# BINARY VAPOR-LIQUID PHASE EQUILIBRIUM FOR CARBON MONOXIDE + HYDROCARBONS

By

XIAOHUA YI <sup>#</sup> Bachelor of Engineering Tianjin University Tianjin, P. R. China

1984

Submitted to the Faculty of the Graduate College of the Oklahoma State University in partial fulfillment of the requirements for the Degree of MASTER OF SCIENCE May, 1992



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Thesis Approved:

19 Thesis Adviser

Dean of the Graduate College

#### PREFACE

This study deals the binary vapor-liquid phase equilibrium for carbon monoxide in selected hydrocarbons. The solubilities of carbon monoxide in these hydrocarbons were measured at temperatures from 50°C to 160°C and pressures to 102 bar. Binary interaction parameters in the Soave-Redlich-Kwong equation and the Peng-Robinson equation were regressed from the experimental results for each of the systems studied; these equations provided excellent representation of the data.

I extend my sincere thanks and gratitude to my adviser, Dr. Robert L. Robinson, Jr.. His intelligent guidance, authoritative knowledge and depth of experience have contributed significantly to the completion of this work.

I would like to thank Dr. K. A. M. Gasem for his direction and supervision in all stages of this work.

Thanks are also extended to Dr. Jan Wagner for his help and encouragement during the course of this work. His service as member of my committee is greatly appreciated.

I would also like to express my thanks to Dr. Zhiqiang Zhu, my previous adviser in Zhejiang University, for his teaching and introducing me into the field of thermodynamics and fluid phase equilibrium.

Dr. Naif Abdelaziz Darwish and Mr. Srinivasa Srivatsan

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deserve special thanks for patiently explaining to me the operation of the equipment and the software packages used in this research.

Finally, I would like to acknowledge the financial support received from the United States Department of Energy.

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### NOMENCLATURE

a	energy parameter in cubic equation of state
b	co-volume parameter in cubic equation of state
В	second viral coefficient
С	constants
F	freedom of a system
f	fugacity
G	Gibbs free energy
H <sub>i,j</sub>	Henry's constant for solute i in solvent j
I	number of solute injections
k,1	interaction parameters in equation of state.
n	number of moles
Ν	number of components
c,d,o,p	constants used to correlate the temperature and density dependence of interaction parameters
max	maximum
min	minimum
Р	pressure
q	constants in mixing rule
R	ideal gas constant
Т	temperature
v	molar volume
x	liquid composition
У	vapor composition

compressibility factor or independent variables

Greek Symbols

 $\mathbf{z}$ 

-

Ŷ	activity coefficient
φ	fugacity coefficient
ρ	density
ω	acentric factor
Δ	constant in Vidal's mixing rule
μ	chemical potential
ε	instrumental error
σ	standard variance
a	ratio of parameter a to parameter b in cubic equation of state
Δ	change in property
π	number of phases

## Subscripts

i,j,k	component or data point or solute injection index
b	property at bubble point
с	critical state
cal	calculated value
exp	experimental value
m	property of mixture
t	total property
ω	infinite dilution or infinite pressure

Superscripts

l liquid phase

xi

- s saturated state
- v vapor phase
- standard state
- E excess property
- ^ property of component in a mixture
- partial molar property
- ★ asymmetric property used in solution model

### CHAPTER I

### INTRODUCTION

A knowledge of the phase behavior of carbon monoxide + hydrocarbons is useful in the design and optimization of processes for the conversion of coal to fluid fuel. Limited experimental data on these mixtures have been reported.

The equation-of-state method has proved an effective means to describe the phase behavior of mixtures containing supercritical components at high pressures if an appropriate equation of state and corresponding mixing rule are applied. Thus, the purpose of this work was to determine the solubilities of carbon monoxide in n-hexane, benzene, ndecane, n-octacosane and n-hexatriacontane and to correlate the data using the Peng-Robinson (PR) and the Soave-Redlich-Kwong (SRK) equations of state.

### CHAPTER II

# FUNDAMENTAL PRINCIPLES OF VAPOR-LIQUID EQUILIBRIUM

For a closed equilibrium system composed of a vapor and a liquid phase, the temperature and pressure of both phases are identical. The chemical potential of a given component i in the two phases is also identical; this condition results from applying the first and the second laws of thermodynamics to the system. Therefore the following equations describe the equilibrium of such a system:

$$T^{\mathbf{v}} = T^{1} \tag{1}$$

$$P^{\mathbf{v}} = P^{1} \tag{2}$$

$$\hat{\boldsymbol{\mu}}_{i}^{v} = \hat{\boldsymbol{\mu}}_{i}^{l} \qquad (i=1,2,\ldots,N) \tag{3}$$

The equal chemical potential criterion can be expressed in terms of fugacity which is defined as follows. For a pure component at constant temperature,

$$d\mu_i = RTdlnf_i \tag{4}$$

$$\lim_{p \to 0} \frac{f_i}{p} = 1 \tag{5}$$

For a component in a mixture at constant temperature and composition,

$$d\hat{\mu}_{i} = RTdln\hat{f}_{i} \tag{6}$$

$$\lim_{p \to 0} \frac{\hat{f}_i}{Px} = 1 \tag{7}$$

Equation 1-3 can be reexpressed as:

$$T^{\mathbf{v}} = T^{\mathbf{l}} \tag{8}$$

$$P^{v} = P^{1} \tag{9}$$

$$\hat{f}_{i}^{v} = \hat{f}_{i}^{1}$$
 (*i*=1,2,...,*N*) (10)

Equation 8-10 are the starting equations to solve problems of vapor-liquid equilibrium.

Component fugacity is a function of temperature, pressure and composition. To obtain the relation between fugacity and these measurable variables, two approaches are generally used. The first is the equation-of-state method, as described by the following relation (1):

$$\ln \frac{f_i}{x_i p} = \int_0^\infty \left( \left( \frac{\partial p}{\partial n_i} \right)_{T, v_t, n_{j \neq i}} - \frac{RT}{v} \right) dv - RT \ln Z$$
(11)

To apply this equation, knowledge of the P-v-T-x behavior of the mixture is required.

The second method, which is called the activity coefficient method, is based on the following definition:

$$\hat{f}_i = \gamma_i x_i f_i^0 \tag{12}$$

or

$$\hat{f}_{i} = \gamma_{i}^{\star} x_{i} H_{i,m} \tag{13}$$

Where f° is fugacity of pure specie i at a specified "standard state", e.g, the pure species i at the system temperature and pressure.  $H_{i,m}$  is the Henry's constant for solute i in mixture. The activity coefficient  $\gamma_i$  itself is also a function of temperature, composition and pressure (although the pressure dependence at low and medium pressures is negligible). Classical thermodynamics use another property of mixtures, the excess Gibbs free energy, to express  $\gamma_i$ :

$$\frac{G^{E}}{RT} = \sum_{i} x_{i} \ln \gamma_{i}$$
(14)

Or equivalently,

$$\ln \gamma_{i} = \left(\frac{\partial (n_{t}G^{E}/RT)}{\partial n_{i}}\right)$$
(15)

Unfortunately, classical thermodynamics does not provide either the P-v-T-x behavior of the mixture or an analytical expression for the Gibbs free energy of the mixture. Statistical thermodynamics can give some guidance but still cannot yet give a perfect model for P-v-T-x or Gibbs free energy. Therefore semi-empirical methods are widely used to study both equations of state and Gibbs free energy models.

By using the so-called semi-empirical method, numerous equations of state have been developed, such as the equations of Van der Waals, Viral, BWR, Martin-Hou, RedlichKwong, and more recently SRK and PR equation. Numerous mixing rules (which are required when applying an equation of state to mixtures) have been proposed, such as the classic quadratic mixing rule, Vidal type mixing rules and the density dependent local composition (DDLC) mixing rules (2).

Many Gibbs free energy models have also been established such as the model of Van Laar, Wilson, NRTL, and group contribution models such as ASOG and UNIFAC. An interesting review of the historical development of the modeling of fluid phase equilibrium has been given by Knapp (3).

