# A MODIFIED EBULLIOMETER FOR MEASUREMENT

## OF INFINITE DILUTION ACTIVITY

## COEFFICIENTS

By

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Thesis Approved:

Thesis Adviser C Dean of the Graduate College

#### PREFACE

An existing design of an ebulliometer was modified by the addition of a vane-type magnetic pump to improve the performance of the ebulliometer at low pressures. The modified apparatus was used to determine the activity coefficients at infinite dilution of seven binary systems at three operating pressures each by the technique known as differential ebulliometry. The experimental results can be used to determine predictive model parameters which are of significant interest because of their ease of use and general applicability in phase equilibrium calculations.

I wish to express my very sincere appreciation to my advisor, Dr. Robert L. Robinson, Jr., for his constant support throughtout the course of this study. I thank Mr. M. Wayne Adkins for his contribution to this work through the construction of all glassware including the modified ebulliometer. I thank Mr. Heinz Hall for his contribution in fabricating the magnetic vane-type pump. I would also like to thank Mrs. Kelly Whitfield for her contribution in typing the entire thesis.

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## LIST OF SYMBOLS

Symbol

А	Rackett Equation Constant
A <sub>1</sub> , A <sub>2</sub> , A <sub>3</sub>	Antoine Equation Constants
К <sub>і</sub>	equilibrium ratio, y/x, for component "i"
m	mass
MW	molecular weight
n	number of moles
Ρ	pressure
P <sub>i</sub> s	vapor pressure of component "i"
P <sub>c</sub>	critical pressure
R	universal gas constant
Т	temperature
T <sub>c</sub>	critical temperature
V	volume
v	specific volume
x	liquid mole fraction
у	vapor mole fraction
Z	overall mole fraction

## Greek symbols

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Ϋ́	infinite dilution activity coefficient
φ	fugacity coefficient
μ	chemical potential

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Λ	Wilson	Equation	parameter
λ	Wilson	Equation	parameter

# Subscripts

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1	solute
2	solvent
C	critical property
i	component number
x	liquid property
у	vapor property
Z	liquid property during experimental run

# Superscripts

ω	infinite dilution property
L	liquid phase property
٧	vapor phase property
S	saturated
р	reference pressure

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#### CHAPTER I

#### INTRODUCTION

The behavior of a compound in a mixture is different from its behavior in the pure state. This difference in behavior is due to the interactions among the unlike molecules which comprise the mixture. If there were no differences in the interactions among different molecules, the properties of a compound would be the same in both the pure state and the mixture, and the mixture would be "ideal". However, in reality there are differences in interactions among different molecules. The activity coefficient is a measure of the extent of non-ideality in the behavior of a liquid in a solution. When the activity coefficient is determined at infinite dilution (i.e., one component is present in very small quantity) in a binary system, the only kind of interaction that exists for a dilute molecule is with the solvent molecules surrounding it.

Infinite dilution activity coefficients can be determined in several ways, one of which is differential ebulliometry. Recently, ebulliometers have been used extensively to obtain infinite dilution activity coefficients (4, 13, 30). However, problems have been encountered in the operation of these devices at low pressures (30). The primary objective of this work was to develop and test a new ebulliometer for improved operation at low pressures. Operation at pressures below atmospheric will enable one to study heavy compounds.

These heavy compounds have very high boiling point temperatures at atmospheric pressure and are unsuitable for the ebulliometric technique.

The plan of attack in the present work was to perform an initial study of mixtures of light hydrocarbons, representing normal paraffins, naphthenes and aromatics. If these results were satisfactory, heavier solute-solvent pairs would be studied in future work.

Ebulliometry is the measurement of the boiling point temperature of a liquid. Differential ebulliometry is the measurement of the change (elevation or depression) in the boiling point of a pure substance as a result of the addition of a small amount of solute.

A knowledge of the rate of change of the boiling point temperature with solute composition enables one to calculate the infinite dilution activity coefficient. These infinite dilution activity coefficients fix the end points of the activity coefficient-composition relationship.

Once infinite dilution activity coefficients have been determined experimentally, they can be compared to values obtained using models for activity coefficient behavior. For example, equations of state could be "tuned" using interaction parameters  $(k_{ij})$  in order to predict mixture properties more accurately. One such equation of state is Soave's modification of the Redlich-Kwong Equation of State:

$$P = \frac{RT}{v-b} - \frac{a(T)}{v(v+b)}$$
(1)

where

$$a = \sum_{i j} \sum_{j} x_{i} x_{j} a_{ij}$$
(2)

$$a_{ij} = (1-k_{ij}) a_i^{1/2} a_j^{1/2}$$
 (3)

$$a_{i}(T) = a_{ci}\alpha_{i}(T)$$
(3.a)

$$a_{ci} = 0.42747 \frac{R^2 T_{ci}^2}{P_{ci}}$$
 (3.b)

$$\alpha_{i}(T) = [1 + m_{i}(1 - T_{Ri}^{0.5})]^{2}$$
 (3.c)

$$m_i = 0.480 + 1.574\omega_i - 0.176\omega_i^2$$
 (3.d)

$$b = \sum_{i} x_{i} b_{i}$$
(4)

$$b_i = 0.08664 \frac{RT_{ci}}{P_{ci}}$$
 (4.a)

The fugacity and activity coefficient are then calculated using the following general thermodynamic relationship:

$$RT \ln\left(\frac{f_i}{Py_i}\right) = \int_V^{\infty} \left[\left(\frac{\partial P}{\partial n_i}\right)_{V,T,n_j} - \frac{RT}{V}\right] dV - RT \ln Z$$
(5)

$$\ln(\frac{f}{P}) = \int_0^P \left(\frac{v}{RT} - \frac{1}{P}\right) dP$$
 (5.a)

$$\gamma_{i} = \frac{f_{i}}{x_{i}f_{i}^{0}}$$
(6)

f<sub>i</sub> = fugacity of component i

 $f_i^0$  = standard state fugacity of i.

In Equation (6),  $f_i$  is calculated using Equation (5) and  $f_i^0$  is calculated using Equation (5.a), where  $f = f_i^0$ , i.e., the standard state fugacity is defined as the condition where only pure liquid i is present  $(x_i = 1)$ .  $x_i$  in Equation (6) is the liquid mole fraction of component i.

Infinite dilution activity coefficients can be used to predict vapor-liquid equilibrium constants which are of great importance in the separation and mass transfer industrial processes. They are also of interest to statistical thermodynamicists by allowing them to neglect the order-disorder problems which are caused by the solute-solute interactions which are absent in the infinitely dilute region (30).

The greatest importance of infinite dilution activity coefficients is that they can be used to optimize activity coefficient models such as the Wilson Equation, given below.

$$\ln \gamma_{k} = -\ln \begin{bmatrix} c \\ \Sigma \\ j=1 \end{bmatrix} \begin{pmatrix} x_{j} & \Lambda_{kj} \end{bmatrix} - \begin{bmatrix} c \\ \Sigma \\ i=1 \end{bmatrix} \begin{bmatrix} x_{i} & \Lambda_{ik} \\ c \\ \Sigma \\ j=1 \end{bmatrix}$$
(7)

which for a binary system can be written as:

$$\ln \gamma_1 = -\ln(x_1 + \Lambda_{12}x_2) + x_2 \left[\frac{\Lambda_{12}}{x_1 + \Lambda_{12}x_2} - \frac{\Lambda_{21}}{\Lambda_{21}x_1 + x_2}\right]$$
(8.a)

$$\ln \gamma_2 = -\ln(x_2 + \Lambda_{21}x_1) - x_1 \left[ \frac{\Lambda_{12}}{x_1 + \Lambda_{12}x_2} - \frac{\Lambda_{21}}{\Lambda_{21}x_1 + x_2} \right]$$
(8.b)

x<sub>i</sub> = liquid mole fraction of component "i"

 $\gamma_k$  = activity coefficient of component "k"

 $\Lambda_{i,j}$  = binary interaction parameters

When insufficient data are available to determine binary Wilson parameters from a best fit of activity coefficients over the entire range of composition, infinite dilution activity coefficients can be used. At infinite dilution, Equations (8.a) and (8.b) reduce to:

$$\ln \gamma_1^{\infty} = 1 - \ln \Lambda_{12} - \Lambda_{21} \quad (x_1 = 0, x_2 = 1)$$
(9.a)

$$\ln \gamma_2^{\infty} = 1 - \ln \Lambda_{21} - \Lambda_{12} \quad (x_1 = 1, x_2 = 0)$$
(9.b)

If infinite dilution activity coefficients are known, Equations (9.a) and (9.b) can be solved simultaneously to obtain the binary interaction parameters  $\Lambda_{12}$  and  $\Lambda_{21}$ .

If temperatures corresponding to  $\gamma_1^{\infty}$  and  $\gamma_2^{\infty}$  are not equal,  $\Lambda_{ij}$  (which are used only under isothermal conditions) are corrected for the effect of temperature and new parameters ( $\lambda_{12} - \lambda_{11}$ ) and ( $\lambda_{12} - \lambda_{22}$ ) are obtained as follows:

$$(\lambda_{12} - \lambda_{11}) = -RT \ln(\frac{v_1^L}{v_2^L} \Lambda_{12})$$
 (10.a)

$$(\lambda_{12} - \lambda_{22}) = -RT \ln(\frac{v_2^L}{v_1^L} \Lambda_{21})$$
 (10.b)

In the present study, since the data obtained were isobaric and not isothermal, Equations (10.a) and (10.b) would be used to obtain the binary interaction parameters  $(\lambda_{12} - \lambda_{11})$  and  $(\lambda_{12} - \lambda_{22})$ , which are generally regarded as being temperature independent.

### CHAPTER II

#### LITERATURE REVIEW

#### 1. Previous Apparatus

Ebulliometers are used for very accurate measurements of the boiling points of liquids. In 1925, Swietoslawski (28) used the Cotrell pump principle to design a very accurate ebulliometer. The Cotrell pump is a small diameter (about 5 mm) glass tube connected to the boiler of an ebulliometer through which the boiling liquid and vapor mixture passes as a result of thermosiphon action. Ebulliometers have been used for determining the boiling point, the molecular weight, mutual solubilities, sample purities, and infinite dilution activity coefficients.

Eckert et al. (30) at the University of Illinois designed a new ebulliometer by modifying the Swietoslawski model. The advantage of the new design was that it prevented the volatile components from escaping from the system and it also reduced pressure fluctuations. This was accomplished by connecting the modified Swietoslawski ebulliometers (through their condensers) to a common manifold.

The design of the present work is a modification of Thomas' (one of Eckert's coworkers) design (30). Thomas' design is shown in Figure 1. As can be seen in the figure, there are six sections in the ebulliometer.

Section D is the boiler where the solution boils in the annulus between the inner and outer tubes. Crushed glass is fused to the inside

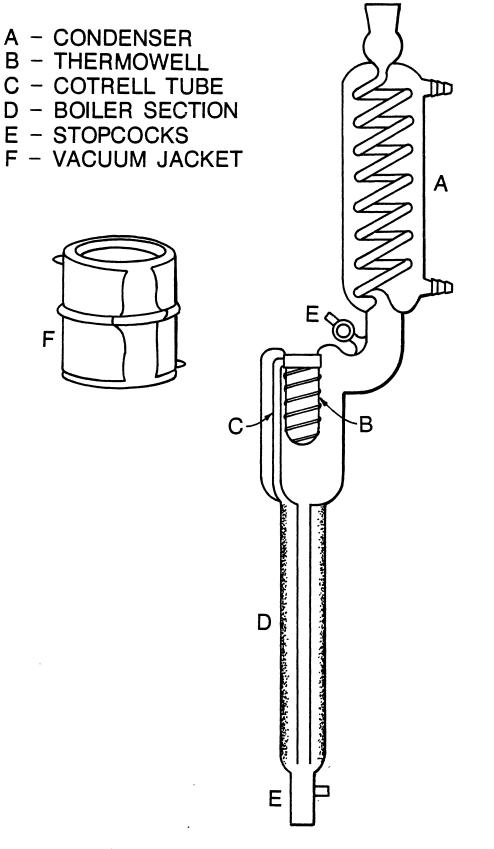


Figure 1. Ebulliometer of Eckert

of the outer tube to provide nucleation sites where the boiling can take place. A stopcock at the bottom of the boiler allows for drainage.

Section C is the Cotrell pump through which the boiling liquid and vapor travel by the boiling action and are delivered to the glass spiral on the thermometer well.

Section B is the thermometer well which is filled with mineral oil and houses the thermometer probe. The glass spiral on the outside of the thermometer well causes the liquid to be distributed evenly on the wall of the thermometer well. This provides enough contact area for the evaporating liquid to reach equilibrium with its vapor. The equilibrium temperature is the temperature registered by the thermometer probe inside the well.

Section A is the condenser where the vapor is totally condensed. There is a Teflon stopcock at the bottom of the condenser where solute is injected and carried to the bulk of the solution by the condensing vapor. This stopcock is also used for (gravimetric) filling of the ebulliometer.

Sections C and B of the ebulliometer are enclosed in a glass vacuum jacket F, which insulates the thermometer well from the surroundings. Section E consists of the stopcocks used for loading and drainage.

Heating is accomplished by a 5 ft long  $(15-20 \Omega)$  nichrome wire wrapped around the boiler in a helix and connected to a powerstat. The cooling in the condenser is provided by a water-ethylene glycol solution in a refrigerated circulating bath.

#### 2. Theory

Gautreaux and Coates (9) derived the mathematical expression for the experimental determination of the infinite dilution activity coefficient. The derivation by Null (16) is reproduced in Appendix B.

$$\gamma_{1}^{\infty} = \frac{\phi_{1}^{(P_{2}^{S})} \{P_{2}^{S} - [1 - \frac{P_{2}^{S} V_{2}}{RT} + \frac{P_{2}^{S}}{\phi_{2}^{S}} (\frac{\partial \phi_{2}}{\partial P})][\frac{dP_{2}^{S}}{dT}][\frac{dT}{dx_{1}}]_{P}^{\infty} \}}{P_{1}^{S} \phi_{1}^{S} \exp[(P_{2}^{S} - P_{1}^{S})\dot{V}_{1}/RT]}$$
(11)

where "1" refers to the infinitely dilute solute component and "2" refers to the solvent.

Neglecting the fugacity coefficient and Poynting correction effects, which are very close to unity at low pressures (in most cases), the above expression reduces to:

$$\gamma_{1}^{\infty} = \frac{P_{2}^{S} - (\frac{dP_{2}^{S}}{dT})(\frac{dT}{dx_{1}})_{p}^{\infty}}{P_{1}^{S}}$$
(12)

The saturated vapor pressures  $P_1^{s}$  and  $P_2^{s}$  and the derivative of the saturated vapor pressure of the solvent with respect to temperature are all pure component properties and can be determined experimentally and represented by an empirical equation such as the Antoine Equation.

The derivative of the boiling point temperature with respect to solute composition at constant total pressure and at infinite dilution is the only term in the above expression which requires binary data. This is the quantity which is determined experimentally using differential ebulliometry. In practice, small incremental amounts of solute are added to the solvent, and the change in the boiling point temperature is measured. The slope of these data at  $x_1 = 0$  represents the  $(\frac{dT}{dx_1})_p^{\infty}$  term, which permits calculation of  $\gamma_1^{\infty}$  from Equation (12).

Equation (12) can be used to examine the sensitivity of  $\gamma_1^{\infty}$  to errors in the measured limiting slope and thus provide a criterion for the applicability of the ebulliometric method to a given binary system (6). The limiting activity coefficient is essentially the algebraic sum of two terms and may become the difference between two much larger numbers. In this case, extremely high accuracy is needed in the data. The fractional change in  $\gamma_1^{\infty}$  with the equivalent change in limiting slope can be considered as a measure of sensitivity:

$$\frac{\left(\frac{\partial T}{\partial x_{1}}\right)^{\infty}_{P}}{\gamma_{1}^{\infty}} \frac{d\gamma_{1}^{\infty}}{d\left(\frac{\partial T}{\partial x_{1}}\right)^{\infty}_{P}} = 1 - \frac{P_{2}^{S}}{P_{1}^{S}\gamma_{1}^{\infty}}$$
(12.a)

The second term on the right side of the equation is essentially the relative volatility at infinite dilution, since the activity coefficient of the solvent is unity at this limit and the total pressure is the vapor pressure of the solvent. According to Equation (12.a), systems where the solute has a much higher vapor pressure than the solvent ( $P_2^S << P_1^S$ ), and systems of similar vapor pressures ( $P_1^S \sim P_2^S$ ) but very high  $\gamma_1^\infty$  values are unsuitable for the ebulliometric method. The ebulliometric technique is, thus, seen to be best suited to systems of similar volatility (6).

