolson) or a backward/forward in time (implicit) scheme can be used (Van Genuchten, 1977). The former results in a second-order accuracy while the latter yields only a first-order accuracy. If higher order basis functions are used, it may be important to obtain a higher order approximation of the time derivatives. This might not be necessary if lower basis functions such as the linear types are employed.

## Numerical Properties of FEM for Non-Linear Systems

Numerical properties of the finite element method, such as the stability, convergence, and accuracy, unlike those of the finite difference methods have not been established adequately although many intuitive proofs and conclusions have been stated (Desai and Christian, 1977). The study of the numerical procedures has often been made in a pragmatic manner. When a given scheme is used for a number of problems and it is found satisfactory, it is considered acceptable. The major criticism of this approach is that it may not yield a general scheme (Desai and Christian, 1977).

Error analysis associated with the solution of the non-linear hyperbolic open channel flow equations may be grouped as: finite element approximation errors; temporal approximation errors; and errors due to any iterative non-linear equation solver, such as the Newton-Raphson, pre-dictor-Corrector, and others. At present, no theoretical finite element error estimates are available for the unsteady nonlinear two-variable equations. However, it is possible to perform error analysis due to
temporal operators, together with the iterative equation solver in a restricted sense. Adopting a procedure described by Chung (1978, pp. 227), by holding the non-linear terms constant during the iteration cycle, it is then possible to generate an approximate amplification property matrix, a technique reported by Lax-Richtmyer (1956) and Richtmyer-Morton (1967). Every eigenvalue $\lambda_{i}$ of the amplification matrix, if made smaller than unity, automatically insures stability. The largest eigenvalue, called the spectral radius of the amplification matrix which governs the stability, and the limiting value of the time step, $\Delta t$, can be determined. As the non-linear terms are updated, the amplification changes, thus altering the stability criteria as calculations progress. This increases the difficulty in the stability analysis of an unsteady non-linear system.

A slightly different approach to the error analysis for non-linear hyperbolic equations (Oden and Fost, 1973) requires the finite element basis functions to satisfy the convergence and completeness criteria as for linear elliptic problems. The study yields a stability estimate that is considered to be consistent with the well-known Von Neumann linear stability criterion which requires the discrete system to propagate information at a rate greater than or equal to the speed of propagation of the actual system. The above approach is too narrow in concept, a linearization technique drawn from elliptic type of problems, and limits the use to the special class of hyperbolic equation studied, the onedimensional homogeneous hyperelastic bodies.

With regard to the finite difference method, Fread (1974) studied the numerical properties of the St. Venant equations for a four-point implicit scheme using the Von-Neumann technique. Since this technique
is only applicable to linear differential equations, linearization of the governing equations is adopted with certain terms omitted on the basis of their relatively small magnitude in order to facilitate the stability analysis. On the other hand, the convergence criterion was analyzed by expanding each term in the Taylor series expansion about the point at which the differential equation is computed. The study concludes that the implicit four-point method is unconditionally stable provided the time weighting factor, $\theta \geq 0.5$, and has a second order accuracy since time step, $\Delta t$, and distance step, $\Delta x$, are quadratic.

Cooley and Moin (1976) studied the numerical properties of the St. Venant equations using the finite element method* and the predictorcorrector iterative solving scheme. They adopted the linearization technique similar to those used by Strelkoff (1970) for stability analysis. Their concluding remarks are identical to those of Fread (1974). Error analyses for other classes of differential equations are reported in literature (Kreig and Key, 1971; Fujii, 1972; Desai, Oden and Johnson, 1975; Desai and Lytton, 1975; and Chung, 1978).

## Finite Element Versus Finite Difference

The purpose of resorting to the Numerical Methods is to be able to solve problems either for which there is no analytical solution or for which the analytical solution is too hard to obtain. For the last two decades, attention has been drifting from the finite difference method

[^1]to finite element method in hydrology and water resources. The search for the most, efficient and accurate simulation model has continued to be the center of inspiration for this change of attention. The question often raised is "Is there any real reason for this change of attention?" Review of major studies in fluid dynamics involving the FDM and FEM sheds some light on answering the question.

In a very vigorous classification of trial functions, Zienkiewics (1974) states that within a broad definition, the finite difference technique falls into a "subclass" of the general finite element methodology, which indeed embraces many other classical approximation procedures. Nevertheless, both techniques can be considered as distinct in a much narrower perspective. For instance, the two methods differ in a manner in which the element equations are generated from the governing equation. While the two adopt the principle of discretization as the initial step in the numerical procedure, the way the concept of discretization is implemented varies. In the FDM, the governing equation is discretized whereas in the $F E M$, the region or continuum of the system is discretized. In other words, in the FEM the problem is formulated as an integral to be minimized, and we use a numerical approximation of the integral to obtain a solution. This step is necessary regardless of the kind of FEM adopted --variational principle or method of weighted residuals (Myers, 1971).

Another distinguishing feature of the FEM is the difference in the grid and element numbering system. For instance, a typical element, e, is the interval between nodal points, $i$ and $j$. This numbering scheme is slightly different from the FDM where nodal-point number is also used to
designate the region surrounding the nodal points. In the FEM, the numbering of the nodal points is entirely separate from the numbering of the elements.

Though it has not been proved technically that any one method is superior to the other, what seems obvious is that the finite element method may prove more advantageous for some classes of problems--those with extremely complex geometry--than FDM. A supporting viewpoint can be drawn from Myers (1971). In a one-dimensional steady state heat transfer problem for a thin rod, solution is sought by the variational principle (FEM) and FDM using the same number of nodal points. It was observed that the FEM solution falls below the exact nodal values by about the same amount that the FEM are above the exact values. This was explained by the fact that FEM was generated by minimizing the integral (Myers, 1971).

Some numerical studies performed by Pinder and Gray (1976) using the equation governing the convective -diffusion transport of a conservative contaminant help to illustrate the relationship between FEM and FDM. By using the Galerkin approximation of the space derivatives and the finite difference approximation of the time derivatives, they observed that the FEM can be expressed in terms of weighted average finite difference approximations. However, this observation had been reproted by others earlier (Myers, 1971; Finlayson, 1972).

Advantages of one method over the other in terms of the numerical properties such as convergence and stability may depend on the nature of the problem as well as the solution technique adopted, Newton-Raphson or Predictor-Corrector method for non-linear problems. For simulation of floods of long duration, a stable algorithm with large distance and
time steps is needed. As Cooley and Moin (1976) and Manan et al. (1977) indicated, the FEM has some answer.

## MATHEMATICAL STATEMENT

Introduction

The mathematical expressions of the unsteady gradually varied streamflow hydraulics are afforded by the well-known "Saint Venant Equations," named after Barre de Saint Venant (1871) who first derived them. The original forms of these equations as presented in Chapter II have been modified to include the lateral flow term. These equations are onedimensional non-linear hyperbolic, initial as well as boundary value partial differential equations, which may be derived from the laws of conservation of mass and momentum.

No attempt is made to re-derive these equations herein, rather interested readers are referred to any basic text on open channel hydraulics, such as Chow (1959), Henderson (1966), Viessman et al. (1972), Wylie and Streeter (1978) and many others.

## Governing Differential Equations

The distribution of depth of flow and velocity of flow and discharge in a stream are represented in Figure 2, following. The mathematical model that predicts the flow on a space and time basis can be represented by the following equations, Viessman et al. (1972):

Equation for conservation of mass

$$
\begin{equation*}
\frac{\partial y}{\partial t}+y \frac{\partial v}{\partial x}+v \frac{\partial y}{\partial x}-q(x, t)=0 \tag{3.1}
\end{equation*}
$$



Figure 2. Streamflow Element

Equation for conservation of momentum

$$
\begin{equation*}
\frac{\partial v}{\partial t}+v \frac{\partial v}{\partial x}+\frac{v}{-y} g(x, t)+g\left(\frac{\partial y}{\partial x}+s_{f}-S_{0}\right)=0 \tag{3.2}
\end{equation*}
$$

where:
$q=$ lateral inflow in the channel reach, $\Delta x$, ft per sec;
$S_{f}=\frac{n 7^{2} v^{2}}{2.2082 R^{4 / 3}} \begin{aligned} & \text { (friction slope, derived from Manning's EQ.), } \\ & \mathrm{ft} / \mathrm{ft} ;\end{aligned}$
$\mathrm{n} 1=$ Manning's roughness factor, sec per ft ${ }^{-1 / 3}$;
$R=$ hydraulic radius, ft;
Other terms are as defined for EO. (2.1) and (2.2).
The two dependent variables in Equations (3.1) and (3.2) are the depth of flow, $y(x, t)$, and the velocity of flow, $v(x, t)$. The channel geometry is specified by the area of flow, $A(x)$, the hydraulic width, $B(x)$, (where $A(x)=B(x) \cdot \partial y / \partial x$ ) and the slope, $S_{0}=S_{0}(x)$. The lateral inflow $q(x, t)$ has about three possible sources of contribution, namely, the rainfall on the stream, overland flow, and the subsurface inflow.

The conservation of mass and momentum equations presented above are classified as one-dimensional in the sense that flow characteristics such as depth and velocity are considered to vary only in the longitudinal $X$ direction of the channe1. Other simplifying assumptions inherent in their derivation are as follows: (1) the velocity is constant and the water surface is horizontal across any section perpendicular to the longitudinal axis; (2) the flow is gradually varied with hydrostatic pressure prevailing at all points in the flow such that the vertical acceleration of water particles may be neglected; (3) the longitudinal axis of the channel can be approximated by a straight line; (4) the bottom slope of the channel is small; (5) the bed of the channel is fixed, i.e. no scouring
or deposition is assumed to occur; (6) the resistance coefficient for steady uniform turbulent flow is considered applicable, and an empirical resistance equation such as the Manning equation describes the resistance effects; and (7) the flow is incompressible and homogeneous in density (Fread, 1976; Freeze, 1972).

Once the velocity of flow and depth of flow are computed from Equations (3.1) and (3.2), the discharge can be computed from the following equation:

$$
\begin{equation*}
Q=v y \tag{3.3}
\end{equation*}
$$

where $Q=$ the streamflow volumetric flow rate; cubic feet per sec per channel width of flow.

## Initial and Boundary Conditions

The essential requirements to initiate any hydraulic routing, be it open channel or overland flows, are the initial and boundary conditions. The distinction between initial and boundary conditions is merely one of position on the $x$ plane at the commencement of the solution procedure, (Viessman et al., 1972). The initial condition on one hand describes the flow depth, velocity, or discharge at all points in space at time, $t=0$. If flow is assumed uniform and steady before any flood wave reaches the point of interest upstream of the entire channel, then either the Manning or Chézy's equation is employed to calculate the initial flow parameters.

On the other hand, the boundary condition refers to the depth, velocity, or discharge at the up- and down-stream points or other point(s) of interest on the river reach at all times, $t>0$. Examples of boundary conditions are discussed elegantly by Fread (1976) and summarized in the following equation:

$$
\begin{equation*}
M y+N v=P \tag{3.4}
\end{equation*}
$$

where
$M, N$ and $P=$ known functions of either $y$ or $v$ or both.
Either $M$ or $N$ is zero at the upstream boundary, and $M$, $N$, and $P$ are segments of a rating curve for the downstream boundary.

Important points to keep in mind in boundary condition specification are as follows: (1) if discharge hydrographs are used for both the upstream and downstream boundary conditions, any error in the initial conditions (the initial depth of flow and velocities at all computational nodes along the stream between the up- and down-stream boundaries when the simulation is started) will be perpetuated in the computations (Fread, 1976). This is not the case when other possible combinations of the boundary conditions (specified depths or discharges upstream and rating curve downstream) are used. (2) Associated with the channel hydraulics are two interacting phenomena, namely, the state of flow (subcritical, critical, and supercritical) and the boundary conditions (Viessman et al., 1972). For subcritical flow, boundary conditions are required at both up- and down-stream of the river reach whereas only two upstream boundary conditions are necessary in supercritical flow. This is because downstream effects can not be propagated backward.

## Simplified Models

The solution of the complete one-dimensional unsteady flow Equations (3.1) and (3.2) oftentimes results in enormous computer time and storage, particularly for floods of long durations. In essence, this has attracted significant interest in the use of simplified models, such as the kinematic and diffusion flow models. The mathematical justification in the use of these simplified models is provided by the slope approximation
analysis and Froude number order of magnitude analysis (Henderson, 1966). While the continuity equation is completely retained, the simplifying assumptions are made in the momentum equation. If Equation (3.2) is re-arranged with the friction slope, $s_{f}$, being the subject of the formula and letting $q(x, t)$ equal zero, the resulting equation is:


The volumetric flow rate, $Q$, is obtained by combining Equation (3.5) and Manning's formula as:

$$
\begin{equation*}
Q=\frac{1.486}{n T} A R^{2 / 3} \sqrt{S_{0}-\frac{\partial y}{\partial x}-\frac{v}{g} \frac{\partial v}{\partial x}-\frac{1}{g} \frac{\partial v}{\partial t}} \tag{3.6}
\end{equation*}
$$

From Equation (3.6), if the last three slope terms are small compared with $s_{0}$, the discharge, $Q$, can be computed as in uniform flow, and it is dependent on depth only. The resulting relationship is known as the kinematic model, and its momentum equation is expressed as:

$$
s_{f}=s_{0}=\frac{n l^{2} v^{2}}{2.22 R^{4 / 3}}
$$

or

$$
Q=\frac{1.486}{n 1} A R^{2 / 3} S_{0}^{\frac{1}{2}}
$$

The kinematic model has been successfully applied in simulating flows in
natural floods in steep river slopes of the order of 10 feet per mile or more, overland flows, and slow-rising hydrographs (Henderson, 1966).

If the longitudinal streambed slope, $S_{0}$, is very flat, the $\partial y / \partial x$ term in Equation (3.6) may well be of the same order as $S_{0}$. In this case the Froude number ${ }^{ \pm}$, $F$, will be very low, so that the third term in Equation (3.6) will be negligible. In fact the third and fourth terms can be shown to be of the same order of magnitude. Details of the mathematical proof are discussed by Henderson (1966). However, for $F^{2} \ll 1$, the terms $v / g \partial v / \partial x$ and $1 / g \partial v / \partial t$ are of the same order of magnitude. This flow condition yields the diffusion flow model. The momentum equation yields:

$$
\begin{equation*}
S_{f}=S_{o}-\frac{\partial y}{\partial x} \tag{3.8}
\end{equation*}
$$

Indeed, the kinematic and diffusion flow models are two extreme cases of slopes--steep and flat--which are frequently encountered in overland flow on watersheds and natural routing of a flood wave in streams. Conceivably, there are possible intermediate values of slope for which all the four slope terms in Equation (3.6) would be appreciable. This is a case where the complete flow model is employed.

[^2]
## CHAPTER IV

## FINITE ELEMENT FORMULATION

## Introduction

The finite element method selected here is the Galerkin's weighted residual principle. This is an excellent choice for the solutions of the unsteady open channel flow equations that are characterized by the nonlinear hyperbolic behavior. This class of equations cannot easily be expressed in the variational form because the governing differential equations are not self-adjoint. Thus, the weighted residual principle, such as the Galerkin's principle, is employed to solve the governing equations of unsteady flow. In its final form, the method generates a system of ordinary differential equations in time for transient problems.

The weighted residuals utilize the concept of orthogonal projections of a residual of a differential equation onto a subspace spanned by certain weighting function. A discussion as it applies to some finite element problems is given by Chung (1978), Norrie and De Vries (1973), Martin and Carey (1975), and Zienkiewicz (1977). The implementation of the finite element formulation of the flow equations is carried out in four basic steps--(1) channel discretization and selection of approximation functions, (2) derivation of element equations, (3) assembly of element equations, (4) transient solution of the system of equations. For sake of clarity, the continuity Equation (3.1) is chosen to illustrate these steps.

Channel Discretization and Selection<br>Of Approximation Function

The natural channel shown in Figures 3 a and 3 b is idealized as a straight line as presented in Figure 3c because the flow equations are one-dimensional. The channel is divided into ( $\mathrm{N}-1$ ) small segments called elements or reaches where $N$ is the total number of nodes for which the solution of the dependent variables is sought. Each element will be modeled with the same flow equation but with different channel geometry and hydraulic properties. The element equations are later assembled into global matrix equations for solution.

To initiate the element equations, the approximation of the dependent variables, such as the velocity of flow, $v(x, t)$, and the depth of flow, $y(x, t)$, that form continuous functions over the infinite distance into discrete variables for a finite distance is necessary in the finite element method. Approximation functions, also known as shape or basis functions, include linear, quadratic, higher order polynomials or spline functions. The linear shape function is utilized to keep calculations simple.

It is important to note that a single function approximating the entire flow domain is difficult to find. The finite element method simplifies the procedure by breaking down or discretizing the function and domain into the elements shown in Figure 4. The characteristics of a shape function are summarized as follows:

1. Each function denoted as $N_{k}^{e}$ is zero, except within the element $e$, and $k$ must be a node of $e$.
2. The function $N_{k}^{e}$ is defined as a continuous function of the independent variable $x_{k}$ over the element $e$ in such a manner that the value


Figure 3. Natural - Idealized Flow Sections.


Figure 4. Approximation of Flow Domain by a Piecewise Continuous Function.
at the nodal point $k$ is unity, and the values at the other nodal points of the element are zero.
3. The function $y(x, t)$, or $y$ for simplicity (depth of flow), is allowed to vary linearly in each element:

$$
\begin{equation*}
y=A+B x \tag{4.1}
\end{equation*}
$$

where

$$
A \text { and } B=\text { constants }
$$

For determining the values of $A$ and $B$, consider Figure 5 and Equation (4.1). Two simultaneous equations are generated by substituting into Equation (4.1) the corresponding values at points $x_{k}$ and $x_{k+1}$ respective1y. These equations are:

$$
\begin{align*}
y_{k} & =A+B x_{k}  \tag{4.2}\\
y_{k+1} & =A+B x_{k+1} \tag{4.3}
\end{align*}
$$

Solve for $A$ and $B$ :

$$
\begin{aligned}
y_{k+1}-y_{k} & =B\left(x_{k+1}-x_{k}\right) \\
B & =y_{k+1} / x_{k+1}-x_{k} \\
y_{k} & =A+\frac{y_{k+1}-y_{k}}{x_{k+1}-x_{k}} x_{k} \\
A & =\frac{y_{k}\left(x_{k+1}-x_{k}\right)}{x_{k+1}-x_{k}}-\frac{x_{k}\left(y_{k+1}-y_{k}\right)}{x_{k+1}-x_{k}} \\
A & =\frac{x_{k+1} n_{k}-x_{k} n_{k+1}}{x_{k+1}-x_{k}}
\end{aligned}
$$

Thus, the linear shape function becomes:

$$
\begin{equation*}
y=\frac{x_{k+1} y_{k}-x_{k} y_{k+1}}{x_{k+1}-x_{k}}+\frac{y_{k+1}-y_{k}}{x_{k+1}-x_{k}} x \tag{4.4a}
\end{equation*}
$$



Figure 5. Linear Shape Function.

$$
\begin{align*}
& =\frac{x_{k+1}-x}{x_{k+1}-x_{k}} y_{k}+\frac{x-x_{k}}{x_{k+1}-x_{k}} y_{k+1}  \tag{4.4b}\\
& =N_{k}^{e} y_{k}+N_{k+1}^{e} y_{k+1} \tag{4.4c}
\end{align*}
$$

By adjusting the coordinate system in Figure 5 such that the origin is at $x_{k}$ and the distance from the new origin to $x_{k+1}$ is $L$, Equation (4.4b) reduces to

$$
\begin{equation*}
y=(1-s) y_{k}+s y_{k+1} \tag{4.5}
\end{equation*}
$$

where

$$
N_{k}^{e}=(1-s), N_{k+1}^{e}=s \text { and } s=x / L
$$

when $x=x_{k}, N_{k}^{e}=1$ and $N_{k+1}^{e}=0$

$$
x=x_{k+1}, N_{k}^{e}=0 \text { and } N_{k+1}^{e}=1
$$

as required part of the characterisitics of the shape function.

## Derivation of Element Equations

The Galerkin's weighted residual method is the basis of the element derivation equations. The method requires that errors or residual between the approximate solution and the true solution be orthogonal to the functions used in the approximation. The principle is expressed mathematicalTy by Segerlind (1976):

$$
\begin{equation*}
\int_{R} N_{\beta} L(\phi) d R=0 \quad \beta=i, j, k, \ldots \tag{4.6}
\end{equation*}
$$

where

$$
\begin{align*}
N_{3} & =\text { shape function; } \\
\phi & =\text { unknown parameter and is approximated by } \\
\phi & =\left[N_{i}, N_{j}, N_{k}, \ldots .\right]\{\phi\} ; \tag{4.7}
\end{align*}
$$

$$
\begin{aligned}
\dot{L}(\phi) & =\text { differential equation governing } \phi ; \text { and } \\
R & =\text { region of interest. }
\end{aligned}
$$

Equation (4.6) implies that the shape function $N_{B}$ must be orthogonal to the residual between the approximate solution and the true solution over the region R. Inserting the continuity Equation (3.1) into Equation (4.6) yields:

$$
\begin{equation*}
\sum_{1}^{k-1} \int_{x_{k}}^{x_{k+1}} N^{\top}\left(\frac{\partial y}{\partial t}+y \frac{\partial v}{\partial x}+v \frac{\partial y}{\partial x}-q(x, t)\right) d x=0 \tag{4.8}
\end{equation*}
$$

where

$$
\begin{aligned}
& \sum_{1}^{k-1}= \text { expression for summing individual element equations from } 1 \text { to } \\
&(k-1) \text { elements; } \\
& N^{\top}= \text { transpose to the shape function, and other terms are as defined } \\
& \text { previously. }
\end{aligned}
$$

Using the shape function, Equation (4.5) into Equation (4.8) gives

$$
\begin{align*}
& \text { Term (1) (2) (3) } \\
& \sum_{1}^{k-1} \int_{0}^{1} N^{\top}\left(\frac{\partial y}{\partial t}+y \frac{\partial v}{\partial x}+v \frac{\partial y}{\partial x}-q(x, t)\right) L d s=0 \tag{4.9}
\end{align*}
$$

Contribution of terms from left to right in Equation (4.9) is given be10w:

Term (1):

$$
\begin{aligned}
\int_{0}^{1} N^{\top}\left(\frac{\partial y}{\partial t}\right) L d s & =\int_{0}^{1}\left[\begin{array}{cc}
(1-s) \\
s
\end{array}\right]\left[\begin{array}{ll}
(1-s) & s
\end{array}\right]\left[\begin{array}{l}
\dot{y}_{1} \\
\dot{y}_{2}
\end{array}\right\} L d s \\
& =\frac{L}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
\dot{y}_{1} \\
\dot{y}_{2}
\end{array}\right\}
\end{aligned}
$$

where

$$
\dot{y}=\frac{\partial y}{\partial t}, \text { time derivative of } y
$$

## Term (2):

For the second term, first consider the following analysis:
Let $v=N_{k} V_{k}+N_{k+1} V_{k+1}=[N]\{V\}$
then $\frac{\partial V}{\partial x}=\frac{\partial}{\partial x}([N]\{V\})$

$$
=\left[\begin{array}{ll}
\frac{\partial N_{k}}{\partial x} & \frac{\partial N_{k+1}}{\partial x}
\end{array}\right]\left\{\begin{array}{l}
V_{k} \\
V_{k+1}
\end{array}\right\}
$$

Thus

$$
\left.\begin{array}{rl}
\int_{0}^{1} N^{\top}\left(y \frac{\partial v}{\partial x}\right) L d s & =s_{0}^{1}\left[\begin{array}{c}
(1-s) \\
s
\end{array}\right][(1-s) s
\end{array}\right]\left\{\begin{array}{l}
y_{1} \\
y_{2}
\end{array}\right\}\left[\begin{array}{ll}
\frac{\partial N_{1}}{\partial x} & \frac{\partial N_{2}}{\partial x}
\end{array}\right]\left\{\begin{array}{l}
v_{1}  \tag{4.9b}\\
v_{2}
\end{array}\right\}
$$

Term (3):
The third term is

$$
\int_{0}^{1} N^{\top}\left(v \frac{\partial y}{\partial x}\right) L d s=1 / 6 \quad\left[\begin{array}{l}
\left(2 v_{1}+v_{2}\right)\left(y_{2}-y_{1}\right)  \tag{4.9c}\\
\left(v_{1}+2 v_{2}\right)\left(y_{2}-y_{1}\right)
\end{array}\right]
$$

Term (4):
The last term is

$$
\begin{align*}
\int_{0}^{1} N^{\top}(q(x, t)) L d s & =q(x, t) s_{0}^{1}\left[\begin{array}{c}
(1-s) \\
s
\end{array}\right] L d s \\
& =\frac{L}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\} q(x, t) \tag{4.9d}
\end{align*}
$$

Combining each of the evaluated terms yields the following element equation:

$$
\begin{align*}
& \frac{L}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
\dot{y}_{1} \\
\dot{y}_{2}
\end{array}\right\}+1 / 6\left[\begin{array}{l}
\left(2 y_{1}+y_{2}\right)\left(v_{2}-v_{1}\right) \\
\left(y_{1}+2 y_{2}\right)\left(v_{2}-v_{1}\right)
\end{array}\right]+1 / 6\left[\begin{array}{l}
\left(2 v_{1}+v_{2}\right)\left(y_{2}-y_{1}\right) \\
\left(v_{1}+2 v_{2}\right)\left(y_{2}-y_{1}\right)
\end{array}\right] \\
&-\frac{9 L}{2}\left\{\begin{array}{l}
7\} \\
7
\end{array}\right\}=0 \tag{4.10}
\end{align*}
$$

Multiplying the Equation (4.10) by a factor of 6 and adding up the two middle terms, we obtain:

$$
L\left[\begin{array}{ll}
2 & 1  \tag{4.11}\\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
\dot{y}_{1} \\
\dot{y}_{2}
\end{array}\right\}+\left[\begin{array}{cc}
v_{2}-4 v_{1} & 2 v_{2}+v_{1} \\
-v_{2}-2 v_{1} & 4 v_{2}-v_{1}
\end{array}\right]\left\{\begin{array}{l}
y_{1} \\
y_{2}
\end{array}\right\}-L q\left\{\begin{array}{l}
3 \\
3
\end{array}\right\}=0
$$

In a manner analogous to the above procedure, the momentum Equation (3.2) for an element can be derived as:

$$
\begin{align*}
& \frac{1}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
\dot{v}_{1} \\
v_{2}
\end{array}\right\}+1 / 12\left[\begin{array}{rr}
-2 v_{1}-v_{2} & -v_{1}-2 v_{2} \\
2 v_{1}+v_{2} & v_{1}+2 v_{2}
\end{array}\right]\left\{\begin{array}{l}
v_{1} \\
v_{2}
\end{array}\right\}+9 / 2\left[\begin{array}{cc}
-1 & 1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
y_{1} \\
y_{2}
\end{array}\right\} \\
& +\frac{L q}{2}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
(v / y)_{1} \\
(v / y)_{2}
\end{array}\right\}+\frac{g L}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
s_{f 1} \\
s_{f}
\end{array}\right\}-\frac{g s_{0} L}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}=0 \tag{4.12}
\end{align*}
$$

## Assembly of Element Equations

The element properties originally expressed in local coordinates need to be transformed into global coordinates before solution algorithm is initiated. Based on the node-to-node relationship (Figure 3c), it is possible to generate an overall element property matrix for the entire domain, a process called assembling of element equations.