#### CHAPTER III

#### LITERATURE REVIEW

#### Experimental Data

An extensive literature survey concerning the vaporliquid equilibrium of mixtures containing carbon monoxide was conducted. The sources searched include Chemical Abstracts, a solubility data compilation (4), a vapor-liquid equilibrium data compilation (5) and specialized journals. The available literature data for carbon monoxide + n-hexane, benzene, n-decane, n-octacosane, and n-hexatriacontane are presented in Table I.

#### Correlation Methods

As described in Chapter II, there are two main methods for phase equilibrium calculations, the activity coefficient method and equation-of-state method. Each has advantages and disadvantages, and a clear comparison has been made by Prausnitz et al.(6). In general, the equation-of-state method is preferred for high pressure problems, especially when the system contains supercritical components.

Much of the work regarding the equation-of-state method focuses on the development and modification of mixing rules for equations of state, which are required when applying an

TA	BL	ĿΕ	Ι

SYSTEM	TEMPERATURE RANGE, K	PRESSURE RANGE, atm	REFERENCE NUMBER
$CH_4 + n-C_{10}$	310.9 - 410.9	1.5 - 340	36
	310.9 - 410.9	10 - 87	37
CO + n−C <sub>6</sub>	298.2	1	39
CO + Benze	ne 433.2 - 533.2	7 - 68	40
CO + n-C <sub>10</sub>	283.2 - 313.4	1	39
CO + n-C <sub>28</sub>	373.2 - 573.2	10 - 50	41
CO + n-C <sub>36</sub>	373.2 - 573.2	10 - 50	41

### AVAILABLE LITERATURE DATA FOR SYSTEMS STUDIED IN THIS WORK

equation of state to multicomponent systems. The following section is devoted to a brief review and discussion of mixing rules for cubic equations of state.

The most common and also the most widely used mixing rule is the classic quadratic mixing rule, in which there is a second order dependence of the parameters of upon compositions. It can be written as following (7):

$$a = \sum_{i j} \sum_{i j} x_i x_j a_{ij}$$
(16)

$$b = \sum_{i j} \sum_{i j} x_i x_j b_{ij} \tag{17}$$

$$a_{ij} = (1 - k_{ij}) (a_i a_j)^{\frac{1}{2}}$$
(18)

$$b_{ij} = (1 - l_{ij}) \frac{b_i + b_j}{2} \tag{19}$$

Where  $k_{i,j}$  and  $l_{i,j}$  are binary parameters used to account for deviations from the simple rules suggested by Van der Waals (8). If both  $k_{i,j}$  and  $l_{i,j}$  are set to be zero, the above equations can be reduced to:

$$a = \sum_{i} x_{i} a_{i}^{\frac{1}{2}}$$
 (20)

$$b = \sum_{i} x_{i} b_{i} \tag{21}$$

The quadratic mixing rule is strongly supported by statistical thermodynamics, since there exists a simple relation between the second viral coefficient and the parameters of cubic equations,

$$B = a - \frac{b}{RT}$$
(22)

and the quadratic dependence of B on composition has been derived theoretically from statistical mechanics (9).

Interaction parameters  $k_{i,j}$  and  $l_{i,j}$  are generally considered to be temperature dependent. Some researchers have studied the generalized relations for  $k_{i,j}$ ,  $l_{i,j}$  as function of temperature and characteristic properties such as acentric factor and carbon number of hydrocarbons (10-12).

The quadratic mixing rule has been shown to give good results for systems consisting of components which are similar both in size and chemical nature. But for mixtures of components which are dissimilar in size or in chemical nature, the correlation results are not as satisfactory.

Some investigators have modified the quadratic mixing rule to make it applicable to polar and asymmetric systems. For example, Panagiotopoulos and Reid (13) proposed the following mixing rule:

$$a = \sum_{i j} x_i x_j (a_i a_j)^{\frac{1}{2}} (1 - k_{ij} + (k_{ij} - k_{ji}) x_i)$$
(23)

$$b = \sum_{i} x_{i} b_{i} \tag{24}$$

where  $k_{i,j}$  does not equal  $k_{j,i}$ . By applying this mixing rule along with a new method of estimating pure component parameters in the equation of state from vapor pressure information (14), they achieved improved results for some binary mixtures containing polar compounds such as water + ethanol. The predictions of vapor-liquid equilibrium of carbon dioxide + ethanol + water and carbon dioxide + acetone + water (15) using binary parameters regressed from binary experimental data are quite acceptable. One of its disadvantages is that it does not reduce to the classic quadratic form at low densities, which has been proved theoretically correct. But the authors show that the deviation from quadratic behavior is not significant. Similar work has also been done by Stryjek (16), Adachi (17), and others.

In 1978 Vidal (18) published his new method to obtain mixing rules from solution models. It is done simply by matching the excess Gibbs free energy derived from an equation of state to that obtained from a solution model at infinite pressure. In fact, a similar idea had been used to obtain the first solution model from Van der Walls equation by Van Laar and Van der Waals in 1910 (19). For cubic equations, Vidal's mixing rule for parameter a can be written as follows if a linear mixing rule for b is assumed,

$$\frac{a}{b} = \sum_{i} x_{i} \frac{a_{i}}{b_{i}} - \frac{G_{\infty}^{B}}{\Lambda}$$
(25)

Where  $\pmb{\Lambda}$  is a constant for a given cubic equation. For example,  $\pmb{\Lambda}$  equal ln2 for the SRK equation.

Huron and Vidal (20) applied this new mixing rule, incorporating the one-parameter Redlich-Kister model and a modified NRTL model for  $G^{B}$ , to correlate vapor-liquid equilibrium data of some binary mixtures containing polar components. Remarkable improvements were obtained (even at low pressures) compared to the results given by the quadratic mixing rule. After that, many investigators used this mixing rule for polar and asymmetric systems by using different excess Gibbs free energy models such as the Wilson, NRTL, UNIQUAC models and group contribution models such as UNIFAC and ASOG.

Strictly speaking, Vidal's mixing rule is valid only at infinite pressure. Therefore the parameters in a solution model regressed from low pressure data can not be used directly to get the parameters for an equation of state at either high or low pressures. To overcome this shortcoming, Molleurp (21) suggested a modification of Vidal's method by matching the excess Gibbs free energy given by an equation of state and solution model at zero pressure. Following the proposed procedure of Mollerup, Michelsen (22) obtained the following mixing rule:

$$q_1(\alpha - \sum_i x_i \alpha_i) + q_2(\alpha^2 - \sum_i x_i \alpha_i^2) = \frac{G^E}{RT} + \sum_i x_i \ln\left(\frac{b_i}{b}\right)$$
(26)

Where a and  $a_i$  is defined as following:

$$\alpha = \frac{a}{bRT}$$
(27)

$$\alpha_{i} = \frac{a_{i}}{b_{i}RT} \tag{28}$$

 ${\bf q}_1$  and  ${\bf q}_2$  are constants having values of -0.478 and 0.0047

respectively (for the SRK equation). This mixing rule, along with the modified UNIFIC model (23) has been applied to predict vapor-liquid equilibrium (24) and gas solubilities (25) for some polar and asymmetric systems at both low and high pressures by using UNIFAC parameters regressed from low pressure data; quite satisfactory results were obtained for the mixtures investigated.

Vidal's mixing rule still has the deficiency that it does not reduce to the quadratic form at low density conditions, as faced by Panagiotopoulos's modified quadratic mixing rule. Therefore another type of mixing rule was developed which can avoid this shortcoming. It is called the density dependent local composition (DDLC) mixing rule. Molleurp (26) originally proposed the idea for such consideration, but it should be ascribed to Whiting (27) who developed the first DDLC mixing rule. The advantage of a DDLC mixing rule is that it reduces to the quadratic form at low densities. For example, Panagiotopoulos and Reid (28) proposed the following DDLC mixing rule:

$$a = \sum_{i j} \sum_{i j} x_i x_j (a_i a_j)^{\frac{1}{2}} (1 - k_{ij}) + \frac{b}{vRT} \sum_{i j} \sum_{i j} x_i x_j (x_i \lambda_{ij} + x_j \lambda_{ji})$$
(29)

Obviously, the quadratic dependence on composition is satisfied at low densities. But the calculated results do not improve very much compared to the Vidal's mixing rule. Luedecke (29) and Sandler (30) also proposed different DDLC mixing rules. A review of DDLC mixing rules has been given by Danner (31).

