# MULTIPLE SCATTERING SUPPRESSION FOR 

CROSS-CORRELATION OF A FLOWING

## FLUID TO DETERMINE

PARTICLE SIZE
By
RYAN MATTHEW CAMBERN
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Thesis Approved:


Wayne B. Powell
Dean of the Graduate College

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## NOMENCLATURE

| a | radius of a particle (nm) |
| :---: | :---: |
| A | cross-sectional area ( $\mathrm{m}^{2}$ ) |
| $\mathrm{A}_{\mathrm{m}}$ | amount of multiple scattering |
| 2b | height of test cell (mm) |
| $\mathrm{B}_{\mathrm{m}}$ | amount of single scattering |
| 2c | width of test cell (mm) |
| d | 'waist' diameter of the laser beam ( $\mu \mathrm{m}$ ) |
| D | incident laser beam diameter (mm) |
| $\mathrm{D}_{\mathrm{h}}$ | hydraulic diameter of the flowing fluid test cell (mm or m) |
| $\mathrm{D}_{0}$ | diffusion constant of particles ( $\mathrm{m}^{2} / \mathrm{s}$ ) |
| E | magnitude of the electric field ( $\mathrm{N} / \mathrm{C}$ ) |
| $\mathrm{E}_{\text {s }}$ | magnitude of the scattered electric field (N/C) |
| E* | complex conjugate of the electric field (N/C) |
| f | focal length of the lens (mm) |
| $\mathrm{g}^{1}$ | normalized electric field correlation function |
| $\mathrm{G}^{1}$ | electric field correlation function (V/m) |
| $\mathrm{g}^{2}$ | normalized intensity correlation function |
| i | complex number $\sqrt{-1}$ |
| $\mathrm{I}_{\text {s }}$ | scattered intensity (kHz) |
| j | count variable used in summations |
| $\mathrm{k}_{\text {B }}$ | Boltzmann constant ( $1.380658 \times 10^{-23} \mathrm{~J} / \mathrm{K}$ ) |
| $\overline{\mathrm{k}}_{\mathrm{i}}$ | incident beam wave vector ( $\mathrm{m}^{-1}$ ) |
| $\overrightarrow{\mathbf{k}}_{\text {s }}$ | scattered wave vector ( $\mathrm{m}^{-1}$ ) |
| l* | effective transport mean free path in DWS theory ( $\mu \mathrm{m}$ ) |
| L | length a particle travels through the incident laser beam |
| $\mathrm{L}_{\text {e }}$ | length necessary to obtain fully developed flow (mm) |
| n | index of refraction |
| np | number of particles with which the viewed particle interacts |
| N | number of particles in the detection volume |
| P | wetted perimeter (m) |
| $\hat{p}$ | pressure (Pa) |
| $\frac{\mathrm{q}}{\mathrm{q}}$ | magnitude of the scattering wave vector $\left(\mathrm{m}^{-1}\right)$ scattering wave vector $\left(\mathrm{m}^{-1}\right)$ |
| Q | flow rate ( $\mathrm{ml} / \mathrm{min}$ ) |

r
argon ion

Reynolds number ( sec ) ( sec )
magnitude of a spatial position vector (m)
spatial position vector (m) correlation time (msec)
absolute temperature (K)
constant used for a two cumulant fit
fluid velocity vector ( $\mathrm{mm} / \mathrm{s}$ )
fluid velocity ( $\mathrm{mm} / \mathrm{s}$ )
fluid velocity in x -direction ( $\mathrm{m} / \mathrm{s}$ )
constant used for a two cumulant fit
denotes translation of the top fiber mount in the x -direction denotes translation of the top fiber mount in the $y$-direction

## Greek

angle through which the laser and detector may travel (degrees) detection width (mm)
angle used in calculation of $\theta$ for the flowing fluid setup (degrees) scattering volume created by a focused beam in a sample ( $\mathrm{cm}^{3}$ )
signal-to-noise ratio
rotation angle of flowing test cell (degrees)
sample volume ( $\mathrm{cm}^{3}$ )
amplitude of the incident wave vector ( V )
viscosity of the sample solvent ( Pa s )
scattering angle (degrees)
experimental proportionality constant
wavelength of the incident laser beam ( nm )
dynamic viscosity ( Pa s )
numerical constant $=3.14159265$
fluid density ( $\mathrm{kg} / \mathrm{m}^{3}$ ) correlation delay time (s)
time it takes a particle to move one unit length under Brownian motion
time between fluctuations of two particles moving at different velocities
time it takes a particle to move through the focused laser beam (sec) tilt angle (degrees)