The concept of discretization employed earlier is based on the fact that a domain with varying geometric and hydraulic properties can be treated independently as subdomains but systematically from one subdomain to another. Assuming that the elements are of variable lengths and that
there are $N$ nodes, the assembled global matrix equation for the continuity Equation (4.11) becomes:

$$
\begin{aligned}
& {\left[\begin{array}{ccccccc}
2 L_{1} & L_{1} & 0 & 0 & 0 & 0 & 0 \\
L_{1} & 2\left(L_{1}+L_{2}\right) & L_{2} & 0 & 0 & 0 & 0 \\
0 & L_{2} & 2\left(L_{2}+L_{3}\right) & L_{3} & 0 & 0 & 0 \\
- & -- & - & -- & -- & -- & -- \\
0 & 0 & 0 & L_{i} & 2\left(L_{i}+L_{i+1}\right) & L_{i+1} & 0 \\
-- & -- & -- & - & -- & -- & -- \\
0 & 0 & 0 & 0 & 0 & L_{N-1} L_{N-1}
\end{array}\right]\left\{\begin{array}{c}
\dot{y}_{1} \\
\dot{y}_{2} \\
\dot{y}_{3} \\
-- \\
\dot{y}_{i} \\
- \\
\dot{y}_{N}
\end{array}\right\}+}
\end{aligned}
$$

$$
\begin{align*}
& \left\{\begin{array}{c}
q_{1} L_{1} \\
q_{1} L_{1}+q_{2} L_{2} \\
q_{2} L_{2}+q_{3} L_{3} \\
-\cdots-\cdots \\
q_{i} L_{i}+q_{i+1} L_{i+1} \\
--\cdots-\cdots \\
q_{N-1} L_{N-1}
\end{array}\right\}=0 \tag{4.13}
\end{align*}
$$

The general form of the above assembled global continuity equation can be expressed as:
$[A]\{\dot{y}\}+[B]\{y\}-\{C\}=0$
where
$A, B$ are matrices and $C$ is a column vector;
$\dot{y}$ is the time derivative;
$y$ and $v$ are dependent variables.
The momentum equation follows the same pattern of assembly.

Transient Solution Approach

The solution to the time-dependent global matrix Equation (4.13) is sought through a "semi-discrete" approach. This approach requires the time derivative of the dependent variable at each node to be replaced by a finite difference scheme in time domain. A simple illustration of the semi-discrete approach can be demonstrated by considering Equation (4.14). The time derivative, $\dot{y}$, will be replaced by a finite difference scheme, such as the forward, backward, and central difference. These are respectively given below:

Forward Difference $\quad \dot{y}=\frac{y^{K+1}-y^{K}}{\Delta t}$
Backward Difference $\quad \dot{y}=\frac{y^{K}-y^{K-1}}{\Delta t}$
Central Difference $\quad \dot{y}=\frac{y^{K+1}-y^{K-1}}{2 \Delta t}$
where:
$K=$ time level.
Substitution of Equation (4.15a) into Equation (4.14) yields
$[A]\left\{\frac{y^{K+1}-y^{K}}{\Delta t}\right\}+[B]\left\{y^{K}\right\}-\{C\}=0$

An implicit equation can be generated from Equation (4.16) with the aid of the time weighting factor. This subject is discussed elegantly in Chapter V.

## CHAPTER V

NUMERICAL FLOW MODELS

Introduction

Three distinct deterministic streamflow routing models are investigated and are discussed in this chapter in their order of increasing complexity: (1) The kinematic flow model comprises (a) the simplified version of the momentum Equation (3.2) that neglects pressure and inertia terms as compared to friction and gravity terms (see Equation 3.7) and (b) the complete form of continuity Equation (3.1); (2) the diffusion flow model combines (c) the simplified momentum equation that accounts only for pressure, friction, and gravity terms, Equation (3.8), and (d) the complete form of continuity Equation (3.1); and (3) the complete flow model comprises the complete forms of both continuity Equation (3.1) and momentum Equation (3.2).

The kinematic flow model is investigated in both an explicit and implicit sense. The explicit kinematic flow model leads to lifnearequations. They are solved using a direct method similar to the tridiagonal matrix algorithm set-up by Varga (1962). Solution proceeds by matrix reduction similar to Gaussian elimination. In contrast to the explicit model, the weighted implicit kinematic model yields a set of non-linear tridiagonal matrix equations which are solved by the functional Newton-Raphson iterative method. This method is known as implicit because the set of equations are solved by an indirect metnod.

The diffusion flow model, as well as the complete flow model each results in a non-linear bi-tridiagonal matrix equation: The functional Newton-Raphson's method, along with the direct solution algorithm*, triangular decomposition technique that yields a recursion algorithm (Douglas etal., 1959; Von Resenberg, 1975), is utilized to predict depth and velocity of flow for each model.

Explicit Kinematic Finite Element Model, EKFEM

The non-linear continuity Equation (3.1) is easily converted to linear form by use of geometric and flow relations:

$$
\begin{equation*}
\frac{\partial A}{\partial t}+\frac{\partial Q}{\partial x}-q(x, t)=0 \tag{5.1a}
\end{equation*}
$$

where

$$
\begin{aligned}
& A=\text { area of flow }, \mathrm{ft}^{2} ; \\
& Q=\text { volumetric flow rate }, \mathrm{ft}^{3} / \mathrm{sec} .
\end{aligned}
$$

The appropriate simplified momentum equation for coupling with the continuity Equation (5.1a) has been obtained and is presented below:

$$
S_{f}=S_{o}=n 7^{2} v^{2} / 2.22 R^{4 / 3}
$$

or

$$
Q=\frac{1.486}{n T} A R^{2 / 3} S_{0} 1 / 3
$$

Applying the Galerkin's weighted residual method to Equation (5.1a) results in the following linear first order ordinary differential equation (see Equation 4.11):

[^3]\[

\frac{L}{6}\left[$$
\begin{array}{ll}
2 & 1  \tag{5.2}\\
1 & 2
\end{array}
$$\right]\left\{$$
\begin{array}{l}
\dot{A}_{1} \\
\dot{A}_{2}
\end{array}
$$\right\}+1 / 2\left[$$
\begin{array}{ll}
-1 & 1 \\
-1 & 1
\end{array}
$$\right]\left\{$$
\begin{array}{l}
Q_{1} \\
Q_{2}
\end{array}
$$\right\}-q L / 2\left\{$$
\begin{array}{l}
1 \\
1
\end{array}
$$\right\}=0
\]

For the entire channel reach the assembled matrix equation becomes:


$$
\left\{\begin{array}{l}
Q_{2}-Q_{1}  \tag{5.3}\\
Q_{3}-Q_{1} \\
Q_{4}-Q_{2} \\
\cdots-\cdots \\
Q_{i+1}-Q_{i-1} \\
\cdots-\cdots- \\
a_{N}-Q_{N-1}
\end{array}\right\}
$$

$$
-1 / 2\left\{\begin{array}{l}
L_{1} q_{1} \\
L_{1} q_{1}+L_{2} q_{2} \\
L_{2} q_{2}+L_{3} q_{3} \\
\cdots \cdots-\cdots \\
L_{i} q_{i}+L_{i+1} q_{i+1} \\
\cdots \cdots \cdots \cdots \\
L_{N-1} a_{N-1}
\end{array}\right\}=0
$$

Equation (5.3) is equivalently expressed in a matrix form:

$$
\begin{equation*}
[K]\{\dot{A}\}+\{D\}-\{F\}=0 \tag{5.4}
\end{equation*}
$$

The time solution of Equation (5.4) is possible upon implementation of the forward differencing in time domain.

$$
\begin{equation*}
[K]\{A\}^{n+1}=[K]\{A\}^{n}+\Delta t\{F\}^{n}-\Delta t\{D\}^{n} \tag{5.5}
\end{equation*}
$$

The solution of the area of flow at various nodes proceeds forward in time with the right hand side evaluated at a previous time level, n. Thus, the Equation (5.5) can be expressed in a more compact form:

$$
\begin{equation*}
[K]\{A\}^{n+1}=\{X\}^{n} \tag{5.6}
\end{equation*}
$$

where

$$
X=\text { known column vector at previous time level. }
$$

The matrix, $K$, is a linear and tridiagonal type that easily leads to a direct solution algorithm. The computer program solving Equation (5.6) is facilitated by the use of the compact tridiagonal algorithm proposed by Varga (1962). The computed area of flow at current time level, $n+1$, is used to update the volumetric flow rate, Q, Equation (5.1b). The solution cycle is repeated as new time level is reached. The coded explicit finite element scheme exhibits dynamic stability due to restrictions on time step. This drawback inherent in explicit numerical schemes is expected regardless of the finite element approach. However, the stability problem is corrected in the weighted implicit flow model.

```
Weighted Implicit Kinematic Finite Element Model, WIKFEM
```

The implicit kinematic flow model begins by combining the non-linear continuity matrix Equation (4.14) with the modified momentum Equation (5.8) (with velocity the subject of the formulation rather than the volumetric flow rate). The introduction of the dimensionless time weighting factor, $\theta$, Figure 1, and the forward differencing to Equation (4.14) yeilds the following:

$$
\begin{equation*}
[A+\Delta t \theta B]\{y\}^{n+1}-\theta \Delta t\{C\}^{n+1}=[A+\Delta t(1-\theta) B]\{y\}^{n}+\Delta t(1-\theta)\{C\}^{n} \tag{5.7}
\end{equation*}
$$

And the modified momentum equation is repeated here for convenience as:

$$
\begin{equation*}
v=\frac{1.486}{n 1} R^{2 / 3} S_{0} 1 / 2 \tag{5.8}
\end{equation*}
$$

where all terms are as previously defined for Equations (3.1), (3.2), and (4.14).

The expanded form of Equation (5.7) for the upstream, interior, and downstream nodes, respectively, are given below:

$$
\begin{align*}
F_{1} \equiv & {\left[\left(2 L_{1}+\Delta t \theta\left(v_{2}-4 v_{1}\right)\right) y_{1}+\left(L_{1}+\Delta t \theta\left(2 v_{2}+v_{1}\right)\right) y_{2}\right]^{n+1}-\left[3 \Delta t \theta L_{1} q_{1}\right]^{n+1} } \\
& +\left[\left(-2 L_{1}+\Delta t(1-\theta)\left(v_{2}-4 v_{1}\right)\right) y_{1}+\left(-L_{1}+\Delta t(1-\theta)\left(2 v_{2}+v_{1}\right)\right) y_{2}\right]^{n} \\
& -\left[3 \Delta t(1-\theta) L_{1} q_{1}\right]^{n}=0  \tag{5.9a}\\
F_{i} \equiv & {\left[\left(L_{i-1}+\Delta t \theta\left(-v_{i}-2 v_{i-1}\right)\right) y_{i-1}+\left(2\left(L_{i-1}+L_{i}\right)+\Delta t \theta\left(v_{i+1}-v_{i-1}\right)\right) y_{i}\right.} \\
& \left.+\left(L_{i}+\Delta t \theta\left(2 v_{i+1}+v_{i}\right)\right) y_{i+1}\right]^{n+1}-3 \Delta t \theta\left[q_{i-1} L_{i-1}+q_{i} L_{i}\right]^{n+1} \\
& +\left[\left(-L_{i-1}+\Delta t(1-\theta)\left(-v_{i}-2 v_{i-1}\right)\right) y_{i-1}+\left(-2\left(L_{i-1}+L_{i}\right)+\Delta t(1-\theta)\left(v_{i+1}^{-v_{i-1}}\right)\right) y_{i}\right. \\
& \left.+\left(-L_{i}+\Delta t(1-\theta)\left(2 v_{i+1}+v_{i}\right)\right) y_{i+1}\right]^{n}-3 \Delta t(1-\theta)\left[q_{i-1} L_{i-1}+q_{i} L_{i}\right]^{n}=0 \\
F_{N} \equiv & {\left[\left(L_{N-1}-\Delta t \theta\left(v_{N}+2 v_{N-1}\right)\right) y_{N-1}+\left(2 L_{N-1}+\Delta t \theta\left(4 v_{N}-v_{N-1}\right)\right) y_{N}\right]^{n+1} }  \tag{5.9b}\\
& -\left[3 \Delta t \theta L_{N-1} q_{N-1}\right]^{n+1}+\left[\left(-L_{N-1}+\Delta t(1-\theta)\left(-v_{N}-2 v_{N-1}\right)\right) y_{N-1}\right. \\
& \left.+\left(-2 L_{N-1}+\Delta t(1-\theta)\left(4 v_{N}-v_{N-1}\right)\right) y_{N}\right]^{n}-\left[3 \Delta t(1-\theta) L_{N-1} q_{N-1}\right]^{n}=0 \quad \text { (5.9c) } \tag{5.9c}
\end{align*}
$$

The solution of Equation (5.9) is obtained through the generalized functional iterative method known as the Newton-Raphson method, first used by Amein and Fang (1969) and later by Fread (1971, 1976). Equation (5.9) expressed in functional form is as follows:

$$
\begin{align*}
& F_{1}\left(y_{1}, y_{2}\right)=0  \tag{5.10a}\\
& F_{i}\left(y_{i-1}, y_{i}, y_{i+1}\right)=0  \tag{5.10b}\\
& F_{N}\left(y_{N-1}, y_{N}\right)=0 \tag{5.10c}
\end{align*}
$$

The $y$ terms in the parentheses are the depth of flow, the dependent variable to be solved. The subscript associated with each $y$ denotes the nodal location. The computed values of $y$ are utilized in Equation (5.8) to generate the corresponding values of velocity of flow, $v$. For the system of $N$ non-linear equations with $N$ unknowns, computation is initiated by assigning trial values to the $N$ unknowns. The substitution of the trial values into the system of non-linear equations yields a set of $N$ residuals. In fact, the residual is the value of the right-hand side of the equation after the trial values are substituted in Equation (5.10). The final solution is obtained when the residuals are reduced to a suitable tolerance level.

If it is assumed that the computations have been carried through the $j^{\text {th }}$ iteration, in other words the values of the unknowns have been approximated through the $j^{\text {th }}$ iteration, then is possible to estimate the value of the residual as follows:

$$
\begin{align*}
& F_{1}\left(y_{1}^{j}, y_{2}^{j}\right)=R_{1}^{j}  \tag{5.11a}\\
& F_{i}\left(y_{i-1}^{j}, y_{i}^{j}, y_{i+1}^{j}\right)=R_{i}^{j}  \tag{5.11b}\\
& F_{N}\left(y_{N-1}^{j}, y_{N}^{j}\right)=R_{N}^{j} \tag{5.11c}
\end{align*}
$$

where $R_{i}{ }^{j}$ is the residual at the $j^{\text {th }}$ iteration cycle for the $i^{\text {th }}$ node. The Newton-Raphson algorithm ties up the residual and partial derivatives of the system of Equations (5.11) in the following manner:
where $\Delta y=y^{j+1}-y^{j}$ (the difference between current and previous iterates of $y$ ).

The matrix to the left-hand side of Equation (5.12) is the tridiagonal Jacobian matrix of size $(N \times N)$. It is possible to store the matrix in a compact ( $\mathrm{N} \times 3$ ) form as shown in Equation (5.13) following:

$$
\left[\begin{array}{lll}
O & \frac{\partial F_{1}}{\partial y_{1}} & \frac{\partial F_{1}}{\partial y_{2}}  \tag{5.13}\\
\frac{\partial F_{2}}{\partial y_{1}} & \frac{\partial F_{2}}{\partial y_{2}} & \frac{\partial F_{2}}{\partial y_{3}} \\
\frac{\partial F_{3}}{\partial y_{2}} & \frac{\partial F_{3}}{\partial y_{3}} & \frac{\partial F_{3}}{\partial y_{4}} \\
-\frac{\partial F_{i}}{\partial y_{i-1}} & \frac{\partial F_{i}}{\partial y_{i}} & \frac{\partial F_{i}}{\partial y_{i+1}} \\
-\frac{\partial F_{N}}{\partial y_{N-1}} & \frac{\partial F_{N}}{\partial y_{N}} & -\cdots- \\
\Delta y_{1} \\
\Delta y_{3} \\
-\overline{y_{i}} \\
\Delta y_{i} \\
\Delta y_{i+1} \\
-- \\
\Delta y_{N-1} \\
\Delta y_{N}
\end{array}\right\}=\left\{\begin{array}{c}
R_{1} \\
R_{2} \\
R_{3} \\
--- \\
R_{i} \\
-- \\
R_{N}
\end{array}\right\}
$$

The right-hand side of Equation (5.13) comprises the column vector generated upon substitution of the trial values of the unknowns into Equation (5.9). The individual terms of the Jacobian are generated from Equation (5.9) and written below as:

$$
\begin{align*}
& \frac{\partial F_{1}}{\partial y_{1}}=2 L_{1}+\operatorname{FAC}\left(v_{2}-4 v_{1}\right) \\
& \frac{\partial F_{1}}{\partial y_{2}}=L_{1}+\operatorname{FAC}\left(2 v_{2}+v_{1}\right) \\
& \frac{\partial F_{i}}{\partial y_{i-1}}=L_{i-1}+\operatorname{FAC}\left(-v_{i}-2 v_{i-1}\right) \\
& \frac{\partial F_{i}}{\partial y_{i}}=2\left(L_{i-1}+L_{i}\right)+F A C\left(v_{i+1}-v_{i}\right)  \tag{5.14}\\
& \frac{\partial F_{i}}{\partial y_{i+1}}=L_{i}+\operatorname{FAC}\left(2 v_{i+1}+v_{i}\right) \\
& \frac{\partial F_{N}}{\partial y_{N-1}}=L_{N-1}+\operatorname{FAC}\left(-v_{N}-2 v_{N-1}\right) \\
& \frac{\partial F_{N}}{\partial y_{N}}=2 L_{N-1}+\operatorname{FAC}\left(4 v_{N}-v_{N-1}\right)
\end{align*}
$$

where $F A C=\theta \Delta t$.
Equation (5.13) is a linearized form of the non-linear weighted implicit kinematic model similar to Equation (5.6) and is solved in the same manner. The computer program for Equation (5.13) does not require the Jacobian matrix to be up-dated for every iteration, rather after every three iterations. The approach seems reasonable in terms of minimizing the computer time because convergence is achieved with relatively few
iterations for most time steps employed. For the guess values of the dependent variable, $y$, required to initiate the iterative Newton-Raphson equation solver, the initial uniform flow depths are utilized. The initial depths of flow prior to the flood into the channel are the best guess to use. Proper upstream and downstream boundary conditions, such as discharge hydrograph and loop-rating curve, are incorporated in the model. The solution is then sought for a prescribed convergence error criterion.

The effectiveness of the weighted implicit model as compared to the explicit version, along with the other two flow models, is discussed in Chapter VI.