Although Vidal's mixing rule and the DDLC mixing rule generally give better results than the classic quadratic mixing rule for some polar and asymmetric systems, for simpler system such as nitrogen + n-butane the classic quadratic mixing rule proves superior to the Vidal's mixing rules and the DDLC mixing rules according to Shibate (2). Adachi (32) considers that the DDLC mixing rules as unnecessary from a practical point of view. In fact, the quadratic mixing rule is still attractive due to the following advantages:

1. theoretical background.

2. simplicity.

3. flexibility in form for further modification. The last point has been partially demonstrated by Panagiotopoulos and Reid (13,28) and also by De Leeuw (33). In the approach of De Leeuw and et al., the temperature and density effect on the interaction parameters is assumed to have the following forms:

$$k_{ij} = k_{ij}^{0} + c_{ij}T + d_{ij}/v \tag{30}$$

$$l_{ij} = l_{ij}^{0} + o_{ij}T + p_{ij}/v$$
 (31)

Good fitting results are achieved for some nitrogencontaining mixtures at pressures as high as 2,000 bars. Based on considerations described above, the classic quadratic mixing rule is adopted as being adequate for the purposes of the present work.

#### CHAPTER IV

# EXPERIMENTAL APPARATUS AND OPERATING PROCEDURES

#### Experimental Apparatus

The experimental apparatus used in this study was originally set by Gasem (12) in 1984 and has been extensively modified by Bufkin (34), Barrick (35) and Darwish (37) since then. A detailed description of the apparatus is given by Darwish (37). Following is a brief description of the main components of the apparatus, which is illustrated in Figure 1.

A 316ss high pressure tubular reactor (High Pressure Equipment Company, Catalog Number: MS-14) is used as the equilibrium cell (EC). One end of the equilibrium cell is plugged and connected to a drive wheel, which is driven by a 1/50 HP variable speed motor to rock the cell. The another end of the cell is connected to the gas injection valve. Five steel balls are placed inside the cell to promote mixing of the solvent and solute.

Three pumps are used in this apparatus. The first one is a 10 cc positive displacement pump (TEMCO Company, Catalog Number: 10-1-12-H) used mainly for solvent injection and cleaning of the equilibrium cell (P1). The second one is



Figure 1. Schematic Diagram of Experimental Apparatus

a 25 cc positive displacement pump (TEMCO Company, Catalog Number: 25-1-10-HAT) used mainly for solute injection (P2). The third one is a 500 cc positive displacement pump (Ruska Company, Catalog Number: 2210-801) mainly used for cleaning of the solvent storage cell (P3).

Two constant temperature ovens are used in this apparatus. One is used to house the solvent injection pump and the solute injection pump, and the temperature is generally set at 50°C. The another is used to house the equilibrium cell as well as the solvent storage cell, the temperature can be adjusted according to requirement.

Two pressure transducers (Sensotec Company, Catalog Number: TJE/743-03) are used to measure the pressure inside the equilibrium cell and the pressure inside the solute injection pump, respectively. These two traducers are calibrated periodically using a dead weight tester (Ruska Instrument Company, Catalog Number: 2426.1-M288).

A high pressure reactor (High Pressure Equipment Company, Catalog number: OC-1) is used as the solvent storage cell (SC7). Another high pressure reactor (High Pressure Equipment Company, Catalog Number: MS-12) is used as the interstorage cell of mercury (SC5), which is connected to the equilibrium cell.

All fittings, tubings and valves used in this study were supplied by High Pressure Equipment Company. All chemicals used in this study were provided by commercial suppliers. No further purification of the chemicals was

attempted. The suppliers and the claimed purities are presented in Table II.

#### Experimental Procedure

The experimental procedure is quite simple theoretically. A known amount of solvent and a known amount of solute are injected into the equilibrium cell. Then the phase transition from vapor-liquid mixture to liquid mixture is realized by the introduction of mercury and the equilibrium pressure (or bubble point pressure) is determined from a plot of pressure versus volume of mercury injected.

According to the Gibbs phase law, the freedom of a twophase binary system is

$$F = N - \pi + 2 = 2 - 2 + 2 = 2 \tag{32}$$

Therefore at least three variables need to be measured to set up a model describing the phase behavior of such system if the model has adjustable parameters which need to be regressed from experimental data. In the experiment of this work, temperature, liquid composition and equilibrium pressure are measured and only isothermal data is acquired in this work (this means that bubble point pressure is observed as a function of liquid composition at fixed temperature). A brief description of the conceptual steps is given below. A step-by-step operating procedure is also given by Darwish (37). There were no significant modifications during the course of this work.

## CHEMICALS AND PURITIES USED IN THIS WORK

Chemicals	Source	Purity( mol%)
Methane	Big 3 Industries, Inc.	99.97+
Carbon Monoxide	Matheson Gas Products	99.99+
n-Hexane	Aldrich Chemical Company	99+
n-Decane	Aldrich Chemical Company	99+
Benzene	Aldrich Chemical Company	99+
n-Octacosane	Aldrich Chemical Company	98+
n-Hexatriacontane	Aldrich Chemical Company	98+

-

1. The solvent and solute inside the equilibrium cell are emptied from the cell and then the cell is cleaned. Vacuum is applied to the equilibrium cell after cleaning to remove any residual solvent or other volatile components.

2. Known amounts of solvent and solute are injected into the equilibrium cell using the solvent injection pump and the solute injection pump.

3. A specific amount of mercury,  $\Delta v$  (0.01 cc or 0.005 cc) is injected into the equilibrium cell and the equilibrium pressure is recorded. This is done repeatedly and the bubble point pressure is interpreted from the P- $\Delta v$  curve. A typical P- $\Delta v$  curve is shown in Figure 2.

By subsequent solute injections, 2 to 4 data points are measured in each run of a isotherm. Each isotherm 1s measured at least two times to confirm the observations.

At the beginning of the study of a new system, the solvent storage cell is cleaned and vacuumed. Then solvent is transferred into the storage cell and degassed.

As described previously, the operating procedure is quite simple theoretically, but several problems were encountered during this work. One serious problem is the presence of non-condensable gas bubbles in the solvent injection pump, which makes the mercury "compressible" and thus affects the accuracy of bubble point pressure determination. Following is a brief description of the procedures used to remove gas bubbles.

1. Isolate the solvent injection pump and the cleaning



Figure 2. Determination of Bubble Point Pressure

pump by closing  $V_{13}$ ,  $V_1$ ,  $V_7$ ,  $V_6$ ,  $V_{16}$ ,  $V_{24}$ ,  $V_{17}$  as shown in Figure 3.

2. Back off the cleaning pump approximately 100 cc and incline the solvent injection pump to remove the mercury inside it into the cleaning pump.

3. Push all mercury between  $ISC \rightarrow V_5 \rightarrow V_{10}$  into the cleaning pump using helium at about 200 psia.

4. Close  $V_{10}$  and apply vacuum to the cleaning pump about 4 to 8 hours through  $V_{27}$ . Then forward the cleaning pump until the pump is filled with mercury. Close  $V_{27}$ .

5. Apply vacuum to the solvent injection pump for 2 to 4 hours through  $V_{26}$ . Then open  $V_{10}$  and forward the cleaning pump very slowly until the solvent injection pump is filled with mercury till  $V_{26}$ . Close  $V_{26}$ .

6. Test the compressibility of mercury inside the solvent injection pump. A typical result is presented in Table III when the pump is gas-free.

TO THE VACUUM SYSTEM



Figure 3. Removing Non-Condensable Gas Bubbles

TABLE 1		
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### COMPRESSIBILITY TEST OF THE SOLVENT INJECTION PUMP

Pump Position	6.05	6.06	6.07	6.08
Pressure, psia	38	228	900	>1600

.
#### CHAPTER V

#### EXPERIMENTAL RESULTS AND ERROR ANALYSIS

#### Experimental Results

In this work, the solubilities of carbon monoxide in nhexane, n-decane, benzene, n-octacosane and nhexatriacontane were measured at temperatures from 50°C to 160°C and pressures to 1500 psia. The experimental data are presented in Tables IV - VIII.

These experimental data and available literature data appear in Figures 4 - 8, which show  $(P_b-P_s)/x$  as a function of x. The experimental data are plotted in this form since any internal scatter or disagreement with literature data are magnified in such plots. The vapor pressures,  $P_s$  were taken from the literature (42).