## Abbreviations

denotes back micrometer for top fiber mount
correlation transfer theory
diffusing-wave spectroscopy

FM denotes front micrometer for top fiber mount
GRIN graded-index lenses
HCN hydrogen-cyanide
$\mathrm{He}-\mathrm{Ne}$
LDV
MSS
Nd: YAG
PMMA
PMT
PSL
helium neon
laser doppler velocimetry
multiple scattering suppression
neodymium-yttrium-aluminum-garnet
polymethylmethacrylate
photomultiplier tube
RT polystyrene latex

VIDE
radiative transfer theory
visible infrared double extinction

## Subscripts

| A, B | denoted different wavelength lasers and detectors |
| :--- | :--- |
| L | laser arm |
| D | detector arm |

## Superscripts

1 denotes electric field correlation function
2 denotes intensity field correlation function

## CHAPTER I

## INTRODUCTION

### 1.1 Background

In a wide variety of industrial applications, a reliable and accurate method of determining small particle characteristics in suspensions is essential. Companies such as those that design and produce filtering systems are looking for more accurate and faster methods to measure particle size and density. These measurements will enable those companies to test the efficiency of their products. There are four main types of particle testing procedures performed today, off-line testing, on-line testing, in-line testing, and in-situ testing [Dorri-Nowkoorani (1995)].

The first three of these procedures require alteration of the original particle suspension to successfully make measurements. Alterations may include either diluting the sample for in-line testing or static sampling (removing a sample from the flow) for both off- and on-line testing. These alterations have raised questions as to the accuracy of the results as well as their representation of actual flow field behavior. Of these procedures, in-situ testing is the least intrusive method used to measure particle properties
in a section of the process. Because of the nature of this procedure, many companies are interested in developing it into an accurate method of testing.

The use of in-situ measuring techniques many times employs the theory of dynamic light scattering (DLS) or photon correlation spectroscopy (PCS). The use of DLS to determine particle characteristics is based on single scattering data. Single scattering occurs within dilute solutions. In single scattering, photons are scattered only once from a colloidal particle prior to detection. When single scattering occurs, an autocorrelation of the light can be used to determine particle characteristics [Berne and Pecora (1976)]. However, single scattering rarely occurs in industrial processes. Multiple scattering is much more common and occurs when dealing with turbid fluids such as oil and gasoline. Photons that have been scattered by a particle in the solution once and are scattered by at least one more particle before being detected are said to be multiply scattered. Multiple scattering affects the time dependence of the autocorrelation function [Lock (1997a)], which can lead to DLS measurements that are inaccurate and difficult to process.

Several techniques have been developed over the last two decades to deal with particle suspensions that are neither single scattering nor completely opaque. Ideally, these techniques look to eliminate the effects of multiple scattering, but this has proven to be a difficult task that requires expensive equipment. One of these techniques incorporates a two-color system that is complex to align, can handle only moderately dense solutions, and still can not be relied upon to deliver high quality single scattering data [Lock (1997a)]. Recently several techniques have been developed that employ various methods of cross-correlation to suppress the effects of multiple scattering instead
of trying to eliminate those effects. In 1997, Meyer et al. developed a different approach to the problem of multiple scattering suppression. The method uses a single focused beam and two slightly separated detectors to suppress multiple scattering effects by crosscorrelation of the single scattering signals. This method was then to be proven highly effective on static fluid samples by Nobbmann et al. [1997].

### 1.2 Objective

The research discussed in this thesis has two goals. The first goal is to continue the work done by Nobbmann et al. (1997) on static fluid samples. The second and most important goal is to prove that the theory of multiple scattering suppression could be adapted to flowing fluid solutions.

To help in the accomplishment of these goals, a foundation of theory will first be presented briefly discussing dynamic light scattering and multiple scattering suppression (MSS). The major equations and some discussion of these theories are given, while a list of published resources will be used to fill in any gaps and to further explain the derivation of those theories. A detailed explanation of the flow suppression theory will then follow.

With a foundation of theory laid, Chapters IV and V will discuss the experimental setups used in implementing those theories. Chapter IV will discuss the static fluid setup, the procedure used to make samples, the computer correlation package, and the procedure used to collect data. Included in the procedure will be a detailed description of the alignment procedure used for each experiment. The experiments included work done to
determine the effects of different scattering angles, different concentrations of particles and different particle sizes on the accuracy of the data collected.