#### CHAPTER III

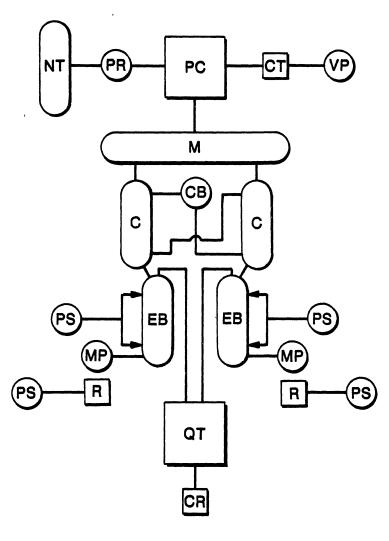
#### EXPERIMENTAL APPARATUS

In order to achieve better operation at low pressures, the ebulliometer of Eckert et al. was modified in this work. The equipment used in the present work for the experimental determination of infinite dilution activity coefficients is shown schematically in Figure 2. The major components are described below.

#### The Ebulliometer

The glass ebulliometer shown in Figure 3 is the central component of the apparatus. The design used in the present work is different from the Thomas design, discussed in the previous section and illustrated in Figure 1, in three areas. First, the Teflon stopcock was replaced by a Teflon septum and its position altered from the bottom of the condenser to the top of the thermometer well encasement, thus ensuring the mixing of the solute and the solvent even at low condensation rates. The condenser in this design is also different in that it has two spherical enlargements in the lower section where 80 to 90 precent of the condensation takes place. (The condensation ring could usually be seen at the top of the lower section during the course of most exprimental runs.)

Second, in this design two Cotrell tubes were used in order to distribute the boiling liquid more evenly across the glass spiral on the thermometer well.



- NT NITROGEN TANK
- PR PRESSURE REGULATOR
- PC PRESSURE CONTROLLER
- CT COLD TRAP
- **VP VACUUM PUMP**
- M MANIFOLD
- CB COOLANT BATH
- C CONDENSER
- **EB EBULLIOMETER**
- PS POWERSTAT
- **MP MAGNETIC PUMP**
- R DRIVE MAGNET ASSEMBLY
- QT QUARTZ THERMOMETER
- **CR CHART RECORDER**

Figure 2. Overall Schematic View of the Apparatus

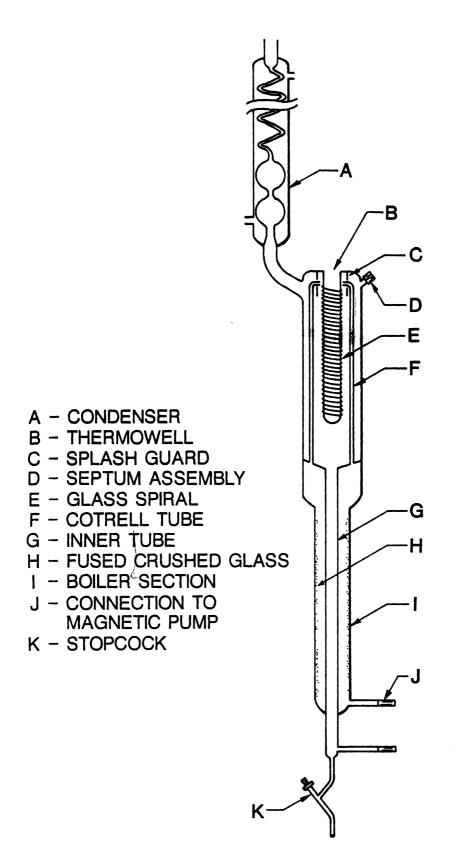


Figure 3. Modified Ebulliometer

The third change, which is also the most important modification, is the addition of a vane-type magnetic pump. This modification was designed to reduce reported problems (30) in operation at low pressures, where boiling becomes unstable (surge boiling) and boiling temperatures fluctuate. The present design replaces the thermosiphon-dependent circulation inside the Cotrell tubes by forced circulation. This leads to better performance of the ebulliometer by constantly covering the thermometer well with a uniform layer of boiling liquid. The new feature enhanced the performance of the ebullimeter, specially at low pressures, as will be discussed in later sections. The result is a decrease in temperature fluctuations caused by the unsteady boiling.

An additional minor change was the deletion of the vacuum jacket. This decision was made because the surrounding room temperature should not have any effect on the boiling point temperature which is measured inside the thermometer well.

The ebulliometer is 3 ft high and the condenser constitutes half of this height. The outlets J (letters refer to Figure 3), which are the connections to the magnetic pump, are 1 1/2" apart and are actually horizontally parallel. If flexible tubing is used to connect the ebulliometer to the pump (as was the case), the distance between the two legs is not critical.

The boiler is 1" in diameter and the inner tube has an ID of 1/4". The thermometer well has an ID of 1/2" and the Cotrell tubes are 6 mm in ID. The encasement is 2" in diameter and the condenser has a diameter of 1 1/2". The septum used was rubber on one side and Teflon on the other and had a diameter of 11 mm. The septa were purchased from Supelco, Inc. (Catalog no. 2-2731). The splash guard, C, guided the

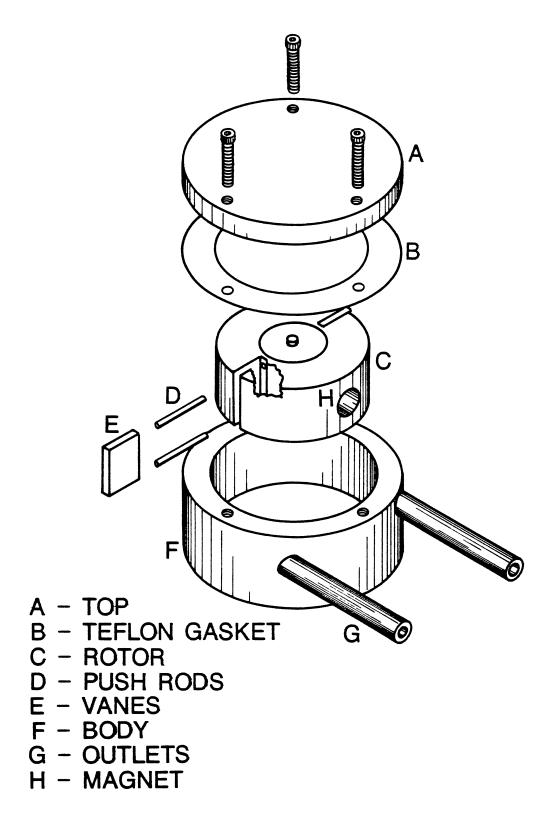
boiling liquid onto the thermometer well spiral. The ground glass fused in the boiling section was 30 mesh and this small size enhanced the boiling characteristics and thus the performance of the ebulliometer.

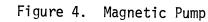
Except for the condenser, which was available commercially, the ebulliometer was assembled by the local glass-blower, Mr. M. Wayne Adkins, including the skillful fusing of the crushed glass to the inside wall of the boiler.

Approximately 5 ft of nichrome ribbon with a resistance of about 17 ohms was wrapped around the boiler in a helix. The ends of the wire were connected through alligator clips to a powerstat. The thermometer well was filled with silicon oil to provide good thermal contact between the thermometer well walls and the thermometer probe.

#### Vane-Type Magnetic Pump

The vane-type pump, Figure 7, is made of brass, is cylindrical in shape, and has a diameter of 35 mm and a height of 15 mm. Inside the pump there is a cylindrical aluminum block, C, in which two rectangular stainless steel vanes, E, are inserted symmetrically. The vanes are 9 mm by 16 mm, and are mutually in contact with two brass "push" rods, D, which are 18 mm in length and 1 mm in diameter. These rods are located horizontally near the top and the bottom of the aluminum block. The magnet, H, is inside the rotor which is mounted slightly off center in the brass housing, F. There is a 5 mm thick Teflon gasket, B, between the top and the body, F, of the pump. The pump is driven by a horse-shoe magnet which is  $2" \times 1 1/2" \times 1 1/2"$  and is mounted on a 0.01 hp motor connected to a powerstat. When the magnet turns, it rotates the cylindrical block inside the pump. The centrifugal force (and the push





rods) drive the vanes out to the wall of the pump and force the fluid to circulate in the manner typical of vane-type pumps. The pump outlets, G, are 5 mm in diameter and 35 mm long. The vane-type magnetic pump was fabricated by Mr. Heinz Hall at the Machine Shop of the Physics Department at Oklahoma State University.

The addition of the magnetic pump improved the performance of the ebulliometer considerably; this can be seen from temperature profiles shown in Figures 5 and 6. These figures demonstrate the boiling point temperature fluctuations of pure n-hexane at 170.53 mm Hg. The thermometer was in the differential mode and the difference in the boiling point of pure hexane in two ebulliometers was measured and recorded. The smallest divisions on the chart paper represent temperature changes of  $10^{-3}$  °C. The chart-recorder speed was 15 mm/min. The temperature changes in Figure 5 are on the order of 2 x  $10^{-2\circ}$ C and those in Figure 6 are on the order of 3 x  $10^{-3\circ}$ C.

#### The Quartz Thermometer

Two ebulliometers are used in differential ebulliometry. One of them contains the pure solvent only and is called the reference ebulliometer. The other one is called a loading ebulliometer. The solute is injected into the loading ebulliometer.

The temperature difference between the reference ebulliometer and the loading ebulliometer was measured by a Hewlett-Packard Quartz Thermometer (Model 2804A) with matching sensor probes capable of resolution to  $10^{-4}$ °C; however, only the medium resolution ( $10^{-3}$ °C) range was used in this work. The probes can be calibrated using either the triple point or freezing point of water. Calibrations were carried out

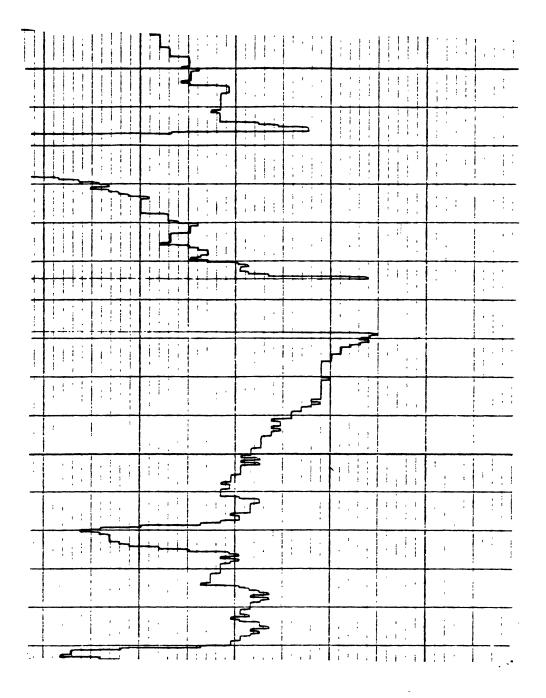


Figure 5. Temperature Fluctuations without the Magnetic Pump

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Figure 6. Temperature Fluctuations with the Magnetic Pump

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in an ice bath in this work. When measuring the boiling point of the pure solvent in the two ebulliometers, both probes should read the same value of temperature. However, due to unmatched probes this difference was usually other than zero; it was read and recorded and subsequently subtracted from other readings. Typical values of this difference in reading were between 0.100 °C and 0.200°C with an average value of 0.150°C.

The probes are 23 cm long and have a maximum diameter of 11 mm. The probes can be used as far as 22 ft away from the thermometer read-out. They are very fragile and sensitive to shock and should be handled with extreme care. An 8" drop onto a hard surface will damage the probes permanently.

The absolute accuracy of the digital thermometer is  $\pm 4 \times 10^{-2}$ °C over the range of -50°C to 150°C. The digital display is pushbutton controlled and there is an internal switch where one can select the readings either in the Celsius or the Fahrenheit temperature scale.

Thumbwheel switches on the front panel can be used to adjust the probes by calibration against the ice point or triple point of water and thus the effects of thermal history on the sensor can be corrected. The thermometer can display the reading of either probe or show the difference between the two readings (differential mode). Three levels of resolution (low, medium, and high) are avilable. The time between display of readings depends on the resolution and mode selected, and ranges from 0.1 second for low resolution and absolute temperature to 20 seconds for high resolution and differential temperature. The overall range of the thermometer is from -80°C to 250°C.

There is an analog output option available on the Quartz Thermometer. The analog output is a d.c. voltage ranging from 0.00 to 9.99 volts and corresponds directly to the value of the selected digits. The desired digits can be selected on the front panel and can be any of the seven digits in display (XXX.XXXX). The analog output is delivered at the rear panel and can be used as input to a recording device.

#### Pressure Controller

A Texas Instrument Precision Pressure Test Set (Model 156) was used to control the pressure. This unit is capable of both reading and regulating pressures. The desired use can be switch-selected at the front panel. The operation of the Test Set is based upon the output of a fused quartz Bourdon capsule. In the Model 156, an absolute Bourdon capsule is installed.

There are three different modes of operation available: Manual Gage, Servo Gage, and Servo Control. In the present study only the Servo Control mode was used. In this mode the counter reading was manually adjusted to correspond to the desired controlled pressure. The error signal from the amplifier drives a pressure regulator. The regulator output pressure is connected to the Bourdon Capsule and the Regulated Pressure Output fitting.

The plumbing at the back of the Test Set consists of three fittings. The first fitting was connected to about 1 ft of stainless steel tubing and then through a vacuum hose to a vacuum pump. The second fitting was connected in a similar manner to the pressure supply source. The third fitting was connected directly to the Bourdon capsule

and contained the regulated pressure. This fitting was connected to the manifold.

The pressure supply can be dry air, nitrogen or any other equivalent inert gas. In the present work nitrogen was used. The pressure at the pressure supply source should be 1.2 times the desired regulated pressure and should be set using a pressure regulator at the output of the gas tank.

The pressure range of the Test Set is from 0 to 1000 mm Hg and the control accuracy, i.e., the regulated pressure output versus counter reading is 0.002% of full scale (or 0.02 mm Hg in the present case).

The deviation meter at the front panel indicated pressure fluctuations in the range of 0.2 to 0.4 mm Hg with an average deviation of 0.3 mm Hg. This corresponded to a maximum temperature fluctuation of  $1 \times 10^{-4}$ °C in this work.

The Test Set is capable of providing output signals at the rear panel for recording purposes, but this feature of the Test Set was not used in the present work.

#### Circulating Bath

An Endocal Model RTE-4 Refrigerated Bath Circulator was used to provide coolant to the condenser. This unit can operate between  $-30^{\circ}$ C to  $100^{\circ}$ C with a  $\pm 0.01^{\circ}$ C temperature stability when water is used as the coolant. It has a 5 liter reservoir. When temperatures below 0°C are desired, a coolant such as ethylene-glycol or ethanol should be mixed in with water to prevent any freezing in the lines.

The unit is equipped with a circulating pump, proportional temperature controller and stainless steel refrigeration coils. The circulator was connected by 3/8" tygon tubing to the two ebulliometer condensers in series.

#### Vacuum Pump

A Sargent-Welch duo seal vacuum pump (Model 1400) was used. The vacuum pump was connected through 1/2 " vacuum tubing to a cold trap which was in turn connected to the vacuum supply port of the pressure controller. The cold trap was not used since no vapor escaped the system.

#### Manifold

The two ebulliometers were connected through their condensers to a manifold using sealed glass ball joints, "O" ring seals, and pinch clamps. The regulated pressure was connected to the top of the manifold through a needle valve. The glass manifold has a diameter of 3" and a length of 18". The manifold was used to provide a large volume of inert gas and thus to dampen pressure fluctuations.

#### Chart Recorder

A Perkin-Elmer Model 561 Recorder was used to record the measured temperature differences. The Model 561 is a straight line potentiometric self-balancing Recorder with stepper motor chart drive. It has chart speeds ranging from 15 to 480 mm/min. It has zero adjust control and an effective chart width of 25 cm. The input voltage is either 0-10 mv or 0-1 v. Therefore, the output signal from the thermometer was reduced by a factor of 10 using a 56 k $\Omega$  and a 6.8 k $\Omega$  resistor in series, to match the recorder input voltage requirement.