As the figures show, the experimental data can be represented simply by a linear expression for each isotherm studied of the form:

$$\frac{p_b - p_s}{x_2} = C_1 + C_2 x_2 \tag{33}$$

Values of  $c_1$  and  $c_2$  and the maximum deviation between correlated and experimental bubble point pressures of each system at each temperature are listed in Table IX. This

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	ΤA	BL	E	IV
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# SOLUBILITY DATA FOR CARBON MONOXIDE + n- HEXANE

Mole Fraction CO	Bubble Point Pres bar (pa	ssure sia)
323.2K (50.0°C,	122.0°F)	
0.0310 0.0549 0.0800 0.1000 0.1346	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	254) 456) 670) 852)
0.1439	86.9 (12 2, 212.0°F)	260)
0.0296 0.0474 0.0807 0.1009 0.1144 0.1408	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	250) 398) 643) 798) 909) 121)
423.2K (150°C,	302°F)	
0.0099 0.0538 0.0799 0.0758 0.1260 0.1466	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	L71) 453) 325) 592) 927) 068)

-

TABLE	V
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# SOLUBILITY DATA FOR CARBON MONOXIDE + BENZENE

Mole	Fraction CO		Bubble bar	Point Pressure (psia)
		323.2K	(50.0°C, 122.0°F)	
0.	.0099		14.7	(213)
0.	0197		29.4	(426)
0.	.0296		44.5	(645)
0.	.0398		60.4	(876)
0.	.0478		73.3	(1063)
0.	0546		84.2	(1221)
0.	.0597		92.6	(1343)
		373.2K	(100.0°C, 212.0°F)	
0.	.0103		14.1	(205)
0.	0150		19.8	(287)
0.	0212		27.1	(394)
0.	0340		42.9	(622)
0.	0371		47.0	(681)
0.	0450		56.8	(824)
0.	0516		65.1	(945)
0.	.0637		80.5	(1167)
		433.2K	(160.0°C, 320.0°F) -	
0.	0148		21.2	(308)
0.	0252		31.3	(445)
0.	0398		45.6	(662)
0.	0428		48.6	(704)
0.	0548		60.5	(877)
Ő.	0590		64.6	(937)

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TABLE V
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## SOLUBILITY DATA FOR CARBON MONOXIDE + n-DECANE

Mole	Fraction CO			Bubble bar	Point Pressure (psia)
		377.6K	(104.4°C,	220.0°F) -	
0 . 0 . 0 . 0 . 0 .	.0388 .0619 .0895 .1197 .1400 .1599			22.2 36.2 53.9 73.7 88.1 102.0	(323) (526) (782) (1070) (1278) (1480)
		310.9K	(37.8°C,	100.0°F)	
0 . 0 . 0 . 0 . 0 .	.0448 .0873 .0953 .1097 .1247 .1399			28.4 58.1 64.5 75.7 87.8 100.0	(412) (842) (936) (1097) (1273) (1451)
		344.3K	(71.1°C,	160.0°F)	
0 . 0 . 0 . 0 . 0 .	0385 0538 0770 0883 1171 1363			23.6 33.0 48.5 56.4 76.9 91.5	(342) (479) (704) (818) (1115) (1327)

TABLE	V	Ι	Ι
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# SOLUBILITY DATA FOR CARBON MONOXIDE + n-OCTACOSANE

Mole Fraction Methane	Bubble bar	Point Pressure (psia)
373.2K (100.0°C,	212.0°F) -	
0.0463	20.1	(291)
0.0577	24.2	(351)
0.1005	44.2	(641)
0.1131	50.2	(728)
0.1463	68.0	(986)
424.2K (150.0°C,	302.0°F)	
0.0551	21.9	(317)
0.1027	41.8	(606)
0.1315	55.9	(811)
0.1623	70.8	(1017)
0.1853	84.1	(1220)

Τź	ABI	$^{1}\mathrm{E}$	V	Ι	Ι	Ι	

# SOLUBILITY DATA FOR CARBON MONOXIDE + n-HEXATRIACONTANE

Mole Fraction Methane	Bubble Poir bar	nt Pressure (psia)
373.2K (100.0°C	2. 212.0°F)	
0.0494	18.0	(261)
0.0638	23.4	(340)
0.0900	34.4	(499)
0.1192	47.5	(689)
0.1603	67.6	(980)
0.2002	89.5	(1299)
423.2K (150.0°	C, 302.0°F)	
0.0568	18.5	(269)
0.0705	23.7	(344)
0.1040	36.0	(522)
0.1289	46.6	(676)
0.1742	66.1	(958)
0.2099	84.0	(1218)

.



Figure 4. Bubble Point Pressure for Carbon Monoxide + n-Hexane



Figure 5. Bubble Point Pressure for Carbon Monoxide + Benzene



Figure 6. Bubble Point Pressure for Carbon Monoxide + n-Decane



Figure 7. Bubble Point Pressures for Carbon Monoxide + n-Octacosane



Figure 8. Bubble Point Pressures for Carbon Monoxide + Hexatriacontane

TABLE I	Χ
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SOLVENT	TEMPERATURE	C <sub>1</sub>	C <sub>2</sub>	Pcal-Pexp  <sub>max</sub>
	°C	(psia)	(psia)	(psia)
$n-C_{f_i}$	50.0 100.0	6426 4343	7769 7137	4
		. 1987	6252	4
$C_{6}H_{6}$	50.0 100.0 160.0	$26894 \\ 11441 \\ 6677$	20827 17186 13760	3 2 6
n-C <sub>10</sub>	37.3 71.1 104.4	12568 9250 8056	8619 8442 7973	6 4 3
n-C <sub>28</sub>	100.0 150.0	7365 6341	5640 5340	6 4
n-C <sub>36</sub>	100.0 150.0	7991 6743	4855 4359	7 6

# EMPIRICAL CORRELATION OF SOLUBILITY DATA

correlation is purely empirical and is required when evaluating the experimental uncertainties of bubble point pressures as described later, which is equivalent to using a second order polynomial to describe the relation between bubble point pressure and solubility:

$$P_b = C_0 + C_1 X_2 + C_2 X_2^2 \tag{34}$$

Obviously  $\textbf{c}_{\emptyset}$  should equal  $\textbf{P}_{g}$  in the above equation.

#### Error Analysis

Two types of errors generally occur in experiments, systematic and random error. Systematic errors are those that have direct relation with the experimental apparatus, operating procedures and other causes that result in consistent deviations of observations from the "true value". Random errors are usually attributed to a combination of disturbances which result in deviations of observations from the "true value". In general, systematic errors affect the accuracy of measurement and random errors affect the precision of measurement.

To eliminate systematic errors in the present work, several efforts were made. One of them was to calibrate the pressure and temperature transducers periodically. The calibration of pressure transducers was done by using a dead weight tester as described in Chapter IV. The calibration of temperature transducers was checked by measuring the icepoint of pure water.

To evaluate the apparatus and procedures used in this

work, initial measurements were made for methane + n-decane and compared with the results given by Reamer (36) and Darwish (37). The experimental data are presented in Table X and the comparison between the results of this work and those of Reamer and Darwish are shown in Figure 9. In the figure, the deviation in solubility predicted by the SRK equation (with the parameters of equation of state regressed from the data of this work) are plotted as a function of bubble point pressure. It shows that the results of this work are in good agreement with those of Darwish (maximum deviation is less than 0.002) and in reasonable agreement with these of Reamer (absolute maximum deviation is less than 0.004).