Chapter V deals with applying the method used in the static fluid setup to a setup where the particle suspension is flowing. As with the static setup, the physical setup and procedure will be discussed. The first flowing fluid experiments dealt with single scattering samples to show how flow rate, particle sizes, and sample alignment affected the accuracy of the data collected. The next experiments performed dealt with the scattering angle, particle concentration, and flow rates. The results of these experiments as well as the results from Chapter IV will be discussed in Chapter VI. The remaining sections of this thesis contain a literature review of work in related areas and some recommendations on how to further advance this project.

## CHAPTER II

## LITERATURE REVIEW

### 2.1 Introduction

The research covered by this thesis is a combination of two relatively new theories. The first theory deals with multiple scattering suppression for a dynamic light scattering analysis of a solution in order to accurately characterize particles, specifically determining their size. The second theory deals with suppression of flow effects in a shearing system to accurately characterize particles. The following literature review is divided into three sections. Section 2.2 deals with the many aspects of light scattering experimentation. The publications presented in this section will give a time line perspective of how light scattering experiments have evolved into the setup used in the research described by this thesis. These publications are not directly related to the current research described within this thesis, but do give the reader background into light scattering as well as knowledge of other techniques used to analyze light scattering data. Section 2.3 of the literature review will discuss the experiments that directly relate to the research described in this thesis, namely single-beam, two-detector multiple scattering
suppression. Publications describing a technique used to suppress flow effects in light scattering experiments are discussed in section 2.4 of the literature review.

### 2.2 Light Scattering Techniques

Studies involving light scattering by small particles have been performed for over a century. Experiments by Tyndall, Rayleigh, Gans, Debye, Mie and others laid the foundation for the light scattering experiments of today. With the invention of the laser, the field of light scattering experimentation has seen great improvements in the accuracy of results as well as an increase in experimental and industrial applications [Berne and Pecora (1976)]. The following review will give some insight into the recent developments in light scattering research. The review is not a complete listing of all of the work done in the field of light scattering, but it does highlight the work that lead to the research covered in this thesis.

George Phillies (1981a) gave a theoretical discussion on quasielastic light scattering spectroscopy. He described how a two-beam, two-detector cross-correlation setup could be more effective in the suppression of multiple scattering effects than a single-beam, single-detector autocorrelation setup. The discussion began with the singlebeam, single detector setup and the intensity equations for one single scattering and one double scattering event as well as the time dependent versions of these equations. These same equations were then developed for the two-beam, two-detector setup. A comparison of the single scattering equations given for both of these setups revealed that they are equivalent as long as the scattering angle remains $90^{\circ}$, while the double
scattering equations are not. Double scattering effects were smaller by a factor of $(\mathrm{np} / \mathrm{N})^{2}$ in the two-beam, two-detector setup than in the single-beam, single-detector setup, where np is the number of particles with which a viewed particle interacts, N is the number of particles in the detection volume, and $\mathrm{np} \ll \mathrm{N}$. Phillies also noted that several conditions must be set for successful data collection. Those conditions are as follows:

1. The wavelength of each laser used must be equal to within one-tenth of the wavelength inside the sample. This means the beams must be aligned within 0.5 mrad of each other.
2. Focusing of any lenses used must be exact, to maintain the planar aspect of the incident and scattered waves.
3. Light scattering by this geometry would be seen by both detectors. At a scattering angle of $90^{\circ}$, this effect was irrelevant; but at other scattering angles, this effect would be cumulative and would make data analysis more difficult.