## Chemicals Used

The ethanol (supplied by U.S. Industrial Chemicals Co.) used in this work was 200 proof and was purchased locally. The other chemicals had minimum purities of 99 mole percent and were purchased from Aldrich Chemical Company, Inc. All the chemicals were used without any further purification.

#### Arrangement of Apparatus

The ebulliometers were clamped to an aluminum rod frame and connected through their condensers to the manifold. The quartz thermometer probes were inserted into the silicon oil in the thermometer wells.

The magnetic pumps were connected to the ebulliometer outlets using thick-walled tygon tubing which would not collapse under vacuum. The regulated pressure line of the pressure controller was connected to the manifold. The pressure supply and the vacuum supply were connected to the controller. The recorder was connected to the quartz thermometer at the rear panel. The magnetic pumps were positioned on the horse shoe magnet assembly. Alligator clips were attached to the nichrome wire heaters on the ebulliometers. The four powerstats were connected to the heating elements and the drive magnet motors.

#### CHAPTER IV

### EXPERIMENTAL PROCEDURE

#### Cleaning the Ebulliometers

The ebulliometers were cleaned using acetone which was pumped through the Cotrell tubes. The acetone was heated slightly so that its vapor would condense inside the condenser and clean that area, too. The acetone was then drained and nitrogen gas at about 20 psia was used to evaporate any remaining acetone from the ebulliometers.

When the ebulliometers became discolored by metallic residue from the aluminum rotors of the magnetic pumps accumulating on the inside walls, the pumps were disconnected and concentrated nitric acid was used to remove the residue. The outlets J(in Figure 3) of the ebulliometers were connected to each other using tygon tubing when the pumps were removed.

Since the ebulliometers were connected to the manifold and were closed at the top, the acetone or other cleaning agents were loaded by removing the septum and pouring the liquid into the ebulliometers through a plastic funnel.

## Charging the Solvent

Approximately 110 cc of the solvent was poured into a plastic bottle and weighed on a 0-100 g analytic balance capable of reading the

mass to 10<sup>-4</sup> g. Since all 110 cc could not be weighed at once using this balance, the process was carried out in two steps. After the liquid had been poured into the loading ebulliometer, the "empty" container was weighed and the mass of the "empty" bottle was subtracted from the mass of the full bottle to determine the exact mass of solvent charged.

The exact mass of the solvent charged to the reference ebulliometer did not need to be determined. 110 cc of solvent was poured into the reference ebulliometer without weighing.

#### Pressure Regulation

After the solvent had been charged, the selected pressure was set on the pressure regulator (using the calibration chart supplied by Texas Instruments). The pressure supply was connected and a pressure 1.2 times the selected value was supplied using the gas tank regulator. The vacuum pump was turned on, the needle valve at the manifold was opened, and the pressure was allowed to reach a stable value as indicated by the deviation meter on the Pressure Test Set. An average deviation of 0.30 mm Hg was considered satisfactory. The pressure fluctuations were very slow at low pressures (150 mm Hg) and rather fast at high pressures (700 mm Hg). Pressure fluctuations were dampened by the ballast in the manifold.

#### Temperature Measurement and Recording

When the pressure stabilized, the magnetic pumps were turned on and set at a sufficiently high flow rate to wet the thermometer well completely. At this point heating was initiated. The voltage selected

depended on the resistance of the nichrome wire used as the heating element. In the present study, the powerstat voltage was set at 35 in order to achieve a boiling point temperature of about  $80^{\circ}$ C, using a  $17_{\Omega}$  resistance.

Once boiling had started and vapor bubbles could be seen in the Cotrell tubes, the temperature was allowed to stabilize in both ebulliometers. The temperature difference between the two ebulliometers when both contained pure solvent (which was close to zero) was read. The temperature differences were recorded for about 40 minutes to make sure that the equilibrium temperature had been reached. A fluctuation of about  $5\times10^{-3\circ}$ C was considered satisfactory; when this level of performance was observed, preparations were made for the first injection. The liquid level in the ebulliometer was observed to be constant when steady state had been reached. The boiling was usually very vigorous and the condensation ring could be seen in the condenser. If periodic fluctuations in the level were observed, temperature stability was not achieved.

#### Solute Injections

At this point about 1 cc of the solute was drawn into a syringe. The syringe was weighed, then the solute was injected through the Teflon septum into the loading ebulliometer. The syringe was immediately reweighed and the exact mass of solute injected was determined. The change in the boiling point temperature was noted and when the temperature reached a new steady state value, the recorder (which had been turned off before the injection) was turned back on and the temperature fluctuations were recorded for about 20 minutes. Again, fluctuations

were typically on the order of 0.005°C. The new boiling point temperature was recorded. A total of 4 or 5 injection was made in a similar manner, and each time the temperature fluctuations were monitored and recorded. The heating was then stopped and the pressure controller was disconnected from the manifold by closing the needle valve. Once the pressure controller, the vacuum pump, the pressure supply and the magnetic pumps were turned off, the ebulliometers were drained and cleaned. The thermometer could be switched off at any time with probes remaining in the thermometer wells.

Each experimental run took about 4 hours and the stabilization of the pressure and temperature took place faster at lower pressures. However, the boiling did not appear as vigorous at lower pressures as it did at higher pressures.

#### CHAPTER V

#### DATA REDUCTION

Since there was condensed vapor holdup in the condenser, liquid holdup in the condenser and on the glass spiral, and vapor holdup in the space between the boiling liquid and the condensing vapor in the ebulliometer, the compositions obtained using the charged masses of liquid needed to be corrected. The procedure for making these corrections and reducing data follows.

The overall mole fraction,  $z_1$ , of the solute in the ebulliometer was calculated as follows.

$$z_1 = \frac{n_1}{n_1 + n_2}$$
(13)

However,  $z_1$  in the above equation is the "overall" mole fraction and is not equal to the liquid mole fraction when any vaporization takes place in the ebulliometer. After boiling begins, some vapor and liquid "holdup" must be taken into account to obtain the liquid composition. When  $z_1$  in the above equation is corrected for vapor and liquid holdups, the corrected liquid mole fraction is referred to as  $x_1$ .

The data obtained experimentally in this work were corrected for vapor and liquid holdups; a computer program (Appendix A) was used to make the corrections (using the following concepts).

The ebulliometer was envisioned to consist of four sections. The first section consisted of the liquid volume during the experimental run. The second section was the volume of the liquid held up on the thermometer well spiral. The third section was the volume of the condensed vapor in the condenser. The fourth section was the vapor volume through the rest of the ebulliometer. These volumes were designated by  $V_z^L$ ,  $V_x^L$ ,  $V_y^L$ , and  $V_y^V$ , respectively. The value of  $V_z^L$  was known to be approximately 106 cc. The values of  $V_y^L$ ,  $V_y^V$ , and  $V_x^L$  were estimated to be 2 cc, 200 cc, and 2 cc, respectively. In order to calculate the vapor composition (mole fraction), equilibrium ratios were needed. These were calculated using the following relationships:

$$K_{i} = \frac{y_{i}}{x_{i}}$$
(14)

K<sub>i</sub> = equilibrium ratio of component "i"

y<sub>i</sub> = vapor mole fraction of component "i"

x<sub>i</sub> = liquid mole fraction of component "i"

$$\hat{f}_i^V = \hat{f}_i^L$$
 (equilibrium criteria) (15)

$$\hat{f}_{i}^{V} = y_{i} \hat{\phi}_{i} P \qquad (16)$$

$$\hat{f}_{i}^{L} = x_{i} \gamma_{i} f_{i}^{0}$$
(17)

$$y_{i}\hat{\phi}_{i}P = x_{i}\gamma_{i}f_{i}^{0}$$
(18)

 $\hat{\phi}_i$  = fugacity coefficient of component "i" in the mixture  $(\hat{\phi}_i = \frac{\hat{f}_i^v}{y_i^p})$ 

 $\gamma_i$  = activity coefficient of component "i" in the mixture  $(\gamma_i = \frac{\hat{f_i}_i^L}{x_i f_i^O})$  $f_i$  = fugacity of component "i" P = total pressure

but 
$$f_i^{0} = P_i^{0} \phi_i^{0}$$
 (19)

$$\hat{y_i \phi_i} P = x_i \gamma_i P_i \phi_i^0$$
(20)

Solving the above equation for  $\gamma_{\mbox{\scriptsize j}}$  and taking logarithms:

$$\ln \gamma_{i} = \ln \frac{y_{i}^{P}}{x_{i}^{P}_{i}} + \ln \hat{\phi}_{i} - \ln \phi_{i}^{0}$$
(21)

At low pressures: 
$$\ln \hat{\phi}_i = \ln \phi_i^0 = 1$$
 (22)

$$\gamma_{i} = \frac{y_{i}^{P}}{x_{i}^{P} p_{i}}$$
(23)

$$K_{i} = \frac{y_{i}}{x_{i}} = \frac{P_{i}^{0} Y_{i}}{P}$$
 (24)

Now using the one constant Margules Equation for  $\gamma_{\mbox{\scriptsize j}}$  :

$$\ln \gamma_1 = A x_2^2 = A (1 - x_1)^2$$
 (25)

$$\ln \gamma_2 = A x_1^2$$
 (26)

and Equation (24) yields

$$K_1 = \frac{P_1^0}{P} \exp [A(1 - x_1)^2]$$
 (27)

$$K_2 = \frac{P_2^0}{P} \exp((Ax_1^2))$$
 (28)

Here  $P_1^0$  and  $P_2^0$  are the vapor pressures of the solute and solvent, respectively. The vapor pressure of the solvent, when the solute is infinitely dilute, is the same as the operating pressure of the system and was known experimentally. The vapor pressure of the solute was calculated using the three constant Antoine Equation:

$$P_i^0 = P_{Ci} \exp [A_1 - \frac{A_2}{A_3 + T}]$$

where  $P_i^0$  is the saturated vapor pressure of component "i" in psia and  $P_{Ci}$  is the critical pressure of the same component in psia. T' is the boiling point temperature in °F.

The Margules constant A is calculated using Equation (25) for the infinitely dilute case  $(x_1 = 0)$ :

$$A = \ln \gamma_1^{\infty}$$
(29)

 $\gamma_1^{\infty}$  in the above equation is calculated using Equation (12). In that equation,  $P_2$  is the solvent vapor pressure and was measured directly.  $\frac{dP_2}{dT}$  was calculated by differentiating the Antoine Equation with respect to temperature:

$$\frac{dP_2^0}{dT} = P_2^0 \left[ \frac{A_2}{(A_3 + T)^2} \right]$$
(30)

The mole fraction versus the boiling point temperature curves were nearly linear over the composition ranges studied, however, there was a slight curvature in some of them. Therefore, both linear and quadratic fits to the experimental data were used, i.e.,

$$\Delta T = a x_1 \tag{31}$$

and

$$\Delta T = ax_1 + bx_1^2$$
 (32)

The slope of either of these curves at infinite dilution  $(x_1 = 0)$  is

 $\frac{d(T)}{dx_1}$  = a . The constants in Equations (31) and (32) were calculated by minimizing the sum of the squares of errors:

$$SS = \Sigma \left( T_{\text{measured}} - T_{\text{calculated}} \right)^2$$
(33)

where the summation extends over all  $\Delta T - x_1$  points in a given experiment. This leads to

$$a = \frac{\sum x_i \Delta T_i}{\sum x_i}, \text{ for the linear case}$$
(34)

$$a = \frac{\sum x_i \Delta T_i - b \sum x_i^3}{\sum x_i^2}, \text{ for the quadratic case}$$
(35)

where

$$b = \frac{\left[\left(\Sigma \times_{i} \Delta T_{i}\right) - \left(\Sigma \times_{i}^{2} \Delta T_{i}\right)\left(\Sigma \times_{i}^{2}\right)\right] / \left(\Sigma \times_{i}^{3}\right)}{\left(\Sigma \times_{i}^{3}\right) - \left[\left(\Sigma \times_{i}^{4}\right)\left(\Sigma \times_{i}^{2}\right) / \left(\Sigma \times_{i}^{3}\right)\right]}$$
(36)

As a first approximation, the overall mole fraction  $z_i$  was used (instead of  $x_i$ ) to get "a" from both (34) and (35). This value of "a" corresponds to the  $\frac{dT}{dx_1}$  term and permits the calculation of  $\gamma_1^{\infty}$  by Equation (12). From  $\gamma_1^{\infty}$ , the Margules constant was calculated using Equation (29). The Margules constant was then used to calculate  $K_1$  and  $K_2$  from Equations (27) and (28). This permitted the calculation of  $y_1$ , the vapor mole fraction of the solute using:

$$y_{i} = K_{i} x_{i}$$
(37)

The following molar material balances were then used to correct  $z_i$  and obtain  $x_i$ :

$$n_1 = n_{10} - n_y^L y_1 - n_y^V y_1$$
 (38)

$$n_2 = n_{20} - n_y^{L}(1 - y_1) - n_y^{V}(1 - y_1)$$
 (39)

where  $n_{10}$  and  $n_{20}$  are the number of moles of the solute and solvent charged to the cell, and  $n_1$  and  $n_2$  are the corrected liquid phase values after some of the solvent and solute have vaporized.

Using the product of the density and volume to give the number of moles in Equations (38) and (39), i.e.,

$$\mathbf{n} = \rho \mathbf{V} \tag{40}$$

Then

$$n_{1} = n_{10} - [\rho^{L} V_{y}^{L} + \rho^{V} V_{y}^{V}]y_{1}$$
(41)

$$n_{2} = n_{20} - [\rho^{L} v_{y}^{L} + \rho^{V} v_{y}^{V}](1 - y_{1})$$
(42)

The vapor densities were calculated using the ideal gas law, since the pressures were low enough to warrant its use. Also, the magnitude of the value of the vapor density was very small compared to the liquid density, and its effect in the material balance equations was minimal. Thus,

$$\rho^{V} = \frac{P}{RT}$$
(43)

The liquid densities of the pure components were calculated using the Rackett Equation (30):

$$\frac{1}{\rho^{L}} = v^{L} = 83.147 \left(\frac{T_{c}}{P_{c}}\right) A^{\left[1 + \left(1 - \frac{T}{T_{c}}\right)\right]^{0.2857}}$$

$$\rho^{L} = \text{liquid density in gmol/cm}^{3}$$

$$v^{L} = \text{liquid specific volume in cm}^{3}/\text{gmol}$$

$$T_{c} = \text{critical temperature in K}$$

$$(44)$$

T = boiling temperature in K

 $P_c$  = critical pressure in atm

A = Rackett Equation Constant

The density of the mixture was calculated assuming ideal mixing as follows:

$$v_{mixture}^{L} = x_1 v_1^{L} + (1 - x_1) v_2^{L}$$
 (45)

$$\rho_{\text{mixture}}^{\text{L}} = \frac{1}{v_{\text{mixture}}^{\text{L}}}$$
(46)

The values of the contants in the above equation and all other constants used in calculations are given in Appendix C.

New mole fractions were then calculated using the following equations:

$$x_1 = \frac{n_1}{n_1 + n_2}$$
(47)

$$x_2 = 1 - x_1$$
 (48)

At this point a convergence criterion was checked and if the new and old values of  $x_i$  agreed within a specified tolerance  $(10^{-4})$ , the calculations for  $x_i$  were stopped. Otherwise, the above procedure was repeated until the convergence condition was satisfied, i.e.,  $x_i$  were used in Equations (34) and (35) to get "a" and "b", "A" was calculated using Equation (29),  $K_i$  were calculated using Equations (27) and (28),  $y_i$  were calculated using Equations (41), (42), (47) and (48). Once the corrected values of the liquid mole fraction were known, the correct infinite dilution activity coefficient was obtained.

#### CHAPTER VI

#### EXPERIMENTAL RESULTS AND DISCUSSION

#### Experimental Results

The experimental results from this study are presented in Tables I through VII. In these tables, the components are labeled by "1" for solute and "2" for solvent. Next, the operating pressure of the system is specified followed by the boiling temperature of the pure solvent and the mass in grams of the solvent charged.