Random errors can be estimated by error propagation. In general, if y is a variable calculated from a set of measurable variables  $z_i$  (i=1,2,...,N), i.e:

$$y = f(z_1, z_2, \dots, z_N)$$
 (35)

The variance of y can be estimated by the following equation provided that  $z_i$  (i=1,2,...,N) are independent variables (38):

$$\sigma_y^2 - \sum_{i=1}^n \left( \frac{\partial f}{\partial z_i} \right)^2 \sigma_{z_i}^2 + \varepsilon_y^2$$
(36)

Where  $\sigma_z$  is the standard derivation of  $z_i$ ,  $\boldsymbol{e}_y$  is instrumental error in y (if y itself is also a measured variable). In this experiment, both bubble point pressure and solubility are determined. For solubility,

ΤA	BI	LΕ	Х

Mole Fraction Methane		Bubble Point bar	Pressure (psia)
	- 344.3K (71.	1°C, 160.0°F)	
0.0521 0.1152 0.1602 0.2000 0.2325 0.2590	377.6K (10	12.5 28.6 41.5 53.2 64.2 72.8 04.4°C, 220°F)	(181) (415) (602) (772) (931) (1055)
0.0466 0.0801 0.1161 0.1498 0.2000 0.2252 0.2502		12.221.031.441.457.666.174.3	(176) (305) (455) (601) (835) (958) (1078)

# SOLUBILITY DATA FOR METHANE IN n-DECANE



Figure 9. Comparsion of Solubilities of Methane in n-Decane

$$x_1 = \frac{n_1}{n_1 + n_2} \tag{37}$$

Where  $n_1$  and  $n_2$  are the number of moles of solvent and solute injected which are determined by the following relations:

$$n_1 = \rho_1 v_{11}$$
 (38)

$$n_2 = \rho_2 \sum_{j=1}^{T} v_{2j}$$
 (39)

Combining Equations 35-39 gives

$$x_{1} = \frac{\rho_{1} v_{11}}{\rho_{1} v_{11} + \rho_{2} \sum_{j=1}^{I} v_{2j}}$$
(40)

$$x_2 = 1 - x_1$$
 (41)

Applying Equation 36 to the above equation, we obtain

$$\sigma_{x_{2}}^{2} = x_{2}^{2} (1 - x_{2})^{2} \left( \left( \frac{\sigma_{\rho_{1}}}{\rho_{1}} \right)^{2} + \left( \frac{\sigma_{\rho_{2}}}{\rho_{2}} \right)^{2} + \left( \frac{\sigma_{v_{11}}}{v_{11}} \right)^{2} + I \sum_{j} \left( \frac{\sigma_{v_{2j}}}{v_{2j}} \right)^{2} \right)$$
(42)

Where I is the number of solute injections. The maximum uncertainty in solubility for each isotherm is listed in Table XI by using Equation 42 along with the following specifications:

$$\sigma_{\rho_1} / \rho_1 = 0.0015$$
 (43)

$$\sigma_{\rho_2}/\rho_2 = 0.0015$$
 (44)

$$\sigma_{v_{11}} = \sigma_{v_{21}} = 0.005CC \tag{45}$$

The uncertainties in bubble point pressure as a

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## TABLE XI

System	Temperature (K)	x <sub>max</sub>	σ <sub>p,max</sub> (psia)	σ <sub>x,max</sub>
CO + n-C <sub>6</sub>	323.2	0.1439	5	0.0004
	373.2	0.1408	4	0.0004
	423.2	0.1466	4	0.0005
CO + C <sub>6</sub> H <sub>6</sub>	323.2	0.0597	4	0.0003
	373.2	0.0637	5	0.0002
	423.2	0.0548	4	0.0002
$CO + n-C_{10}$	310.9	0.1399	6	0.0005
	344.3	0.1172	6	0.0005
	377.6	0.1400	6	0.0005
$CO + n-C_{28}$	373.2	0.1463	6	0.0007
	423.2	0.1623	5	0.0006
$CO + n-C_{36}$	373.2	0.2002	8	0.0009
	423.2	0.2099	7	0.0008

-

## MAXIMUM UNCERTAINTIES IN SOLUBILITY AND BUBBLE POINT PRESSURE

function of temperature and solubility can be obtained from the following equation:

$$\sigma_{p_b}^2 = \varepsilon_{p_b}^2 + \left(\frac{\partial p_b}{\partial x_2}\right)^2 \sigma_{x_2}^2 + \left(\frac{\partial p_b}{\partial T}\right)^2 \sigma_T^2$$
(46)

The temperature effect is generally neglected since it is quite small compared with the first two terms. The partial derivative of pressure with solubility is obtained directly from Equation 34:

$$\frac{\partial p_b}{\partial x_2} = c_1 + 2c_2 x_2 \tag{47}$$

The instrumental error in pressure is expressed as follows according to Darwish (37):

$$\boldsymbol{\varepsilon}_{p} = 0.004p \tag{48}$$

The maximum uncertainties in bubble point pressures for each isotherm are also presented in Table XI. These estimates are in excellent agreement with the actual experimental results given in Table IX for the error in calculated bubble points.

#### CHAPTER VI

#### CORRELATION OF EXPERIMENTAL DATA

SRK and PR Equation Representations

#### Representation

All experimental data were correlated using the SRK and PR equations. The critical properties of the pure components were taken from the literature (43) for those components for which experimental values are available. For those that lacked experimental critical properties, Gasem's method (12) was adopted to obtained the required properties. All parameters used in the correlation are presented in Table XII. Detailed procedures for data reduction are given by Gasem (12).

Table XIII is a summary of correlation results and the regressed binary parameters are summarized in Table XIV. Detailed correlation results are presented in Tables A.1-A.6 in Appendix. Binary interaction parameters  $k_{ij}$  and  $l_{1j}$ , root mean square errors in solubilities

$$RMSE = \left\{ \frac{1}{N} \sum_{i=1}^{N} (x_{cal} - x_{exp})^2 \right\}^{\frac{1}{2}}$$
(49)

and maximum deviations between experimental and calculated solubilities are listed in these tables. Each table contains

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Component	Pressure bar	Temperature K	Acentric Factor	Reference
CH4	46.60	190.5	0.0110	(43)
со	34.94	132.9	0.0490	(43)
n-C <sub>6</sub>	30.30	507.9	0.2980	(43)
$C_6H_6$	48.94	562.1	0.2120	(43)
$n-C_{10}$	21.08	617.6	0.4900	(43)
n-C <sub>28</sub>	8.26	845.4	1.1073	(12)
n-C <sub>36</sub>	6.91	901.1	1.2847	(12)

-

## CRITICAL PROPERTIES AND ACENTRIC FACTORS USED IN THE SRK AND PR EQUATIONS OF STATE

TABLE XII

# TABLE XIII

# SUMMARY OF CORRELATION RESULTS BY USING SRK EQUATION

Solvent	Temperature	RMSE	in Solubilit	y, Mole	Faction	
	К	Case 1	Case 2	Case 3	Case 4	
n-C <sub>6</sub>	323.2	0.0002	0.0002	0.0019	0.0019	
	373.2	0.0002	0.0003			
	423.2	0.0003	0.0005			
Benzene	323.2	0.0000	0.0002	0.0008	0.0008	
	373.2	0.0001	0.0002			
	433.2	0.0000	0.0001			
n-C <sub>10</sub>	310.9	0.0003	0.0003	0.0014	0.0014	
	344.3	0.0002	0.0002			
	377.6	0.0002	0.0003			
	373.2	0.0002	0.0008	0.0029	0.0029	
20	423.2	0.0008	0.0012			
n-C <sub>36</sub>	373.2	0.0002	0.0025	0.0069	0.0084	
	423.2	0.0005	0.0026			

## TABLE XIV

## SUMMARY OF BINARY PARAMETERS IN THE SRK EQUATION

Solvent	m (17)	С	ase 1	Case 2	Case 3		Case 4
Solvenc	I (K)	k <sub>ij</sub>	l <sub>ij</sub>	k <sub>ij</sub>	k <sub>ij</sub>	l <sub>ij</sub>	k <sub>ij</sub>
n-C <sub>6</sub>	323.2 373.2 423.2	0.0608 0.0297 0.0146	0.0006 0.0115 0.0329	0.0626 0.0633 0.1072	0.0850	-0.0065	0.0657
Benzene	323.2 373.2 423.2	-0.0122 -0.0321 -0.0316	0.0326 0.0306 0.0332	0.0953 0.0700 0.0803	0.1646	-0.0230	0.0884
n-C <sub>10</sub>	310.9 344.3 377.6	0.1356 0.1115 0.0800	-0.0031 -0.0004 0.0043	0.1204 0.1094 0.1006	0.1104	0.0007	0.1136
n-C <sub>20</sub> (*)	323.2 373.2 423.2	0.2772 0.2257 0.2292	-0.0113 -0.0105 -0.0140	0.1602 0.1166 0.0873	0.1439	-0.0018	0.1253
n-C <sub>28</sub>	373.2 423.2	0.1498 0.1413	-0.0071 0.0100	0.0438 0.0031	0.0427	-0.0013	0.0246
n-C <sub>36</sub>	373.2 423.2	0.1937 0.1218	-0.0122 -0.0151	-0.0045 -0.1299	0.3063	-0.0219	-0.0542

\*: Experimental data is given by Srivatsan (44).

results for the following four cases:

1. Two interaction parameters,  $k_{ij}$  and  $l_{ij}$  are determined from the experimental data for each system at each temperature.