Weighted Implicit Diffusion Finite<br>Element Model, WIDFEM

Another simplified model is the diffusion flow model. The model is developed by coupling the continuity Equation (3.1) and the simplified momentum Equation (3.8). The finite element transformation procedure for Equation(3.1) is given in Equations:(4.8) through (4.14). The same principles are applied to Equation (3.8), resulting in the following element equation:

$$
1 / 2\left[\begin{array}{ll}
-1 & 1  \tag{5.15}\\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
y_{1} \\
y_{2}
\end{array}\right\}+\frac{L}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\left\{\begin{array}{l}
s_{f 1} \\
s_{f 2}
\end{array}\right\}-\frac{L}{2} s_{0}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}=0
$$

where

$$
\begin{aligned}
& S_{f}=F R v^{2} / R^{4 / 3} \text { and } \\
& F R=n 7^{2} / 2.2082 .
\end{aligned}
$$

The assembled matrix equation of Equation (5.15), along with the dimensionless time weighting factor. and the forward time differencing, becomes:

$$
\begin{align*}
& \Delta t \theta[F]\left\{y^{n+1}\right\}+\Delta t \theta \frac{F R}{\left(R^{n+1}\right)^{4 / 3}}[D]\left\{v^{n+1}\right\}^{2}-\Delta t \theta\left\{M^{n+1}\right\} \\
& =-\Delta t(1-\theta)[F]\left\{y^{n}\right\}-\Delta t(1-\theta) \frac{F R}{\left(R^{n}\right)^{4 / 3}}[D]\left\{v^{n}\right\}^{2}+t(1-)\left\{M^{n}\right\} \tag{5.16}
\end{align*}
$$

where the superscripts $(n+1)$ and ( $n$ ) for the variables $y, v$, and $R$ are the current and previous iterations, respectively. The expanded forms of Equation (5.16) similar to Equation (5.9) for the upstream, interior, and downstream nodes are respectively given as:

$$
\begin{align*}
G_{1} & \equiv L_{1}(F R) \Delta t \theta\left[2 v_{1}{ }^{2} / R_{1}{ }^{4 / 3}+v_{2}{ }^{2} / R_{2}^{4 / 3}\right]^{n+1}+3 \Delta t \theta\left[y_{2}-y_{1}\right]^{n+1} \\
& -3 \Delta t \theta\left[S_{0} L_{1}\right]^{n+1}+L_{1}(F R) \Delta t(1-\theta)\left[2 v_{1}{ }^{2} / R_{1} 4 / 3+v_{2}{ }^{2} / R_{2}^{4 / 3}\right]^{n} \\
& +3 \Delta t(1-\theta)\left[y_{2}-y_{1}\right]^{n}-3 \Delta t(1-\theta)\left[S_{0} L_{1}\right]^{n}  \tag{5.17a}\\
G_{i} & \equiv \Delta t \theta(F R)\left[L_{i-1} v_{i-1}{ }^{2} / R_{i-1}^{4 / 3}+2\left(L_{i-1}+L_{i}\right) v_{i}^{2} / R_{i}{ }^{4 / 3}+L_{i} v_{i+1}^{2} / R_{i+1}^{4 / 3}\right]^{n+1} \\
& +3 \Delta t \theta\left[y_{i+1}-y_{i-1}\right]^{n+1}-3 \Delta t \theta\left[\left(L_{i-1}+L_{i}\right) S_{0}\right]^{n+1} \\
& +\Delta t(1-\theta)(F R)\left[L_{i-1} v_{i-1}{ }^{2} / R_{i-1}^{4 / 3}+2\left(L_{i-1}+L_{i}\right) v_{i}^{2} / R_{i}^{4 / 3}+L_{i} v_{i+1}{ }^{2} / R_{i+1}^{4 / 3}\right]^{n} \\
& +3 \Delta t(1-\theta)\left[y_{i+1}-y_{i-1}\right]^{n}-3 \Delta t(1-\theta)\left[\left(L_{i-1}+L_{i}\right) S_{0}\right]^{n}  \tag{5.17b}\\
G_{N} & \equiv L_{N-1}(F R) \Delta t \theta\left[v_{N-1}^{2} / R_{N-1}^{4 / 3}+2 v_{N}{ }^{2} / R_{N}{ }^{4 / 3}\right]^{n+1}+3 \Delta t \theta\left[y_{N}-y_{N-1}\right]^{n+1} \\
& -3 \Delta t \theta\left[S_{0} L_{N-1}\right]^{n+1}+L_{N-1}(F R) \Delta t(1-\theta)\left[v_{N-1}^{2} / R_{N-1}^{4 / 3}\right]^{n}+3 \Delta t(1-\theta)\left[y_{N}-y_{N-1}\right]^{n} \\
& -3 \Delta t(1-\theta)\left[S_{0} L_{N-1}\right]^{n} \tag{5.17c}
\end{align*}
$$

The simultaneous solution of Equations (5.9) and (15.17) is possible using the generalized functional iterative method, known as the NewtonRaphson method discussed earlier in the implicit kinematic flow model. The functional representations of Equations (5.9) and (15.17) are as follows:
upstream nodes

$$
F_{1}\left(y_{1}, v_{1}, y_{2}, v_{2}\right)=0
$$

$$
G_{1}\left(y_{1}, v_{1}, y_{2}, v_{2}\right)=0
$$

$$
F_{i}\left(y_{i-1}, v_{i-1}, y_{i}, v_{i}, y_{i+1}, v_{i+1}\right)=0
$$ nodes

$$
G_{i}\left(y_{i-1}, v_{i-1}, y_{i}, v_{i}, y_{i+1}, v_{i+1}\right)=0
$$

$G_{i}\left(y_{i-1}, v_{i-1}, y_{i}, v_{i}, y_{i+1}, v_{i+1}\right)=0$
$F_{N}\left(y_{N-1}, v_{N-1}, y_{N}, v_{N}\right)=0$
downstream nodes

$$
G_{N}\left(y_{N-1}, v_{N-1}, y_{N}, v_{N}\right)=0
$$

Similar to the implicit kinematic model, the substitution of the trial values for $v$ and $y$ into the system of non-linear Equations (5.18) yields a set of 2 N residuals. Furthermore, the computations are carried through the $j^{\text {th }}$ iteration cycle; then the estimates of the residuals are as follows:

$$
\begin{align*}
& F_{1}\left(y_{1}^{j}, v_{1}^{j}, y_{2}^{j}, v_{2}^{j}\right)=R_{F_{1}}^{j} \\
& G_{1}\left(y_{1}^{j}, v_{1}^{j}, y_{2}^{j}, v_{2}^{j}\right)=R_{G_{1}}^{j} \\
& F_{i}\left(y_{i-1}^{j}, v_{i-1}^{j}, y_{i}^{j}, v_{i}^{j}, y_{i+1}^{j}, v_{i+1}^{j}\right)=R_{F_{i}}^{j} \\
& G_{i}\left(y_{i-1}^{j}, v_{i-1}^{j}, y_{i}^{j}, v_{i}^{j}, y_{i+1}^{j}, v_{i+1}^{j}\right) R_{G_{i}}^{j}  \tag{5.19}\\
& F_{N}\left(y_{N-1}^{j}, v_{N-1}^{j}, y_{N}^{j}, v_{N}^{j}\right)=R_{F_{N}}^{j} \\
& G_{N}\left(y_{N-1}^{j}, v_{N-1}^{j}, y_{N}^{j}, v_{N}^{j}\right)=R_{G_{N}}^{j}
\end{align*}
$$

Where $R_{F_{i}}^{j}$ and $R_{G_{i}}^{j}$ are the residuals at the $j^{\text {th }}$ interation cycle for the continuity and momentum equations, respectively, at $i^{\text {th }}$ node.

The Newton-Raphson algorithm couples the residuals and the partial derivatives of the systems of Equation (15.19) in the following form:
where $\Delta y=y^{j+1}-y^{j}$ and $\Delta v=v^{j+1}-v^{j}$.
The matrix to the left-hand side of Equation (5.20) containing the partial derivatives of the functions $F$ and $G$ is the bi-tridiagonal Jacobian matrix of size ( $2 \mathrm{~N} x \geq 2 \mathrm{~N}$ ). The maximum non-zero elements in any single row is six. Thus, the Jacobian matrix is stored in a compact ( $2 \mathrm{~N} \times 6$ ) matrix. The individual terms of the Jacobian matrix for the function $G$ are given below whereas those of F are noted in Equation (5.14).

$$
\begin{align*}
& \frac{\partial G_{1}}{\partial y_{1}}=-F A C\left[\frac{8 L_{1}(F R) v_{1}{ }^{2}}{3 R_{1}^{7 / 3}} \frac{\partial R_{1}}{\partial y_{1}}+3.0\right] \\
& \frac{\partial G_{1}}{\partial y_{2}}=4 L_{1}(F R)(F A C) v_{1} / R_{1}^{4 / 3}  \tag{5.21a}\\
& \frac{\partial G_{1}}{\partial y_{2}}=-F A C\left[\frac{4 L_{1}(F R) v_{2}^{2}}{3 R_{2}^{7 / 3}} \frac{\partial R_{2}}{\partial y_{2}}-3.0\right] \\
& \frac{\partial G_{1}}{\partial v_{2}}=2 L_{1}(F R)(F A C) v_{2} / R_{2}^{4 / 3} \\
& \frac{\partial G_{i}}{\partial y_{i-1}}=-F A C\left[\frac{-4 L_{i-1}(F R) v_{i-1}^{2}}{3 R_{i-1}^{7 / 3}} \frac{\partial R_{i-1}}{\partial y_{i-1}}+3.0\right] \\
& \frac{\partial G_{i}}{\partial v_{i-1}}=2 L_{i-1}(F R)(F A C) v_{i-1} / R_{i-1}^{4 / 3} \\
& \frac{\partial G_{i}}{\partial y_{i}}=-F A C\left[\frac{8\left(L_{i-1}+L_{i}\right)(F R) v_{i}^{2}}{3 R_{i} 7 / 3} \frac{\partial R_{i}}{\partial y_{i}}\right]  \tag{5.21b}\\
& \frac{\partial G_{i}}{\partial v_{i}}=4\left(L_{i-1}+L_{i}\right)(F R)(F A C) v_{i} / R_{i}^{4 / 3} \\
& \frac{\partial G_{i}}{\partial y_{i+1}}=-F A C\left[\frac{4 L_{i}(F R) v_{i+1}^{2}}{3 R_{i+1}^{7 / 3}} \frac{\partial R_{i+1}}{\partial y_{i+1}}-3.0\right] \\
& \frac{\partial G_{i}}{\partial v_{i}+1}=2 L_{i}(F R)(F A C) v_{i+1} / R_{i+1}^{4 / 3} \\
& \frac{\partial G_{N}}{\partial y_{N-1}}=-F A C\left[\frac{4 L_{N-1}(F R) v_{N-1}^{2}}{3 R_{N-1}^{7 / 3}} \frac{\partial R_{N-1}}{\partial y_{N-1}}+3.0\right] \tag{5.21c}
\end{align*}
$$

$$
\begin{aligned}
& \frac{\partial G_{N}}{\partial v_{N-1}}=2 L_{N-1}(F R)(F A C) v_{N-1} / R_{N-1}^{4 / 3} \\
& \frac{\partial G_{N}}{\partial y_{N}}=-F A C\left[\frac{8 L_{N-1}(F R) v_{N}^{2}}{3 R_{N-1}^{7 / 3}} \frac{\partial R_{N}}{\partial y_{N}}-3.0\right] \\
& \frac{\partial G_{N}}{\partial v_{N}}=4 L_{N-1}(F R)(F A C) v_{N} / R_{N}^{4 / 3}
\end{aligned}
$$

where

$$
\begin{aligned}
& F A C=\theta \Delta t ; \\
& \frac{\partial R}{\partial y}=\text { rate of change of hydraulic radius with depth. }
\end{aligned}
$$

For a natural channe1, $R=A / P$ (area, A, divided by the wetted perimeter, P), then

$$
\partial R / \partial y=\frac{\left(P \frac{\partial A}{\partial X} y^{-}-\frac{\partial P}{\partial X}\right)_{y}}{P^{2}}
$$

For a rectangular channel, $R=B y / B+2 y$, and

$$
\frac{\partial R}{\partial y}=\frac{B^{2}}{(B+2 y)^{2}}
$$

The solution af Equation (5.20) is initiated iteratively by evaluating the right hand column vector and the Jacobian matrix using the prevous nodal values of depth, $y$, and velocity, $v$. Good starting values for nodal depths and velocities are those of the uniform flow before the flood wave arrives at the upstream section of the channel.

At the upstream boundary node, if the flood discharge hydrograph is imposed as a known condition, then the corresponding upstream velocity, $v_{1}$, at any time level is evaluated as:

$$
\begin{equation*}
v_{1}=Q_{1} / B y_{1} \tag{5.22}
\end{equation*}
$$

On the other hand, the simplified momentum equation adopted for the diffusion model adequately describes the downstream boundary condition as a loop rating curve. Thus, no further modification is necessary.

## Weighted Implicit Complete Finite Element Mode1, WICFEM

Solution of the complete flow model follows the same basic steps as the implicit diffusion flow model. The significant difference between the two models resides in the total number of terms in the momentum Equation (3.2). For this reason, only the manipulation of the complete momentum equation deserves further discussion. The finite element transformed version of the complete momentum Equation (3.2) is presented in Chapter IV as Equation (4.12).

The assembled matrix equation (Equation 4.12), together with the dimensionless time weighting factor, $\theta$, and forward time differencing yields:

$$
\begin{align*}
& {[D+(\theta \Delta t) E]\left\{v^{n+1}\right\}+\theta \Delta t[F]\left\{y^{n+1}\right\}+\theta \Delta t[G]\left\{\frac{v^{n+1}}{y^{n+1}}\right\}+\theta \Delta t[D]\left\{S_{f}{ }^{n+1}\right\}} \\
& -\theta \Delta t\left\{M^{n+1}\right\}=[D-\Delta t(1-\theta) E]\left\{v^{n}\right\}-\Delta t(1-\theta)[F]\left\{y^{n}\right\}-\Delta t(1-\theta)[G]\left\{\frac{v^{n}}{y^{n}}\right\} \\
& -\Delta t(1-\theta)[D]\left\{S_{f} n^{n}\right\}+\Delta t(1-\theta)\left\{M^{n}\right\} \tag{5.23}
\end{align*}
$$

where
$D, E, F$ and $D=$ assembled matrices of Eq. 4.12 from left to right, respectively;
$M=$ assembled vector of Eq. 4.12.
Letting $S_{f}=F R v^{2} / R^{4 / 3}$, where $F R=g n 7^{2} / 2.2082$, then Equation (5.23) becomes:

$$
\begin{align*}
& {[D+(\theta \Delta t) E]\left\{v^{n+1}\right\}+\theta \Delta t[F]\left\{y^{n+1}\right\}+\theta \Delta t[G]\left\{\frac{v^{n+1}}{y^{n+1}}\right\}+\theta \Delta t\left(\frac{F R}{\left(R^{n+1}\right)^{4 / 3}}\right)[D]} \\
& \left\{v^{n+1}\right\}^{2}-\theta \Delta t\left\{M^{n+1}\right\}=[D-\Delta t(1-\theta) E]\left\{v^{n}\right\}-\Delta t(1-\theta)[F]\left\{y^{n}\right\}-\Delta t(1-\theta)[G] \\
& \left\{\frac{v^{n}}{y^{n}}\right\}-\Delta t(1-\theta) \frac{F R}{\left(R^{n}\right)^{4 / 3}}[D]\left\{v^{n}\right\}^{2}+\Delta t(1-\theta)\left\{M^{n}\right\} \tag{5.24}
\end{align*}
$$

The expanded forms of Equation (5.24) for the upstream, interior and downstream nodes are given as:

$$
\begin{align*}
& G_{1} \equiv\left[\left(4 L_{1}+\theta \Delta t\left(-2 v_{1}-v_{2}\right)\right) v_{1}+\left(2 L_{1}+\theta \Delta t\left(-v_{1}-2 v_{2}\right)\right) v_{2}\right]^{n+1} \\
& +2 \theta \Delta t\left[2 L_{1} q_{1} \frac{v_{1}}{y_{1}}+L_{1} q_{1} \frac{v_{2}}{y_{2}}\right]^{n+1}+2 \theta \Delta t\left[2 L_{1} \frac{F R}{R_{1}{ }^{4 / 3}} v_{1}^{2}+L_{1} \frac{F R}{R_{2}^{4 / 3}} v_{2}^{2}\right]^{n+1} \\
& +\theta \Delta t\left[6 g\left(y_{2}-y_{1}\right)\right]^{n+1}-\theta \Delta t\left[6 g S_{0} L_{1}\right]^{n+1}+\left[\left(-4 L+\Delta t(1-\theta)\left(-2 v_{1}-v_{2}\right)\right) v_{1}\right. \\
& \left.+\left(-2 L_{1}+\Delta t(1-\theta)\left(-v_{1}-2 v_{2}\right)\right) v_{2}\right]^{n}+2 \Delta t(1-\theta)\left[2 L_{1} q_{1} \frac{v_{1}}{y_{1}}+L_{1} q_{1} \frac{v_{2}}{y_{2}}\right] n \\
& +2 \Delta t(1-\theta)\left[2 L_{1} \frac{R F}{R_{1}} 4 / 3 v_{1}^{2}+L_{1} \frac{F R}{R_{2}^{4 / 3}} v_{2}^{2}\right]^{n}+\Delta t(1-\theta)\left[6 g\left(y_{2}-y_{1}\right)\right]^{n} \\
& -\Delta t(1-\theta)\left[6 g S_{0} L_{1}\right]^{n}=0 \tag{5.25a}
\end{align*}
$$

$$
\begin{aligned}
G_{i} \equiv & {\left[\left(2 L_{i-1}+\theta \Delta t\left(2 v_{i-1}+v_{i}\right)\right) v_{i-1}+\left(4\left(L_{i-1}+L_{i}\right)+\theta \Delta t\left(v_{i-1}-v_{i+1}\right)\right) v_{i}\right.} \\
& \left.+\left(2 L_{i}+\theta \Delta t\left(-v_{i}-2 v_{i+1}\right)\right) v_{i+1}\right]^{n+1}+2 \theta \Delta t\left[q_{i-1} L_{i-1} \frac{v_{i-1}}{y_{i-1}}+2\left(L_{i-1} q_{i-1}+L_{i} q_{j}\right)\right. \\
& \left.\frac{v_{i}}{y_{i}}+q_{i} L_{i} \frac{v_{i+1}}{y_{i+1}}\right]^{n+1}+2 \theta \Delta t\left[L_{i-1} \frac{F R}{R_{i-1}^{4 / 3}} v_{i-1}^{2}+2\left(L_{i-1}+L_{i}\right) \frac{F R}{R_{i}^{4 / 3}} v_{i}^{2}\right. \\
& \left.+L_{i} \frac{F R}{R_{i+1}^{4 / 3}} v_{i+1}^{2}\right]^{n+1}+\theta \Delta t\left[6 g\left(y_{i+1}-y_{i-1}\right)\right]^{n+1}-\theta \Delta t\left[6 g S_{0}\left(L_{i-1}+L_{i}\right)\right]^{n+1} \\
& +\left[\left(-2 L_{i-1}+\Delta t(1-\theta)\left(2 v_{i-1}+v_{i}\right)\right) v_{i-1}+\left(-4\left(L_{i-1}+L_{i}\right)+\Delta t(1-\theta)\left(v_{i-1}-v_{i+1}\right)\right) v_{i}\right.
\end{aligned}
$$

$$
\begin{align*}
& \left.+\left(-2 L_{i}+\Delta t(1-\theta)\left(-v_{i}-2 v_{i+1}\right)\right) v_{i+1}\right]^{n}+2 \Delta t(1-\theta)\left[q_{i-1} L_{i-1} \frac{v_{i-1}}{y_{i-1}}\right. \\
& \left.+2\left(L_{i-1} q_{i-1}+L_{i} q_{i}\right) \frac{v_{i}}{y_{i}}+q_{i} L_{i} \frac{v_{i+1}}{y_{i+1}}\right]^{n}+2 \Delta t(1-\theta)\left[L_{i-1} \frac{F R}{R_{i-1}^{4 / 3}} v_{i-1}^{2}\right. \\
& \left.+2\left(L_{i-1}+L_{i}\right) \frac{F R}{R_{i}^{4 / 3}} v_{i}^{2}+L_{i} \frac{F R}{R_{i}^{4 / 3}} v_{i+1}^{2}\right]^{n}+\Delta t(1-\theta)\left[6 g\left(y_{i+1}^{-y_{i-1}}\right)\right]^{n+1} \\
& -\Delta t(1-\theta)\left[6 g S_{0}\left(L_{i-1}+L_{i}\right)\right]^{n}=0  \tag{5.25b}\\
& G_{N} \quad\left[\left(2 L_{N-1}+\theta \Delta t\left(2 v_{N-1}+v_{N}\right)\right) v_{n-1}+\left(4 L_{N-1}+\theta \Delta t\left(v_{N-1}+2 v_{N}\right)\right) v_{N}\right]^{n+1} \\
& +2 \theta \Delta t\left[L_{N-1} a_{N-1} \frac{v_{N-1}}{y_{N-1}}+2 L_{N-1} a_{N-1} \frac{v_{N}}{y_{N}}\right] n+1+2 \theta \Delta t\left[L_{N-1} \frac{F R}{R_{N-1}^{4 / 3}} v_{N-1}^{2}\right. \\
& \left.\left.+2 L_{N-1} \frac{F R}{R_{N}^{4 / 3}} v_{N}\right]^{2}\right]^{n+1}+\theta \Delta t\left[6 g\left(y_{N}-y_{N-1}\right)\right]^{n+1}-\theta \Delta t\left[6 g S_{0} L_{N-1}\right]^{n+1} \\
& \div\left[\left(-2 L_{N-1}+\Delta t(1-\theta)\left(2 v_{N-1}+v_{N}\right)\right) v_{N-1}+\left(-4 L_{N-1}+\Delta t(1-\theta)\left(v_{N-1}+2 v_{N}\right)\right) v_{N}\right]^{n} \\
& +2 \Delta t(1-\theta)\left[L_{N-1} q_{N-1} \frac{v_{N-1}}{y_{N-1}}+2 L_{N-1} a_{N-1} \frac{v_{N}}{y_{N}}\right]^{n}+2 \Delta t(1-\theta)\left[L_{N-1} \frac{F R}{R_{N-1}^{4 / 3}} v_{N-1}^{2}\right. \\
& \left.+2 L_{N-1} \frac{F R}{R_{N}^{4 / 3}} v_{N}^{2}\right]^{n}+\Delta t(1-\theta)\left[6 g\left(y_{N}-y_{N-1}\right)\right]^{n}-\Delta t(1-\theta)\left[6 g S_{0} L_{N}-7\right]^{n} \tag{5.25c}
\end{align*}
$$

By replacing all the Jacobian terms associated with the momentum equation for the diffusion flow model in Equation (5.20) with those of the complete flow model, the solution thereafter follows the same routine. However, it is possible to modify the downstream boundary condition for the momentum equation similar to the diffusion flow model as an adequate loop rating curve, (Equation 15.17c). The upstream momentum Equation (5.25a) needs no modification if Equation (5.22) is employed to update the upstream velocity, $v_{1}$.

The required Jacobian-momentum terms to replace those in Equation (5.20) are as follows:

$$
\begin{align*}
& \frac{\partial G_{1}}{\partial y_{1}}=-F A C\left[\frac{16(F R) L_{1} v_{1}{ }^{2}}{3 R_{1}{ }^{7 / 3}} \frac{\partial R_{1}}{\partial y_{1}}+\frac{4 L_{1} q_{1} v}{y_{1}{ }^{2}}+6 g\right] \\
& \frac{\partial G_{1}}{\partial v_{1}}=4 L_{1}+F A C\left[\left(-4 v_{1}-2 v_{2}\right)+\frac{8 L_{1}(F R) v_{1}}{R_{1}{ }^{4 / 3}}+\frac{4 L_{1} q_{1}}{y_{1}}\right] \\
& \frac{\partial G_{1}}{\partial y_{2}}=-F A C\left[\frac{8(F R) L_{1} v_{2}{ }^{2}}{3 R_{2}{ }^{7 / 3}} \frac{\partial R_{2}}{\partial y_{2}}+\frac{2 L_{1} q_{1} v_{2}}{y_{2}{ }^{2}}+6 g\right]  \tag{5.26a}\\
& \frac{\partial G_{1}}{\partial v_{2}}=2 L_{1}+F A C\left[\left(-2 v_{1}-4 v_{2}\right)+\frac{4 L_{1}(F R) v_{2}}{R_{2}}+\frac{2 L_{1} q_{1}}{y_{2}}\right] \\
& -\frac{\partial G_{i}}{\partial y_{i-1}}=-F A C\left[\frac{8(F R) L_{i-1} v_{i-1}}{3 R_{1}{ }^{7 / 3}} \frac{\partial R_{i-1}}{\partial y_{i-1}}+\frac{2 L_{i-1} q_{i-1} v_{i-1}}{y_{i-1}^{2}}+6 g\right] \\
& \frac{\partial G_{i}}{\partial v_{i-1}}=2 L_{i-1}+F A C\left[\left(4 v_{i-1}+2 v_{i}\right)+\frac{4 L_{i-1}(F R) v_{i-1}}{R_{i-1}^{4 / 3}}+\frac{2 L_{i-1} q_{i-1}}{y_{i-1}}\right] \\
& \frac{\partial G_{i}}{\partial y_{i}}=-F A C\left[\frac{16(F R)\left(L_{i-1}+L_{i}\right) v_{i}{ }^{2}}{3 R_{i}^{7 / 3}} \frac{\partial R_{i}}{\partial y_{i}}+\frac{4\left(L_{i-1} q_{i-1}+L_{i} q_{i}\right) v_{i}}{y_{i}^{2}}\right] \tag{5.26b}
\end{align*}
$$

$$
\begin{align*}
& \frac{\partial G_{i}}{\partial v_{i}}=4\left(L_{i-1}+L_{i}\right)+\operatorname{FAC}\left[2\left(v_{i-1}-v_{i-1}\right)+\frac{8\left(L_{i-1}+L_{i}\right)(F R) v_{i}}{R_{i}^{4 / 3}}+\frac{4\left(L_{i-1} q_{i-1}+L_{i} q_{i}\right)}{y_{i}}\right] \\
& \frac{\partial G_{i}}{\partial y_{i+1}}=-\operatorname{FAC}\left[\frac{8(F R) L_{i} v_{i+1}}{3 R_{i+1}^{7 / 3}} \frac{\partial R_{i+1}}{\partial y_{i+1}}+\frac{2 L_{i} q_{i} v_{i+1}}{y_{i+1}^{2}}-6 g\right] \\
& \frac{\partial G_{i}}{\partial v_{i+1}}=2 L_{i}+F A C\left[\left(-2 v_{i}-4 v_{i+1}\right)+\frac{4 L_{i}(F R) v_{i+1}}{R_{i+1}^{4 / 3}}+\frac{2 L_{i} q_{i}}{y_{i+1}}\right] \\
& --\frac{\partial G_{N}}{\partial y_{N-1}}=-\operatorname{FAC}\left[\frac{8(F R) L_{N-1} v_{N-1}}{3 R_{N-1}^{7 / 3}} \frac{\partial R_{N-1}}{\partial y_{N-1}}+\frac{2 L_{N-1} q_{N-1} v_{N-1}}{2}+6 g\right] \tag{5.26c}
\end{align*}
$$

$$
\begin{aligned}
& \frac{\partial G_{N}}{\partial v_{N-1}}=2 L_{N-1}+F A C\left[\left(4 v_{N-1}+2 v_{N}\right)+\frac{4 L_{N-1}(F R) v_{N-1}}{R_{N-1}^{4 / 3}}+\frac{2 L_{N-1} q_{N-1}}{y_{N-1}}\right] \\
& \frac{\partial G_{N}}{\partial y_{N}}=-F A C\left[\frac{16(F R) L_{N-1} v_{N}{ }^{2}}{3 R_{N}{ }^{7 / 3}} \frac{\partial R_{N}}{\partial y_{N}{ }^{2}}+\frac{4 L_{N-1} q_{N-1} v_{N}}{y_{N}{ }^{2}}-6 g\right] \\
& \frac{\partial G_{N}}{\partial v_{N}}=4 L_{N-1}+F A C\left[\left(2 v_{N-1}+4 v_{N}\right)+\frac{8 L_{N-1}(F R) v_{N}}{R_{N}}+\frac{4 L_{N-1} q_{N-1}}{y_{N}}\right]
\end{aligned}
$$

Similar to the weighted implicit kinematic and diffusion flow models, the Jacobian terms of Equation (5.26) are up-dated after every three iterations. The numerical performance of the complete flow model and simplified models in predicting the depth of flow and velocity of flow is the subject of Chapters VI and VII.