Phillies presented experimental proof of the theory previously described in a second publication (1981b). The experimental setup consisted of a $25 \mathrm{~mW} \mathrm{He}-\mathrm{Ne}$ laser, an attenuator, a $50 \%$ reflective plate beamsplitter, two focusing lens, two mirrors, three pairs of irises, two RCA 7265 photomultiplier tubes, and a 64 channel Langley-Ford digital correlator. The single laser beam was split into two equal beams that were routed into opposing sides of a $1-\mathrm{cm}^{2}$ fluorimeter cell with polished sides. The photomultiplier tubes and irises were then positioned across from each other on the other two sides of the test cell. As described in the first publication, the major problem with this experimental setup was exact alignment of all components. A majority of the publication dealt with the alignment procedure for the split beams and the two photomultipliers. Alignment also dealt with the test cell itself. The test cell was tilted less than $0.5^{\circ}$ to eliminate backscattering into the photomultipliers from the opposing walls. Experimental results
were presented for two scattering media: $0.091 \mu \mathrm{~m}$ polystyrene latex spheres (PSL) suspended in $0.4 \mathrm{~g} / \mathrm{l}$ sodium lauryal sulphate at concentrations of $2.1 \times 10^{-5}$ to $1.18 \times 10^{-2}$ percent by volume and bovine serum albumin in 0.15 M NaCl at a concentration of 20 percent by weight. The data showed that the use of a single-detector setup would overestimate the diffusion coefficient of $0.091 \mu \mathrm{~m}$ PSL particles in concentrations greater than $10^{-3}$, while the two-detector setup would be concentration independent. Finally, it was reported that the experiments on serum albumin matched the results of earlier work done by Hall et al. (1980).

Like Phillies (1981a,b), Dhont and de Kruif (1983) published work done on the theory followed by a second publication [Mos et al. (1986)] that described the experimental work done on that theory. In 1983, Dhont and de Kruif described how Phillies had overlooked a second order term in dealing with the theory of double scattering. They further describe how Phillies used an incomplete form of the Siegert relation in his calculations. This oversight is found to be insignificant, however, since the use of a two-laser beam, two-detector setup causes the double scattering effects to be negligible in the measured intensity correlation function. The paper also describes the difficulties associated with collecting and analyzing data from a scattering angle other than $90^{\circ}$. To help eliminate these difficulties, it was suggested that two different wavelength lasers and two different absorbing filters be used. It was also shown that the two detectors should be placed at scattering angles related to each other and to the wavelength of the two lasers used. This relation was given by the following equation:

$$
\begin{equation*}
\frac{\sin \left(\frac{\theta_{\mathrm{A}}}{2}\right)}{\sin \left(\frac{\theta_{\mathrm{B}}}{2}\right)}=\frac{\lambda_{\mathrm{A}}}{\lambda_{\mathrm{B}}}, \tag{2-1}
\end{equation*}
$$

where the subscripts $A$ and $B$ represent two different lasers, $\theta_{A}$ and $\theta_{B}$ are the scattering angles for each detector, and $\lambda_{\mathrm{A}}$ and $\lambda_{\mathrm{B}}$ are the wavelengths for each laser.

In 1986, Mos et al. (including Dhont and de Kruif) described the experimental setup that they used to prove their theory for elimination of multiple scattering effects. The experimental setup described in this paper resembles the one previously described in the review of Phillies' (1981) experiments. Also like the work done by Phillies, this paper describes in detail the alignment procedure involved in collecting reliable data and the difficulties associated with aligning the setup. Experiments were performed on two colloidal systems. The first sample was polystyrene latex spheres (diameter equal to 176 nm ) of differing concentrations in distilled water. The second sample was silica spheres dispersed in toluene or xylene. The silica spheres had radii reported in the paper to be $36.7 \pm 2 \mathrm{~nm}$. Experimental results showed that the diffusion coefficient could be reliably found with the cross-correlation method for various concentrations, while autocorrelation results were not reliable at higher concentrations.

The need to determine particle characteristics is not limited to just liquid media. Gougeon et al. (1987) described a technique called visible infrared double extinction (VIDE) used to characterize a cloud of coal particles (diameters ranged from 20 to 80 $\mu \mathrm{m})$. These clouds had a high optical thickness, but were still in a range considered to be weakly multiple scattering. Two lasers of different wavelengths were used in the experiments: the visible light source was a $\mathrm{He}-\mathrm{Ne}$ laser ( 632.8 nm wavelength) and the infrared light source was a Hydrogen-Cyanide (HCN) laser (337 nm wavelength). A cylinder specially designed for these experiments rotated and vibrated to produce the coal cloud. A photomultiplier tube collected the visible light from the sample, while a moll
thermopile collected the infrared light. Outputs from both detectors were then sent to a two-channel recorder. Experimental results overestimated the diameter of the coal particles compared to the value of the sieve diameter used, and the difference was attributed to the use of non-spherical coal particles in the experiments. They concluded that the technique would still be applicable to diagnose dense laden flows.