2. One interaction parameter,  $k_{1,j}$ , is determined from the experimental data for each system at each temperature  $(1_{j,j}=0)$ .

3. Two interaction parameters,  $k_{ij}$  and  $l_{ij}$  are determined from the experimental data for each system, independent of temperature.

4. One interaction parameter,  $k_{ij}$ , is determined from experimental data for each system, independent of temperature  $(1_{ij}=0)$ .

Table XIII shows that when two temperature dependent binary interaction parameters are used for each isotherm (Case 1), the SRK or PR equation is capable of describing the solubilities with RMSE no more than 0.0008 in mole fraction for each system studied.

When only one temperature dependent interaction parameter  $k_{ij}$  is used for each isotherm (Case 2), nearly the same RMSE is achieved as in Case 1 for carbon monoxide + nhexane, benzene, and n-decane systems. For carbon monoxide + n-octacosane, n-hexatriacontane systems, however, the RMSE obtained are larger than for Case 1. This may be explained briefly from the meaning of the second interaction parameter  $l_{ij}$ , which is used to correct the so called co-volume of the mixture. When  $l_{ij}$  is set to be zero, Equation 17 and Equation 19 reduces to:

$$b = \sum_{i} x_{i} b_{i} \tag{50}$$

For carbon monoxide + n-hexane, benzene, n-decane system, this simplification does not introduce as much error as for carbon monoxide + n-octacosane, n-hexatriacontane systems, since the differences between the solvent and solute in molecular sizes are less than for the later systems.

The RMSE is larger than the experimental uncertainties when lumped  $k_{ij}$  and  $l_{ij}$  (or lumped  $k_{ij}$ ) is used (Case 3 and Case 4) for all systems under study. This suggests that the temperature effect on the interaction parameters should be -considered if better correlation results are required.

Correlation results for Case 1 are also plotted in Figures 10- 14. The comparisons of our results with those of other investigators are also shown in these figures by predicting the solubilities using the binary interaction parameters regressed from our data at each temperature.

Figure 11 shows that for carbon monoxide + benzene system our results are in good agreement with those of Connolly (40) at 160.0°C. The maximum deviation between prediction and measurement is 0.0006. For carbon monoxide + n-octacosane, n-hexatriacontane systems, our solubilities are slightly higher than those reported by Huang et al.(41) at 100.0°C and 150.0°C respectively as indicated by Figure 13 and Figure 14.

There exists correlation between the two interaction



Figure 10. Comparsion of Solubilities of Carbon Monoxide in n-Hexane



Figure 11. Comparsion of Solubilities of Carbon Monoxide in Benzene



Figure 12. Comparsion of Solubilities for Carbon Monoxide in n-Decane



Figure 13. Comparsion of Solubilities of Carbon Monoxide in n-Octacosane



Figure 14. Comparsion of Solubilities for Carbon Monoxide in n-Hexatriacontane

parameters  $k_{ij}$  and  $l_{ij}$  (the correlation coefficients are close to -1 for all isotherms/ systems studied). Although this does not mean that two interaction parameters are not necessary to describe the phase behavior of the systems, it is obvious that for paraffins with carbon numbers less than 10, the improvement in prediction is marginal after the introduction of the second interaction parameter  $l_{ij}$  as clearly indicated by Figure 15. Here the optimized interaction parameter  $k_{ij}$  corresponds to a specified interaction parameter  $l_{ij}$ , and the predicted bubble point pressure in the form of a "reduced" RMSE (the ratio of RMSE in bubble point pressure prediction to that of Case 1) are presented as a function of specified  $l_{ij}$ .

For n-C<sub>6</sub>, the reduced RMSE is not greatly affected by changes in  $l_{ij}$  if  $k_{ij}$  is optimized for a given value of  $l_{ij}$ . For n-C<sub>36</sub>, however, the prediction results become increasingly sensitive to the second interaction parameter  $l_{ij}$  as the carbon number of solvent increases as is also shown in Figure 15. In Figure 16, the effects of  $k_{ij}$  and  $l_{ij}$ on solubility prediction show the same pattern as in bubblepoint pressure predictions presented in Figure 15.

#### Parameter Generalization

To facilitate the interpolation and extrapolation to conditions or systems for which no experimental data exist, parameter generalizations were considered on the basis of the experimental data obtained in this work. Temperature is

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Figure 15. Effect of k;; and l;; on SRK Bubble Point Pressure Prediction



Figure 16. Effects of kij and lijon SRK solubility Prediction

chosen as one of the correlation variables since the temperature dependence is significant, as shown in Table XIII. Another correlation variable chosen is the carbon number of solvent. The acentric factor of the solvent, which is frequently used as a correlation variable in parameter generalization by some investigators (10) (12), is deliberately avoided since the experimental information on critical properties of heavy normal paraffins (and thus acentric factor) is lacking.

The interaction parameters are (simply) expressed as a linear function of both temperature and carbon number of solvent as following:

$$k_{ij} = A_1 + A_2 CN + (A_3 + A_4 CN) T$$
 (51)

$$l_{ii} = B_1 + B_2 CN + (B_3 + B_4 CN) T$$
(52)

It should be mentioned that it is the residual in bubble point pressure between prediction and experimentation (not the obtained interaction parameters) that was used as the criterion for the optimization. Several cases were studied; the detailed descriptions of the cases are listed in Table XV. The summary of the prediction results by different generalization approaches and the optimized parameter constant are presented in Table XVI. In Table XVIII, the calculation results of Case 1 in the previous section are compared with the results corresponding to Case 10 in this section.

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## TABLE XV

# SPECIFIC CASES FOR GENERALIZATION OF INTERACTION PARAMETERS OF SRK EQUATION

Case	e Number	Description
1	$\begin{array}{c} \mathbf{k}_{1j} = 0 \\ 1_{1j} = 0 \end{array}$	
2	$k_{ij}$ =Constant $l_{ij}$ =0	Constant k <sub>ij</sub> is used for application to all binary systems.
3	$k_{ij}$ =Constant $l_{ij}$ =Constant	Constant $k_{ij}$ and $l_{ij}$ are used for application to all binary systems.
4	$k_{ij} = k_{ij} (CN)$ $l_{ij} = 0$	$\mathbf{k}_{ij}$ is correlated as a function of carbon number.
5	$k_{ij} = k_{ij}(CN)$ $l_{ij} = constant$	$\mathbf{k}_{ij}$ is correlated as a function of carbon number and constant $\mathbf{l}_{ij}$ is used for all systems.
6	$k_{ij} = k_{ij} (CN, T)$ $l_{ij} = 0$	${\bf k}_{ij}$ is correlated as a function of both carbon number and temperature.
7	$k_{ij} = k_{ij}(CN,T)$ $l_{ij} = constant$	$\mathbf{k}_{ij}$ is correlated as a function of both carbon number and temperature and constant $\mathbf{l}_{ij}$ is used for all systems.
8	$ \mathbf{k}_{ij} = \mathbf{k}_{ij} (CN) $ $ \mathbf{l}_{ij} = \mathbf{l}_{ij} (CN) $	Both $k_{1j}$ and $l_{1j}$ are correlated as functions of carbon number.
9	$k_{ij}$ =Constant $l_{ij}$ = $l_{ij}$ (CN,T)	Constant $k_{ij}$ is used for all systems and $l_{ij}$ is correlated as a function of temperature and carbon number.
10	$k_{1j} = k_{1j} (CN, T)$ $l_{ij} = l_{ij} (CN, T)$	Both $k_{ij}$ and $l_{ij}$ are correlated as functions of carbon number and temperatures.

.