In the first column, the mass of the solute added in grams is indicated. In the second column, the uncorrected (overall) mole fraction is recorded. These mole fractions were calculated in the following manner:

$$n_1 = \frac{m_1}{MW_1} \tag{49}$$

$$n_2 = \frac{m_2}{MW_2} \tag{50}$$

$$x_1 = \frac{n_1}{n_1 + n_2}$$
(51)

The third column gives the experimental change in the boiling point temperature corresponding to each solute mole fraction. The calculated values of the change in the boiling point temperature based on the linear and quadratic fits are presented in columns four and five. In the fifth and sixth columns, the linearly and quadratically fitted values of the corrected solute mole fraction are recorded.

All tables are based on the point  $x_1 = 0$ ,  $\Delta T \doteq 0$ , i.e., the initial  $\Delta T$  reading from the thermometers before any solute was added has been subtracted from the subsequent readings. Each table contains three sets of data obtained at three operating pressures for the same binary system.

The final calculated results are shown in Table VIII. The binary system is specified in the first column. The operating pressure is indicated in the second column. The boiling point temperature of the pure solvent is given in the third column. The linearly and quadratically fitted values of the infinite dilution activity coefficient are presented in columns four and five. The  $\gamma_1^{\infty}$  values of Eckert et. al are given in column six.

The effects of temperature on infinite dilution activity coefficient are shown in Figures 7 through 13. Each figure contains three data points representing the three pressures at which the binary system data were obtained. Error bars of  $\pm 1\%$  are shown on each figure, except Figure 4 where they would cover the entire figure. These figures are a measure of the linearity of the experimental data. Error bars take into account the non-linearity that might have been introduced by experimental errors. These figures show that the data are consistent. Most plots are linear if the experimental uncertainty is taken into account.

Figures 14 through 20 show typical isobaric boiling point results. These figures correspond to the intermediate pressures at which the experiments were conducted. These plots show the effect of

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Т	A	В	L	E	Ι

ISOLARIC BOILING POINTS FOR TOLUENE (1)-ETHANOL (2) MIXTURES

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Solute	Overall	Boiling Pt.	Error in Ca	alcd. ∆T(°C)	Calcd. Liqui	d Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	Quadratic	Linear	Quadratic
		705 mm Hg, T <sub>b, so</sub>	lvent = 76.36	54°C, m <sub>2</sub> = 81.4	44 gm	
1.173	0.00716	-0.150	-0.01283	-0.00180	0.00703	0.00703
2.387	0.01446	-0.290	-0.01256	-0.01256	0.01422	0.01422
3.227	0.01974	-0.382	-0.00284	-0.00284	0.01943	0.01943
4.212	0.02524	-0.472	+0.01303	+0.01303	0.02486	0.02486
	P =	310 mm Hg, T <sub>b, so</sub>	lvent = 56.81	.5°C, m <sub>2</sub> = 83.9	55 gm	,
1.040	0.00617	-0.144	-0.01149	-0.01149	0.00604	0.00603
2.088	0.01230	-0.279	-0.01426	-0.01426	0.01207	0.01207
3.313	0.01938	-0.423	-0.00535	-0.00535	0.01904	0.01900
4.526	0.02586	-0.554	-0.01328	+0.01328	0.02586	0.02581
	P =	• 180 mm Hg, T <sub>b</sub> , so	lvent = 45.90	)1°C, m <sub>2</sub> = 76.323	3 gm	
1.3267	0.00781	-0.225	-0.02755	-0.00350	0.00763	0.00760
2.5342	0.01481	-0.406	-0.03088	-0.00002	0.01450	0.01445
3.8790	0.02249	-0.591	-0.02035	+0.00170	0.02206	0.02199
4.9128	0.02831	-0.722	-0.00274	+0.00118	0.02780	0.02772
6.2798	0.03591	-0.880	+0.03350	-0.00123	0.03531	0.03521

### TABLE II

### ISOBARIC BOILING POINTS FOR CYCLOHEXANE (1)-BENZENE (2) MIXTURES

Solute	Overall	Boiling Pt.	Error in (	Calcd. ∆T(°C)	Calcd. Liquid	d Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	Quadratic	Linear	Quadratic
	P =	669 mm Hg, T <sub>b, so</sub>	lvent = 75.5	507°C, m <sub>2</sub> = 92.483	2 gm	
1.216	0.01196	-0.163	-0.01808	-0.01201	0.01185	0.01185
2.306	0.02243	-0.271	+0.00104	+0.00737	0.02225	0.02224
3.354	0.03230	-0.391	+0.00084	+0.00310	0.03204	0.03204
4.497	0.04283	-0.516	+0.00387	+0.00285	0.04251	0.04251
	Ρ =	418 mm Hg, T <sub>b, so</sub>	lvent = 61.7	76°C, m <sub>2</sub> = 90.13	5 gm	
1.029	0.01052	-0.150	-0.00758	-0.00018	0.01042	0.01041
1.989	0.02014	-0.280	-0.00725	+0.00119	0.01995	0.01994
3.044	0.03049	-0.418	-0.00478	-0.00136	0.03022	0.03020
4.162	0.04123	-0.550	+0.00900	+0.00047	0.04089	0.04086
	P =	193 mm Hg, T <sub>b, so</sub>	lvent = 41.3	86°C, m <sub>2</sub> = 90.619	) gm	
1.044	0.01061	-0.162	-0.00524	+0.00286	0.01049	0.01048
2.077	0.02088	-0.321	-0.01218	-0.00331	0.02066	0.02064
3.093	0.03079	-0.451	-0.0151	-0.00150	0.03045	0.03045
4.105	0.04045	-0.590	+0.00881	-0.00019	0.04003	0.04003

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## TABLE III

ISOBARIC BOILING POINTS FOR BENZENE (1)-CYCLOHEXANE (2) MIXTURES

Solute	Overall	Boiling Pt.	Error in (	Calcd. ∆T(°C)	Calcd. Liqui	d Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	Quadratic	Linear	<u>Quadratic</u>
	Р	= 727 mm Hg, T <sub>b,so</sub>	lvent = 78.69	90°C, m <sub>2</sub> = 79.545	gm	
1.348	0.01549	-0.228	-0.02103	-0.01025	0.01535	0.01534
2.505	0.03138	-0.411	-0.00332	+0.00817	0.03112	0.03109
3.418	0.04553	-0.537	-0.00346	-0.00021	0.04512	0.04512
4.687	0.05841	-0.709	+0.01005	-0.00151	0.05798	0.05798
	Р	= 389 mm Hg, T <sub>b, so</sub>	lvent = 59.8	874°C, m <sub>2</sub> = 84.649	9 gm	
1.360	0.01701	-0.226	-0.00294	+0.00249	0.01685	0.01685
2.246	0.02779	-0.371	-0.00651	-0.00108	0.02755	0.02754
3.311	0.04043	-0.534	-0.00340	-0.00140	0.04011	0.04010
4.504	0.05421	-0.705	+0.00679	+0.00082	0.05380	0.05379
	Р	= 170 mm Hg, T <sub>b, so</sub>	lvent = 37.6	51°C, m <sub>2</sub> = 80.80	L gm	
1.179	0.01696	-0.199	-0.01184	+0.00075	0.01696	0.01680
2.427	0.03107	-0.364	-0.01485	-0.00272	0.03107	0.03079
3.573	0.04192	-0.527	-0.00229	+0.00269	0.04192	0.04155
4.647	0.05661	-0.662	+0.01338	-0.00072	0.05661	0.05614

## TABLE IV

ISOBARIC BOILING POINTS OF CYCLOHEXANE (1)-HEXANE (2) MIXTURES

Solute	Overall	Boiling Pt.	Error in	Calcd. ∆T(°C)	Calcd. Liqu	id Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	<u>Linear</u>	Quadratic	Linear	Quadratic
	Ρ	= 550 mm Hg, T <sub>b, so</sub>	lvent = 58.	.550°C, m <sub>2</sub> = 67.517	gm	
0.947	0.01612	0.112	+0.00046	-0.00004	0.01612	0.01623
2.053	0.03021	0.249	+0.00040	-0.00012	0.03021	0.03040
3.001	0.04289	0.357	+0.00041	+0.00021	0.04289	0.04315
3.906	0.05783	0.458	-0.00064	-0.00008	0.05783	0.05817
	Р	= 300 mm Hg, T <sub>b</sub> , so	lvent = 41.	.924°C, m <sub>2</sub> = 69.547	gm	
0.860	0.01250	0.095	+0.00033	+0.00144	0.01250	0.01250
1.686	0.02422	0.182	-0.00253	-0.00115	0.02437	0.02437
2.770	0.03943	0.298	-0.00049	+0.00015	0.03943	0.03942
3.890	0.05417	0.414	+0.00141	+0.00008	0.05450	0.05449
	Р	= 171 mm Hg, T <sub>b, so</sub>	lvent = 27.	958°C, m <sub>2</sub> = 68.172	gm	
1.090	0.01422	0.115	-0.00330	-0.00292	0.01422	0.01623
2.074	0.03033	0.215	+0.00152	+0.00193	0.03052	0.03040
2.983	0.04372	0.305	+0.00028	+0.00040	0.04400	0.04315
4.086	0.05616	0.410	-0.00021	-0.00062	0.05652	0.05818

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ISOBARIC BOILING POINTS OF HEXANE (1)-CYCLOHEXANE (2) MIXTURES

Solute	Overall	Boiling Pt.	Error in (	Calcd. ∆T(°C)	Calcd. Liquio	l Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	Quadratic	Linear	Quadratic
	P	= 700 mm Hg, T <sub>b so</sub>	lvent = 77.90	63°C, m <sub>2</sub> = 83.973	gm	
0.684	0.00789	-0.131	+0.00686	+0.00399	0.00779	0.00789
1.374	0.01572	-0.271	+0.00482	-0.00090	0.01553	0.01554
2.202	0.02496	-0.436	+0.00047	-0.00238	0.02466	0.02467
3.097	0.03476	-0.612	-0.00407	+0.00120	0.03435	0.03437
	P :	= 350 mm Hg, T <sub>b</sub> , so	olvent = 56.9	961°C, m <sub>2</sub> = 81.15	7 gm	
0.908	0.01081	-0.187	-0.00259	+0.00183	0.01081	0.01066
1.858	0.02187	-0.380	-0.00668	-0.00183	0.02187	0.02156
2.707	0.03155	-0.540	-0.00126	+0.00058	0.03155	0.03115
3.646	0.04203	-0.713	+0.00508	+0.00005	0.04203	0.04152
	P =	• 150 mm Hg, T <sub>b, so</sub>	lvent = 35.0	)45°C, m <sub>2</sub> = 80.54	7 gm	
0.988	0.01184	-0.204	-0.00567	-0.00030	0.01167	0.01184
2.298	0.02711	-0.460	-0.00550	+0.00012	0.02675	0.02674
3.195	0.03730	-0.627	-0.00142	+0.00016	0.03682	0.03680
4.105	0.04741	-0.790	-0.00567	-0.00012	0.04683	0.04681

## TABLE VI

### ISOBARIC BOILING POINTS OF BENZENE (1)-HEXANE (2) MIXTURES

Solute	Overall	Boiling Pt.	Error in Ca	alcd. ∆T(°C)	Calcd. Liquid	d Comp., × <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	<u>Quadratic</u>	Linear	Quadratic
	P	= 700 mm Hg, Tb, sc	lvent <sub>.</sub> = 66.09	95°C, m <sub>2</sub> = 72.65	3 gm	
1.141	0.01703	-0.020	-0.00704	-0.00124	0.01703	0.01703
2.157	0.03172	-0.027	-0.00287	+0.00127	0.03172	0.03172
3.309	0.04785	-0.032	+0.00441	-0.00040	0.04785	0.04785
	P :	= 350 mm Hg, T <sub>b, sc</sub>	olvent = 45.92	20°C, m <sub>2</sub> = 69.78	5 gm	
1.267	0.01963	-0.016	-0.00291	-0.00187	0.01963	0.01963
2.498	0.03798	-0.027	-0.00168	-0.00061	0.03798	0.03798
3.738	0.05578	-0.034	-0.00319	+0.00339	0.05578	0.05578
4.708	0.06962	-0.047	-0.00083	-0.00186	0.06926	0.06926
	Р	= 150 mm Hg, T <sub>b, sc</sub>	olvent = 24.95	50°C, m <sub>2</sub> = 66.72	9 gm	
1.341	0.02170	-0.035	-0.01186	-0.00219	0.02170	0.02170
2.559	0.04061	-0.052	-0.00869	+0.00111	0.04061	0.04061
3.724	0.05802	-0.064	-0.00211	+0.00096	0.05802	0.05802
4.948	0.07565	-0.071	+0.00969	-0.00071	0.07565	0.07565

### TABLE VII

ISOBARIC BOILING POINTS OF HEXANE (1)-BENZENE (2) MIXTURES

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Solute	Overall	Boiling Pt.	Error in	Calcd. ∆T(°C)	Calcd. Liqui	id Comp., x <sub>1</sub>
Added, m <sub>l</sub> (gm)	Mol. Frn., z <sub>l</sub>	Change, ∆T(°C)	Linear	Quadratic	Linear	Quadratic
	Ρ =	700 mm Hg, T <sub>b</sub> , so	lvent = 77.	347°C, m <sub>2</sub> = 90.961	gm	
1.015	0.01001	-0.397	-0.01504	+0.00779	0.00973	0.00971
1.874	0.01832	-0.724	-0.02384	-0.00196	0.01785	0.01780
2.664	0.02585	-1.001	-0.01224	-0.00698	0.02520	0.02514
3.400	0.03277	-1.227	+0.02753	-0.00422	0.03197	0.03190
	P =	: 350 mm Hg, T <sub>b</sub> , so	lvent = 56.	642°C, m <sub>2</sub> = 91.055	gm	
0.962	0.00948	-0.412	-0.02375	-0.00198	0.00919	0.00916
1.797	0.01757	-0.743	-0.02255	+0.00127	0.01704	0.01700
2.581	0.02504	-1.040	-0.01170	+0.00033	0.02433	0.02427
3.684	0.03537	-1.429	+0.02577	-0.00034	0.03442	0.03435
	P =	150 mm Hg, T <sub>b, so</sub>	lvent = 35.	378°C, m <sub>2</sub> = 91.662	gm	
0.986	0.00966	-0.462	-0.02966	-0.00002	0.00931	0.00927
2.075	0.02011	-0.935	-0.03245	+0.00095	0.01943	0.01937
3.011	0.02892	-1.315	-0.01452	-0.00124	0.02800	0.02791
4.073	0.03873	-1.710	+0.03496	+0.00043	0.03757	0.03747

## TABLE VIII

### EXPERIMENTAL INFINITE DILUTION ACTIVITY COEFFICIENTS

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				ite Dilutio oefficient	on Activity
Solute/Solvent	Pressure mm Hg	Temperature <u>°C</u>	This Linear	Work E Quadratic	ckert (30)
Tolune/Ethanol	705.38	76.364	4.91	5.22	5.14
	309.55	56.815	5.28	5.61	5.34
	180.00	45.901	5.88	6.23	5.60
Cyclohexane/Benzene	669.44	75.507	1.42	1.45	1.45
	417.50	61.776	1.48	1.52	1.52
	192.70	41.386	1.57	1.62	1.61
Benzene/Cyclohexane	726.70	78.690	1.37	1.40	1.35
	388.60	59.870	1.43	1.45	1.41
	169.50	37.651	1.50	1.54	1.48
Cyclohexane/Hexane	550.00	58.550	1.07	1.07	1.07
	300.00	41.924	1.06	1.07	1.09
	170.53	27.958	1.07	1.06	1.09
Benzene/Hexane	700.00	66.095	1.46	1.49	
	350.00	45.920	1.52	1.53	
	150.00	24.950	1.62	1.67	
Hexane/Benzene	700.00	77.347	1.59	1.68	
	350.00	56.642	1.76	1.85	
	150.00	35.378	1.92	2.05	
Hexane/Cyclohexane	700.00	77.963	1.08	1.14	
	350.00	56.966	1.10	1.11	
	150.00	35.045	1.12	1.06	