Pine et al. (1988) described the use of a new technique in quasielastic light scattering. The new technique was called diffusing-wave spectroscopy (DWS), and it was meant to extend the use of quasielastic light scattering to densely multiple scattering media. Their theory dealt with backscattering and transmission autocorrelation functions for a diffusive media. For transmission the only parameter remaining as a variable during experiments was the transport mean free path, $l^{*}$. Experiments used a laser with a wavelength of 488 nm on an optically thick media of $0.497 \mu \mathrm{~m}$ polystyrene latex particles and were found to match their model if $l^{*}$ was set at $1.43 \mu \mathrm{~m}$. For comparison with the transmission model, the backscattering model used the same 1*. For an l* value of 1.43 $\mu \mathrm{m}$, the experimental data from the backscattering model fit well when a second variable parameter, $\kappa$, was set to 2.0 . Further experiments were performed on other particle sizes and fit to the backscattering model if k was varied by $\pm 10 \%$. Finally, experiments were conducted on mixtures of two different interacting and non-interacting particle sizes ( $0.312 \mu \mathrm{~m}$ and $0.497 \mu \mathrm{~m}$ spheres). It was found that the correlation function for the interacting particles decayed slower in both transmission and backscattering.

Wiese and Horn (1991) described the use of a fiber optic as a different approach to quasielastic light scattering detection. Their experimental setup consisted of a $\mathrm{He}-\mathrm{Ne}$ laser ( 632.8 nm wavelength), a single mode fiber optic, a coupler that allowed the fiber
optic to illuminate the sample and simultaneously collect the backscattering signal, a photomultiplier, and a correlator. They performed experiments to compute the autocorrelation functions of polymer latex particles (diameters ranging from 41 to 326 nm ) at various concentrations (ranging from 1 to 40 percent by weight). The experimental results show that multiple scattering effects have a negligible impact on the autocorrelation function even at the higher concentrations. The claim that negligible multiple scattering effects were seen at such high concentrations was intriguing.

Ackerson et al. (1992) and Dorri-Nowkoorani et al. (1993) and (1994) discussed an alternative to the multiple scattering suppression method for particle characterization. This method utilized the Correlation Transfer Theory (CT) to determine particle size in media ranging from single to highly multiple scattering. The main focus of these papers was that the CT equation and the Radiative Transfer (RT) equation were similar. Therefore, RT solution techniques could be applied to the CT equation in order to obtain a solution.

Segrè et al. (1995) described the use of a two-color dynamic light scattering system for multiple scattering suppression. A detailed description of the theory behind this two-color system is provided in their paper. The system consisted of two lasers of differing wavelengths ( 488 nm and 514.5 nm ) and two photomultipliers with filters designed to detect only one color. A German commercial company called ALV provided this setup. The lasers were generally operated at 100 mW and experiments were performed over a range of scattering angles from $20^{\circ}$ to $140^{\circ}$. Experiments were conducted on two suspensions: polystyrene (PS) spheres ( 56 nm radius) in water and polymethylmethacrylate (PMMA) particles (185 nm radius) in dodecane. The
concentrations of the two suspensions tested ranged from 0.05 to $1.00 \%$ volume fraction for polystyrene and from 0.01 to $0.25 \%$ volume fraction for PMMA. Results showed that the two-color system worked effectively to characterize particles in multiple scattering suspensions. Preliminary work using the two-color technique described above can also be seen in Schätzel et al. (1990).

In 1998, Aberle et al. discussed the use of a 3-D cross-connotation method to suppress multiple scattering effects in turbid media. The experimental setup utilized a single laser beam that had been split into two parallel beams. Those two beams were then directed by a lens to cross in the center of the sample cell. This was done to produce two equal scattering vectors that were detected by two monomode fibers. Like the setups described by Phillies (1981a,b) and Mos et al. (1986), the alignment of the laser beams and the detection fibers were critical to the collection of reliable data. Experiments were performed over a range of scattering angles from $10^{\circ}$ to $135^{\circ}$ on $69,107 \pm 10.5,236 \pm$ $6.58,453 \pm 9.0$, and $481 \pm 1.8 \mathrm{~nm}$ diameter latex spheres in deionized water. The concentrations of the samples were based on optical transmission and ranged from 0.7 to 99.3 \%. Experimental results showed that the technique was excellent at suppressing multiple scattering effects for a wide range of particle sizes, a wide range of concentrations and was easy to use over a wide range of scattering angles. Experimental results also showed that this technique was useful in both static and dynamic light scattering applications.