## TABLE XVI

### SUMMARY OF RESULTS OF PARAMETER GENERALIZATION

Case	1	2	3	4	5	6	7	8	9	10
					PRESS	SURE				
RMSE	8.96	5.62	5.04	4.92	4.68	3.62	3.40	2.30	1.49	1.23
BIAS	-5.79	0.24	-0.27	-0.20	0.31	-0.09	0.31	-0.01	0.19	-0.04
AAD	7.02	4.19	3.62	3.51	3.50	2.84	2.68	1.60	1.16	0.93
AAD%	13.08	8.44	7.15	7.04	7.46	5.60	5.73	3.22	2.46	1.86
NRMSE <sup>a</sup>	7.28	4.57	-4.10	4.00	3.80	2.94	2.76	1.87	1.21	1.00
					CONST	ANTS				
A <sub>1</sub>	0	0.0832	0.1312	0.1343	0.1301	-0.4136	-0.3710	-0.0664	-0.0207	-0.1604
A <sub>2</sub>	0	0	0	-0.0032	-0.0160	0.0474	0.0338	0.0112	0	0.0393
A <sub>3</sub>	0	0	0	0	0	0.0015	0.0013	0	0	3.09E-4
A <sub>4</sub>	0	0	0	0	0	-1.3E-4	-1.2E-4	0	0	-7.90E-5
B <sub>1</sub>	0	0	-0.0061	0	0.0269	0	2.2E-2	0.0300	0.0217	-0.0202
B <sub>2</sub>	0	0	0	0	0	0	0	0.0015	9.42E-4	-5.48E-5
B <sub>3</sub>	0	0	0	0	0	0	0	0	4.27E-5	1.22E-4
B <sub>4</sub>	0	0	0	0	0	0	0	0	-5.32E-6	-3.20E-6

NRMSE<sup>a</sup>=RMSE/RMSE for Case 10
#### TABLE XVII

Т(К)	k <sub>1j</sub>	$k^{g}_{1j}$	l <sub>ij</sub>	l <sup>g</sup> <sub>ij</sub>	¦∆ P ¦ %	¦∆P ¦ <sup>8</sup> %	NP
			CO + n	-C <sub>6</sub>			
323.2 373.2 423.2	$0.0608 \\ 0.0297 \\ 0.0146$	0.0218 0.0135 0.0052	$0.0006 \\ 0.0115 \\ 0.0329$	$0.0128 \\ 0.0179 \\ 0.0231$	$0.3 \\ 0.4 \\ 0.3$	$0.7 \\ 0.9 \\ 2.5$	6 6 6
			CO +n-	$\mathbf{C}_{10}$			
310.9 344.3 377.6	$0.1356 \\ 0.1115 \\ 0.0800$	0.0826 0.0665 0.0505	-0.0031 -0.0004 0.0043	$0.0073 \\ 0.0103 \\ 0.0133$	$0.3 \\ 0.3 \\ 0.3 \\ 0.3$	$0.7 \\ 1.9 \\ 2.5$	6 6 6
			CO +n-	·C <sub>20</sub>			
323.2 373.2 423.2	0.2772 0.2257 0.2292	$0.2193 \\ 0.1503 \\ 0.0867$	-0.0113 -0.0105 -0.0140	-0.0025 0.0004 0.0034	0.2 0.1 0.1	1.4 0.7 1.9	6 6 8
		-	CO + n	-C <sub>28</sub>			
$373.2 \\ 423.2$	0.1498 0.1413	0.2285 0.1333	-0.0071 -0.0100	-0.0096 -0.0079	$0.1 \\ 0.3$	5.5 2.1	5 5
			CO +n-	·C <sub>36</sub>			
$373.2 \\ 423.2$	0.1937 0.1218	0.3067 0.1798	-0.0122 -0.0151	-0.0196 -0.0192	0.2 0.4	2.4 1.6	6 6
					0.2	1.8	88

### COMPARISON OF INTERACTION PARAMETERS AND CORRESPONDING DEVIATIONS

g: The properties corresponding to the generalized parameters.

# Krichevsky-Kasarnovsky Equation Representation

The Krichevsky-Kasarnovsky (K-K) equation was also used to analyze the solubility data of carbon monoxide + neicosane, n-octacosane and n-hexatriacontane to get the Henry's constants of carbon monoxide in these solvents. The procedure for deriving the K-K equation is described below.

From the pressure dependence of component fugacity in a mixture

$$\left(\frac{\partial \ln f}{\partial p}\right)_{T,x} = \frac{\overline{v}_1}{RT}$$
(53)

we can get:

$$\left(\frac{\partial \ln\left(\hat{f}_{i}/x_{i}\right)}{\partial p}\right)_{T,x} = \frac{\overline{v}_{i}}{RT}$$
(54)

Combining the definition of Henry's constant

$$H_{i,s} = \lim_{X_i \to 0} \hat{f}_i / X_i \tag{55}$$

and Equation 54 gives the pressure dependence of Henry's constant:

$$\left(\frac{\partial \ln H_{i,s}}{\partial p}\right)_{T} = \frac{\overline{v_{1}^{\infty}}}{RT}$$
(56)

Integrating the above equation with the assumption that the infinite dilute partial molar volume of the solute does not change with pressure gives:

$$lnH_{i,s}(T,P) = lnH_{i,s}(T,P^{r}) + \frac{\overline{v_{i}^{\infty}}(P-P^{r})}{RT}$$
(57)

Where  $P^{r}$  is a arbitrary reference pressure and  $H_{1,s}(T, P^{r})$  is the Henry's constant at system temperature and reference

pressure.

Assuming that the solute obeys Henry's law:

$$\hat{f}_{i} = H_{i,s}(T, P) x_{i} \tag{58}$$

Choosing the saturation pressure of the solvent as the reference pressure and substituting Equation 58 into Equation 57 gives:

$$\ln \frac{f_{i}}{x_{i}} = \ln H_{i,s}(T, P^{s}) + \frac{\overline{v_{i}^{\infty}}(p - p^{r})}{RT}$$
(59)

This equation is called the K-K equation.

To apply the K-K equation, the fugacity of the solute in the vapor mixture is needed. For carbon monoxide + neicosane, n-octacosane and n-hexatriacontane, no significant errors will result if the fugacity of pure solute at system temperature and pressure is substituted with the assumption that the vapor phase is composed of pure solute. For carbon monoxide + n-hexane, benzene, n-dencane, however, errors introduced by such a simplification are not negligible. Therefore Krichevsky-Kasarnovsky analysis has been applied only for carbon monoxide + n-eicosane, n-octacosane and nhexatriacontane.

Goodwin (45) has given a precise but complicate empirical equation to correlate the properties of carbon monoxide. Figure 17 shows that the difference in fugacity of carbon monoxide given by Goodwin's correlation and the PR equation is negligible. Therefore, the PR equation was used to calculate the fugacity of pure carbon monoxide.



Figure 17. Comparsion of Fugacity of Carbon Monoxide Predicted by the PR Equation and Goodwin's Correlation

The parameters regressed from the experimental data and the RMSE in solubilities from this correlation are presented in Table XVIII. The RMSE in solubility predictions is quite close to the experimental uncertainties and comparable to the prediction results from the equation of state representations (when two temperatere depandent parameters are used), as Table XVIII shows.

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## TABLE XVIII

solvent	т	$H_{i,s}(T,P^s)$	v <sub>i</sub>	RMSE
	(K)	(psia)	cm <sup>3</sup> /gmol	mole frac.
$n-C_{20}$	348.2	7233 (72) <sup>‡</sup>	76 (5)	0.0006
	373.2	6843 (60)	76 (6)	0.0004
	423.2	6252 (32)	94 (3)	0.0003
n-C <sub>28</sub>	373.2	5934 (120)	64 (14)	0.0010
	423.2	5503 (110)	81 (11)	0.0008
n-C <sub>36</sub>	373.2	4999 (18)	102 (2)	0.0004
	423.2	4513 (29)	121 (4)	0.0006

PARAMETERS IN KRICHEVSKY-KASARNOVSKY CORRELATION

\*: standard errors in the tabulated values are given in parentheses.

#### CHAPTER VII

#### CONCLUSIONS AND RECOMMENDATIONS

#### Conclusions

1. The solubilities of carbon monoxide in n-hexane, ndecane, benzene, n-octacosane, n-hexatriacontane have been measured at temperature from 50°C to 160°C and pressures up to 1500 psia. No experimental data have been found for the solubilities of carbon monoxide in n-hexane and n-decane in the literatures searched in this work.