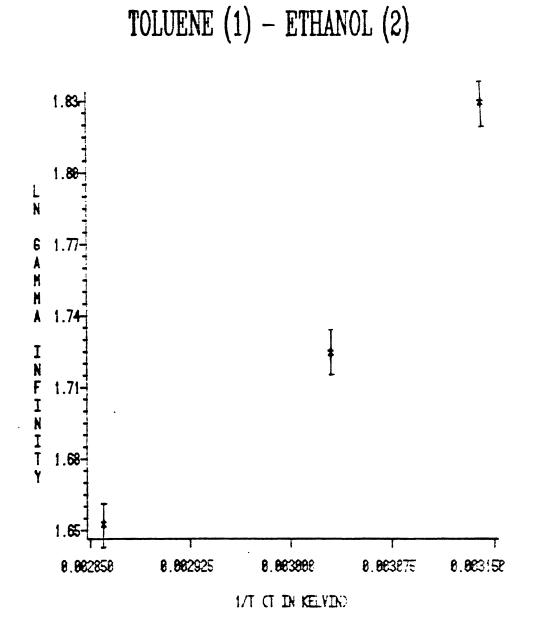


Figure 7. Infinite Dilution Activity Coefficient for Toluene in Ethanol

CYCLOHEXANE (1) - BENZENE (2)

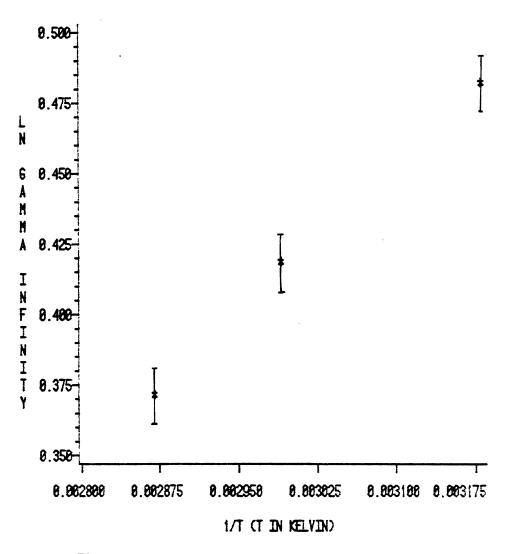


Figure 8. Infinite Dilution Activity Coefficient for Cyclohexane in Benzene

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# BENZENE (1) - CYCLOHEXANE (2)

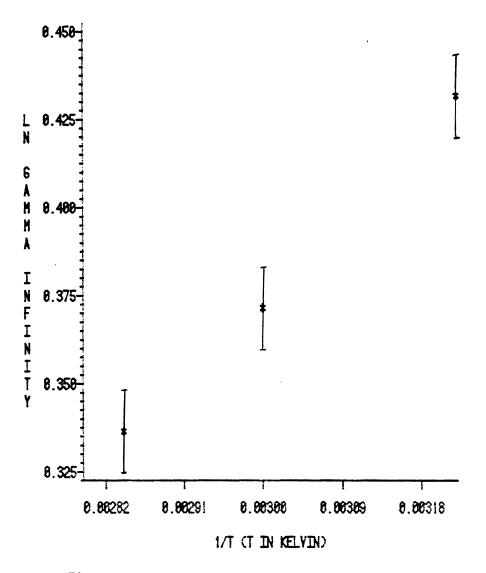


Figure 9. Infinite Dilution Activity Coefficient for Benzene in Cyclohexane

## CYCLOHEXANE (1) - HEXANE (2)

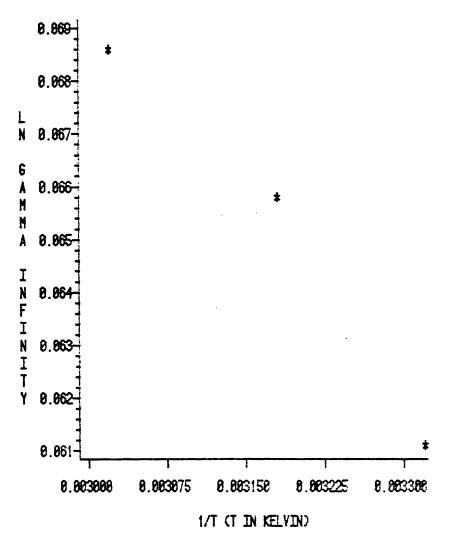


Figure 10. Infinite Dilution Activity Coefficient for Cyclohexane in Hexane

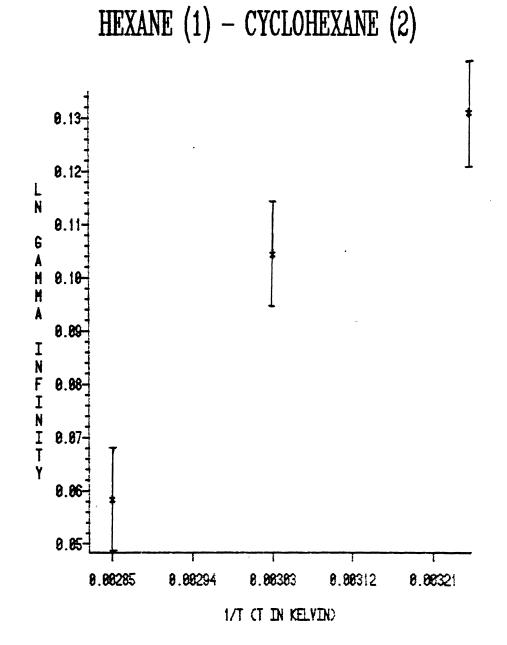


Figure 11. Infinite Dilution Activity Coefficient for Hexane in Cyclohexane

BENZENE (1) - HEXANE (2)

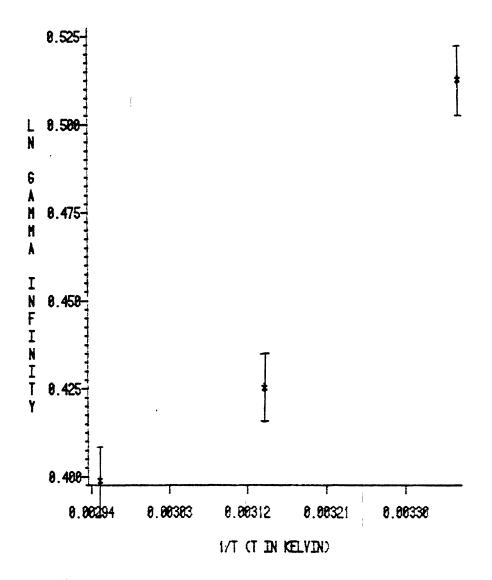
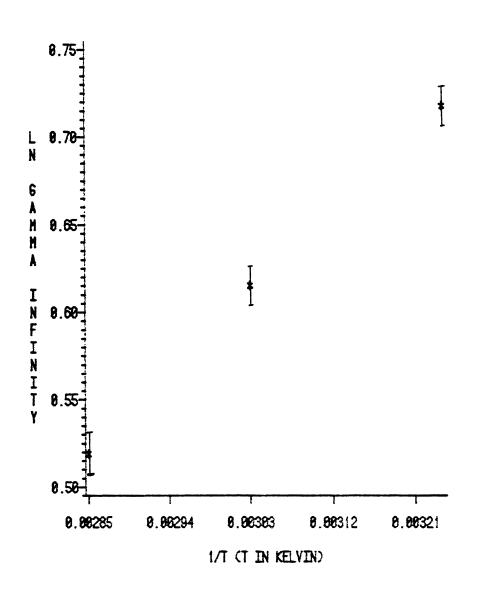


Figure 12. Infinite Dilution Activity Coefficient for Benzene in Hexane



HEXANE (1) – BENZENE (2)

Figure 13. Infinite Dilution Activity Coefficient for Hexane in Benzene

# TOLUENE (1) - ETHANOL (2) AT 310 MMHG

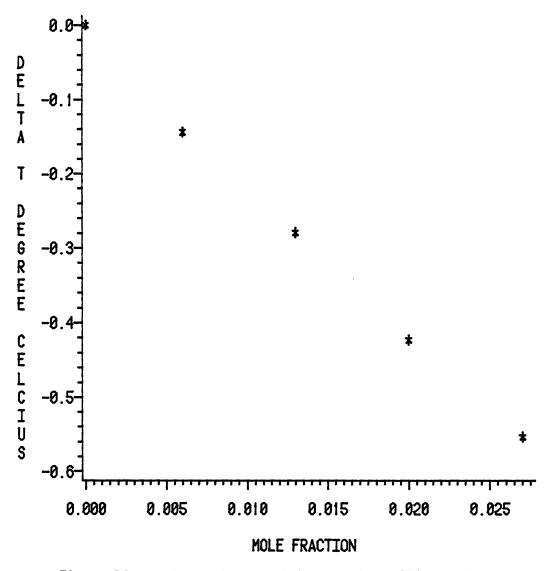


Figure 14. Effect of Composition on the Boiling Point Temperature of the Binary System Toluene (1) - Ethanol (2)

## CYCLOHEXANE(1)-BENZENE(2) AT 418 MMHG

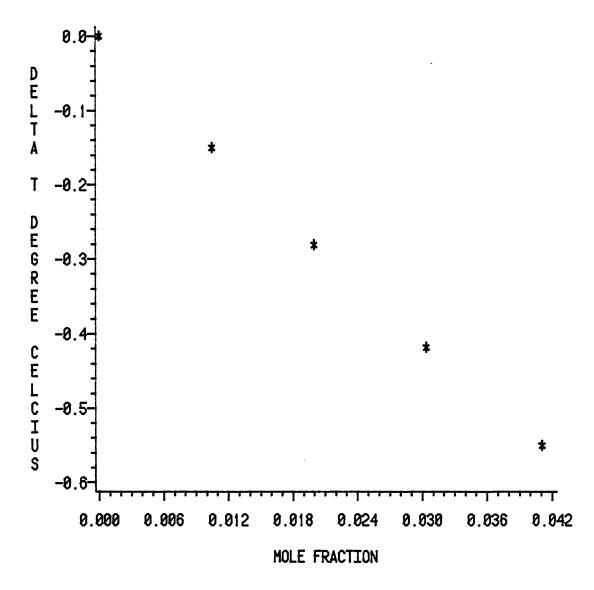


Figure 15. Effect of Composition on the Boiling Point Temperature of the Binary System Cyclohexane (1) - Benzene (2)

## BENZENE(1)-CYCLOHEXANE(2) AT 388.6 MMHG

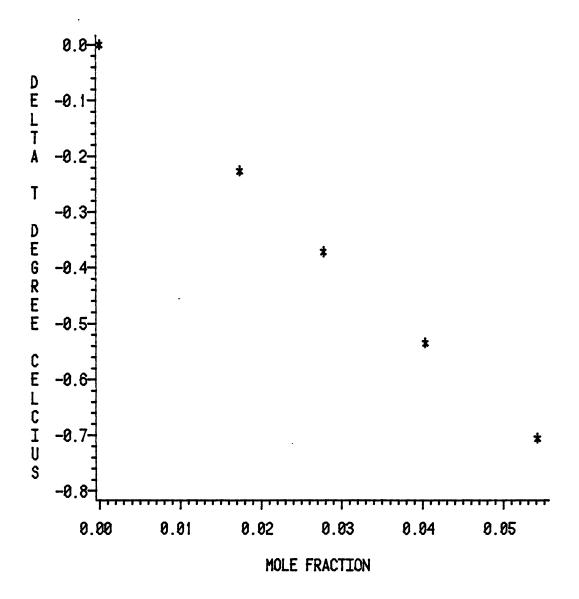


Figure 16. Effect of Composition on the Boiling Point Temperature of the Binary System Benzene (1) - Cyclohexane (2)



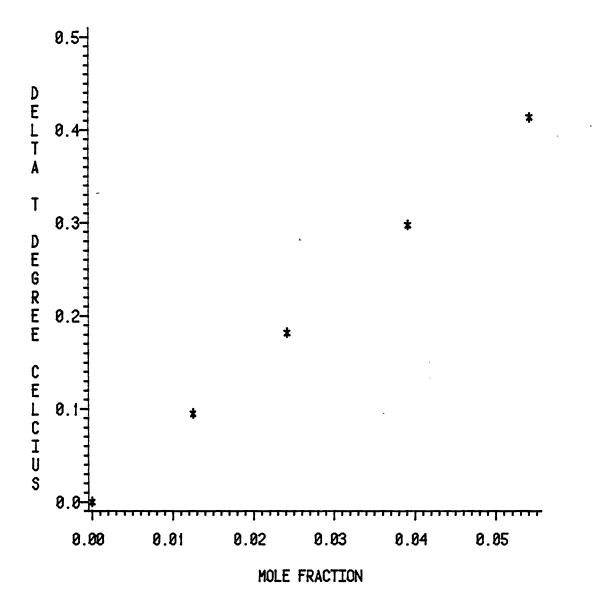


Figure 17. Effect of Composition on the Boiling Point Temperature of the Binary System Cyclohexane (1) - Hexane (2)

# HEXANE(1)-CYCLOHEXANE(2) AT 350 MMHG

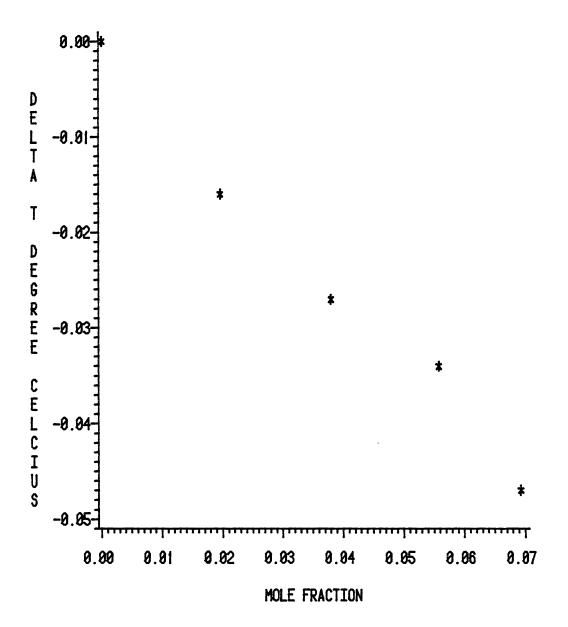


Figure 18. Effect of Composition on the Boiling Point Temperature of the Binary System Hexane (1) - Cyclohexane (2)

# BENZENE(1)-HEXANE(2) AT 350 MMHG

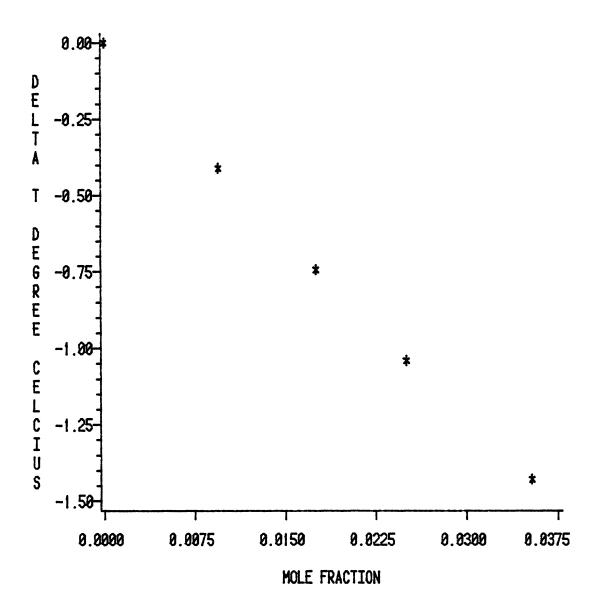


Figure 19. Effect of Composition on the Boiling Point Temperature of the Binary System Benzene (1) - Hexane (2)