## CHAPTER VI

## VERIFICATION OF MODELS

## Introduction

Although the use of numerical methods for unsteady flow investigations has increased tremendously in recent years, most of the investigations are still exploratory, and serious attempts at making them accessible to the users in the field have not yet been made. The responsibility for developing an efficient numerical model that needs minimal or no modification, except the insertion of input data, is partly the objective of this study. However, further testing of the models with different problems and with a variety of boundary conditions is necessary for general use.

The performance of each model, particularly the weighted implicit diffusion and complete flow models in predicting flows in a natural channel, was assessed by comparing simulated and observed hydrographs. Possible discrepancies between simulated and observed flows are attributable to the following sources: errors in field measurements of the flows, survey errors in the measurement of channel sections, errors in estimating resistance coefficients, and, most importantly, changes in the channel properties before and during the unsteady flow event (Amein and Fang, 1969). Other sources of errors are associated with the numerical method itself, namely: finite element approximation errors, temporal approximation errors, and errors due to any iterative non-linear
equation solver (Chung, 1978). To eliminate the contribution of the first kind of sources of errors, it was necessary to verify the models in two parts.

The first part involves simulation of flow in an idealized channel of rectangular geometry. Simulated hydrographs were compared with similar results from those predicted using an explicit finite difference scheme (Viessman et al., 1972). This approach helps to explore the basic principles in the numerical development of the individual models. As a result, any discrepancies observable on application to natural channels should not be all blamed on the mathematical model development. The second part involves simulation of a flood in a natural channel, the Illinois River located in Oklahoma. Limited data in other major streams in Oklahoma, mostly cross sections, roughness coefficients and recorded floods made the Illinois River the best choice.

## Application to Idealized Channel

The computer programs of the explicit, weighted implicit kinematic, diffusion, and complete flow models have been written in FORTRAN IV for an IBM 360 model 75. The models were applied separately to simulate the hypothetical flood in a rectangular channel presented by Viessman et al. (1972) using the explicit finite difference scheme. The example problem considers a $2-\mathrm{mile}$ long and $2-\mathrm{ft}$ wide rectangular channel having a depth of flow of 6 ft . It is subjected to an upstream increase in flow to 2000 cfs in a period of 20 minutes, and then it decreases uniformly to the initial depth of flow in an additional period of 40 minutes. The channel has a bottom slope of $0.0015 \mathrm{ft} / \mathrm{ft}$ and an estimated Manning co-
efficient, $n$, of 0.02 .
Similar to the example of Viessman et al. (1972), a distance step of 528 ft was used in the simulation, although the four models can accept variable distance steps.. Also, the weighted implicit kinematic model has a built-in option to route the flood in a trapezoidal, triangular, or rectangular channel. For the first two geometries, the right- and left-side slopes captioned as ZRS and ZLS should have assigned values other than zeros, except for a rectangular channel. The triangular geometry will have zero width for input value.

## Hydrographs from EKFEM and WIKFEM

The flow hydrographs for upstream, midreach, and downstream sections are predicted by the explict and weighted implict kinematic models. These hydrographs are plotted along with those predicted by Viessman et al. (1972), shown in Figure 6. It should be noted that the explicit difference scheme of Viessman et al. (1972) solves the continuity and momentum equations completely. The kinematic flow models depict attenuations in the peak flows at midreach as well as the downstream section. This performance is acceptable since the longitudinal channel slope utilized for the simulation falls within 10 feet per mile ( $10.3 \%$ ) for which the use of kinematic approximation is justified. Details of the slope approximation for use of simplified models are discussed by Henderson (1966).

While the explicit kinematic finite element model is limited to a time step of 2 seconds because of stability considerations, the weighted implicit scheme appears to be unconditionally stable. The influence of the time weighting factor, $\theta$, on the numerical distortion (dispersion and


Figure 6. Comparison of Dishcarge Hydrographs for EKFEM and WIDFEM and Viessman's Solution.
attenuation of computed stage or discharge hydrographs) is shown in Figure 7. The plotted discharge hydrograph indicates that the lower range of the allowable $\theta$ values, such as 0.75 , as compared to the upper limiting value of 1.00 minimized the attenuation of the peak flow which results from the use of a large time step, $\Delta t$ of 300 seconds. This observation is not unique but confirms that of Fread (1973). The weighted implicit kinematic flow model was run for $\Delta t=180,300$, and 600 seconds with various values of the weighting factor $\theta$, such as $0.55,0.75$, and 1.0 , respectively. In all time steps, fastest convergence was obtained with $\theta=1.0$, and it is recommended for use with this routine. It is not surprise that $\theta=1.0$ affords rapid convergence because the scheme becomes fully implicit. However, instability results in WIKFEM with $\theta<0.55$. Thus, the allowable range of the time weighting factor, $\theta$, is $0.55 \leq \theta \leq 1.0$.

## Hydrographs from WIDFEM

Simulated discharge hydrographs for a time step of 60 seconds and time weighting factors of 0.55 and 1.0 , respectively, along with those of Viessman et al. (1972) are compared in Figure 8. The predicted hydrographs denoted as plots $B$ and $C$ in the figure are in close agreement with those of Viessman et al. (1972). However, the slight influence of the time weighting factor in the predicted peak flows at mid-reach and down-stream locations can be observed.

Though the difference in peak flows with $\theta$ values of 0.55 and 1.0 is minimal for a time step of 60 seconds, significant differences for larger time steps such as 300 secondsor more are apparent. Figure 9 illustrates very clearly the iteractive effect of the time weighting


Figure 7. Weighted Implicit Kinematic Finite Element Model Simulation at 300 Seconds for $\theta$ of 0.75 and 1.00.


Figure 8. Comparison of Hydrographs from WIDFEM at $\Delta t$ of 60 Seconds and $\theta$ of 0.55 and 1.00 and Viessman's Solution.


Figure 9. Weighted Implicit Diffusion Finite Element Model Simulation at 300 Seconds for $\theta$ of 0.55 and 1.00.
factor, $\theta$, and the numerical dispersion resulting from use of large time steps. Values of $\theta$ greater than 0.55 tend to attenuate the peak discharge. This observation equally validates that of the weighted implicit kinematic finite element model discussed earlier.

Like the WIKFEM, the weighted implicit diffusion is unconditionally stable for the time weighting factor in the range of $0.55 \leq \theta \leq$ 1.0. In spite of the numerical distortion associated with the use of large time steps, only a $\theta$ value of 0.55 predicted hydrographs identical to Viessman et al. (1972). As a result, subsequent simulations of the WIDFEM for time steps large than 60 seconds were executed with a $\theta$ value of 0.55 .

## Hydrographs from WICFEM

Applications of the weighted implicit complete finite element model to the idealized channel using a time step of 60 seconds and a $\theta$ value of 0.55 predicted the discharge hydrographs shown in Figure 10. Hydrograph results are identical to those of Viessman et a1. (1972) on the rising limbs but differ slightly on the receding limbs. On the average this difference is insignificant. The WICFEM affords an unconditionally stable solution for the time weighting factor in the range of $0.55 \leq \theta \leq 1.0$. Also the model shares the same basic characteristic as the WIDFEM discussed earlier.

## Flow Simulation in a Natural Channel

The second test analyzed flow through a natural river channel. The Illinois River between Watts and Tahlequah gaging stations (Sta. 1955 and 1965, respectively) in Oklahoma, shown in Figure 11, was chosen.


Figure 10. Comparison of Hydrographs from WICFEM at $\Delta t$ of 60 Seconds and $\theta$ of 0.55 and Viessman's Solution.


Figure 11. Map of Illinois River Basin.

Geometric cross-sectional data (developed from topographic maps) were collected from Weigant (1982) along with the flood data of April 10, 1979, and were used in predicting the flow hydrograph at the Tahlequah station. Figure 11 shows the Illinois River and locations of the stations.

Owing to the nature of the available topographic data, simulation was executed using a composite channel section. The changes from section to section in some locations are significant enough that smaller distance steps are necessary to adequately represent them in the model. Thus, the channel sections were averaged with a single longitudinal bottom slope of 4.5 feet per mile (Weigant, 1982).

## Initial and Boundary Conditions

Initial depths of flow were generated by backwater calculation starting from a downstream depth. Discharge values at intermediate nodes were estimated by linear interpolation applied to the two initial discharges at up-and down-stream locations. Nodal velocities corresponding to initial depths are calculated by dividing the nodal discharge by corresponding cross section. At the upstream point, the discharge was prescribed as a function of time. At the downstream boundary, a loop rating curve was imposed.

For the 1979 flood, the initial discharge values are given by the unsteady nonuniform flow of 482 cfs at the Watts station and 596 cfs at the Tahlequah station at time $t=0$. The discharge hydrograph at the Watts Station increased from 482 to 22980 cfs in 28 hr and then decreased to 1722 cfs in additional 68 hr . Figures 12 and 13 show the observed discharge hydrographs at Watts and Tahlequah and the rating curves at the


Figure 12. Observed Discharge Hydrographs at Watts, Tahlequah, and Flint Creek Stations for April 10, 1979, Flood, Illinois River, Oklahoma.


Figure 13. Rating Curves at Watts and Tahlequah, Illinois River, Oklahoma, for April 10, 1979.
stations, respectively. Computed flow at Tahlequah, 50.4 miles from Watts, was compared to the observed flow at the same station.

## Determination of Flow Parameters

The flow parameters necessary for simulation in natural channels are the channel cross sections, A; top widths, B; Manning's roughness coefficient, n ; and lateral inflow, q .

Average cross-sectional and top width data were utilized to generate a fourth-order polynomial equation, using a least square fitting program (Davis, 1973). A fourth-order polynomial yielded the best fit from the analysis of variance. By increasing the order of the polynomial beyond fourth, it was necessary to see if the increase in the degree of the polynomial significantly improved the fit of the regression. Such statistics as the sum of square due to deviation defined as the difference between total sum of square $\left(S S_{T}\right)$ and sum of square due to regression $\left(S S_{R}\right)$ and the goodness-of-fit defined as $S S_{R} / S S_{T}$ were used for assessment. The general form of the equation adopted to model the averaged cross section areas and top width ${ }^{\text { }}$ is represented as:

$$
\begin{align*}
& A(Y)=b_{0}+b_{1} Y+b_{2} Y^{2}+b_{3} Y^{3}+b_{4} Y^{4}  \tag{6.1a}\\
& B(Y)=c_{0}+c_{1} Y+c_{2} Y^{2}+c_{3} Y^{3}+c_{4} Y^{4} \tag{6.1b}
\end{align*}
$$

where $A(Y)$ and $B(Y)$ implies that the area and top width are functions of depth of flow only. Figure 14 illustrates a typical cross section geometry of the Illinois River as given in Equation (6.1). Results are included in the computer sample output.

The initial estimated Manning's roughness coefficient variation


Figure 14. Typical Cross-section of Illinois
River, Oklahoma.
against discharge for the Illinois River as provided is plotted in Figure 15. Also shown are the fitted third-order polynomial regression equations of the initial estimated roughness coefficient and of the modified coefficient values. The fitted regression curves are necessary because the plot of the initial estimates of roughness coefficient versus discharge depicts shape variations in some adjoining corners. Thus, a smooth curve was deemed necessary to better represent actual roughness coefficient variations Equation (6.2).

$$
\begin{equation*}
n=0.03713+0.14097 E-05 Q+0.41739 E-10 Q^{2}-0.23004 E-14 Q^{3} \tag{6.2}
\end{equation*}
$$

The modified initial estimates of the roughness coefficient variation are represented as:

$$
\begin{equation*}
n=0.02615+0.42801 E-05 Q-0.21618 E-09 Q^{2}+0.39355 E-14 Q^{3} \tag{6.3}
\end{equation*}
$$

The lateral inflow hydrograph at Flint Creek, a tributary of the Illinois River 13.2 miles downstream of the Watts Station, recorded during the same date, was imposed as a function of time, Figure 12. The main channel reach corresponding to 13.2 miles from Watts was allowed to receive the lateral inflow from Flint Creek. The inflow is represented in cubic feet per second per area of reach.

## Hydrographs from WIDFEM and WICFEM

Application of the flow models to the Illinois River was limited to WIDFEM and WICFEM because of the inherent flat slope of the channel. Use of the kinematic flow models would not be adequately justified in this particular example based on the slope approximation analysis (Henderson, 1966).


Figure 15. Estimated and Fitted Manning's Roughness Coefficient Variation with Discharge, Illinois River, Oklahoma.

The depth of flow at the Tahlequah Station for the WICFEM using a time step of 30 minutes and a time weighting factor of 0.55 is predicted and compared with that measured in Figure 16. Simulated results are in excellent agreement with observed flow. The marginal difference between computed and actual depths at the early portion of the rising limb and at the tailing edge of the hydrograph reflects the uncertainity of the input data. Among other things, the models are sensitive to variations of the Manning's roughness coefficients in predicting flows. Higer roughness coefficients imply reduced flows and vice versa. Thus, close predictions are possible as long as the roughness coefficient and other input data are accurate.

Figure 16 also illustrates the response of the WICFEM to a modified Manning's roughness coefficient regression Equation (6.3) in predicting depths of flow. The predicted stage hydrograph indicates a slightly high peak, at six hours earlier than the previous prediction using equation (6.2). Indeed, the simulated depths of flow using Equation (6.2) yielded the time of the peak that are more similar to those observed than to those from Equation (6.3). Thus, Equation (6.2) is more representative of the actual roughness variation in the Illinois River. Figure 17 shows the depth of flow predicted using the weighted implicit diffusion finite element mode1, WIDFEM. The same observations are valid as discussed above using WICFEM. Comparison of the computed flows from WIDFEM and WICFEM using Manning's regression Equation (6.2) against the observed records at the Tahlequah Station is provided in Figures 18 and 19. Discharge hydrographs depict a compounded error of the computed depth of flow and velocity of flow for a given location in the stream. For instance, the difference in peak flows as indicated in Figure 16 is


Figure 16. Observed and Predicted Stage Hydrographs at Tahlequah Station from Weighted Implicit Complete Finite Element Model for $\Delta t$ of 1800 Seconds.


Figure 17. Observed and Predicted Stage Hydrographs at Tahlequah Station from Weighted Implicit Diffusion Finite Element Model for $\Delta t$ of 1800 Seconds.


Figure 18. Observed and Simulated Discharge Hydrographs at Tahlequah Station for $\Delta t$ of 1800 Seconds.


Figure 19. Observed and Simulated Discharge Hydrographs at Tahlequah Station for $\Delta t$ of 900 Seconds.
about $7 \%$ considering prediction with Equation (6.2), while Figure 18 shows an error of $15 \%$ for WICFEM. Apparently, the error distribution amongst the predicted depth and velocity as illustrated in the discharge hydrograph are bound to be uneven.

Simulated results of the WIDFEM and WICFEM as shown in Figure 19 for a time step of 15 minutes and weighting factor 0.55 are exactly the same. However, a comparison of the two models for $\Delta t$ of 30 minutes and $\theta$ of 0.55 indicates a slight difference only at the peaks, Figure 18. Invariably, the WICFEM seems to sustain lesser numerical distortion for larger time steps than the WIDFEM. Still, there is much to be gained in the use of WIDFEM. Hydrograph results from WIDFEM are more comparable to those from WICFEM. In addition, the computer time and cost are slightly less for WIDFEM. Appendix $K$ compares computer CPU time and cost of models.

## CHAPTER VII

## SUMMARY AND CONCLUSIONS

## Numerical Performance of Models

The numerical properties of the flow models--EKFEM, WIKFEM, WIDFEM, and WICFEM--such as rate of convergence, accuracy, and stability, need to be assessed through well established mathematical relations. For instance, the Courant condition is employed in the explicit finite difference technique to evaluate the dynamic stability condition arising from the size of the time steps. Since similar conditions in the finite element techniques are not versatile and few in use are formulated under limited assumptions, we are therefore encouraged to draw comparisons from documentations established for the finite difference schemes at least for the time being.

The convergence criterion is a condition in which the solution of the finite element equation for a finite grid size approaches the true solution of the original partial differential equation. For the weighted implicit finite difference scheme proposed by Fread (1974), the convergence criterion was developed by determining the functional form of the truncation error through the Taylor series expansion about the point at which the difference equation is computed. The truncations error, $T R$, can be expressed as:

$$
\begin{equation*}
T R=(2 \theta-1) 0(\Delta t)+0\left(\Delta t^{2}\right)+0\left(\Delta x^{2}\right) \tag{7.1}
\end{equation*}
$$

where 0 indicates "order of", and when $\theta=1$, the truncation error is:

$$
\begin{equation*}
T R=0(\Delta t)+0\left(\Delta t^{2}\right)+0\left(\Delta x^{2}\right) \tag{7.2}
\end{equation*}
$$

Equation (7.2) shows that the fully implicit difference scheme is only first order accurate due to $\Delta t$ term. However, when $\theta=0.5$, the error shows a second order accuracy for $\Delta t$ and $\Delta x$.

The WIKFEM, WIDFEM, and WICFEM converge to the true solution for various values of the weighting factor ranging from 0.55 to 1.00 . For $\theta$ less than 0.55 , the models are completely unstable and invariably do not converge. This leads to the concept of numerical stability, defined as a condition whereby the numerical round-off errors introduced in a computational procedure fail to be amplified into an unlimited error. If errors generated at time level $(t+\Delta t)$ are smaller than the errors at time $t$ and not vice versa, the solution is said to be stable.

Stability of the non-linear difference equations of Saint Venant has been investigated by fourier analysis (Fread, 1973, 1974). This analysis is known as the Von Neumann method. In general, results indicate that an implicit difference formulation of the unsteady flow equation is unconditionally stable for any ratio of $\Delta x / \Delta t$, when the weighting factor, $\theta$, is restricted to the range $0.5<\theta<1.0$. The analysis proves also that stability of the implicit difference equation does not depend on the ratio $\Delta x / \Delta t$ like the explicit method and method of characteristics. The weighted implicit finite element flow mod-els--WIKFEM, WIDFEM, and WICFEM--are found to be unconditionally stable for the weighting factor in the range of $0.55 \leq \theta \leq 1.0$. However, rapid convergence for weighting factorof unity is observed only for :NKEEM.

The EKFEM bears similar restrictions as the explicit finite difference scheme. Numerical stability is conditional as defined by the Courant condition. Also, the WIKFEM, WIDFEM, and WICFEM reflect similar numerical properties as the implicit finite difference routine. The concept of explicit and implicit schemes applied to the finite element, FE, and finite difference, FD, formulations tends to tie the FE and FD in the same numerical subset.

Conclusions

Based on the results of the finite element modeling of the streamflow routing for idealized and natural channels, the following conclusions can be drawn:

1. Explicit and weighted implicit kinematic, weighted implicit diffusion, and complete flow models have been developed to predict the velocity of flow, depth of flow, and discharge in a stream.
2. The explicit kinematic finite element model, EKFEM, solves the flow routing problems, having a maximum time step of two seconds.
3. The weighted implicit kinematic finite element model, WIKFEM yields accurate results, with a maximum time interval of ten minutes and weighting factor in a range of 0.55 to 1.00 for a rectangular channel.
4. Both the weighted implicit diffusion and complete finite element models yield accurate and unconditionally stable solutions.
5. A11 the models--EKFEM, WIKFEM, WIDFEM, and WICFEM--have been tested against a problem presented by Viessman et a1. (1972). The comparisons of the flood hydrographs are in close agreement, and the observed difference resides on the speed and stability. In this regard, the weighted implicit models excel.
6. Use of a simplified model such as WIKFEM and WIDFEM in terms of computer storage and cost will be preferred provided good engineering judgement is exercised in their application. For this reason, these models will be favored over the complete solution of the unsteady flow equations.
7. Only the weighted implicit diffusion and complete finite element models were applied to a natural channel, the Illinois River in Oklahoma, for a flood observed on April 10, 1979. Simulated discharge hydrographs at the Tahiequah station, 50.4 miles downstream from the Watts Station, with time steps of 15 and 30 minutes and a weighting factor of 0.55 are in close agreement with the observed flows. A discrepancy of 8 percent in the maximum stage and 15 percent in the maximum discharge is attributed to the degree of accuracy of the input data, especially the roughness coefficient.
8. Not much observable difference exists between the simulated results of the WIDFEM and WICFEM for the natural channel flood routing test. For these particular test results, there is more to be gained in using the simplified diffusion model as discussed earlier in (6) above.

## CHAPTER VIII

SUGGESTIONS FOR FUTURE STUDY

The following suggestions for future study would be helpful in using the flow models for predicting the depth, velocity, and volumetric flow rate in a natural channel.

1. Modify the present flow models to incorporate boundary geometry at bridges showing contracting and expanding flow. In addition, field surveys of the hydraulic roughness values for various channel reaches are vital. Variation should be indicated in terms of longitudinal channel distance as well as the depth of flow or volumetric flow rate. Roughness coefficient values imposed on each cross section are usually helpful in locating where a cross section should be subdivided to determine distributed properties. For instance, values of 0.3 and 0.1 are assumed for the expansion and contraction coefficients, respectively.
2. Determine what portion of the cross section conveys flow and what portion stores water, particularly for smaller flood events. This might not be necessary for a very large flood wave. For instance, in the present study, it was assumed that the entire cross section conveyed flow for the flood of April 10, 1979, in the Illinois

River. For smaller events, it is assumed that all conveyance occurs in or near the main channel (Thomas, 1975).
3. Study other possible forms of modeling the flow cross sections and the corresponding top width besides using higher order polynomial curve fitting methods. Clearly, there are many possible approaches such as: (a) higher order spline function (cubic spline), (b) logarithmic or exponential regression equations, and (c) simple averaging and interpolation of the input data for intermediate values.

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APPENDIX A

NEWTON-RAPHSON FUNCTIONAL ITERATIVE METHOD FOR SOLUTION OF NON-LINEAR SYSTEM(S)

## A. Single Variable Equation

Consider a single non-linear variable equation expressed in functional form as follows:

$$
\begin{equation*}
f(x)=0 \tag{A.1}
\end{equation*}
$$

where

$$
x=a \text { real variable; }
$$

$f(x)=$ any reasonably well-behaved function.
The solution of the variable x of Equation (A.1) is obtained in an iterative manner, proceeding from the first solution estimate, $x^{n}$, towards the succeeding improved estimate, $x^{n+1}$, which tends to converge toward the solution variable $x$. The orderly procedure by which the improved solution estimate $x^{n+1}$ is sought, such that it converges to the true solution $x$, is known as Newton-Raphson Iteration and is described as follows.