2. The solubilities of carbon monoxide in benzene agree with those of Connolly within 0.0008 mole fraction.

3. The solubilities of carbon monoxide in n-octacosane, n-hexatriacontane from this work are in agreement with literature data within 0.004 mole fraction.

4. Good representations are achieved for all the binary systems investigated using the SRK or PR equation along with the classic quadratic mixing rule if two temperaturedependent binary interaction parameters are used; the differences between representations and experiment are with experimental uncertainties.

5. Good representations are achieved for all the binary systems investigated except carbon monoxide + n-hexatriacontane using the SRK or PR equation along with

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the classic quadratic mixing rule if only one temperature dependent binary interaction parameters is used; the differences between representations and experiment are within experimental uncertainties.

6. The Henry's constants of carbon monoxide in neicosane, n-octacosane, n-hexatriacosane have been obtained by using the Krichevky-Kasarnovsky equation.

#### Recommendations

1. The solvent injection pump should be replaced by one with larger capacity to facilitate studies of lower solubility systems such as carbon monoxide + naphthalene.

2. Further studies are recommended on carbon monoxide + cyclic hydrocarbons (naphthionic and aromatic) that are found in coal to better define the behavior of these systems.

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#### APPENDIX

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# SRK AND PR EQUATION OF STATE REPRESENTATION OF CARBON MONOXIDE + HYDROCARBONS

#### SRK AND PR EQUATION OF STATE REPRESENTATIONS OF SOLUBILITY OF CARBON MONOXIDE IN n-HEXANE

Temperature K (°F)	Soave Par (PR Par <sup>k</sup> ij	rameters rameters) L <sub>1j</sub>	Error in Mole RMS	Fraction* ¦MAX¦
323.2 (122.0)	0.0608 (0.0716)	0.0006 (-0.0001)	0.0002	0.0003
	0.0624 (0.0857)		0.0002	0.0004
373.2 (212.0)	0.0297 (0.0518)	0.0115 (0.0044)	0.0002	0.0004
	0.0633 (0.0633)		0.0003	0.0006
373.2 (302.0)	0.0146 (0.0495)	0.0329 (0.0151)	0.0005	0.0008
	0.1072		0.0003	0.0006
323.2 373.2 423.2	0.0850 (0.0860)	-0.0065 (-0.0059)	0.0019	0.0046
420.2	0.0657 (0.0705)		0.0019	0.0049

SRK	AND	$\mathbf{PR}$	EQUA	TIC	DN	OF	$S_{1}$	ΓΑΤΕ	REP	RES	EN	TATI	DNS
$\mathbf{OF}$	SOLU	JBIL	ITY	OF	CA	RBO	N	MONC	DXID	ΕI	Ν	BENZI	ENE

Temperature K (°F)	(°F) Soave Parameters (°F) (PR Parameters) $k_{1j}$ $l_{ij}$		Error in Mole RMS	Fraction* ¦MAX¦	
323.2 (122.0)	-0.0122 (0.0046)	0.0326 (0.0375)	0.0000	0.0001	
	0.0953 (0.1157)		0.0002	0.0003	
373.2 (212.0)	-0.0321 (0.0060)	0.0306 (0.0280)	0.0001	0.0001	
	0.0700 (0.0886)	-	0.0002	0.0003	
423.2 (320.0)	-0.0316 (0.0125)	0.0332 (0.0251)	0.0000	0.0000	
	0.0803 (0.0886)		0.0001	0.0002	
323.2 373.2 423.2	0.1465 (-0.0998)	-0.0230 (-0.0699)	0.0008	0.0024	
	0.0884 (0.1070)		0.0008	0.0023	

# SRK AND PR EQUATION OF STATE REPRESENTATIONS OF SOLUBILITY OF CARBON MONOXIDE IN n-DECANE

Temperature K (°F)	Soave Pa: (PR Pa: <sup>k</sup> ij	rameters rameters) $1_{ij}$	Error in Mole RMS	Fraction*
310.9 (100.0)	0.1356 (0.1378)	-0.0031 (-0.0023)	0.0003	0.0005
	0.1204 (0.1280)		0.0003	0.0007
344.3 (160.0)	0.1115 (0.1220)	-0.0004 (-0.0019)	0.0002	0.0003
	0.1094 0.1137		0.0002	0.0003
377.6 (220.0)	0.0800	0.0043 (0.0004)	0.0002	0.0004
	0.1006 (0.1008)		0.0003	0.0005
310.9 344.3 377.6	0.1104 (0.0676)	0.0007 (0.0118)	0.0014	0.0032
511.0	0.1136 (0.1184)		0.0014	0.0032

Temperature K (°F)	Soave Pa (PR Pa <sup>k</sup> 1]	rameters rameters) l <sub>lj</sub>	Error ın Mole RMS	Fraction* ¦MAX¦
373.2(212.0)	0.1498 (0.1701)	-0.0071 (-0.0084)	0.0002	0.0004
	0.0438 (0.0699)		0.0008	0.0018
423.2(302.0)	0.1413 (0.1711)	0.0100	0.0008	0.0011
	0.0031 (0.0239)		0.0012	0.0019
373.2 423.2	0.0427 (-0.0142)	-0.0013 (0.0051)	0.0028	0.0050
	0.0246 (0.0448)		0.0030	0.0051

#### SRK AND PR EQUATION OF STATE REPRESENTATIONS OF SOLUBILITY OF CARBON MONOXIDE IN n-OCTACOSANE

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Soave Parameters Error in Mole Fraction\* Temperature K (°F) (PR Parameters) k<sub>ij</sub> l<sub>ij</sub> RMS MAX 323.2 (212.0) 0.1937 -0.0122 0.0001 0.0003 (0.2162) (-0.0134)-0.0045 0.0025 0.0029 (0.0333)-----373.2 (302.0) 0.1218 -0.0151 0.0005 0.0008 (0.1593) (-0.0175)-0.12990.0026 0.0035 (-0.0788)\_\_\_\_\_ 323.2 0.3063 -0.0219 0.0069 0.0120 373.2 (0.2149) (-0.0175)-0.05420.0084 0.0150 (-0.0128)

SRK AND PR EQUATION OF STATE REPRESENTATIONS OF SOLUBILITY OF CARBON MONOXIDE IN n-HEXATRIACONTANE

SRK	AND	$\mathbf{PR}$	EQUATIO	N O	F	STATE	REF	PRESENTATIONS
	OF	SOL	UBILITY	OF	ME	THANE	IN	n-DECANE

Temperature K (°F)	Soave Par (PR Par <sup>K</sup> ij	rameters rameters) l <sub>ij</sub>	Error in Mole RMS	Fraction* ¦MAX¦
344.3 (160.0)	0.0590 (0.0575)	-0.0123 (-0.0115)	0.0005	0.0007
	0.0285 (0.0308)		0.0015	0.0025
377.6 (220.0)	0.0529 (0.0516)	-0.0103 (-0.0107)	0.0005	0.0007
	0.0273 (0.0268)		0.0011	0.0018
344.3 377.6	0.0564 (0.0548)	-0.0114 (-0.0111)	0.0005	0.0011
	0.0280 (0.0291)		0.0016	0.0023

\* The RMS and maximum errors in CH4 mole fraction are essentially the same for both the SRK and PR equations of state.

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#### XIAOHUA YI

#### Candidate for the Degree of

Master of Science

#### Thesis: BINARY VAPOR-LIQUID PHASE EQUILIBRIUM FOR CARBON MONOXIDE + HYDROCARBONS

Major Field: Chemical Engineering

Biographical:

- Personal Data: Born in Shimen, Hunan, China, February 14, 1964, the son of Ziwen Yi and Tintin Chen. Married to Yanfang Liu on July 20, 1990.
- Education: Graduated from the First High School, Shimen, Hunan, China, in August 1980; received Bachelor of Engineering Degree in Chemical Engineering from Tinjing University, Tinjing, China, in August 1984. Completed requirements for the Master of Science in May, 1992.
- Professional Experience: Research Assistant, School of Chemical Engineering, Oklahoma State University, January, 1991 to May 1992.