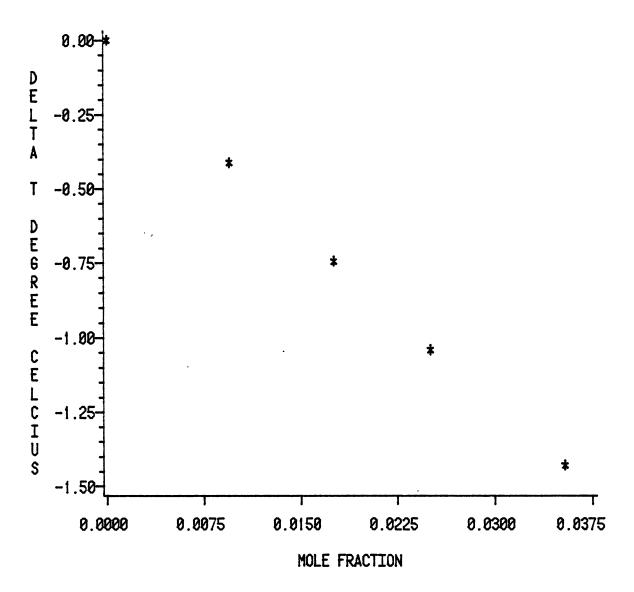


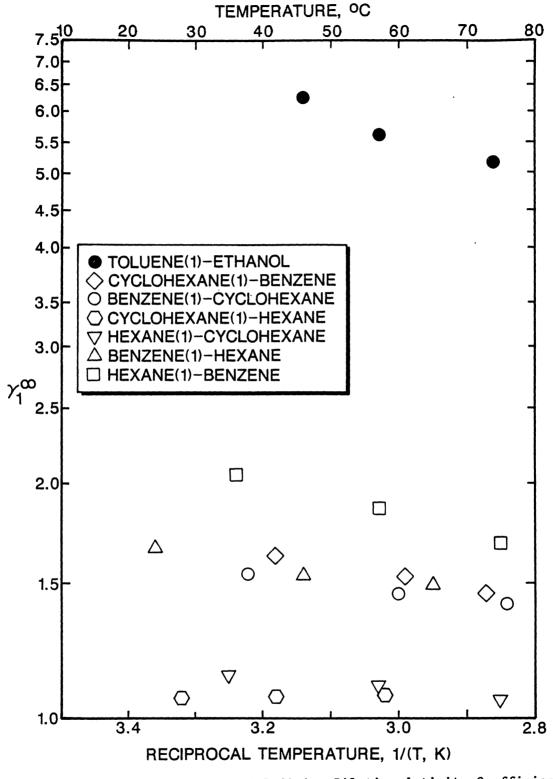
Figure 20. Effect of Composition on the Boiling Point Temperature of the Binary System Hexane (1) - Benzene (2)

composition on boiling point temperature. Again the linearity of the experimental data demonstrates the reliability of the experimental technique. Nearly the same change in the boiling point temperature is observed for a given amount of solute added in all of the binary systems studied. Figure 21 shows infinite dilution activity coefficients for the seven systems studied in this work. This plot is included to show relative magnitudes of  $\gamma_1^{\infty}$  for the various systems. The activity coefficients range between 1.06 and 6.23 for a temperature range of 24.950 C to 78.690 C. The largest activity coefficients correspond to the toluene (1) - ethanol (2) system. This means that the above system is a very non-ideal mixture. This can be attributed to the polarity of ethanol through its -OH bond. Ethanol was the only alcohol studied in this work. The hydrocarbons demonstrate similar non-idealities proven by the fact that the activity coefficients for all three groups: paraffins, naphthenes and aromatics are in the range of 1.06 to 2.05. In this range, paraffins and naphthenes exhibit similar non-idealities, while the aromatics are more non-ideal in a given mixture.

#### Discussion

Differential ebulliometry was used to obtain the data presented in Table VIII. The infinite dilution activity coefficients of all the systems studied are shown in Figure 21. The temperature fluctuations were reduced to about  $\pm 5 \times 10^{-3}$ °C by using a magnetic pump to provide forced circulation of the boiling liquid-vapor mixtures. Light hydrocarbons representing normal paraffins, naphthenes, and aromatics were used in this study because these compounds are important in chemical and petroleum processes and, primarily, because data are

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Figure 21. Comparison of Infinite Dilution Activity Coefficients for the Systems Studied

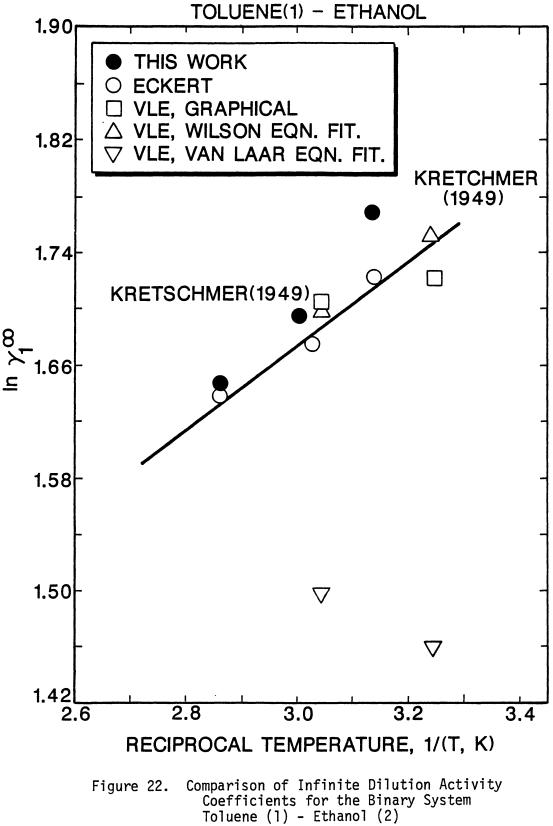
available in the literature for  $\gamma_1^{\infty}$  values which provide a means to assess the results from the modified ebulliometer.

Of the two models used to fit the data, the quadratic model gave better results, as expected. The differences between the experimental boiling point temperatures and the calculated values were on the order of 0.01 C for the linear case and 0.005 C for the quadratic case. The agreement is well within the experimental accuracy of the data. Since the quadratic model gives results which are more consistent with the experimental uncertainties, there is no reason why this model should not be used exclusively.

Plots of the effect of temperature on  $\gamma_1^{\infty}$  are a measure of the consistency of data. Figures 22 and 23 show a comparison of the measured  $\gamma_1^{\infty}$  values with data from the literature. The data of this work, the experimental ebulliometric data of Eckert et al., the data obtained by Kretchmer, Scatchard, and Nagata (30) from the graphical extrapolation of vapor-liquid equilibrium data, and the Wilson and Van Laar fits of the vapor-liquid equilibrium data are all shown in Figures 22 and 23. These figures were taken from reference (30) directly and the experimental data of this work were simply superimposed on Eckert's plot. The line drawn through the data points is based on the data of Eckert only.

Figure 22 shows the case where this work differs most with the Eckert data. The trend of the data points of this work have a different slope than the line based on Eckert's data; however, the other points are also scattered. In Figure 23, much better agreement is observed between the data of this work and those of Eckert et al. This is a more typical representation of the agreement observed between the two

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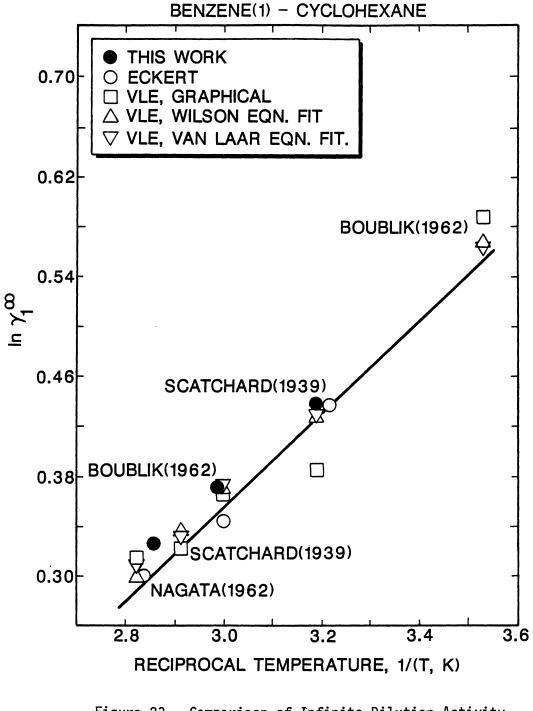


Figure 23. Comparison of Infinite Dilution Activity Coefficients for the Binary System Benzene (1) - Cyclohexane (2)

works. In fact, the present data appear to be in marginally better agreement with other workers' data than do the data of Eckert et al.

The sensitivity of the aquired data to the estimated amount of liquid holdup is illustrated in Table IX. The estimated value of liquid holdup used in the calculations was 2 cc. Results are shown for the systems with the lowest  $\gamma_1^{\infty}$  (hexane (1)-cyclohexane (2) system), and the highest  $\gamma_1^{\infty}$  (toluene (1)-ethanol (2) system). Values of 0.2 cc and 5 cc were used to test the sensitivity. As can be seen from this table, the smaller the holdup, the smaller the correction in the liquid mole fraction, which results in a smaller value for the infinite dilution activity coefficient. The effect of the liquid holdup is (as expected) greater on the larger  $\gamma_1^{\infty}$ 's. (The quadratically fitted  $\gamma_1^{\infty}$  was used in these comparisons.) The five percent change in the values of infinite dilution activity coefficient corresponding to a holdup range of 0.2 cc to 5 cc suggests that the reduced data obtained in this work with an estimate of liquid holdup equal to 2 cc would not be far from the actual values.

An attempt was made to study heavier hydrocarbons such as decane and 1-methylnaphthalene as a test of the performance of the apparatus for such components. When heavy hydrocarbons were used as solutes in light hydrocarbons, the performance of the apparatus was quite satisfactory, but Equation (12) failed to produce reliable results. Eckert (6) argues that when the solute is much less volatile than the solvent, as in the case of decane (1)-hexane (2) studied in this work, the ebulliometric method of determining infinite dilution activity coefficients may not be applicable. This was proven to be the case in the decane (1)-hexane (2) system, where a negative value of  $\gamma_1^{\infty}$  was

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## TABLE IX

## EFFECT OF LIQUID HOLDUP ON CALCUATED INFINITE DILUTION ACTIVITY COEFFICIENTS

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Liquid Holdup, cc	Calculated Value of <sub>۲1</sub> <sup>∞</sup> Hexane(1)-Cyclohexane(2) Toluene(1)-Ethanol(2		
0.2	1.055	6.096	
2	1.060	6.226	
5	1.067	6.399	

obtained using Equation (12). There is a 105°C difference between the boiling points of hexane and decane, and this could explain the failure of Equation (12) and the ebulliometric technique.

When heavy hydrocarbons were used as both solute and solvent, e.g., decane (1)-1-methylnaphthalene (2), the performance of the apparatus was not satisfactory. The higher temperatures at which the heavier hydrocarbons boiled resulted in the flaking off of the fused crushed glass inside the annulus section of the ebulliometer. This glass found its way into the magnetic pump and stopped its operation. The thick-walled tygon tubing was also attacked by the high boiling chemicals (specially aromatics) and collapsed.

Since the infinite dilution activity coefficient was calculated in an iterative manner, a one constant Margules Equation was deemed sufficient. Additional constants would not have much effect on the final  $\gamma_{1}^{\infty}$  values, but they would have made the calculations more difficult.

The linear and quadratic statistical fits gave infinite dilution activity coefficients which differed from one another by at most 6%, as shown in Table VIII.

The operating pressure in this work ranged between 150.0 mmHg and 726.7 mmHg. Typical values of the low, medium, and high pressures at which the experiments were conducted were 150 mmHg, 350 mmHg, and 700 mmHg. The above pressures corresponded to boiling point temperatures between 24.95°C to 78.69°C.

The first objective of this work which was to build and test a modified ebulliometer for better operation at low pressures was accomplished successfully. However, the second objective of this work which was to study heavy solvent-solute pairs was not achieved due

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mainly to mechanical problems such as the plugging of the magnetic vanetype pump with flaked off crushed glass, and the collapse of tygon tubing at the high temperatures at which heavier hydrocarbons such as normal decane and 1-methylnaphthalene boiled.

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#### CHAPTER VII

#### CONCLUSIONS AND RECOMMENDATIONS

#### Conclusions

The compact ebulliometer (which requires only about 110 cc of solvent) designed by Eckert and coworkers (6) was modified in this work to replace the thermosiphon-dependent circulation in the ebulliometer by a forced convection operation. This modification enhanced appreciably the performance of the ebulliometer at low pressures by improving the boiling behavior as a result of the constant flow of boiling liquid and vapor provided to the glass spiral on the thermometer well.

The system was tested using light hydrocarbons, and the performance was quite satisfactory. The data obtained were compared to those of other investigators (Table VIII, Figures 22 and 23) and were shown to be in good agreement (maximum deviation of 5%).

The performance of the apparatus was unsatisfactory when heavy hydrocarbons were used. The ebulliometric technique was shown not to be suitable for mixtures with widely varying boiling points (high relative volatility), as expected from the analysis given by Eckert et al. (6).

#### Recommendations

The apparatus of this work demonstrated excellent performance at low pressures with light hydrocarbons and should continue to be used. A few minor modifications could render it suitable for heavy hydrocarbons,

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as well. The recommended modifications follow.

- A more powerful magnetic pump could be used to reduce temperature fluctuations at high temperatures (greater than 150°C) by providing a higher flow rate of the boiling mixture to the thermometer well.
- Installation of a mesh screen at the pump inlet would eliminate the possibility of the flaked-off crushed glass entering the pump and stopping its operation.
- 3. Use of Teflon tubing which is less sensitive to the attack of highboiling aromatics could prevent the problems caused by the failure of the tygon tubing.

The quartz thermometer can operate at temperatures to 250°C. An attempt should be made to test heavier chemicals (which have boiling points below 250°C) in the apparatus once the above modifications have been implemented. A glass vacuum jacket should be used at high temperatures to avoid heat loss to the surroundings.

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APPENDIXES

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#### APPENDIX A

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#### COMPUTER PROGRAM FOR DATA REDUCTION

The following program was developed to reduce the experimental data to yield infinite dilution activity coefficients. It calculates a Margules Constant (Equation (29)) using the overall uncorrected mole fractions (x = z) based on a linear and a quadratic model:

$$\Delta T = ax \qquad (linear) \qquad (31)$$

$$\Delta T = ax + bx^2 \qquad (quadratic) \qquad (32)$$

It then uses the Margules constant to make corrections for the vapor and liquid holdups using molar material balances (Equations (38), (39)). Once the corrected values of the mole fractions are known, the program uses the same linear and quadratic functions to calculate "a" in the above equations and consequently the infinite dilution activity coefficient (Equation (12)).