Let the non-linear equation $f(x)$ be expanded using its Taylor series for an initial iterate $x^{0}$.
i.e.
$f(x)=f\left(x^{0}\right)+\frac{\left(x-x^{0}\right)}{1!} f^{\prime}\left(x^{0}\right)+\frac{\left(x-x^{0}\right)^{2}}{2!} f^{\prime \prime}\left(x^{0}\right)+\frac{\left(x-x^{0}\right)^{3}}{3!} f^{\prime \prime \prime}\left(x^{0}\right)(A .2)$
The linear function of $x^{0}$ that best approximates the non-linear function $f(x)$, evaluated at $x^{0}$, is obtained by retaining only the first order terms of Equation (A.2) such as:

$$
\begin{equation*}
f(x)=f\left(x^{0}\right)+\Delta x f^{\prime}\left(x^{0}\right) \tag{A.3}
\end{equation*}
$$

where:
$\Delta x=x-x^{0}$ (correction value);
$f^{\prime}\left(x^{0}\right)=\frac{\partial f\left(x^{0}\right)}{\partial x^{0}} \quad$ (Jacobian term evaluated at $x^{0}$ ).
An iteration procedure is desired which will cause the function $f\left(x^{0}\right)$ to approach zero as $\Delta x$ approaches zero. Thus, theleft-hand side of Equation (A.3) is made equal to zero with the following resulting generalized
iteration algorithm:

$$
\begin{equation*}
f^{\prime}\left(x^{n}\right)\left(x^{n+1}-x^{n}\right)=-f\left(x^{n}\right) \tag{A.4}
\end{equation*}
$$

where:
n and $\mathrm{n}+1$ are previous and current iterates respectively.
The Jacobian $f^{\prime}\left(x^{n}\right)$ needs to be updated at every iteration cycle. However, the initial Jacobian can be kept and used for all cycles or updated at selected iteration cycles at the expense of slow convergence. The iteration process is stopped when convergence is achieved. This can be checked in two ways--the absolute and relative tests. The former requires the absolute difference between the current and previous iterates to be less or equal to a specified value called error criterion.

$$
\begin{equation*}
\left|x^{n+1}-x^{n}\right| \leq \varepsilon_{1} \tag{A.5}
\end{equation*}
$$

where:
$\varepsilon_{p}=$ error criterion
The relative test is expressed as:

$$
\begin{equation*}
\frac{\left|x^{n+1}-x^{n}\right|}{\operatorname{MAX}\left(\left|x^{n}\right|,\left|x^{n+1}\right|\right)} \leq \varepsilon_{2} \tag{A.6}
\end{equation*}
$$

The relative error test is usually preferred to the absolute test because while the latter requires the knowledge of the size of $x^{n}$, the former takes that already into account.
B. Multi-Variable Equation

For a system of non-linear multi-variable cquations, the Fewton?aphson method is equally efficient in providing the roots or solution o.: such a system (Amein and Fang, 1969; Fread, 1976). Consider the following $N$-dimensional system of non-linear algebraic equations:

$$
\begin{align*}
& f_{1}\left(x_{1}, x_{2}, x_{3}, \cdots, x_{N}\right)=0 \\
& f_{2}\left(x_{1}, x_{2}, x_{3}, \cdots, x_{N}\right)=0 \\
& \vdots  \tag{A.7}\\
& f_{N}\left(x_{1}, x_{2}, x_{3}, \cdots, x_{N}\right)=0
\end{align*}
$$

or in a vector notation:

$$
\begin{equation*}
f_{i}(x)=0 \tag{A.8}
\end{equation*}
$$

where:
subscript i denotes a particular equation.
In a manner analogous to the steps discussed for a single variable equation (EQ. A. 1 through A.6), the linearized form of equation $A .8$ is as follows (see EQ. A.4):
$f_{i}^{\prime}\left(x^{n}\right)\left(x^{n+1}-x^{n}\right)=-f_{i}\left(x^{n}\right)$
Express EQ. (A.9) in a more concise form as:

$$
\begin{equation*}
f_{i}^{\prime}\left(x^{n}\right) \Delta x=-f_{i}\left(x^{n}\right) \tag{A.10}
\end{equation*}
$$

where:

$$
\Delta x=\left(\begin{array}{c}
x_{1}{ }^{n+1}-x_{1}^{n} \\
x_{2}^{n+1}-x_{2}^{n} \\
\cdots \cdots \cdots \cdots \\
\cdots \cdots \cdots \cdots \\
\cdots \cdots \cdots \cdots \\
x_{N}^{n+1}-x_{N}^{n}
\end{array}\right) \quad, f_{i}\left(x^{n}\right)\left(\begin{array}{c}
f_{1}\left(x^{n}\right) \\
f_{2}\left(x^{n}\right) \\
\cdots \cdots \\
\cdots \cdots \\
\cdots \cdots \cdots \\
f_{N}\left(x^{n}\right)
\end{array}\right)
$$

$$
f_{i}^{\prime}\left(x^{n}\right)=\left[\begin{array}{ccccc}
\frac{\partial f_{1}}{\partial x_{1}} & \frac{\partial f_{1}}{\partial x_{2}} & \frac{\partial f_{1}}{\partial x_{3}} & \cdots & \frac{\partial f_{1}}{\partial x_{N}} \\
\frac{\partial f_{2}}{\partial x_{1}} & \frac{\partial f_{2}}{\partial x_{2}} & \frac{\partial f_{2}}{\partial x_{3}} & \cdots & \\
\cdot & \cdot & \cdot & \cdot & \frac{\partial f_{2}}{\partial x_{N}} \\
\cdot & \cdot & \cdot & \cdot & \cdot \\
\frac{\partial f_{N}}{\partial x_{1}} & \frac{\partial f_{N}}{\partial x_{2}} & \frac{\partial f_{N}}{\partial x_{3}} & \cdots & \cdot \\
\frac{\partial f_{N}}{\partial x_{N}}
\end{array}\right]
$$

The solution of the linear system of equation represented in vector form by Equation (A.10) is sought for the unknown linear correction vector $\Delta x$ by a suitable matrix solution technique. For a system of ( $N \times N$ ) matrix equation, Gaussian elimination may be employed. However, the most efficient triangular decomposition solution technique for a compact bitridiagonal matrix (Douglas et al., 1959) is presented in Appendix B.

The convergence of the iteration process, Equation (A.6), depends on a good initial solution vector estimate $x^{\circ}$. If the initial iterates are sufficiently close to $x$, convergence is attained at a quadratic rate since the iterative procedure is second order, that is, involves the first derivative.

APPENDIX B

COMPACT BI-TRIDIAGONAL SOLUTION ALGORITHM

Consider the following system of linear algebraic equations generated by weighted implicit diffusion- or complete-flow models of the finite element approximations of Saint-Venant equations and the Newton-Raphson iterative method as a bi-tridiagonal system:

$$
\begin{align*}
& { }^{(1)} y_{i} y_{i-1}+{ }_{a}^{(2)} v_{i-1}+b_{i}^{(1)} y_{i}+b_{i}^{(2)} v_{i}+c_{i}^{(1)} y_{i+1}+c_{i}^{(2)} v_{i+1}=d_{i}^{(1)} \\
& a_{i}^{(3)} y_{i-1}+a_{i}^{(4)} v_{i-1}+b_{i}^{(3)} y_{i}+b_{i}^{(4)} v_{i}+c_{i}^{(3)} y_{i+1}+c_{i}^{(4)} v_{i+1}=d_{i}^{(2)} \\
& \quad \text { for } 1 \leq i \leq N \\
& \text { with } a_{1}^{(m)}=c_{N}^{(m)}=0 \text { for } 1 \leq m \leq 4
\end{align*}
$$

Equation (B.1) is an equivalent form of Equation (A.10), Appendix $A$, and can be conveniently expressed in a compact (2N×6) matrix form as follows:

The compact solution algorithm developed by Douglas et al (1959) and later used by Von Rosenberg (1969) is a direct solution technique for a system of linear equations. The algorithm is an efficient triangular decomposition method that yields a recursion equation, thus substantially reducing computations and computer core storage.

The algorithm is as follows:
First Computer

$$
\begin{align*}
& \beta_{i}^{(1)}=b_{i}^{(1)}-a_{i}^{(1)} \lambda_{i-1}^{(1)}-a_{i}^{(2)} \lambda_{i-1}^{(3)} \\
& \beta_{i}^{(2)}=b_{i}^{(2)}-a_{i}^{(1)} \lambda_{i-1}^{(2)}-a_{i}^{(2)} \lambda_{i-1}^{(4)}  \tag{B.3}\\
& \beta_{i}^{(3)}=b_{i}^{(3)}-a_{i}^{(3)} \lambda_{i-1}^{(1)}-a_{i}^{(4)} \lambda_{i-1}^{(3)} \\
& \beta_{i}^{(4)}=b_{i}^{(4)}-a_{i}^{(3)} \lambda_{i-1}^{(2)}-a_{i}^{(4)} \lambda_{i-1}^{(4)}
\end{align*}
$$

and

$$
\begin{aligned}
& \delta_{i}^{(1)}=d_{i}^{(1)}-a_{i}^{(1)} \gamma_{i-1}^{(1)}-a_{i}^{(2)} \gamma_{i-1}^{(2)} \\
& \delta_{i}^{(2)}=d_{i}^{(2)}-a_{i}^{(3)} \gamma_{i-1}^{(1)}-a_{i}^{(4)} \gamma_{i-1}^{(2)} \\
& \text { with } \delta_{i}^{(1)}=d_{i}^{(1)} \text { and } \delta_{i}^{(2)}=d_{i}^{(2)} \\
& \text { and } \mu_{i}=\beta_{i}^{(1)} \beta_{i}^{(4)}-\beta_{i}^{(2)} \beta_{i}^{(3)}
\end{aligned}
$$

The $\beta_{i}^{(m)}, \delta_{i}^{(m)}$, and $\mu_{i}$ are computed to aid in the computation of the following functions and need not be stored after the computation of

$$
\begin{align*}
& \lambda_{i}^{(1)}=\left(\beta^{(4)} c_{i}^{(1)}-\beta_{i}^{(2)} c_{i}^{(3)}\right) / \mu_{i} \\
& \lambda^{(2)}=\left(\beta^{(4)} c_{i}^{(2)}-\beta_{i}^{(2)} c_{i}^{(4)}\right) / \mu_{i}  \tag{B.4}\\
& \lambda_{i}^{(3)}=\left(\beta_{i}^{(1)} c_{i}^{(3)}-\beta_{i}^{(3)} c_{i}^{(1)}\right) / \mu_{i}  \tag{B.4}\\
& \lambda^{(4)}=\left(\beta_{i}^{(1)} c_{i}^{(4)}-\beta_{i}^{(3)} c_{i}^{(2)}\right) / \mu_{i}
\end{align*}
$$

and

$$
\begin{aligned}
& \gamma_{i}^{(1)}=\left(\beta_{i}^{(4)} \delta_{i}^{(1)}-\beta_{i}^{(2)} \delta_{i}^{(2)}\right) / \mu_{i} \\
& \gamma_{i}^{(2)}=\left(\beta_{i}^{(1)} \delta_{i}^{(2)}-\beta_{i}^{(3)} \delta_{i}^{(1)}\right) / \mu_{i}
\end{aligned}
$$

The values of $\lambda_{i}^{(m)}$ and $\gamma_{i}^{(m)}$ must be stored as they are used in the back solution. This is

$$
\begin{align*}
& y_{N}=\gamma_{N}^{(1)} \\
& v_{N}=\gamma_{N}^{(2)} \tag{B.5}
\end{align*}
$$

and

$$
\begin{gathered}
y_{i}=\gamma_{i}^{(1)}-\lambda_{i}^{(1)} y_{i+1}-\lambda_{i}^{(2)} v_{i+1} \\
v_{i}=\gamma_{i}^{(2)}-\lambda_{i}^{(3)} y_{i+1}-\lambda_{i}^{(4)} v_{i+1} \\
\text { for }(N-1) \geq i \geq 1
\end{gathered}
$$

APPENDIX C

DESCRIPTION OF SUBROUTINES

The list of subroutines and their corresponding functions is given below. The list of major variables and symbols used in the computer program is provided in the comment page of computer program listing, Appendix E. Any temporary storage variables are not included because their definitions are obvious.

## Subroutines

MAIN Coordinates the functions of the other subprograms, and prints converged solutions for prescribed time increments.

READW Reads and echo-checks all the input data.
JACOBI Evaluates and updates the Jacobian matrix of (2Nx6) terms.
VECTR Evaluates and updates the column vector of size ( $2 \mathrm{~N} \times 1$ ).
GEOMTR Evaluates and updates the nodal flow area, wetted perimeter, variation of Manning's roughness coefficient with discharge, and the change of hydraulic radius with respect to depth of flow.

BTRIDG
Solves the compact ( $2 \mathrm{~N} \times 6$ ) bi-tridiagonal matrix equations.

APPENDIX D

GUIDE FOR DATA INPUT

The formats for entering the data are given below. The same format statements for READ and :URITE are applicable to both diffusion and complete flow models as provided in the subprogram READW. The data deck for the particular example of the Illinois River flood of April 10, 1979, is presented. However, data for the Flint Creek as lateral inflow for a single reach is entered via the MAIN program as a DATA STATEMENT. The reader should refer to Appendix $E$ for the definition of the variables.

| CARD | COLUMNS | FORMAT | VARIABLE |
| :---: | :---: | :---: | :---: |
| 1 | 1-10 | F 10.4 | TPRINT |
|  | 11-20 | F 10.4 | TTA |
|  | 21-30 | F 10.4 | TSUM |
|  | 31-40 | F 10.4 | T |
| 2 | 1-10 | F 10.4 | TETHA |
|  | 11-20 | F 10.4 | dETA |
|  | 21-30 | F $10 \cdot 4$ | DETV |
|  | 31-40 | F $10 \cdot 4$ | So |
|  | 41-50 | I10 | IMAX |
|  | 51-60 | I10 | N1 |
| 3 | 1-72 | 6F12.5 | YO(J) |
| 4 | 1-72 | 6F12.5 | QRE (J) |
| 5 | 1-72 | 6F72.5 | QLAT (J) |
| 6 | 1-72 | 6F12.5 | XL (J) |
| 7 | 1-5 | I5 | JORD |
| 8 | 1-50 | 5F10-5 | ASF(J) |


| CARD | COLUMN | FORMAT | VARIABLE |
| :---: | :---: | :---: | :--- |
| 9 | $1-50$ | $5 F 10.5$ | $\operatorname{PSF}(\mathrm{~J})$ |

Note: $\operatorname{QSTR}(\mathrm{J})$ and $\operatorname{TRS}(\mathrm{J})$ are included in the main program in DATA STATEMENTS.

## APPENDIX E

COMPUTER PROGRAM LISTING FOR WICFEM

```
    SOOB TIME=5
```



```
    * *
    * - DIMENSIONAL STREAMFLOM ROUTING MODEL *
```




```
    *
    COMPLETE FLOM FINITE ELEMENT STREAMFLOM ROUTING
    * *
    * SOLYED IMPLICITLY BY ITERATIVE NEWTON-RAPHSON MTD. *
```



```
    *
    * *.OEFINITION OF TERMS \bullet.. *
    * O
    *UARIABLES UNITS ARE AS FOLLOWS: TIME{SECIgLENGTMEFT)
    *DETPH(FTI, VELOCITY&FT. PER SECI,OISCHARGEICFS)
    *ACC IS THE ACCELERATION OF GRAVITY,32.2FT. PER SEC PER SEC
    #ACF,PCF ARE POLYNO. COEFF. FOR AREA E METTEO PERIMETER
    *ANSPM ARE THE AREA E GETTED PERIMETER OF FLON RESPECTIYELY
    *HYD 15 THE RATE OF CHANGE OF HYOR.RADIUS WITH DEPTH
    *QSTR IS THE UPSTREAM INFLON DISCHARGE HYDRCGRAPM
    *OLAT IS THE LATERAL FLOW TERMoFT. PER SEC
    *QFL IS THE LATERAL IMFLOW HYOROGRAPH AT FLINT CREEK *
    #SO IS THE CONSTANT CHANNEL SLOPE
    *RN IS THE MANNING ROUGHNESS COEFF.
    *XL IS THE NODAL SPACING.
    *YO IS THE INITIAL UNIFORM NORMAL DEPTH.
    *VO 15 THE INITIAL UNIFORM NORMAL VELOCITY *
    *N2 IS THE TOTAL NUMBER OF NODES *
    * IS THE TIME STEP (SECONDSI *
    *TSUR IS THE ENTIRE FLOOD DURATION IN SECONOS. *
    *TSR IS THE TIME FOR UPSTREAM INFLOM HYOROGFAPM *
    *TFL IS THE TIME FCR LATERAL INFLOH HYOROGRAPH AT FLINT *
    *TPRINT IS THE TIME FOR INITIAL PRIATING (SECONOS: *
    *TTA IS THE INCREMENTAL PRINTING TIME (SECONOS) *
    *-\infty-IF TIME STEP IS GREATER THAN TTA PRINTIAG MILL BE -\infty- *
    * PERFORMED AT THE INCREMENT OF THE TIME STEPgT-\infty-* *
    *IMAX IS THE MAX. ITERATION LIMIT
    *NT IS THE NUMBER OF POINTS FOR UPSTREAM INFLON HYDROGRAPN *
    *NTP IS THE UPSTREAM IMFLOK HYOROG POINT FOE THE PEAK FLOE *
    *JGOJGP ARE SAME AS NTONTP FOR LATERA INFLO# HYDROG FOR FLINT*
    - JORC IS THE CRDER OF POLYNOMIAL EQ. FOR AREA E W. PERIMETER *
    #DETA IS THR CONVERGENCE CRITERIA FGR DEPTH
    *TETHA IS THE TIME WEIGHTING FACTOR. *
    * YN IS THE CALCULATED DETH OF FLON *
    *VN IS THE CCRRESPONDIAG VELOCITY OF FLOM
    *日SR IS THE JACOSIAN MATRIX OF DIMENSION ENI X 6 %
    *CXVI E CXYZ ARE THE (AIXII COLUMN VECTORS EVALUATED AT *
    *(1-TETHA) E TETHA RESPECTIVELY.
    *LDIM,LOIN ARE THE VARIABLE DIMENSIONING PARAMETERS
    *REAOM IS THE SUBPROGRAM TO READ E EGMOE CMECK INPUT DATA
    GEOMTR IS THE SUBPROGRAM TO UPOATE FLON AREAgMETTED,
    - MAANING*S ROUGHNESS COEFF. E RATE OF CHAAGE HYOR. RACIUS
    *
```



```
    C
        DIMENSION ACF(5),AN(26), BSR(52,6%,CSV1(52),CSV2(521,FCF(5).
        1 OTPH{26),HYD(26):QLAT(25),QRE{261,PN(26):VELY(26):
```

```
    2 VO{26),YN{26:,XL{25),YO{26},YN(26)
        DIMENSION GFL(7),QSTR(11):TFL(7),TSR&11),RN(26)
        OATA GFL,JG,JGP/81.0,408.0.629.,9.31.,566.,318..9164.,94.7/
        DATA TFL/0.0.43200.,50400.,57600.,86400.,136800.,345600./
        DATA ACC,LDIN,LDIN,NR,LP,LK/32.2,52,26,5,2*%/
        DATA NT,NTP/11,%6/
    C
    C
    C
    C
    C
    OD-CAL CULATE INITIAL GEGMETRIC PARAMETERS FROM SUBPRCGRAM.
        CALL GEOMTR\ACF,AN,HYO,QRE,PCF,PN:YO,RN,AI,LDIN:
        PR = 2.13.
        0O 100 J = L,N1
    100 VO&J) = ORE{J)/AMPJ\
        OO = QRE{1%
        WRITELLP,140)
```



```
        HRITE(LP,145)
```



```
        1 'OISCHARGE*,4X, DDEPTH* 3X, "VELOCITY*, 6X, 'DISCHARGE*,AX,*DEPTH*,
        3X:*VELOCITY*)
    C
    C -OUSE INITIAL VELOCITY E DEPTH GF FLON AS GUESS VALLES
        TO INITIATE SIMULATION.
        OO 150 K=1%N1
        YN(K) = YO(K)
        150 VN(K) = VO(K)
        FAC1 = TETHA*T
        FAC2 = 1&- TETHAIOT
        OO 155 L = 1,N1
        155 RN(L) = ACC*RN(L)**2/2.2082
        -\infty SET LCOP FOR TIME SIMULATION.
        OR = 00
        00 900 JL =1,JSILE
        TIPE=FLOAT\JLIFT
        ——OUPDATE THE LATERAL INFLOM HYOROGRAPH FOR REACH,7
        TCK = TIME - TFLIJGI
        IF(TCK)160,160,200
    160 00 180 KC = 2&JG
        IF&TIME - TFL(KC))190:190.180
    180 CONTINUE
    190 GF% = GFL(KC-1) + (GFL{KC) - OFLSKC-1)|)
    1 (TFL{KC) - TFL(KC-1)]*{TIME - TFL(XC-1)!
        GO 10 320
200 IF&TCK - IFL(JGP),220,220,300
220 00 230 KS = MSOJGP
        IF&TIME - TFL\KSJ1240.240.230
230 CONTINUE
2*O OFG = OFLIKS-1) - (OFLEKS-1) - QFL(KSII)
```

```
            1 (TFL(KS)- TFL(KS-1)\#(TIME - TFL(XSO1))
            GOTO 320
    300 QFM = QFL(JGP)
    320 QLATMT)= GFH&\XL(7)FPM(7)I
C
C O- UPDATE UPSTREAM BQUNOARY GOMOITION.
        TO = TIME - TSRENTPI
        IF{TO\350,350,380
    350 DO 360 LG = 2,NTP
        IFITIME - TSR(LC)I370.370,360
    360 COMTINUE
    370 QR = QSTR{LC-1) + OSTRILC) - OSTR(LC-1)]/
        1(TSR{LC)- TSR(LC-1)\#{TIME - TSR(LC-1)!
        GOTO 520
    380 IFETD - ISRSNT\1400,A00,500
    400 DO 420 JC = MC,NT
        IF(TIME - TSR(JCJ)450.450.420
    420 CONTINUE
    450 QR = OSTR(JC-1) = ESTR(JC-1) - OSTR(JCJI/
        1 (TSR(JC)-TSR(JC-1)I##TIME - TSR{JC-1))
        60 10 520
    SOO OR = OSTR&NTI
    520 CONTINUE
C ---CALL SUBROUTINE TO GENERATE GOLUMN YECTOR (2N X 1I
        JSBTCH = 1
        CALL VECTR&CSVI,FAC2,QLAT,YO,VO,SO,ACC,XL,OR,RNON1,
        1 AN,PN, QRE,LDN,LDIN,LOIM,JSWTCHI
    C
C
        -\infty-GENERATE JACOBIAN MATRIX.
        LUP=0
        CALL JACCEICBSR,FACIFYN,YN,XL,OLAT,OR,RM,
        1 AN:PN, QRE,HYO,SO,ACC,LDIM,LDN,LOIN,N1,LKI
    C
        -\infty- ITERATE TO CONYERGENCE FOR EACH TIME STEP.
        OO 599 LL =1,IMAX
        LAST = 2*N1
        JSWTCH = 2
        LUP = LUP + 1
        CALL Y ECTR&CSV2,FAC1,OLAT,YN:VN,SO,ACC,XL,QR,RN,N1s
            1 AN&PN, QRE,LON,LDIN,LDIM,JSWTCHS
        DO 530 K = 1,LAST
        CSV2(K) = CSV2(K)-CSV1(K)
    530 CONTIMUE
    C - -OBTAIN SOLUTION YIA TRI-DIAGONAL SUBPROGRAM.
        CALL 3TRIDGICSV2, BSR,DTPH&VELY,LDIM,LDIN,N1,LK)
        ---UPJATE THE NODAL GEOMETRIC PARAMETERS.
        CALL GEOMTR\ACF,AN,HYO,QRE,PCF,PN,YN&RN,NI,LDIN)
        DO 550 L = 1,N1
        550 RN(L) = ACC *RN(L) **2/2.2082
        JS = N1 - 1
        VEC = OR/AN&11
        VELY(1) = VN(1) - YEC
    C
    C
    C
    -\infty GHECK FOR RELATIYE CONVERGENCE FOR ALL VARTABLES.
    JERR = 0
    DO 56) J =1.N1
    Y&1 = ABSIDTPHC JO)
    VB1 = ABS(VELYTJ))
    YB2 = YN(J) - DTPH\J)
    VB2 = VN(J) - VELYCJI
```



```
126
127
128
129
1 3 0
131
132
133
134
135
136
137
138
139
140
141
142
143
144
145
146
147
148
149
150
```