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C	THIS PROGRAM MAKES CORRECTIONS FOR THE VAPOR AND LIQUID	00000130
С	HOLD-UPS BASED ON LINEAR AND QUADRATIC MODELS OF DELTA T VS X	00000140
С	AND THEN CALCULATES THE INFINITE DILUTION ACTIVITY COEFFICIENT	00000150
	REAL MW1, MW2, P1, P2, M2, M20, Y1(10), Y2(10)	00000160
	REAL K1(10),K2(10),M1(10),M(10),T(10),VLG(10),ROLG(10)	00000170
	DIMENSION X(10),Y(10),Z(10),ZC(10),T1(10)	00000180
	DIMENSION XC(10),XT(10),X2(10),AZ(10),AM1(10)	00000190
	DIMENSION ROL(10), ROV(10), VL(10), VV(10)	00000200
	DIMENSION TCALC(10), DELTAT(10), X2T(10), X3(10), X4(10)	00000210
	WRITE (6,79)	00000220
79	FORMAT (1H1,////, ' SYSTEM: TOLUENE(1)-ETHANOL(2)')	00000230
	READ (9,*) N,P,R,TC,P3,AM20	00000240
С	N IS THE NUMBER OF INJECTIONS	00000250
С	P IS THE SYSTEM PRESSURE IN PSIA	00000260
С	R IS THE GAS CONSTANT	00000270
C	TC IS THE BOILING POINT OF PURE SOLVENT IN DEGREES C	00000280
c	P3 IS THE VAPOR PRESSURE OF THE SOLUTE IN PSIA	00000290
Ċ	AM20 IS THE GRAMS OF SOLVENT CHARGED	00000300
•	WRITE (6,*) ' VALUES OF N.P.R'	00000310
	WRITE (6,*) N.P.R	00000320
	WRITE (6,*) ' VALUES OF TC,P3,AM20'	00000330
		00000340
	WRITE (6,*) TC,P3,AM20 WRITE (6,*) ' VALUES OF T(1) THROUGH T(N)' DD 1001 L=1,N	00000350
	D0 1001 L=1.N	00000360
	READ (9,+) T1(L)	00000370
С	T(1) THROUGH T(N) ARE THE EXPERIMENTAL DELTA T IN DEGREE C	00000380
U	WRITE (6.*) T1(L)	00000390
1001	CONTINUE	00000390
1001	WRITE (6,*) ' VALUES OF AM1(1) THROUGH AM1(N)'	
		00000410
	D0 1002 L=1,N	00000420
•	READ (9,*) AM1(L)	00000430
С	AM1(1) THROUGH AM1(N) ARE THE GRAMS OF SOLUTE INJECTED	00000440
	WRITE (6,*) AM1(L)	00000450
1002	CONTINUE	00000460
	WRITE (6,*) ' VALUES OF PC1,PC2,TC1,TC2'	00000470
_	READ (9,+) PC1,PC2,TC1,TC2	00000480
C C	PC1 IS THE CRITICAL PRESSURE OF COMPONENT 1 IN PSIA	00000490
С	PC2 IS THE CRITICAL PRESSURE OF COMPONENT 2 IN PSIA	00000500
C	TC1 IS THE CRITICAL TEMPERATURE OF COMPONENT 1 IN K TC2 IS THE CRITICAL TEMPERATURE OF COMPONENT 2 IN K	00000510
С		00000520
	WRITE (6,*) PC1,PC2,TC1,TC2	00000530
	WRITE (6,*) ' VALUES OF VVY,VLY,VX,VZ'	00000540
	READ (9,*) VVY,VLY,VX,VZ	00000550
C C	VVY IS THE VAPOR VOLUME IN THE EBULLIOMETER IN CC	00000560
С	VLY IS THE CONDENSED VAPOR VOLUME IN CC	00000570
С	VX IS THE LIQUID VOLUME IN CC	00000580
С	VZ IS THE INITIAL VOLUME IN CC	00000590
	WRITE (6,*) VVY,VLY,VX,VZ	00000600

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	WRITE (6,*) ' VALUES OF A11,A21,A31'	00000720
-	READ (9,*) A11,A21,A31	00000730
С	A11, A21, AND A31 ARE ANTOINE CONSTANTS OF COMPONENT 1	00000740
	WRITE (6.*) A11.A21.A31	00000750
	WRITE (6,*) ' VALUES OF A12,A22,A32'	00000760
_	READ (9,*) A12,A22,A32	00000770
С	A12,A22,AND A32 ARE THE ANTOINE CONSTANTS OF COMPONENT 2	00000780
	WRITE (6.*) A12,A22,A32	00000790
	WRITE (6,*) ' VALUES OF AR1,AR2,MW1,MW2'	00000800
	READ (9,*) AR1,AR2,MW1,MW2	00000810
С	AR1 AND AR2 ARE RACKETT EQUATION CONSTANTS	00000820
С	MW1 AND MW2 ARE THE MOLECULAR WEIGHTS OF SOLUTE AND SOLVENT	00000830
	WRITE (6,*) AR1.AR2.MW1.MW2	00000840
	LQ=1	00000850
	TK=TC+273.15	00000860
С	TK IS THE BOILING POINT TEMPERATURE OF THE SOLVENT IN KELVIN	00000870
	TF=TC+1.8+32	00000880
С	TF IS THE BOILING POINT TEMPERATURE OF THE SOLVENT IN DEGREE I	F 00000890
	M2C=AM20/MW2	00000900
С	GRAMS OF SOLVENT CONVERTED TO MOLES OF SOLVENT	00000910
	DO 112 I=1.N	00000920
	M1(I)=AM1(I)/MW1	00000930
С	GRAMS OF SOLUTE CONVERTED TO MOLES OF SOLUTE .	00000940
J	Z(I) = M1(I)/(M1(I) + M2O)	
С	MOLES OF SOLUTE ARE CONVERTED TO MOLE FRACTION	00000950
112	CONTINUE	00000960
112		00000970
	DO 123 I=1,N	00000980
	X(I)=Z(I)	00000990
•	AZ(I)=Z(I)	00001000
C	THE LIQUID MOLE FRACTION IS SET EQUAL TO THE INITIAL VALUE	00001010
123	CONTINUE	00001020
	Q=1	00001030
	GO TO 134	00001040
8	Q = 1	00001050
41	DO 16 I=1,N	00001060
1	DO 977 K=1,N	00001070
	T(K)=T1(K)	00001080
977	CONTINUE	00001090
	P1=PC1*EXP(A11-(A21/(A31+TF+(T(I)*1.8))))	00001100
С	P1 IS THE VAPOR PRESSURE OF THE SOLUTE CALCULATED BY ANTOINE (	
•	P2=PC2+EXP(A12-(A22/(A32+TF+(T(I)+1.8))))	00001120
С	P2 IS THE VAPOR PRESSURE OF THE SOLVENT CALCULATED BY ANTOINE I	
-		00001140
	X(I)=Z(I)	00001150
	VL1=83.147*TC1*(14.696/PC1)*AR1**(1+(1-((TK+T(I))/TC1))**0.28	
+	1)	00001170
•	VL2=83.147*TC2*(14.696/PC2)*AR2**(1+(1-((TK+T(I))/TC2))**0.28	
-		
+	1)	00001190

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С	HERE THE LIQUID MOLAR VOLUMES OF COMPONENTS 1 AND 2 ARE	00001310
С	CALCULATED USING THE RACKETT EQUATION	00001320
15	VL(I)=VL1+X(I)+VL2+(1-X(I))	00001330
С	THIS IS THE MIXTURE MOLAR VOLUME IN CC/GMOL	00001340
	J+1	00001350
	ROI(T) = 1/VI(T)	00001360
С	POLISTIC INCLUD DENSITY IN CMOL/CC	00001380
57		00001370
57	$(1)^{-}$	00001380
•	CALCULATED USING THE RACKETT EQUATION VL(I)=VL1*X(I)+VL2*(1-X(I)) THIS IS THE MIXTURE MOLAR VOLUME IN CC/GMOL J=J+1 ROL(I)=1/VL(I) ROL IS THE LIQUID DENSITY IN GMOL/CC K1(I)=(P1/P)*EXP(A*(1-X(I))**2) K2(I)=(P2/P)*EXP(A*X(I)**2)	00001390
С	THESE ARE THE EQUILIBRIUM CONSTANTS OF COMPONENTS 1 & 2	00001400
	Y1(I)=K1(I)*X(I)	00001410
	Y2(I)=K2(I)+(1-X(I))	00001420
С	THESE ARE THE VAPOR MOLE FRACTIONS	00001430
	Y(I) = Y1(I)/(Y1(I) + Y2(I))	00001440
С	Y IS NORMALIZED HERE	00001450
	VLG(I) = VL(I + Y(I) + VL(2 + (1 - Y(I)))	00001460
	Y2(I)=K2(I)+(1-X(I)) THESE ARE THE VAPOR MOLE FRACTIONS Y(I)=Y1(I)/(Y1(I)+Y2(I)) Y IS NORMALIZED HERE VLG(I)=VL1+Y(I)+VL2+(1-Y(I)) ROLG(I)=1/VLG(I)	00001470
С	VLG AND ROLG ARE THE CONDENSED VAPOR MOLAR VOLUME AND DENSITY	00001470
U	ROV(I) ≥P/(R*(TK+T(I))*14.696)	
101	$K_0 \vee (1)^{-r} / (K^* (1K^* (11))^* 14.050)$	00001490
101	M(1) = M(1) - ((VV) + ROV(1) + VL + ROLG(1)) + Y(1))	00001500
•	M2=M2O-((VVY+ROV(I)+VLY+ROLG(I))+(1-Y(I)))	00001510
С	M(I)=M1(I)-((VVY*ROV(I)+VLY*ROLG(I))*Y(I)) M2=M2O-((VVY*ROV(I)+VLY*ROLG(I))*(1-Y(I))) THESE ARE THE COMPONENT MOLE BALANCES XC(I)=M(I)/(M(I)+M2)	00001520
	XC(I)=M(I)/(M(I)+M2)	00001530
С	THIS IS THE NEW VALUE OF X USING THE MATERIAL BALANCE	00001540
	IF (ABS(XC(I)-X(I)).LE.O.0001) GO TO 16	00001550
С	THIS IS THE CONVERGENCE CRITERION	00001560
	X(I)=XC(I)	00001570
	GO TO 15	00001580
16	CONTINUE	00001590
	IF (LO.FO.2) GO TO 1340	00001600
С	THESE ARE THE COMPONENT HOLE BALANCES xC(I)=M(I)/(M(I)+M2) THIS IS THE NEW VALUE OF X USING THE MATERIAL BALANCE IF (ABS(XC(I)-X(I)).LE.O.0001) GD TO 16 THIS IS THE CONVERGENCE CRITERION x(I)=xC(I) GD TO 15 CONTINUE IF (LQ.EQ.2) GD TO 1340 THIS SECTION IS USED IF A LINEAR FIT IS DESIDED	*00001610
č	THIS SECTION IS USED TO A LINEAD FIT IS DESIDED	00001620
134		00001630
134	$\frac{1}{2} \frac{1}{2} \frac{1}$	
		00001640
	X2(I)=X(I)++2	00001650
154	CONTINUE	00001660
	DO 19 I=2,N	00001670
	XT(I)=XT(I)+XT(I-1)	00001680
	X2(I)=X2(I)+X2(I-1)	00001690
С	THESE ARE SUMMATIONS USED TO GET THE DT/DX TERM	00001700
19	CONTINUE	00001710
	DT = XT(N)/X2(N)	00001720
С	THIS IS THE DT/DX TERM	00001730
-	DO 185 I=1.N	00001740
	TCALC(I) = DT + X(I)	00001750
с	THIS SECTION IS USED IF A LINEAR FIT IS DESIRED DO 154 I=1,N XT(I)=X(I)+T1(I) X2(I)=X(I)++2 CONTINUE DO 19 I=2,N XT(I)=X2(I)+X2(I-1) THESE ARE SUMMATIONS USED TO GET THE DT/DX TERM CONTINUE DT=XT(N)/X2(N) THIS IS THE DT/DX TERM DO 185 I=1,N TCALC(I)=DT+X(I) THIS IS THE LINEARLY FITTED VALUE OF DELTA T DELTAT(I)=T1(I)-TCALC(I) THIS IS THE DIPERDENCE RETWEEN THE EXPERIMENTAL AND THE	00001760
-	DELTAT(I)=T1(I)-TCALC(I)	00001770
С	THIS IS THE DIFFERENCE BETWEEN THE EXPERIMENTAL AND THE	00001780
c	STATISTICALLY OBTAINED DELTA T	00001790
L.	STATISTICALLY UDIAINED DELTA I	00001790

	DP2=P*(A22/((TF+A32)**2))	
С	THIS IS THE DP2/DT TERM	00001920
•	G1 = (DP2 + DT + 1, B))/P3	00001920
С	THIS THE INFINITE DILUTION ACTIVITY COEFFICIENT	
J	IF (Q.EQ.2) GO TO 43	00001940
	AG1=G1	00001950
	WRITE (6.79)	00001960
	WRITE (6.998)	00001970
998	FORMAT (///,2X,' LINEAR MODEL:')	00001980
000	WRITE (6,731) AG1	00001990
731	FORMAT (///, ' THE ORIGINAL VALUE OF G1 IS: ', F10.5)	00002000
	Q=Q+1	00002010
	A=ALOG(G1)	00002020
С	THIS IS THE MARGULES CONSTANT	00002030
C	GO TO 41	00002040
43	AC=ALDG(G1)	00002050
40	IF (ABS(AC-A).LE.0.001) GO TO 420	00002060
	A=AC	00002070
	DO 146 J=1.N	00002080
		00002090
	X1(U)=0	00002100
146	X2(J)=0 CONTINUE	00002110
140	GO TO 41	00002120
420		00002130
420	WRITE (6, 1840)	00002140
1840	FORMAT (///,6X,'INJECTION NO',6X,'Z(I)',10X,'X(I)',10X,'T(I)', BX,'TCALC(I)',8X,'DELTAT(I)')	
Ŧ	WRITE (6,1860) I,Z(I),X(I),T1(I),TCALC(I),DELTAT(I)	00002160
1860	FORMAT (12X, 15, 2X, 5(F12, 5, 2X))	00002170
185	CONTINUE	00002180
100	GO TO 1000	00002190
с	**************************************	00002200
c	THIS SECTION IS USED IF A QUADRATIC FIT IS DESIRED	
1340	DO 1440 I=1.N	00002220
1340	XT(I)=X(I)+T1(I)	00002230
	$x_2(1) = x(1) + 1$	00002240
	X3(I)=X(I)++3	00002250
	$X_{4}(I) = X(I) + 4$	00002260
	$x_4(1) = x_1(1) = x_4$ $x_2T(1) = x_2(1) = T_1(1)$	00002270
1440		00002280
1440	CONTINUE	00002290
	D0 190 I=2,N	00002300
	XT(I)=XT(I)+XT(I-1)	00002310
	$X_2(I) = X_2(I) + X_2(I-1)$	00002320
	X3(I)=X3(I)+X3(I-1)	00002330
	X4(I)=X4(I)+X4(I-1)	00002340
•	X2T(I) = X2T(I) + X2T(I-1)	00002350
C	THESE ARE SUMMATIONS USED TO GET THE DT/DX TERM	00002360
190		00002370
•	B=(XT(N)-(X2T(N)+X2(N))/X3(N))/(X3(N) -(X4(N)+X2(N)/X3(N)))	00002380
+	~(X4(N)*AZ(N)/AJ(N)))	00002390 00002400
		0002400

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С	THIS IS THE COEFFICIENT OF X**2 IN THE QUADRATIC FIT	00002510
	DT=(XT(N)-B+X3(N))/X2(N)	00002520
С	THIS IS THE DT/DX TERM	00002530
	DO 1850 I=1,N	00002540
	TCALC(I)=DT+X(I)+B+(X(I)++2)	00002550
С	THIS IS THE CALCULATED VALUE OF T	00002560
	DELTAT(I)=T1(I)-TCALC(I)	00002570
С	THIS IS THE DIFFERENCE BETWEEN THE EXPERIMENTAL AND THE	00002580
С	STATISTICALLY FIT VALUE OF DELTA T	00002590
	DP2=P*(A22/((TF+A32)**2))	00002600
С	THIS IS THE DP2/DT TERM	00002610
	G1=(P-(DP2+DT+1.8))/P3	00002620
С	THIS IS THE INFINITE DILUTION ACTIVITY COEFFICIENT	00002630
	IF (Q.GE.2) GO TO 430	00002640
	AG1=G1	00002650
	WRITE (6.79)	00002660
	WRITE (6,999)	00002670
999	FORMAT (///,2X,' QUADRATIC MODEL:')	00002680
	WRITE (6,737) G1	00002690
737	FORMAT (///, ' THE ORIGINAL VALUE OF G1 IS: '.F10.5)	00002700
	Q=Q+1	00002710
	A=ALOG(G1)	00002720
С	THIS IS THE MARGULES CONSTANT	00002730
-	GO TO 41	00002740
430	AC=ALOG(G1)	00002750
	IF (ABS(AC-A), LE.0.001) GO TO 42	00002760
	A=AC	00002770
	DO 1980 J=1.N	00002780
	XT(J)=0	00002790
	X2(J)=0	00002800
	X3(J)=0	00002800
	X4(J)=0	00002810
	X2T(J)=0	00002820
1980	CONTINUE	00002830
1980	GO TO 41	00002840
42	WRITE (6,184)	
184	FORMAT (///.6X.'INJECTION NO'.6X.'Z(I)'.10X.'X(I)'.10X.'T(I)'.	00002860
+	8X,'TCALC(I)',8X,'DELTAT(I)') WRITE (G.186) I,Z(I),X(I),T1(I),TCALC(I),DELTAT(I)	00002880
400		00002890
186	FORMAT (12X, I5, 2X, 5(F12.5, 2X))	00002900
1850	CONTINUE	00002910
-	***************************************	00002920
1000	WRITE (6,23)	00002930
23	FORMAT('O', 5X, 'TC', 12X, 'P', 11X, 'G1', 10X, 'DT', 10X, 'DP2', 12X, 'A'	
	WRITE (6,20) TC,P,G1,DT,DP2,A	00002950
		00002960