    50 FORMAT (4F10.2)
    ```
    50 FORMAT (4F10.2)
    READ&NR,80ITETHA,DETA,DETV,SO, IMAX,M1
    READ&NR,80ITETHA,DETA,DETV,SO, IMAX,M1
    80 FORMAT 4F10.4,2I1O:
    80 FORMAT 4F10.4,2I1O:
    REAO&MR,90)(QSTR{L),L = 1,NT)
    REAO&MR,90)(QSTR{L),L = 1,NT)
    READ&NR,90I|TSR|L|,L = L,NT)
    READ&NR,90I|TSR|L|,L = L,NT)
    -O-READ INITIAL DEPTHS OF FLOW,DISCHARGE&LATERAL FLOME SPACING.
    -O-READ INITIAL DEPTHS OF FLOW,DISCHARGE&LATERAL FLOME SPACING.
        LDN = NL - L
        LDN = NL - L
        READ{NR,90 I{YOR J),J = 1,N1}
        READ{NR,90 I{YOR J),J = 1,N1}
        READ(YR,90)(QRE(J);J = 1;N1)
        READ(YR,90)(QRE(J);J = 1;N1)
        READ(NR,90)&OLAT(J)&J = 1:LON)
        READ(NR,90)&OLAT(J)&J = 1:LON)
        READINR,90){XL(J),J = 1.LDN)
        READINR,90){XL(J),J = 1.LDN)
        FORMAT ( GF12.5)
        FORMAT ( GF12.5)
    C ---READ OROER OF POLYNOMIAL EQ.
    C ---READ OROER OF POLYNOMIAL EQ.
        READINR,110JJORD
        READINR,110JJORD
    110 FORMAT&ISJ
    110 FORMAT&ISJ
        LR = JORD + 1
        LR = JORD + 1
    C ---REAO AREA E mETTED PERIMETER POLYNO. COEFF. MATRICE
    C ---REAO AREA E mETTED PERIMETER POLYNO. COEFF. MATRICE
    C ONE POMAT A TIME.
    C ONE POMAT A TIME.
        READ|NR,140)(ASF(d)sJ = 1,LR)
        READ|NR,140)(ASF(d)sJ = 1,LR)
        READ(NR=140)&PSF(d)sd = 1,LR)
        READ(NR=140)&PSF(d)sd = 1,LR)
    140 FORMAT (5F10.5)
    140 FORMAT (5F10.5)
C
C
C
C
    -\triangle-PRINT OUT INPUT DATA.
    -\triangle-PRINT OUT INPUT DATA.
        GRITESLP,150)
        GRITESLP,150)
    150 FORMAT (1H1:
    150 FORMAT (1H1:
        WRITEILP,160)NL,T,SO
        WRITEILP,160)NL,T,SO
    160 FORMAT ////20X, *TOTAL NO. OF NODES =4,I5//20X,
    160 FORMAT ////20X, *TOTAL NO. OF NODES =4,I5//20X,
        1 TIME STEP =*,F10.3,1X,*SEC."//20X, "CHANMEL BOTTOM SLOPE =* *
        1 TIME STEP =*,F10.3,1X,*SEC."//20X, "CHANMEL BOTTOM SLOPE =* *
        2 F10.41
        2 F10.41
        WRITESLP,170ITETHA, IMAX, DETA,DETV
        WRITESLP,170ITETHA, IMAX, DETA,DETV
    170 FORMAT|/f20x,*TIME WEIGHTING FACTOR =*,F10.4//20X,
    170 FORMAT|/f20x,*TIME WEIGHTING FACTOR =*,F10.4//20X,
        1 *MAX. ITERATION LIMIT =*,I5//20X,
        1 *MAX. ITERATION LIMIT =*,I5//20X,
        2 CONVERGENCE CRITERIA FOR DEPTH =0,F10.4//20X:
        2 CONVERGENCE CRITERIA FOR DEPTH =0,F10.4//20X:
        3 'CONVERGENCE CRITERIA FOR VELOCITY =',F10.4;
        3 'CONVERGENCE CRITERIA FOR VELOCITY =',F10.4;
    C
```

    C
    ```


```

    180 FORMAT///////23X, UPSTREAM DISCHARGE HYDROGRAPH=//20 X:
    ```
    180 FORMAT///////23X, UPSTREAM DISCHARGE HYDROGRAPH=//20 X:
        1 'J",5X, "TIME PERIOD*,5X, 'MEASURED FLOW'//(20X,I2,4X,
        1 'J",5X, "TIME PERIOD*,5X, 'MEASURED FLOW'//(20X,I2,4X,
        2F10.1,3X,F10.1):
        2F10.1,3X,F10.1):
    C
    C
        WRITE(LP,190){K,YO{K), QRE{K),K=1,N1)
        WRITE(LP,190){K,YO{K), QRE{K),K=1,N1)
        FORMAT (//////20X,*NODE*,5X:*INITIAL DEPTH*:5X,*INITIAL DISCHARGE*
        FORMAT (//////20X,*NODE*,5X:*INITIAL DEPTH*:5X,*INITIAL DISCHARGE*
        1//{18X,I3,8X,F10.3,12X,F10.3)1
        1//{18X,I3,8X,F10.3,12X,F10.3)1
        WRITEILP,200)(J,XL(J), QLATRJI,J =1, LDNI
        WRITEILP,200)(J,XL(J), QLATRJI,J =1, LDNI
    200 FORMAT / //////20X,*REACH*,9X; LENGTH*,9X, 'LATERAL FLOM*//& 20X,
    200 FORMAT / //////20X,*REACH*,9X; LENGTH*,9X, 'LATERAL FLOM*//& 20X,
        1 I3,8X,F10.2,5x:F10.6)1
        1 I3,8X,F10.2,5x:F10.6)1
        WRITEILP,220;
        WRITEILP,220;
        FORMAT (//////20X,*FOURTH-ORDER REGRESSION COEFF. FOR AREA*//23X,
```

        FORMAT (//////20X,*FOURTH-ORDER REGRESSION COEFF. FOR AREA*//23X,
    ```


```

            WRITELLP,300I(ASF(J);J = 1:LR)
    ```
            WRITELLP,300I(ASF(J);J = 1:LR)
            WRITEILP,2401
            WRITEILP,2401
    240 FORMAT///////20X,* FOURTH-ORDER REGRESSION COEFF. FOR EETTED
```

    240 FORMAT///////20X,* FOURTH-ORDER REGRESSION COEFF. FOR EETTED
    ```


```

        NRITEILP,300)|PSF{J|OJ = 1:LR)
    ```
        NRITEILP,300)|PSF{J|OJ = 1:LR)
        FORMAT (/20X,5F10.51
        FORMAT (/20X,5F10.51
        RETURY
        RETURY
            END
```

            END
    ```


```

    C
    ```
```

    C
    ```




\begin{tabular}{|c|c|c|}
\hline 320 & & \(00203 J=1, N 1\) \\
\hline 321 & & IFIJ = 1)100:100, 150 \\
\hline 322 & 100 & BETAIS \(=\) VELA1.3) \\
\hline 323 & & BETAS2) \(=\) VEL(1.4) \\
\hline 324 & & BETA(3) = VEL(2,3) \\
\hline 325 & & BETAIA) \(=\) VEL(2.4) \\
\hline 326 & & DETAS1) \(=\) COL(1) \\
\hline 327 & & DETA(2) \(=\) COL(2) \\
\hline 328 & & 2U = BETA(1)*BETA(4) - BETA(2)*EETA(3) \\
\hline 329 & & IFSZU.EO. \(0.012 \mathrm{CO}=0.001\) \\
\hline 330 & &  \\
\hline 331 & &  \\
\hline 332 & &  \\
\hline 333 & & SAC(1, 4 ) \(=(B E T A(1) * V E L(2,6)-B E T A(3) * V E L(1,6)) / 2 U\) \\
\hline 334 & & G0 TO 180 \\
\hline 335 & 150 & \(K=k+2\) \\
\hline 336 & & \(M=K+1\) \\
\hline 337 & &  \\
\hline 338 & &  \\
\hline 339 & &  \\
\hline 340 & &  \\
\hline & \(c\) & \\
\hline 341 & & DETA(1) = COL(K)-VEL(K, 1) \#GAMA(J-1, 1)-VEL(K,2)*GAMA(J-1,2) \\
\hline 342 & &  \\
\hline & \(c\) & \\
\hline 343 & & \(2 U=8 E T A(1) * \theta E T A P 4)-\operatorname{BETA}(2) * B E T A(3)\) \\
\hline 344 & & IFILU - EQ. \(0.012 \mathrm{C}=0.001\) \\
\hline 345 & &  \\
\hline 346 & &  \\
\hline 347 & &  \\
\hline 348 & &  \\
\hline & c & \\
\hline 349 & 180 & GAMA(J, 1) = (8ETA(4) DEDA(1)-BETA(2)* DETA(2))/ZU \\
\hline 350 & &  \\
\hline 351 & \(c^{200}\) & ```
CONTIMUE ---COMPUTE SOLUTION vIA RECURSIVEEE.
``` \\
\hline 352 & & LIMIT \(=\) N1 - 1 \\
\hline 353 & & JK \(=\) LIMIT \\
\hline 354 & & YXAN1) = GAMACLDIN: 11 \\
\hline 355 & & vx(N1) = GAMA(LDIN:23 \\
\hline 356 & & DO 305 L ( L,LIMIT \\
\hline 357 & &  \\
\hline 358 & &  \\
\hline 359 & & JK \(=\) JK - 1 \\
\hline 360 & 300 & cont in ue \\
\hline 361 & & Retury \\
\hline 362 & & END \\
\hline
\end{tabular}

\section*{APPENDIX F}

SAMPLE OUTPUT FOR WICFEM

The sample output listed in the pages following is the format with which the input parameters are reprinted for correction and referral. A clear illustration is drawn from a natural channel simulation of Illinois River, using flood of April 10, 1979, and the complete flow model.
```

TOTALND. OF NODES = 26
TIME STEP = 1800.000 SEG.
CHANNEL BOTTOM SLOPE = 0.0009
TIME WEIGHTING FACTOR = 0.5500
MAX. ITERATION LIMIT= }=6
CONVERGENCE CRITERIA FOR DEPTH = 0.0100
COYVERGENCE GRITERIA FOR YELOCITY = 0.1000

```
    UPSTREAM OISCHARGE HYDROGRAPH
\begin{tabular}{rrr} 
& TIME PERIOD & MEASURED FLOW \\
& & \\
1 & 0.0 & 482.0 \\
2 & 50400.0 & 757.0 \\
3 & 64800.0 & 5590.0 \\
4 & 79200.0 & 7710.0 \\
5 & 86400.0 & 11000.0 \\
6 & 100800.0 & 22980.0 \\
7 & 129600.0 & 11320.0 \\
8 & 158400.0 & 5100.0 \\
9 & 172800.0 & 4110.0 \\
10 & 208800.0 & 3104.0 \\
11 & 345600.0 & 1722.0
\end{tabular}
\begin{tabular}{lcr} 
NODE & IMITIAL DEPTH & INITIAL DISCHARGE \\
1 & 3.340 & \\
2 & 3.340 & 482.000 \\
3 & 3.340 & 482.000 \\
4 & 3.340 & 482.000 \\
5 & 3.340 & 482.000 \\
6 & 3.340 & 482.000 \\
7 & 3.340 & 482.000 \\
8 & 3.340 & 482.000 \\
9 & 3.340 & 482.000 \\
10 & 3.340 & 482.000 \\
11 & 3.340 & 482.000 \\
12 & 3.340 & 482.000 \\
13 & 3.340 & 482.000 \\
14 & 3.340 & 482.000 \\
\hline & & 482.000
\end{tabular}
\begin{tabular}{llr}
15 & 3.340 & 482.000 \\
16 & 3.340 & 482.000 \\
17 & 3.340 & 482.000 \\
18 & 3.340 & 482.000 \\
19 & 3.340 & 482.000 \\
20 & 3.340 & 482.000 \\
21 & 3.340 & 482.000 \\
22 & 3.340 & 482.000 \\
23 & 3.340 & 482.000 \\
24 & 3.340 & 482.000 \\
25 & 3.340 & 482.000 \\
26 & 3.340 & 482.000
\end{tabular}
\begin{tabular}{lll} 
REACH & LENGTH & LATERAL FLOM \\
& & \\
1 & 10560.00 & 0.000000 \\
2 & 10560.00 & 0.000000 \\
3 & 10560.00 & 0.000000 \\
1 & 10560.00 & 0.000000 \\
5 & 10560.00 & 0.000000 \\
6 & 10560.00 & 0.000000 \\
7 & 10560.00 & 0.000050 \\
\hline 8 & 10560.00 & 0.000000 \\
9 & 10560.00 & 0.000000 \\
12 & 10560.00 & 0.000000 \\
11 & 10560.00 & 0.000000 \\
12 & 10560.00 & 0.000000 \\
13 & 10560.00 & 0.000000 \\
11 & 10560.00 & 0.000000 \\
15 & 10560.00 & 0.000000 \\
15 & 10560.00 & 0.000000 \\
17 & 10560.00 & 0.000000 \\
18 & 10560.00 & 0.000000 \\
17 & 10560.00 & 0.000000 \\
23 & 10560.00 & 0.000000 \\
21 & 10560.00 & 0.000000 \\
22 & 10560.00 & 0.000000 \\
23 & 10560.00 & 0.000000 \\
24 & 10560.00 & 0.000000 \\
25 & 12672.00 & 0.000000
\end{tabular}
\begin{tabular}{ccccc} 
FOURTH-ORDER REGRESSION COEFF. FOR AREA \\
O-TH & \(15 T\) & 2NO & 3RD & 4TH \\
-32.01320 & 84.60530 & 5.47340 & 0.91215 & -0.00452
\end{tabular}

FOURTH-ORDER REGRESSION COEFF. FOR WETTED PERIMETER
\begin{tabular}{ccccc}
\(0-\) TH & \(15 T\) & \(2 N 0\) & \(38 D\) & 4 TH \\
10.082 .35 & 57.67010 & -4.90130 & 0.34127 & -0.00601
\end{tabular}
\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|}
\hline & \multicolumn{3}{|c|}{upstream} & \multicolumn{3}{|c|}{midstream} & \multicolumn{2}{|l|}{downstream} & \multirow[b]{2}{*}{VELCCITY} \\
\hline TINE(HR.1 & dis charge & DEPTH & velocity & DISCHARGE & DEPTH & velocity & DISCHARGE & DEPTH & \\
\hline 1.00 & 501.643 & 2.854 & 1.807 & 565.813 & 3.340 & 1.640 & 546.570 & 3.317 & 1.618 \\
\hline 2.00 & 521.285 & 2.895 & 1.845 & 606.491 & 3.340 & 1.758 & 618.323 & 3.296 & 1.850 \\
\hline 3.00 & 540.928 & 3.060 & 1.767 & 628.423 & 3.340 & 1.821 & 656.780 & 3.327 & 1.921 \\
\hline 4. 00 & 560.571 & 3.148 & 1.774 & 644.660 & 3.349 & 1.860 & 669.159 & 3.343 & 1.937 \\
\hline 5.00 & 583. 214 & 3.132 & 1.845 & 665.587 & 3.384 & 1.891 & 664.443 & 3.346 & 1.920 \\
\hline 6.00 & 599.857 & 3.199 & 1.859 & 699.077 & 3.460 & 1.923 & 651.312 & 3.344 & 1.913 \\
\hline 7.00 & 619.500 & 3.292 & 1.842 & 740.050 & 3.560 & 1.955 & 659.136 & 3.343 & 1.908 \\
\hline -. 00 & 639.142 & 3.331 & 1.863 & 778.815 & 3.652 & 1.980 & 657.959 & 3.342 & 1.906 \\
\hline 9.00 & 659.785 & 3,377 & 1.883 & 806.762 & 3.719 & 1.996 & 657.274 & 3.341 & 1.904 \\
\hline 10.00 & 679.428 & 3.449 & 1.882 & 826.363 & 3.765 & 2.007 & 656.856 & 3.341 & 1.903 \\
\hline 11.00 & 698.071 & 3.501 & 1.892 & 837.034 & 3.793 & 2.013 & 656.619 & 3.340 & 1.903 \\
\hline 12.00 & 717.714 & 3.539 & 1.914 & 840.218 & 3.799 & 2.013 & 656.492 & 3.340 & 1.902 \\
\hline 13.00 & 737.357 & 3.589 & 1.927 & 8.38 .264 & 3.794 & 2.012 & 656.418 & 3.340 & 1.902 \\
\hline 14.00 & 757.000 & 3.652 & 1.928 & 042.535 & 3.801 & 2.015 & 658.431 & 3.340 & 1.502 \\
\hline 15.00 & 1965.249 & 5.979 & 2.306 & 064.539 & 3.847 & 2.032 & 656.830 & 3.341 & 1.903 \\
\hline 16.00 & 3173.499 & 7. 102 & 2. 760 & 908.835 & 3.939 & 2.061 & 658.747 & 3.346 & 1.504 \\
\hline 17.00 & 4381.742 & 8.519 & 2.691 & 983.721 & 4.092 & 2. 106 & 664.227 & 3.359 & 1.909 \\
\hline 18.00 & 5587.992 & 9.264 & 2.907 & 1095.981 & 4.311 & 2.166 & 676.106 & 3.387 & 1.920 \\
\hline 19.00 & 6113.992 & 9.512 & 3.037 & 1235.558 & 4.570 & 2.230 & 695. 399 & 3.436 & 1.934 \\
\hline 20.00 & 6649.992 & 9.790 & 3.167 & 1365.632 & 4.007 & 2.277 & 122.138 & 3.502 & 1.952 \\
\hline 21.00 & 7179.992 & 10.503 & 2. 913 & 1466.053 & 4.983 & 2.307 & 752.192 & 3.575 & 1.972 \\
\hline 22.00 & 7703.992 & 10.741 & 3.008 & 1591.610 & 5.199 & 2.329 & 779.893 & 3.644 & 1.996 \\
\hline 23.00 & 9356.992 & 11.755 & 2.948 & 2041.995 & 5.766 & 2.529 & 802.375 & 3.703 & 1.998 \\
\hline 24.00 & 10999.990 & 12.304 & 3.264 & 2905.542 & 6.715 & 2.766 & 821.332 & 3.748 & 2.008 \\
\hline 25.00 & 13994.990 & 13.543 & 3. 307 & 3825.629 & 7.613 & 2.906 & 834.800 & 3.781 & 2.014 \\
\hline 26.00 & 16983.990 & 14.924 & 3.267 & 4662.191 & 8.349 & 2.980 & 850.800 & 3.818 & 2.023 \\
\hline 27.00 & 19984.990 & 16.389 & 3.006 & 5413.211 & 8.958 & 3.019 & 873.799 & 3.876 & 2.037 \\
\hline 28.00 & 22979.990 & 17.102 & 3.235 & 6137.531 & 9.494 & 3.056 & 925.853 & 3.974 & 2.072 \\
\hline 29.00 & 21522.490 & 16.609 & 3.165 & 6763.393 & 10.000 & 3.035 & 1002.750 & 4.117 & 2.116 \\
\hline 30.00 & 20064.990 & 16.373 & 3. 041 & 7510.277 & 10.505 & 3.048 & 1125.035 & 4.336 & 2.189 \\
\hline 31.00 & 18607.490 & 16.201 & 2.917 & 8483.727 & 11.031 & 3.113 & 1323.278 & 4.697 & 2.280 \\
\hline 32.00 & 17159.000 & 15.590 & 2.981 & 9468.043 & 11.582 & 3.137 & 1680.693 & 5.246 & 2.430 \\
\hline 33.00 & 15692.500 & 14.766 & 3.063 & 10539.300 & 12.153 & 3.152 & 2272.846 & 5.902 & 2.642 \\
\hline 34.00 & 14235.000 & 14.530 & 2.861 & 11656.490 & 12.732 & 3.154 & 2980.063 & 6.748 & 2.009 \\
\hline 35.00 & 12777.500 & 13.486 & 2.998 & 12632.130 & 13.303 & 3.107 & 3649.134 & 7.374 & 2.902 \\
\hline 36.00 & 1132).000 & 12.740 & 3. 059 & 13835.280 & 13.835 & 3.119 & 4363.063 & 8.068 & 2.973 \\
\hline 37.00 & 10542.490 & 12.447 & 3.033 & 14774.170 & 14.301 & 3.091 & 4932.383 & 8.595 & 3.023 \\
\hline 38.00 & 9764.996 & 12.134 & 2. 887 & 15524.680 & 14.671 & 3.069 & 5637.734 & 9.049 & 3.084 \\
\hline 39.00 & 8987.496 & 11.436 & 3.032 & 15938.890 & 14.932 & 3.032 & 6221.613 & 9.452 & 3.104 \\
\hline 40.00 & 8203.996 & 11.010 & 3.070 & 16219.390 & 15.060 & 3.019 & 6783.055 & 9.960 & 3.106 \\
\hline 41.00 & 7432.496 & 10.601 & 2.935 & 16113.430 & 15.077 & 2.995 & 7471.414 & 10.340 & 3.135 \\
\hline 42.00 & 6655.000 & 9.830 & 3.080 & 15850.500 & 15.002 & 2.979 & 8095.895 & 10.741 & 3.139 \\
\hline 43.00 & 5877.500 & 9.427 & 3.017 & 15419.780 & 14.050 & 2.963 & 6579.770 & 11.098 & 3.143 \\
\hline 44.00 & 5103.000 & 8. 925 & 2.848 & 14851.550 & 14.635 & 2.952 & 9299.203 & 11.42\% & 3.171 \\
\hline 45.00 & 4852.496 & 8.695 & 2.840 & 14168.860 & 14.360 & 2.939 & 9919.813 & 11.798 & 3.152 \\
\hline 46.00 & 4604.996 & 8.461 & 2.829 & 13430.570 & 14.035 & 2.933 & 10523.770 & 12.111 & 3.172 \\
\hline 47.00 & 4357.496 & 8.045 & 2.935 & 12682.270 & 13.658 & 2.944 & 11093.890 & 12.383 & 3.180 \\
\hline 40.00 & 4109.996 & 0.015 & 2.830 & 11765.590 & 13.226 & 2.932 & 11604.490 & 12.668 & 3.159 \\
\hline 49.00 & 4009.399 & 7.827 & 2.881 & 10807.480 & 12.737 & 2.923 & 12074.970 & 12.882 & 3.179 \\
\hline 50.00 & 3908.799 & 7.836 & 2. 805 & 9824.559 & 12.194 & 2.919 & 12464.820 & 13.103 & 3.165 \\
\hline
\end{tabular}
\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|}
\hline 51.00 & 3809． 299 & 7.673 & 2.848 & 8829.344 & 21.612 & 2.910 & 12821.030 & 13.280 & 3． 163 \\
\hline 52.00 & 3707.600 & 7.614 & 2.801 & 7882.184 & 11.006 & 2.906 & 13075．450 & 13.444 & 3.144 \\
\hline 53.00 & 3607.000 & 7.509 & 2．807 & 7011.211 & 10.412 & 2.896 & 13275.830 & 13.567 & 3.130 \\
\hline 54.00 & 3505.400 & 7.399 & 2.793 & 6257．238 & 9.857 & 2.887 & 13356.730 & 13.661 & 3.10 ＊ \\
\hline 55.00 & 3405.800 & 7.324 & 2.771 & 5638.133 & 9.373 & 2.875 & 13365.330 & 13.710 & 3.083 \\
\hline 56．00 & 3305.200 & 7.195 & 2.772 & 5154.316 & 8.976 & 2.862 & 13246.800 & 13.722 & 3.053 \\
\hline 57.00 & 3204.599 & 7.123 & 2.740 & 4790.555 & 8． 666 & 2.848 & 13044.680 & 13.684 & 3.026 \\
\hline 58.00 & 3104．000 & 6.995 & 2．740 & 4520.584 & 8.431 & 2.833 & 12714.070 & 13.604 & 2．991 \\
\hline 59.00 & 3067．631 & 7.019 & 2.698 & 4317.703 & 8． 253 & 2.819 & 12294．370 & 13.472 & 2.957 \\
\hline 60.00 & 3031． 262 & 6．940 & 2.720 & 4158.090 & 8.112 & 2.805 & 11749.540 & 13.293 & 2.914 \\
\hline 61.00 & 2994．894 & 6.933 & 2.693 & 4036.201 & 7.995 & 2.799 & 11076.960 & 13.000 & 2.875 \\
\hline 62.00 & 2958．526 & 6.897 & 2.680 & 3919.003 & 7.887 & 2.787 & 10216.430 & 12.505 & 2．841 \\
\hline 63.00 & 2922.157 & 6.872 & 2.664 & 3306．999 & 7.784 & 2.775 & 9293.941 & 12.085 & 2．818 \\
\hline 64.00 & 2885．788 & 6.835 & 2.654 & 3699．744 & 7.683 & 2.762 & 6346．797 & 11.521 & 2.799 \\
\hline 65.00 & 2849．420 & 6． 802 & 2.643 & 3599.432 & 7.588 & 2.750 & 7437．984 & 10.922 & 2.767 \\
\hline 66.00 & 2813.052 & 6． 764 & 2.635 & 3509.756 & 7.501 & 2.739 & 6611．473 & 10.326 & 2.718 \\
\hline 67.00 & 2776．683 & 6． 726 & 2.626 & 3431．770 & 7.425 & 2.729 & 5900.816 & 9.765 & 2．74 \\
\hline 68.00 & 2743．315 & 6.686 & 2.619 & 3365.245 & 7.359 & 2．720 & 5319.344 & 9.271 & 2．772 \\
\hline 69.00 & 2703.947 & 6.617 & 2．614 & 3302.410 & 7.301 & 2.707 & 4865.309 & 8.883 & 2.764 \\
\hline 70.00 & 2667．579 & 6.581 & 2.626 & 3251.927 & 7.252 & 2.699 & 4543.902 & 8.578 & 2．764 \\
\hline 71.00 & 2632． 210 & 6.537 & 2.626 & 3207．462 & 7.207 & 2.692 & 4283.523 & 8.329 & 2.756 \\
\hline 72.00 & 2594．841 & 6.490 & 2.619 & 3166.590 & 7.166 & 2.686 & 4083.307 & 0.130 & 2．750 \\
\hline 73.00 & 2558．473 & 6.451 & 2.612 & 3127.312 & 7.126 & 2.680 & 3915．717 & 7.963 & 2.741 \\
\hline 74.00 & 2522.105 & 6.402 & 2.607 & 3088． 362 & 7.086 & 2.673 & 3778．913 & 7.823 & 2.734 \\
\hline 75.00 & 2485．737 & 6.368 & 2.587 & 3058．964 & 7.044 & 2.661 & \(4-371.3 .560\) & 7.697 & 2.121 \\
\hline 76.00 & 2449．368 & 6.286 & 2.565 & 3023.510 & 7.002 & 2.659 & 3616.291 & 7.608 & 2．720 \\
\hline 77.00 & 2413．000 & 6.251 & 2.602 & 2985.725 & 6.958 & 2.655 & 3517.497 & 7.519 & 2.708 \\
\hline 78．00 & 2376．632 & 6.264 & 2.554 & 2946．345 & 6.914 & 2.649 & 3434．561 & 7.437 & 2．700 \\
\hline 7500 & 2343.262 & 6.131 & 2.607 & 288．7．579 & 6.870 & 2．643 & 3343．439 & 7.358 & 2.699 \\
\hline S 0.00 & 2303．894 & 6.186 & 2.549 & 2876.619 & 6.826 & 2.646 & 3284.313 & 7.304 & 2.683 \\
\hline 81.00 & 2267．526 & 6.063 & 2.565 & 2810.322 & 6.781 & 2.631 & 3233.575 & 7.237 & 2.691 \\
\hline 82.00 & 2231．157 & 6.037 & 2.547 & 2797.600 & 6.736 & 2.634 & 3176.162 & 7.191 & 2.673 \\
\hline 83.00 & 219．789 & 6.013 & 2.520 & 2746．666 & 6.691 & 2.617 & 3157.211 & 7.144 & 2． 675 \\
\hline 84.00 & 2158． 120 & 5.924 & 2．552 & 2688.710 & 6.644 & 2.609 & 3092.535 & 7.097 & 2.650 \\
\hline 85.00 & 2122.052 & 5.940 & 2.498 & 2674.288 & 6.599 & 2.611 & 3060.258 & 7.066 & 2.657 \\
\hline 86.00 & 2065．634 & 5．826 & 2.534 & 2611.124 & 6.556 & 2.594 & 3020．498 & 7.014 & 2.650 \\
\hline 87.00 & 2049.315 & 5.837 & 2．484 & 2599.255 & 6.515 & 2.596 & 2971.886 & 6.975 & 2.643 \\
\hline 88.00 & 2012．947 & 5.742 & 2．507 & 2539．855 & 6.473 & 2.580 & 2939.010 & 6.926 & 2.647 \\
\hline 09.00 & 1976．578 & 5.732 & 2.460 & 2526.038 & 6.430 & 2.581 & 2889.070 & 6.885 & 2.629 \\
\hline 90.00 & 1943． 210 & 5.640 & 2．495 & 2463.275 & 6.383 & 2.564 & 2856.548 & 6.838 & 2.632 \\
\hline 91.00 & 1903．842 & 5.653 & 2．431 & 2448．794 & 6.336 & 2.565 & 2808.084 & 6.796 & 2.614 \\
\hline 92.00 & 1867.473 & 5.560 & 2．440 & 2389．665 & 6.289 & 2.536 & 2787．404 & 6.745 & 2．61 2 \\
\hline 93.00 & 1834.105 & 5.488 & 2．435 & 2357．700 & 6.243 & 2.534 & 2744．664 & 6.709 & 2.607 \\
\hline 94.00 & 1796．736 & 5.456 & 2． 428 & 2320.382 & 6.197 & 2.526 & 2711.566 & 6.673 & 2.599 \\
\hline 95.00 & 1754.368 & 5．410 & 2．395 & 2283.573 & 6.151 & 2.510 & 2674.237 & 6.629 & 2.592 \\
\hline 96.00 & 1722．000 & 5.324 & 2.436 & 2238.257 & 6.103 & 2.509 & 2622．735 & 6.583 & 2.586 \\
\hline
\end{tabular}

\section*{APPENDIX G}

COMPUTER PROGRAM LISTING FOR WIDFEM

 －
IMPLICIT DIFFUSION FINITE ELEMENT METHOD SOLYED＊
＊BY ITERATIVE NEMTONGRAPHSOM TECHNIOUE＊
 ＊
* \(\bullet\) DEFINITION OF TERMS ••• *
*VARIABLES UNTTS ARE AS FOLLONS: TIMEUSECIOLENGTHIFTI *
    GARIABLES UNITS ARE AS FOLLONS\& TIMERSECJILENGTHIFTS *
    *DETPH(FTI, VELOCITY\&FT. PER SECI,DISCHARGEGCFSI *
    *ACC IS THE ACCELERATION OF GRAVITY,32.2FT. PER SEC PER SEC
    *ACF, PCF ARE POLYNO. COEFF. FOR AREA E WETTED PERIMETER
    *AN:PN ARE THE AREA E WETTED PERIMETER OF FLOM RESPECTIVELY
    *HYD IS THE RATE OF CHANGE OF HYDR,RADIUS MITH DEPTH
    *OSTR IS THE UPSTREAM INFLOW DISCHARGE HYDROGRAPH
    *QLAT IS THE LATERAL FLON TERM,FT• PER SEC
    \#OFL IS THE LATERAL INFLOW HYDROGRAPH AT FLINT CREEX
    *SO IS THE CONSTANT CHANMEL SLOPE
    *RN IS THE MANNING ROUGHNESS COEFF. *
    *XL IS THE NODAL SPACING.
    *YO IS THE INITIAL UNIFORM NORMAL DEPTH. *
    *VO IS THE INITIAL UNIFORM NORMAL VELOCITY *
    *N1 IS THE TOTAL NUMBER OF NODES *
    *T IS THE TIME STEP (SECONDS)
*
    *TSUM IS THE ENTIRE FLODO DURATION IN SECONDS.
    \#TSR IS THE TIME FOR UPSTREAM INFLOW HYOROGRAPM
    -TFL IS THE TIME FOR LATERAL INFLOM HYOROGRAPH AT FLINT *
    *TPRINT IS THE TIME FOR INITIAL PRINTING ISECONDS) *
    *TTA IS THE INCREMENTAL PRINTING TIME 《SECONDS: *
    *- - IF TIME STEP IS GREATER THAN TTA PRINTING WILL BE - -
    * PERFORMED AT THE INCREMENT OF.THE TIME STEPgTーm *
    -ImAX IS THE MAX. ITERATION LIMIT
    *NT IS THE NUMBER OF POINTS FOR UPSTREAM INFLOW HYOROGRAPH *
    *NTP IS THE UPSTREAM INFLOH HYDROG POINT FOR THE PEAK FLOW *
    *JGOJGP ARE SAME AS NT, NTP FOR LATERA INFLOM HYORDG FOR FLINT*
    FJORD IS THE ORDER OF POLYNOMIAL EQ. FOR AREA E M. PERIMETER *
    *DETA IS THR CDMVERGENCE CRITERIA FOR DEPTH
    *TETHA IS THE TIME WEIGHTIMG FACTOR. *
    *YN IS THE CALCULATED DETH OF FLOW *
    FYN IS THE CORRESPONDING YELOCITY OF FLOM .
    *日SR IS THE JACOBIAN MATRIX OF DIMENSION (2NI X 6
    \#CXVI E EXV2 ARE THE (NIXI) COLUMN VECTORS EVALUATED AT *
    * (I-TETHA) \(E\) TETHA RESPECTIVELY.
        -
    -LDIM,LOIN ARE THE VARIABLE DIMENSIONING PARAMETERS
    *READW IS THE SUBPROGRAM TO READ E ECHOE CHECK INPUT DATA
    *GEOMTR IS THE SUBPROGRAM TO UPDATE FLOW AREA, NETTED,
    * MANNING*S ROUGHNESS COEFF. E RATE OF CMANGE HYDR. RADIUS
    *

DIMENS ION ACF（5），AN（26），BSR（52，6），CSV1（52），CSV2（52），PCF（5）： 1 DTPH（26），HYD\｛26！：QLAT（25），QRE（26），PN（26），VELY（26），
        --ФREAD E ECHOE-CHECK INPUT DATA FROM SUBPROGRAM.
    LDN = LDIN - 1
    CALL READW(ACF, DETA,DETY, QLAT, QRE, PCF, OSTR,TSR,NT,
    1 TPRINT, TTA, TSUM,T,TETHA,SO,IMAX,XL,YO,N1,NR,LP,LDN,LDIN!
    JSIZE = TSUM/T
    \(M S=J G+1\)
    MC \(=\) NTP +1
\(C\)
\(c\)
C
    TWRIT = TPRIMT
    JSTP \(=N_{1}-2\)
    - \(-C\) CALULATE INITIAL GEOMETRIC PARAMETERS FROM SUBPROGRAM.
    CALL GEOMTRIACF,AN, HYD, QRE, PCF, PN, YO, RN, NI, LDIM
    \(P R=2.13\).
    DO 100 J=1,N1
    100 VOKJI \(=\) QREUJMAN(J)
        \(00=Q R E(1)\)
        WRITE LP:140)
        140 FORMAT \(1 / / / / 28 x,{ }^{*} U P S T R E A M *, 24 X,{ }^{\circ}\) MIDSTREAM*, \(24 X,{ }^{\circ}\) DOMNSTREAM*)
    MRITE LP,145)


        \(23 X,{ }^{\circ} \mathrm{VELOCITY}\) •
\(C\)
\(C\)
\(C\)
    150 VN(K) \(=\) VORKI
    150 VN(K) \(=\) VORKI
    FACI = TETHA*T
    FAC2 = 1. - TETHAD*T
        DO \(155 L=1, N 1\)
        155 RN\{L) \(=\) RN(L)* \(2 / 2.2082\)
C
C - - UPDATE THE LATERAL INFLOM HYOROGRAPH FOR REACH, 7
        TCK = TIME - TFLIJG1
        IF(TCX) \(160,160,200\)
    \(16000180 \mathrm{KC}=2, J G\)
        IFITIME-TFL(KC);190:190:180
    180 CONTINUE

    1 (TFL(KC)-TFL(KC-1) \#\# (TIME - TFLSKC-1)
        GOTO 320
    200 IFITCX - TFLIJGPI; \(220,220,200\)
    220 DO \(230 \mathrm{KS}=M S, J G P\)
        IF\&TIME - TFLEKSI1240,240.230
    230 CONTINUE
    240 QFM \(=\) QFL(KS-1) - \(\operatorname{OFL}(K S-1)\) - QFLSKSI)

    OIMENSION OFL(7), QSTR(11),TFL\{7), TSR(11), RN(26)

    DATA TFL/O.0.43200., 50400.957600..86400.. \(136800.9345600 . /\)
    DATA ACC,LDIM,LOINsNR,LP,LK/32.2,52,26:5:2*6/
    DATA NT, NTP/11,6/
7
8
    c
10
10
12
13
14
15
16
17
18
19
20
20
22
22
23
24
25
26
    \(c\)
\(C\)
\(C\)
    - - USE. INITIAL VELOCITY E DEPTH OF FLOW AS GUESS VALUES
    TO INITIATE SIMULATION.
    DO \(150 \mathrm{~K}=1, \mathrm{NI}\)
    YN(K) \(=\) YOAK)
    -- SET LOOP FOR TIME SIMULATION.
    \(Q R=00\)
    DO 900 JL =1, JSIZE
    TIME = FLOATSNLまT


88 89 90 91 92 93
```

        YB3= MAX1{ABS{YB2),ABS{YN&J\)!
        VB3 = MAXI(ABS(VB2),ABS(VNSJ)\)
        IF{YB3 .LE. O.0.OR. VB3 .LE. 0.0JG0 TO 570
        YERROR = YB1/YB3
        VERRCR = YB1/VB3
        IF\YERROR -LE. DETA -AND. VERROR -LE. DETYJJERR=JERR+1
        560 CONTINUE
    C
    C}57
    -\infty SWITCH CURRENT VALUES OF DEPTH OF FLOM TO OLD ONES.
    5 8 0
    C
    C
    C
    585 CALL JACOBIIBSR,FACI,YN,VN,XL,QLAT,QR,RN,
            1 AN,PN, QRE,HYD,SO,ACC,LDIM,LDN,LDIN,N1,&K)
            LUP = 0
        590 CONTINUE
    C OD DATE DEPTHS & VELOCITIES OF PREYIOUS TIME STEP.
        600 DO 680 J=1,N1
        YO(d) = VNR\!
        YO&d) = YNPJS
        680 CONTIMUE
    C
    700 TPRINT = TPRINT + TTA
    TM = IIME/TMRIT
    00710 J = 1,N1
    QRE{J) =AN{J\*VN(J)
        710 CONTINUE
        WRITE(LP,720)TM,QRE(1):YN(1), YN(1), QRE(13), YN(13), VN(13),
        QRE(N1), YN(N1),VN{N1)
    720 FORMAT( 2X,F10.2,5X,3F10.3.5X,3F10.3,5X,3F10.3)
    750 IFETIME - TSUMI900,950.950
    C -- ADVANCE THE TIME STEP.
    900 CONTINUE
    920 WRITEILP,930I
    930 FORMAT////10X,*MAX. ITERATION LIMIT EXCEEDED*':
    950 STOP
    END
    ```

```

    C SUBPRTGRAM TO READ AND ECHDE INPUT OATA *
    C SUBPRJGRAM TO READ AND ECHOE INPUT DATA
        *
    C *
            SUBROUTINE READMSASF,DETA,DETV,QLAT,QRE,PSF,OSTR,TSR,
            1 NT,TPRINT,TTA,TSUM&T,TETHA,SO,IMAX,XL,YO,NI,NR&LP&LDN,LDIN&
    C *
    ```

```

                DIMENSION ASF(5), PSF(5):QLAT(LONI, QRE|LDIN), XLILDNI, YOILDINSO
            1 OSTR(NTJITSR&MT)
    C
    C ---READ TIME PARAMETERS.
        REAO&NR,SOITPRINT:TTA,TSUAOT
    ```


\begin{tabular}{|c|c|c|}
\hline 188 & \multirow[t]{2}{*}{} & \multirow[t]{2}{*}{} \\
\hline 189 & & \\
\hline 190 & &  \\
\hline 191 & & G1 = XL(1)*RX(1)*FAC*(2.*VX(1)**2/RY1 + vxi2)**2/RY2) \\
\hline 192 & &  \\
\hline 193 & & S1 = 3- \#FAC*XL (1)*S0 \\
\hline 194 & & IF\&JSM TCH - 1) 50,50,80 \\
\hline 195 & 50 &  \\
\hline 196 & & B = (XL(1) - FAC* (2, \#VX(2)t SPK) ) \#YX(2) \\
\hline 197 & & CXV(1) \(=A+8+C\) \\
\hline 198 & & CXY(2) \(=\) S1-G1-P1 \\
\hline 199 & & G0 1090 \\
\hline 200 & 80 & CXV(1) \(=A+8-C\) \\
\hline 201 & & CXV(2) \(=\) G1 + P1-S1 \\
\hline & C & --INTERIOR NOOAL CALCULATION. \\
\hline 202 & 90 & \(00200 \downarrow=1\) LSTP \\
\hline 203 & & \(K=K+2\) \\
\hline 204 & & \(m=K+1\) \\
\hline 205 & & RO1 \(=(A \times(J) / P \times(J)) * * P S\) \\
\hline 206 & & RO2 \(=(A \times(J+1) / P \times(J+1)\rangle\) \#*PS \\
\hline 207 & & ROJ \(=(A \times(J+2) / P \times(J+2): * * P S\) \\
\hline 208 & &  \\
\hline 209 & &  \\
\hline 210 & &  \\
\hline 211 & &  \\
\hline 212 & &  \\
\hline & 1 &  \\
\hline 213 & &  \\
\hline 214 & &  \\
\hline 215 & & 1FIJSMTCH 1)100,100,150 \\
\hline 216 & 100 &  \\
\hline 217 & &  \\
\hline 218 & &  \\
\hline 219 & &  \\
\hline 220 & & CXVAM) \(=51-\mathrm{HI}\) - P1 \\
\hline 221 & & GO 10 200 \\
\hline 222 & 150 & Cxy(k) = 11 + 8i + CI OI \\
\hline 223 & & CXY(M) \(\quad\) CI + PI SI \\
\hline 224 & 200 & COATINUE \\
\hline & \(C\) & - - OOMSSTREAM NCDAL CALCULATIOA. \\
\hline 225 & & HT1 = AX(A1-1)/fX(A1-1) \\
\hline 226 & & HT2 = AXANIMPX(N1) \\
\hline 227 & & RYJ \(=\) HT1**PS \\
\hline 228 & & RY4 \(=\) M T2**PS \\
\hline 229 & &  \\
\hline 230 & &  \\
\hline 231 & &  \\
\hline 232 & &  \\
\hline 233 & &  \\
\hline 234 & & SN = 3-*FAC*XL(A1-1)*S0 \\
\hline 235 & & IFIJSWTCH - 13300,3C0.400 \\
\hline 236 & 300 &  \\
\hline 237 & &  \\
\hline 238 & &  \\
\hline 239 & & CXY(LOIM) \(=\) SN-GN-PN \\
\hline 240 & & 6010560 \\
\hline 241 & 400 & CXY(LOIP-1) \(=\) AN + EN-CN \\
\hline 242 & &  \\
\hline 243 & 500 & REIURN \\
\hline 244 & & EMO \\
\hline
\end{tabular}

\begin{tabular}{|c|c|c|}
\hline 288 & \multicolumn{2}{|r|}{\multirow[t]{2}{*}{```
ROA2=1X(A1)/PX(M1)
BTf(LB,1)= XL(N1-1)= FAC*{YX(N1)+2.#V\&N1-1)]
```}} \\
\hline 289 & & \\
\hline 290 & \multicolumn{2}{|r|}{BTF（LB， 2 ）－FAC＊（2．＊YX（N1－1）＋YX\＆M1）} \\
\hline 291 & \multicolumn{2}{|r|}{} \\
\hline 292 & \multicolumn{2}{|r|}{BTF\｛L8，4）＝FAC＊（4．＊YX（N1）－Yx（A1－1））} \\
\hline 293 & \multicolumn{2}{|r|}{} \\
\hline & \multicolumn{2}{|r|}{1 RCA1＊＊P\％301} \\
\hline 294 & \multicolumn{2}{|r|}{EIf（L）IF，} \\
\hline 295 & \multicolumn{2}{|r|}{} \\
\hline & \multicolumn{2}{|r|}{1 ROA2＊＊M－3．1} \\
\hline 296 & \multicolumn{2}{|r|}{} \\
\hline 297 & \multicolumn{2}{|r|}{} \\
\hline 298 & \multicolumn{2}{|r|}{ETf《LE， \(61=0.0\)} \\
\hline 299 & \multicolumn{2}{|r|}{ETF（L）IMg \(=0.0\)} \\
\hline 300 & \multicolumn{2}{|r|}{ETK\｛LOIM，E）\(=0.0\)} \\
\hline 301 & \multicolumn{2}{|r|}{RETERN} \\
\hline 302 & \multicolumn{2}{|r|}{ENC} \\
\hline & \multicolumn{2}{|r|}{} \\
\hline & \(C\)＊ & －\({ }_{\text {－}}\) \\
\hline & \(C\)－ & － \\
\hline & \(C\)－ & ＊SUEPROGFAP TO SCLYE THE BI－TRIOIAGOMAL MATRIX＊ \\
\hline & C＊ & ＊ \\
\hline & \(C\)＊ & ＊＊ \\
\hline 303 & \multicolumn{2}{|r|}{SUEROUTIAE ETRICGICOL，VEL， \(\mathrm{YX,VX,LOIM,LDIA,NISLK)}\)} \\
\hline & \(C\)－ & －\({ }^{\text {－}}\) \\
\hline & \multicolumn{2}{|r|}{} \\
\hline 304 & \multicolumn{2}{|r|}{\multirow[t]{2}{*}{DIFENSICA VEL\｛LCIMg（K），YX（LDIN）gYX（LDIN）gCOL（LDIMI？ 1 日ETA（4），CETAC2），SAC（26，4J，GAMA（26．2）}} \\
\hline & & \\
\hline & \multicolumn{2}{|l|}{C} \\
\hline & \multicolumn{2}{|l|}{\multirow[t]{2}{*}{C \(\quad \begin{aligned} & \text {－}-P E R F C G F ~ M A T R I X ~ R E D U C T I O N ~ O P E R A T I O N . ~\end{aligned}\)}} \\
\hline 305 & & \\
\hline 306 & \multicolumn{2}{|r|}{DC \(200 \downarrow=1, \mathrm{NL}\)} \\
\hline 307 & \multicolumn{2}{|r|}{1FsJ－11100，100．150} \\
\hline 308 & 100 & BETAS1）\(=\) VELILO31 \\
\hline 309 & \multicolumn{2}{|r|}{QETA\＆2）＝VEL\｛1：4）} \\
\hline 310 & \multicolumn{2}{|r|}{BETA（3）\(=\) YEL（2，3）} \\
\hline 311 & \multicolumn{2}{|r|}{EETA（4）＝YEL（2，4）} \\
\hline 312 & \multicolumn{2}{|r|}{DETA（1）＝COL（1）} \\
\hline 313 & \multicolumn{2}{|r|}{DETA（2）＝COL 21} \\
\hline 314 & \multicolumn{2}{|r|}{ZU＝EETA（1）＊日ETA（4）－日ETA（2）\％8ETA（3）} \\
\hline 315 & \multicolumn{2}{|r|}{IFR2U．EG＊0．0）2U \(=0.002\)} \\
\hline 316 & \multicolumn{2}{|r|}{} \\
\hline 317 & \multicolumn{2}{|r|}{} \\
\hline 318 & \multicolumn{2}{|r|}{} \\
\hline 319 & \multicolumn{2}{|r|}{} \\
\hline 320 & \multicolumn{2}{|r|}{GO 10180} \\
\hline 321 & \multicolumn{2}{|l|}{\(150 K=K+2\)} \\
\hline 322 & \multicolumn{2}{|r|}{\(m=x+1\)} \\
\hline 323 & \multicolumn{2}{|r|}{} \\
\hline 324 & \multicolumn{2}{|r|}{} \\
\hline 325 & \multicolumn{2}{|r|}{} \\
\hline 326 & \multicolumn{2}{|r|}{} \\
\hline & \multicolumn{2}{|l|}{C} \\
\hline 327 & &  \\
\hline 328 & \multicolumn{2}{|l|}{\multirow[t]{2}{*}{}} \\
\hline & & \\
\hline 329 & \multicolumn{2}{|r|}{} \\
\hline 330 & \multicolumn{2}{|r|}{IFEZU．EQ． \(0.012 \mathrm{U}=0.001\)} \\
\hline 331 & \multicolumn{2}{|r|}{} \\
\hline
\end{tabular}
```

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348

```


```

    SAC(d)3) = (EETA(1:QVEL(M,S) -EETA(3)&VEL{KgS)\/LU
    ```
    SAC(d)3) = (EETA(1:QVEL(M,S) -EETA(3)&VEL{KgS)\/LU
    SAC(dy4) = (8ETA&1)&VEL(M,6) -EETA&3)&VEL&Kg6)\/2U
    SAC(dy4) = (8ETA&1)&VEL(M,6) -EETA&3)&VEL&Kg6)\/2U
    C
    C
        180 GAFASJ,1: = (BETA&4I*OETA(1)-BETA(2)*OETA(2):/2U
```

        180 GAFASJ,1: = (BETA&4I*OETA(1)-BETA(2)*OETA(2):/2U
    ```


```

    200 CONTINUE
    ```
    200 CONTINUE
    C
    C
    ---COMPUTE SOLUTIOA VIA GECURSIVE EQE
    ---COMPUTE SOLUTIOA VIA GECURSIVE EQE
        LINIT = M1-1
        LINIT = M1-1
        」K = LIm1T
        」K = LIm1T
        YX{A1) = CAMA&LDIN,1:
        YX{A1) = CAMA&LDIN,1:
        yX(N1)=GAMA(LCIN:2)
        yX(N1)=GAMA(LCIN:2)
        00 300 L E LSLIMIT
```

        00 300 L E LSLIMIT
    ```




```

        JK = JK - 1
    ```
        JK = JK - 1
        300 COATINUE
        300 COATINUE
            RETUAN
            RETUAN
            EMO
            EMO
SEMTRY
```