		00003130
20	FORMAT (4X,6(F10.5,2X))	00003140
	EP=((G1-AG1)/G1)*100	00003150
	WRITE (6,937) EP	00003160
937	FORMAT (///,2X,' THE PERCENT CHANGE IN G1 IS: ',F10.5)	00003170
	IF (LQ.EQ.2) GO TO 657	00003180
	LQ=LQ+1	00003190
	GO TO 8	00003200
657	STOP	00003210
	END	00003220

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SYSTEM: TOLUENE(1)-ETHANOL(2) VALUES OF N,P,R 82.0500000 5 3.4820000 VALUES OF TC, P3, AM20 45.6000000 1.4829990 84.3314900 VALUES OF T(1) THROUGH T(N) -0.2250000 -0.4060000 -0.5910000 -0.7220000 -0.8800000 VALUES OF AM1(1) THROUGH AM1(N) 1.3267000 2.5341990 3.8789990 4.9127990 6.2798000 VALUES OF PC1, PC2, TC1, TC2 587.8000000 925.3000000 591.7900000 516.2600000 VALUES OF VVY,VLY,VX,VZ 200.0000000 2.0000000 2.0000000 106.0000000 VALUES OF A11, A21, A31 5.9442500 5836.2850000 374.7451000 VALUES OF A12, A22, A32 7.4358190 6162.3590000 359,3825000 VALUES OF AR1, AR2, MW1, MW2 0.2646000 0.2520000 92.1139900 46.0690000 SYSTEM: TOLUENE(1)-ETHANOL(2)

LINEAR MODEL:

THE ORIG	INAL	VALUE OF G1 IS:	5.30019			
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	1	0.00781	0.00763	-0.22500	-0.19745	-0.02755
INJECTION	N0	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	2	0.01481	0.01450	-0.40600	-0.37512	-0.03088
INJECTION	ND	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	3	0.02249	0.02206	-0.59100	-0.57065	-0.02035
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	4	0.02831	0.02780	-0.72200	-0.71926	-0.00274
INJECTION	N0	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	5	0.03591	0.03531	-0.88000	-0.91350	0.03350
TC 45.60001		P G1 9.48200 5.35352	DT 2 -25.86975	DP2 0.09572	A 1.67758	

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THE PERCENT CHANGE IN G1 IS: 0.99625

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SYSTEM: TOLUENE(1)-ETHANOL(2)

QUADRATIC MODEL:

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THE ORIGINAL VALUE OF G1 IS: 5.85281

INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	1	0.00781	0.00760	~0.22500	-0.22150	-0.00350
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	2	0.01481	0.01445	-0.40600	-0.40598	-0.00002
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	3	0.02249	0.02199	-0.59100	-0.59270	0.00170
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	4	0.02831	0.02772	-0.72200	-0.72318	0.00118
INJECTION	NO	Z(I)	X(I)	T(I)	TCALC(I)	DELTAT(I)
	5	0.03591	0.03521	-0.88000	-0.87877	-0.00123
TC	Р	G1	DT	DP2	A	
45.60001	3.48	200 5.8664	0 -30.28419	0.09572	1.76919	

THE PERCENT CHANGE IN G1 IS: 0.23169

84

.

#### APPENDIX B

## DERIVATION OF $\gamma_1^{\infty}$ EXPRESSION

The criterion for equilibrium between two phases "I" and "II" is that their chemical potentials be equal:

$$\mu_{i}^{I} = \mu_{i}^{II} \quad (i = 1, N) \tag{1}$$

if "I" is taken as the vapor phase and "II" is taken as the liquid phase, the vapor-liquid equilibrium is demonstrated by:

$$\nu_{i} V = \nu_{i} L$$
 (2)

In a given phase, the quantity  $\mu_{i}$  for any component can be expressed as the sum of two terms:

$$\mu_{i} = \mu_{i}^{0} + \Delta \mu_{i}$$
 (3)

The term  $\mu_i^0$  is the chemical potential of component "i" in the standard state, where standard state refers to the property of the pure component at the system temperature and total pressure. The quantity  $\Delta\mu_i$  is known as the partial molal free energy of mixing and is evaluated in terms of the activity coefficient,  $\gamma$ , as follows.

$$\Delta \mu_{i} = RT \ln(\gamma_{i} x_{i})$$
(4)

Therefore Equation (3) can be rewritten as:

$$\mu_{i} = \mu_{i}^{0} + RT \ln(\gamma_{i} x_{i})$$
(5)

When Equation (5) is applied to two phases (vapor and liquid) in equilibrium, using Equation (2):

$$(\mu_{i}^{0})^{V} + RT \ln(\gamma_{i}^{V} x_{i}^{V}) = (\mu_{i}^{0})^{L} + RT \ln(\gamma_{i}^{L} x_{i}^{L})$$
 (6)

where  $x_i^L = x_i$  is the liquid composition, and  $x_i^V = y_i$  is the vapor composition. Upon rearranging:

RT ln 
$$\frac{x_i \gamma_i^L}{y_i \gamma_i} = (\mu_i^0)^V - (\mu_i^0)^L$$
 (7)

The equilibrium ratio is defined as:

$$K_{i} = \frac{y_{i}}{x_{i}} = \frac{\gamma_{i}^{L}}{\gamma_{i}^{V}} \exp\left[\frac{(\mu_{i}^{0})^{L} - (\mu_{i}^{0})^{V}}{RT}\right]$$
(8)

Using the following thermodynamic relationship:

$$d\mu_{i} = -s_{i}dT + v_{i}dP \qquad (9)$$

and integrating between the total system pressure and the vapor pressure of component "i", at constant temperature (dT = 0):

$$\Delta \mu_{i} = \int_{P_{i}}^{P} (v_{i})^{L} dP$$
 (10)

$$\Delta \mu_{i} = \int_{P_{i}}^{P} \frac{z_{i}RT}{P} dP$$
(11)

$$\Delta \mu_{i} = RT \ln \frac{P}{P_{i}s} - \int_{P_{i}s}^{P} [\frac{RT}{P} - (v_{i})^{V}]dP \qquad (12)$$

$$(\mu_{i}^{0})^{V} - (\mu_{i}^{0})^{L} = (\nu_{i}^{1})^{L}(P^{S} - P) + RT \ln \frac{P}{P^{S}} - \int_{P}^{P} [\frac{RT}{P} - (\nu_{i}^{1})^{V}]dP$$
 (13)

Substituting in Equation (9):

$$K_{i} = \frac{y_{i}}{x_{i}} = \frac{\gamma_{i}^{L}}{\gamma_{i}^{V}} \frac{P_{i}^{S}}{P} \exp\left[\frac{(v_{i})^{L}(P - P_{i}^{S})}{RT}\right] \exp\left\{\frac{1}{RT}\int_{P_{i}^{S}}^{P} (v_{i})^{V}\right] dP\}$$
(14)

The fugacity of component "i" is defined as follows:

$$f_{i}^{s} = P_{i}^{s} \exp\{\frac{1}{RT} \int_{0}^{P_{i}^{s}} [(v_{i})^{V} - \frac{RT}{P}] dP\}$$
 (15)

$$f_i^{p} = P \exp \left\{\frac{1}{RT} \int_0^P \left[\left(v_i\right)^V - \frac{RT}{P}\right] dP\right\}$$
(16)

$$\frac{f_i^{S}}{f_i^{P}} = \frac{P_i^{S}}{P} \exp \left\{\frac{1}{RT} \int_{P_i^{S}}^{P} \left[\frac{RT}{P} - (v_i^{V})^{V}\right] dP\right\}$$
(17)

Substituting Equation (17) in Equation (14):

$$\frac{y_{i}}{x_{i}} = \frac{\gamma_{i}^{L}}{\gamma_{i}^{V}} \frac{f_{i}^{S}}{f_{i}^{p}} \exp\left[\frac{(v_{i})^{L} (P - P_{i}^{S})}{RT}\right]$$
(18)

For a binary system:

$$f_{1}Py_{1} = \frac{\gamma_{1}^{L}}{\gamma_{1}} f_{1}^{S} x_{1}exp[\frac{(v_{1})^{L} (P - P_{1}^{S})}{RT}]$$
(19)

$$f_{2}Py_{2} = \frac{\gamma_{2}^{L}}{\gamma_{2}^{V}} f_{2}^{S} x_{2} \exp \left[\frac{(v_{2})^{L} (P - P_{2}^{S})}{RT}\right]$$
(20)

Using the definition of the fugacity coefficient:

$$\phi_{i} = \frac{f_{i}}{P} = \exp \left\{\frac{1}{RT} \int_{0}^{P} \left[\left(v_{i}\right)^{V} - \frac{RT}{P}\right] dP\right\}$$
(21)

Equations (19) and (20) can be rewritten as:

$$Py_{1\phi_{1}}^{p} = \frac{\gamma_{1}^{L}}{\gamma_{1}} P_{1}^{s} \phi_{1}^{s} x_{1} \exp[\frac{(v_{1})^{L} (P - P_{1}^{s})}{RT}]$$
(22)

$$Py_{2} \phi_{2}^{P} = \frac{\gamma_{2}^{L}}{\gamma_{2}^{V}} P_{2}^{S} \phi_{2}^{S} x_{2} \exp[\frac{(v_{2})^{L} (P - P_{2}^{S})}{RT}]$$
(23)

For a binary system:

$$P = P(y_1 + y_2)$$
(24)

Solving Equation (22) and (23) for  $\text{Py}_1$  and  $\text{Py}_2$  and adding:

$$P = \frac{\gamma_{1}^{L}\phi_{1}^{s}}{\gamma_{1}^{V}\phi_{1}^{p}} P_{1}^{s}x_{1}exp[\frac{(v_{1})^{L}(P-P_{1}^{s})}{RT}] + \frac{\gamma_{2}^{L}\phi_{2}^{s}}{\gamma_{2}^{V}\phi_{2}^{p}} P_{2}^{s}x_{2}exp[\frac{(v_{2})^{L}(P-P_{2}^{s})}{RT}]$$
(25)

Assuming ideal gas phase solution:  $\gamma_1^V = \gamma_2^V = 1$ , and differentiating with respect to solute composition:

$$\frac{dP}{dx_{1}} = P_{1}^{S} \phi_{1}^{S} exp[\frac{(v_{1})^{L}(P-P_{1}^{S})}{RT}][\frac{\gamma_{1}x_{1}(v_{1})^{L}}{\phi_{1}^{P}RT}\frac{dP}{dx_{1}} - \frac{\gamma_{1}x_{1}}{(\phi_{1}^{P})^{2}}\frac{d\phi_{1}^{P}}{dP}\frac{dP}{dx_{1}} + \frac{\chi_{1}}{\phi_{1}^{P}}\frac{dP}{dx_{1}} + \frac{\gamma_{1}}{\phi_{1}^{P}}] + P_{2}^{S} \phi_{2}^{S} exp[\frac{(v_{2})^{L}(P-P_{2}^{S})}{RT}]\frac{\gamma_{2}x_{2}(v_{2})^{L}}{\phi_{2}^{P}RT}\frac{dP}{dx_{2}} - \frac{\varphi_{1}^{2}x_{1}}{\phi_{2}^{P}RT}\frac{dP}{dx_{2}}]$$

$$\frac{\gamma_{2}x_{2}}{\left(\phi_{2}^{P}\right)^{2}} \frac{d\phi_{2}^{P}}{dP} \frac{dP}{dx_{2}} + \frac{x_{2}}{\phi_{2}^{P}} \frac{d\gamma_{2}}{dx_{2}} + \frac{\gamma_{2}}{\phi_{2}^{P}} \frac{dx_{2}}{dx_{1}}$$
(26)

where the superscripts of  $\gamma_1^L$  and  $\gamma_2^L$  were dropped. Letting  $x_1$  approach zero (infinitely dilute solution):

$$\frac{dx_2}{dx_1} = -1$$

$$\frac{d\gamma_1}{dx_1} = \text{finite at } x_1 = 0$$

$$\begin{array}{c}
\gamma_{2} \neq 1 \\
x_{2} \neq 1 \\
\frac{d\gamma_{2}}{dx_{2}} \neq 0 \\
P \neq P_{2}^{S} \\
(\frac{dP}{dx_{1}})^{\infty} = \frac{\gamma_{1}^{\infty}P_{1}^{S}\phi_{1}^{S}}{(P_{2}^{S})} \exp[\frac{(v_{1})^{L}(P_{2}^{S}-P_{1}^{S})}{RT} + P_{2}^{S}\{[\frac{(v_{2})^{L}}{RT} - \frac{1}{\phi_{2}^{S}}\frac{d\phi_{2}^{P}}{dP}](\frac{d\phi_{2}^{P}}{dP})^{\infty}-1\} \\
\phi \qquad (27)
\end{array}$$

Solving for 
$$\gamma_1^{\infty}$$
:

$$\gamma_{1}^{\infty} = \frac{\phi_{2}^{S}}{(\phi_{1}^{P})^{S}} \frac{P_{2}^{S} [1 - \frac{(v_{2})^{L}}{RT} + \frac{1}{\phi_{2}^{S}} \frac{d\phi_{2}^{P}}{dP} (\frac{dP}{dx_{1}})^{\infty}] + (\frac{dP}{dx_{1}})^{\infty}}{P_{1}^{S} exp[\frac{(v_{1})^{L} (P_{2}^{S} - P_{1}^{S})}{RT}]}$$
(28)

If the gas phase is ideal gas:  $\phi_1 = \phi_2 = 1$  and  $\frac{d\phi}{dP} = 0$ ,

$$\gamma_{1}^{\infty} = \frac{P_{2}^{S} + [1 + \frac{P_{2}^{S}(v_{2})^{L}}{RT}](\frac{dP}{dx_{1}})_{T}^{\infty}}{P_{1}^{S}} \exp \left[\frac{(v_{1})^{L}(P_{1}^{S} - P_{2}^{S})}{RT}\right]$$
(29)

If the liquid compressibility factor and the Poynting effect are neglected:

$$\gamma_{1}^{\infty} = \frac{P_{2}^{S} + (\frac{dP}{dx_{1}})^{\infty}}{P_{1}^{S}}$$
(30)

but 
$$\left(\frac{dP}{dx_1}\right)_T = - \left(\frac{dP}{dT}\right)_{x_1} \left(\frac{dT}{dx_1}\right)_P$$
 (31)

and for 
$$x_1 = 0$$
:  $\left(\frac{dP}{dx_1}\right)_T^{\infty} = -\left(\frac{dP}{dT}\right) \left(\frac{dT}{dx_1}\right)_T^{\infty} P$  (32)

.

.

.

$$\gamma_{1}^{\infty} = \frac{P_{2}^{S} - (\frac{dP_{2}^{S}}{dT})(\frac{dT}{dx_{1}})_{P}^{\infty}}{P_{1}^{S}}$$
(33)

#### APPENDIX C

#### TABLE OF THE CONSTANTS USED

In the following table all the constants used to obtain and reduce the data are listed. The Antoine Equation constant  $A_1$ , and the Rackett Equation constant A are unitless. The units for the other constants are specified in parentheses after the symbol.

Antoine Constants:				
Compound	A_1	A <sub>2</sub> (°F)	A3(°F)	
Benzene	5.658375	5307.813	379.456	
Cyclohexane	5.473055	5030.253	371.2755	
Ethanol	7.43437	6162.36	359.3826	
Hexane	6.039243	5085.758	382.794	
Toluene	5.944251	5836.287	374.745	
Hexane	6.039243	5085.758	382.794	

These constants are consistent with the following Equations:

$$\ln(\frac{P_{i}}{P_{ci}})^{0} = A_{1} - \frac{A_{2}}{A_{3} + T}$$

where T' is in °F

Critical Properties, Rackett Equation constants, and Molecular Weights

P <sub>c</sub> (psia	<u>A</u>	Mol. Wt.
6 714.2	0.2696	78.12
4 591.5	0.2729	84.16
6 925.3	0.2520	46.07
3 440.0	0.2635	86.18
9 587.8	0.2646	92.14
	6 714.2 4 591.5 6 925.3 3 440.0	6       714.2       0.2696         4       591.5       0.2729         6       925.3       0.2520         3       440.0       0.2635

The gas constant:  $R = 82.05 \frac{\text{cm}^3 \text{ atm}}{\text{gmole K}}$ 

Values obtained from:

Henley, E.J., Seader, J.D., "Equilibrium Stage Separation Operations in Chemical Engineering", Wiley, New York, 1981, 714-724.

# VITA

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