SEMTRY

```

APPENDIX H

SAMPLE OUTPUT FOR WIDFEM
```

TOTAL NO. OF NODES = 26
TIME STEP = 1800.000 SEC.
CHAMAEL BOTTCM SLOPE = 0.0009
TIME MEIGHTING FACTOR=0.3500
MAX. ITERATION LIMIT= 60
CONYERGENCE CRITERIA FOR DEPTH = 0.0100
CONVEREENCE CRITERIA FOR vELOCITY: 0.1000

```

\begin{tabular}{|c|c|c|}
\hline NODE & INITIAL DEPTH & INITIAL OISCHARGE \\
\hline 1 & 3. 340 & 482.000 \\
\hline 2 & 3.340 & 482.000 \\
\hline 3 & 3.340 & 482.000 \\
\hline 4 & 3. 340 & 482.000 \\
\hline 5 & 3.340 & 482.000 \\
\hline 6 & 3.340 & 482.000 \\
\hline 7 & 3.340 & 482.000 \\
\hline 8 & 3. 340 & 482.000 \\
\hline 8 & 3.340 & 482.000 \\
\hline 10 & 3.340 & 482.000 \\
\hline 11 & 3.340 & 482.000 \\
\hline 12 & 3.340 & 482.000 \\
\hline 13 & 3.340 & 482.000 \\
\hline 14 & 3.340 & 482.000 \\
\hline
\end{tabular}
\begin{tabular}{lll}
15 & 3.340 & 482.000 \\
16 & 3.340 & 482.000 \\
17 & 3.340 & 482.000 \\
18 & 3.340 & 482.000 \\
19 & 3.340 & 482.000 \\
20 & 3.340 & 482.000 \\
21 & 3.340 & 482.000 \\
22 & 3.340 & 482.000 \\
23 & 3.340 & 482.000 \\
24 & 3.340 & 482.000 \\
25 & 3.340 & 482.000 \\
26 & 3.340 & 482.000
\end{tabular}
\begin{tabular}{lll} 
REACM & LEFGTH & LATERAL FLOZ \\
& & \\
1 & 10560.00 & 0.000000 \\
2 & 10560.00 & 0.000000 \\
3 & 10560.00 & 0.000000 \\
1 & 10560.00 & 0.000000 \\
5 & 10560.00 & 0.000000 \\
6 & 10560.00 & 0.000000 \\
7 & 10560.00 & 0.000050 \\
8 & 10560.00 & 0.000000 \\
9 & 10560.00 & 0.000000 \\
10 & 10560.00 & 0.000000 \\
11 & 10560.00 & 0.000000 \\
12 & 10560.00 & 0.000000 \\
13 & 10560.00 & 0.000000 \\
14 & 10560.00 & 0.000000 \\
15 & 10560.00 & 0.000000 \\
16 & 10560.00 & 0.000000 \\
17 & 10560.00 & 0.000000 \\
18 & 10560.00 & 0.000000 \\
19 & 10560.00 & 0.000000 \\
20 & 10560.00 & 0.000000 \\
21 & 10560.00 & 0.000000 \\
22 & 10560.00 & 0.000000 \\
23 & 10560.00 & 0.000000 \\
24 & 10560.00 & 0.000000 \\
25 & 12672.00 & 0.000000
\end{tabular}
\begin{tabular}{|c|c|c|c|c|}
\hline FOURIH-CRDER & REGRESSION & N COEFF. & FOR MaEA & \\
\hline O-TH & 151 & 2N0 & 380 & 4TH \\
\hline -32.01320 & 84.60530 & 5.47340 & 0.91215 & -0.0045 \\
\hline FOURIM-CRDER & REGRESSICM & N COEFF. & FOR METTEO & FERIMETER \\
\hline O-IK & 151 & 2M0 & 380 & 47 H \\
\hline 10.0 2 235 & 57.67010 & -4.90130 & 0.34127 & -0.00601 \\
\hline
\end{tabular}












○~~No






 ~~




\section*{APPENDIX I}

COMPUTER PROGRAM LISTING FOR WIKFEM

The computer program for the weighted implicit kinematic model is for an idealized channel . It has a built-in option to route flood in a trapezoidal, triangular, or rectangular channel. For the first two geometries, the right- and left-side slopes, captioned as ZRS and ZLS should have assigned values other than zeros, except for rectangular channel. The triangular geometry will have zero width for input value.

The definition of the variables and symbols used in the computer program is provided in the comment page of the program listing. Any temporary storage variables are not included because their definitions are obvious. Instruction for the input data is provided in the MAIN program for each READ STATEMENT and is self-explanatory.

510 C
, \(\operatorname{TIME}=10040\) :
C
C *


*
IMPLICIT KINEMATIC FINITE ELEMENT METHOD SOLVED
*
* BY ITERATIVE NEWION-RAPHSON TECHNIQUE
* *

* \(\bullet\).. DEFINITION OF TERMS ...
*VARIABLES UNITS ARE AS FOLLOWS: TIME\&SECSOLENGTHIFTB *
- DEPTHIFT VELCCITYIFT PER SECI, DTSCHAPCERCES) © DEPTH\{FTJ, VELCCITY(FT. PER SEC), DISCHARGE (CFS)
-ACC IS THE ACCELERATION OF GRAVITY,32.2FT. PER SEC PER SEC - QLAT IS THE LATERAL FLOW TERM,FT. PER SEC *TLL IS THE TOTAL LENGTH OF CHANNEL REACH BEING INYESTIGATED * * SO IS THE CONSTANT CHANNEL SLOPE
-RN IS THE MANNING ROUGHNESS COEFF.
*XL IS THE NODAL SPACING.
*YO IS THE INITIAL UNIFORM NORMAL DEPTH.
*VO IS THE INITIAL UNIFORM NORMAL VELOCITY
*NI IS THE TOTAL NUMBER OF NODES
*QMAX IS THE PEAK FLOOD DISCHARGE OF INFLOW HYDROGRAPH * T IS THE TINE STEP (SECONDSI
- TMX IS THE TIME PERIOD BETMEEN OO E GMAX
-TAP IS THE TIME PERIOD AFTER OMAX UNTIL QD - TPRINT IS THE TIME FOR INITIAL PRINTING (SECONDSA -TTA IS THE INCREMENTAL PRINTING TIME \&SECONDS:
*---IF TIME STEP IS GREATER THAN TTA PRINTING HILL BE - -
* PERFCRMED AT THE INCREMENT OF THE TIME STEP:T-D. * *
*TETHA IS THE TIME WEIGHIING FACTOR.
- YN IS THE CALCULATED DETH OF FLOW
* UN IS THE CORRESPUNDING VELOCITY OF FLOW
- ZSR IS THE TRAPEZOIDAL CHANNEL RIGHT SIOE SLOPE
-ZSL IS THE TRAPEZOIDAL CHANNEL LEFT SIDE SLOPE
*ZSL IS THE TRAPEZOIDAL CHANNEL LEFT SIDE SLOPE
*IMAX IS THE MAX. ITERATION LIMIT
*DETA IS THR CONVERGENCE CRITERIA FOR DEPTH
*BSR IS THE JACOBIAN MATRIX OF DIMENSION (NIX3) \#CXVI E CXVZ ARE THE (NIXI) CDLUMN VECTORS EVALUATED AT *\{1-TETHA) E TETHA RESPECTIVELY. *LDIMgLDIN ARE THE VARIABLE DIMENSIONING PARAMETERS +

DIMENSION BSR(21:3),CSVI(21), CSV2\&21), OTPH\& 21 ) QLAT(20 de
1 VELY(21),YO(21):VO(21),YN(21), VN(21), XL(20), QRE(21)

DATA LDIN,LK/21,3/
OATA NR,LP/5,6/
C
- - -READ DATA AND ECHOE CHECK
READ\&NR, SO) TPRINT, TTA, TSUM, T, YB, ZSR,ZSL
READ\&NR 5 50)TPR
FORMATS 7F10.2)
REAOKNR, 60, OMAX,TMX, TAP, TLL, B1
```

    60 FORMATS5F10.21
        READINR, &OITETKA,DETA,IMAX,SO:RN:N1,NN
        FORMAT(2F10.4:I10.2F10.4:2I10)
        \SIZE = TSUM/T
        LDN = LDIN - 1
    c
        90
        FORMAT:IHI)
    c
    WRITEILP,100INI,T,SO,RN
    FORMATY///2OX,'TOTAL NO. OF NODES =*,I5//20X.
        1 -TIME STEP =0,F10.3,1X,*SEC.*//20X,0CHANNEL BOTTON SLOPE =0,
        2 F10.4,//20x,4MANNING ROUGHNESS COEFF. =0,F10.41
        WRITEILP,12OITETHA,IMAX,DETA,LSR,ZSL
    120 FORMATS//20X,昂IME WEIGHTING FACTOR =*&F10.4//20X,
    1. MAX. ITERATION LIMIT =4,I5//20X,
    2 - CONVERGENCE CRITERIA FOR DEPTH =0,FIO.5//20X,
        3.-TRAPEZOIDAL CHANNEL SIDE SLOPES: RIGHT =0,FIO.2:
        5X,'LEFT =':F10.2)
        WRITE(LP,125|XL(1), OLAT(1)
    ```

```

        1 CLATERAL FLOW =*,F10.3:1X,0FT.PER SEC.";
        TWRIT = TPRINT
        JSTP = NL - 2
    C --- CALGULATE THE INITIIAL MORHAL OISCHARGE & VELOCITY.
        PR = 2.13.
        CM = 1.486/RN*SORTSSO:
        2PP = SQRTI1. + 2SR**21 + SQRT(1. + 2SL**2I
        AE=81*YB + 5*VB**2*RZSSR + ZSLI
        PE = Bl + YB=ZPP
        VB=CAF(AE/PE)**PR
        OB=AE*VB
        - DO 230 J=1.NI
        YO(J)= %B
        VO(J) = VB
        ORE(J)=0B
    230 CONTINUE
        QO = ORE(1)
        WRITEILP,2501YO(1), VO& 1H:QRE{1)
    250 FORMATI//20X.0INITIAL OEPTH =*,F10.3.3X.
        1.INITIAL VEL, =0,FIO.3.3X,0 INITIAL DISCH. =0,F1O.3)
    c
        WRITEILP:140)
    ```

```

        WRITE(LP,145)
    ```

```

        | 'DISCHARGE*,4X,'DEPTH*,3X,*VELOCITY*,6X,'DISCHARGE*,4X,*DEPTH**
        2 3X,`VELOCITY`)
    C
        ---USE INITIAL vElocity e depth of flom as guesS values
        TO INITIATE SIMULATION.
        OO 280 K = 1,N1
        YN(K) = YO(X)
        VN(X) = VO(K)
    280
        CONTINUE
        FAC1 = TETHA*T
        FAC2 = (1.-TETHA)*T
    C. --- SET LOOP FOR TIME SIMULATIDN.
        QR = OO
        DO 900 JL =1gJSIIE
    ```
\begin{tabular}{|c|c|c|}
\hline 50 & C & TIME F FLOATY JLJFT \\
\hline 51 & & TD = TIME - TMK \\
\hline 52 & & IFCTDI4 50:450.480 \\
\hline 53 & 450 & QR = 00 - (QMAX - UOJ/TMX TIME \\
\hline 54 & & GOTO 520 \\
\hline 55 & 480 & IF \({ }^{\text {S }}\) - TAP)490.490.500 \\
\hline 56 & 490 & QR = QMAX - (QMAX - GOI/TAP \# TD \\
\hline 57 & & 6010520 \\
\hline 58 & 500 & QR = U0 \\
\hline 59 & 520 & continue \\
\hline & C & ---CALL SUBROUTINE TO GENERATE COLUMA VECTOR [N1 \(x\) 1) \\
\hline 60 & & JSmTCH = 1 \\
\hline 61 & &  \\
\hline & 1 & LDN,LDIN,LKOJSWTCHI \\
\hline & 6 & \\
\hline & \(C\) & ---GEMERATE JACO8LAN MATRIX. \\
\hline 62 & & LUP \(=0\) \\
\hline .63 & &  \\
\hline & \(c^{1}\) & ```
LONOLDIN,N1,LKI
-\infty- ITERATE TO CONVERGENCE FOR EACH TIME STEP.
``` \\
\hline 64 & & 00590 LL 31, IMAX \\
\hline 65 & & JSWTCH \(=2\) \\
\hline 66 & & LUP = LUP + 1 \\
\hline 67 & &  \\
\hline & 1 & LON=LDIN,LX, LSATCHI \(^{\text {d }}\) \\
\hline 68 & & DO 530 K = 1, N1 \\
\hline 69 & & CSV2(K) = CSV2イK) - CSVI(K) \\
\hline 70 & 330 & CONTINUE \\
\hline & \(c\) & --DOBTAIN SOLUTION VIA TRI-DIAGONAL SUBPROGRAM. \\
\hline \(71^{\circ}\) & &  \\
\hline 72 & & \(J S=N_{1}-1\) \\
\hline 73 & & DO \(535 \mathrm{~L}=1, N 1\) \\
\hline 74 & & DPI = YN\{L) - DTPH(L) \\
\hline 75 & &  \\
\hline 76 & &  \\
\hline 77 & &  \\
\hline 78 & 535 & CONTINUE \\
\hline & \(c\) & -- CHECK FIOR RELATIVE CONYERGENCE FOR ALL VARIABEES. \\
\hline 79 & & JERR = 0 \\
\hline 80 & & 00560 J \(=1, N 1\) \\
\hline 81 & & YB1 = ABSIOTPHCJ): \\
\hline 82 & & YB2 = YN\&J) - DTPHCJ) \\
\hline 83 & &  \\
\hline 84 & & IF(YB3 LLE. O.O)GO TO 570 \\
\hline 85 & & YERROR = Y81/YE3 \\
\hline 86 & & IFSYERROR -LE DETAJJERR = JERR * 1 \\
\hline 87 & 560 & CONTINUE \\
\hline & \[
\begin{aligned}
& C \\
& C
\end{aligned}
\] & -- SHITCH CURRENT VALUES OF DEPTH OF FLOM T.O OLD ONES. \\
\hline 88 & 570 & 00580 L=1.N1 \\
\hline 89 & & YNイL) = YN(L) - OTPHEL! \\
\hline 90 & 580 & CONTINUE \\
\hline 91 & & IF\&JERR - EQ. N1) GO TO 600 \\
\hline & \(c\) & -- CHECX IF SPECIFIED ITERATION LIMIT IS EXCEEDED. \\
\hline & \(C\) & \\
\hline \multirow[t]{3}{*}{92} & & IFIJERR -LT. N1. AMD. LL .GE. IMAXI GO TO 920 \\
\hline & \(c\) & \\
\hline & 6 & ---UPDATE THE JACOBIAN MATRIX AT EVERY 3 ITERATIOS. \\
\hline 93 & & IF6LUP - 31590.585.585 \\
\hline
\end{tabular}
\begin{tabular}{|c|c|c|}
\hline 94 & 585 &  LDN,LDIN,NI,LKI \\
\hline 95 & & LUP \(=0\) \\
\hline 96 & \(c^{590}\) & CONTINUE \\
\hline 97 & 600 & 00 6a0 J=1, \({ }^{\text {di }}\) \\
\hline 98 & & void) = VN(J) \\
\hline 99 & & YO(J) = YN(J) \\
\hline 100 & 680 & continue \\
\hline & \(c\) & --- PRINT OUT RESULTS. \\
\hline 101 & & IFITIME - TPRINT + .0038750.700.700 \\
\hline 102 & 700 & TPRINT = TPRINT + PTA \\
\hline 103 & & TM = TIMESTWRIT \\
\hline 104 & & DO \(720 \mathrm{~J}=10 \mathrm{Ns}\) \\
\hline 105 & &  \\
\hline 106 & & QREJJ = AREAFVNTJ \\
\hline 107 & 710 & continue \\
\hline 108 & &  \\
\hline & & GRESN1),YN(NLIOUN(NI) \\
\hline 109 & 720 &  \\
\hline 110 & 750 & IFITIME - TSUMD900.950.950 \\
\hline & C & --- advance the time step. \\
\hline 111 & 900 & continue \\
\hline 112 & 920 & WRITEILP,9301 \\
\hline 113 & 930 & FORMATS/I/10X.A MAX. ITERATION LIMIT EXCEEDED.*) \\
\hline 11. & 950 & Stop \\
\hline 115 & & ENO \\
\hline &  &  \\
\hline & c & SUBPROGRAM TO GENERATE COLUNMN YECTOR (NX1) \\
\hline & C & * * * * * * \\
\hline 116 & 1 & SUBROUTINE VECTREBI,CXV,FAC, QLAT,YX,VX,XL, OT. RN,N1, ZSR,ZSL,LDN:LOIN,LKOJSNTCH: \\
\hline & 6 & . \({ }^{\text {a }}\) \\
\hline & \(c\) &  \\
\hline 117 & &  \\
\hline & c & \\
\hline & \(C\) & -a UPSTREAM NOOAL CALCULATION \\
\hline 118 & & LSTP = N1-2 \\
\hline 119 & &  \\
\hline 120 & & SPK = QT/AREA \\
\hline 121 & &  \\
\hline 122 & &  \\
\hline 123 & & \(C=3 . * F A C=G L A T(1)=X L\{1)\) \\
\hline 124 & & IFSJSwTCH - 1) 50.50 .80 \\
\hline 125 & 50 &  \\
\hline 126 & & \(8=(X L(1)-F A C F(20 * V \times(2)+S P K) 1\) \#YX(2) \\
\hline 127 & & Cxv(1) \(=A+8+C\) \\
\hline 128 & & GO TO 90 \\
\hline 129 & 80 & CXV(1) \(=1+B-C\) \\
\hline & \(c\) & ---INTERIOR NODAL CALCULATION. \\
\hline 130 & 90 & 00 200 J = 1,LSTP \\
\hline 131 & & \(\kappa=\downarrow+1\) \\
\hline 132 & & \(A I=\) (XL(J) - FACE\{Vx(J+1) + 2e*vx(J)) ) =rx(J) \\
\hline 133 & &  \\
\hline 134 & . &  \\
\hline 135 & & OI = 3. FFAC=(OLAT(J) \\
\hline 136 & & IFIJSWTCH - 11100,100,150 \\
\hline
\end{tabular}


```

176
C

```

```

    OIMENSIQN STIFF{LDINOLKI,RH{LDINI:YR(LDINO
    UIMENSION G(5OIGM\SO)
    --OEEGIN TRIANGULAR REDUCTION OPERATION.
    n(1) = STIFF(1:3)/STIFF{1,2)
    G(1) = RH(1)/STIFF(1,2)
    00 100 J = 2,NI
    SAVI = STIFFIJ.2) - STIFF(d,1)*W(J-1)
    SAV2 = RH(J) - STIFF(JO1)=G(J-1)
    W(J) = STIFFIJ.3)/SAVI
    GIJ) = SAV2/SAV1
    100 CONTINUE
    C --- OETAIN SỌlUTION VIA RECURSIVE EO.
LIMIT = NI - L
K = LIMIT
YR(N1) = G(N1)
00 200 L=1sLIMI
YR(K) = G(K) - W{K)=YR{K+1)
K=K - 1
200 CONTINUE
RETURN
END
SENTRY

```

APPENDIX J

SAMPLE OUTPUT FOR WIKFEM

Two sample outputs for a rectangular channel flood routing using time steps of 300 and 600 seconds respectively are included. Input data are drawn from the example problem given by Viessman et al. (1972). The computer print-out includes the input data and the simulated flow parameters--depth of flow, velocity of flow, and the volumetric flow rate.

```

TIME STEP = 300.000 SEC.
ChanNEL BOTTOM SLOPE = 0.0015
TIME WEIGHIING FACIOR = 1.0000

```
CONVERGENCE CRITERGA fOR DEPTH = 0.01000

TOTAL NO. OF NODES \(=21\)
TIME STEP \(=600.000\) SEC.
CHANNEL BOTTGM SLGPE \(=0.0015\)
manning roughness coeff. = 0.0200

TIME WEIGHTING FACTOR = 1.0000
MAX. ITERATION LIMIT \(=50\)
CONVERGENCE CRITERIA FOR DEPTH \(=0.01000\)
TRAPEZOIDAL CHANNEL SIOE SLOPESE RIGHT = 0.00 LEFT = 0.00

NOUAL SPACING \(=528.000\) FT. LATERAL FLOK \(=0.000\) FT.PER SEC.

INITIAL DEPTH \(=6.000\) INITIAL VEL. \(=6.946\) INITIAL DISCH. \(=833.495\)
\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|}
\hline & \multicolumn{3}{|c|}{UPSTREAM} & \multicolumn{2}{|c|}{Mrostream} & \multicolumn{4}{|c|}{donnstream} \\
\hline timeimin) & discharge & DEP TH & velocity & discharge & DEPTH & velocity & discharge & depth & velocity \\
\hline 10.00 & 1417.652 & 8.802 & 8.053 & 1080.712 & 7.221 & 7.483 & 936.491 & 6.517 & 7.185 \\
\hline 20.00 & 2004.586 & 11.415 & 8.781 & 1559.296 & 9.445 & 8.253 & 1237.687 & 7.968 & 7.767 \\
\hline 30.00 & 1710.933 & 10.125 & 8.449 & 1727.718 & 10.200 & 8.470 & 1529.685 & 9.312 & 8.213 \\
\hline 40.00 & 1418.977 & 8.808 & 8.055 & 1596.775 & 9.615 & 8.304 & 2611.924 & 9.683 & 8.324 \\
\hline 50.00 & 1127.654 & 7.446 & 7.572 & 1349.300 & 8.487 & 7.949 & 1497.784 & 9.168 & 8.169 \\
\hline 60.00 & 836.238 & 6.014 & 6.953 & 1073.936 & 7.189 & 7.470 & 1280.156 & 6.166 & 7.838 \\
\hline
\end{tabular}

APPENDIX K

COMPUTER CPU TIME AND COST

The numerical computation associated with each of the flow models; EKFEM, WIKFEM, WIDFEM, and WICFEM respectively is a direct function of computer CPU time and cost. Iterative solution algorithm with a prolonged convergence will translate into enormous computer CPU time and cost.

Comparisons of models and their corresponding CPU time and cost are presented below for a given time weighting factor, time step, channel geometric, and hydraulic data. Simplified models are expected to have less CPU time and cost as compared to the complete flow model.

\section*{Models Versus CPU Time and Cost}
a) Idealized Channel
\begin{tabular}{|c|c|c|c|c|}
\hline Time Weighting & \multicolumn{4}{|l|}{Factor, \(\theta=0.55\)} \\
\hline \multicolumn{5}{|c|}{Time step, \(\Delta t\)} \\
\hline MODELS & \multicolumn{2}{|l|}{60 seconds} & \multicolumn{2}{|l|}{300 seconds} \\
\hline WIKFEM & \[
\begin{gathered}
\mathrm{CPU}^{+} \\
\mathrm{COST}^{F}
\end{gathered}
\] & \[
\begin{aligned}
& =3.59 \\
& =0.91
\end{aligned}
\] & CPU
COST & \[
\begin{aligned}
& =1.44 \\
& =0.40
\end{aligned}
\] \\
\hline \multirow[t]{2}{*}{WIDFEM} & CPU & \(=9.15\) & CPU & \(=2.67\) \\
\hline & COST & \(=2.23\) & COST & \(=0.69\) \\
\hline \multirow[t]{2}{*}{WICFEM} & CPU & \(=10.68\) & & \(=4.54\) \\
\hline & COST & \(=2.60\) & COST & \(=1.14\) \\
\hline EKFEM & \multicolumn{4}{|l|}{\[
\begin{aligned}
\Delta T & =2 \text { seconds } \\
\text { CPU } & =20.87 \\
\text { COST } & =5.01
\end{aligned}
\]} \\
\hline
\end{tabular}

\footnotetext{
+ Unit of CPU time is seconds
F Unit of cost is dollars
}
b) Natural Channel
\begin{tabular}{|c|c|c|c|c|}
\hline Time Weighting & \multicolumn{4}{|l|}{Factor, \(\theta=0.55\)} \\
\hline & \multicolumn{4}{|l|}{Time step, \(\Delta t\)} \\
\hline & 900 & conds & 1800 & seconds \\
\hline \multicolumn{5}{|l|}{MODELS} \\
\hline \multirow[t]{2}{*}{WIDFEM} & CPU & \(=72.90\) & & \(=52.58\) \\
\hline & COST & \(=17.35\) & COST & \(=12.53\) \\
\hline \multirow[t]{2}{*}{WICFEM} & CPU & \(=86.50\) & & \(=66.94\) \\
\hline & COST & \(=20.57\) & COST & \(=15.94\) \\
\hline
\end{tabular}

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\section*{Thesis: FINITE ELEMENT MODELING OF STREAMFLOW ROUTING}

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Membership in Honorary Societies: Chi Epsilon, Tau Beta Pi.```


[^0]:    The author was not able to discover any documentation of FEM in time derivatives applied to unsteady flows in open channelmodeling.

[^1]:    * The author observed very astonishingly, the constant use of finite differencing for time derivatives and finite element for space derivatives (mostly linear basis functions). This seems to explain the formidable difficulties associated with a complete finite element error analysis of unsteady non-linear hyperbolic equations.

[^2]:    $\pm$ Froude number is a dimensionless flow parameter utilized to characterize the state of flow. If the Froude number is less, equal, or greater than unity, the flow is subcritical, critical, and supercritical, respectively.

[^3]:    * Direct solution algorithm for a linearized bi-tridiagonal matrix equation stored in compact (2Nx6) matrix, where $N$ is total number of nodes, was originally developed by Douglas et al. (1959) and is presented in Appendix B.