### 2.3 Multiple Scattering Suppression: Single Beam - Two Detector Setups

Previous to 1997, the majority of work done in the area of multiple scattering had involved two beam, two detector setups for suppression. These setups were effective in suppressing the multiple scattering effects only in semidilute samples. A major complication in dealing with these two-beam, two-detector systems was alignment. Due to this complication, the desire to find a simpler scattering geometry that would allow multiple scattering suppression even in very turbid solutions has been growing. This desire has lead to the design of a system that utilizes a single laser beam and two closely configured detectors to suppress multiple scattering effects. The following articles describe the development and use of this system.

Meyer et al. (1997a) developed an experimental setup incorporating a single laser beam ( 514.5 nm wavelength) and two fiber optic cables with polished ends. The setup was capable of running experiments with the fibers separated at three different vertical distances of $0.25,0.50$, and 0.75 mm . Experiments were performed at scattering angles of $60^{\circ}, 90^{\circ}, 120^{\circ}$, and $135^{\circ}$ on 0.107 and $0.204 \mu \mathrm{~m}$ diameter particles of polystyrene latex spheres for concentrations ranging from 0.0017 to 5.0 percent by weight. It was found in their experiments that decay time had little dependence on the three separation lengths listed above. Therefore results were given for only a fiber separation of $0.25 \mathrm{~mm}(1.5$ mrad) and involved mostly $0.107 \mu \mathrm{~m}$ particles. Normalized autocorrelation and crosscorrelation functions were plotted for various concentrations against delay time so the effect of concentration on each could be seen. As the concentrations increased, the
linearity (a measure of data reliability) of the autocorrelation functions became markedly less while the cross-correlation functions remained steady. They also gave results showing how the calculation of diameter was correct using cross-correlation data for all concentrations while calculations using the autocorrelation data dropped by more than an order of magnitude. A second publication by Meyer et al. describing the same experiment is also listed in the references of this thesis under (1997b).

Theoretical work done by Lock (1997a) accompanied the experimental work done by Meyer et al. (1997a,b) on multiple scattering suppression by a single beam/two detector setup. Lock limited the calculations given in his paper to single and double scattering only, but noted that his work could be applied to higher orders of multiple scattering. Lock first gave a mathematical description of the experimental scattering geometry. Then in terms of the electric field equations, he described single and double scattering by non-interacting spherical particles in a liquid solution. Next, he derived the electric-field correlation functions for a single beam cross-correlation system. Lastly, the intensity cross-correlation function was derived, the level of double scattering suppression was given and the time dependence of the autocorrelation and crosscorrelation functions was derived. Lock concluded with the fact that a focused laser beam produced smaller far-zone coherence in the direction transverse to the beam for multiple scattering than that produced by single scattering. Like the experimental papers that Meyer et al. (1997a,b) published, two papers dealing with this subject were published in 1997 by Lock. Both are listed in the references of this thesis. The difference between Lock's two papers is that the calculations in the (1997a) paper ignore polarization and
angular dependence of scattered light, while the (1997b) paper includes those effects in the derivation of the autocorrelation and cross-correlation functions.

The final publication builds on the work covered by the first four publications cited in this section. In 1997, Nobbmann et al. furthered the experimental work done by Meyer et al. (1997a,b) by developing a setup that allowed examination of how the multiple scattering contribution to the cross-correlation function was reduced with increasing fiber separation. Nobbmann et al. utilized the same basic setup as the one described by Meyer et al., with one major difference. The detector housing in Nobbmann's setup used a beam-splitter to divide the scattered beam equally into two single-mode fibers that use GRIN lenses for better detection. This design allowed the two fibers to be physically separated by $90^{\circ}$ when they are effectively separated by only milliradians. The separation was accomplished by the use of a micrometer mounted to the top fiber holder. Experimental results were given for polystyrene latex spheres with a diameter of $0.107 \mu \mathrm{~m}$ at volume fractions of $0.15 \%$ and $0.25 \%$. These experiments were all performed at a scattering angle of $90^{\circ}$. Nobbmann et al. were successful in showing the ability of their setup to sweep through a measured amount of separation between fibers and to correctly give the particle diameter when that separation became great enough ( $\sim 1 \mathrm{mrad}$ ).

### 2.4 Flow Effect Suppression

The majority of published articles involving light scattering and flowing fluids deal with laser Doppler velocimetry applications. Laser Doppler velocimetry experiments
estimate the velocities of particles within flowing systems. Contrary to LDV experiments, flow effects must be suppressed for experiments where the ability to characterized particles within the flowing fluid is desired. The following review describes two articles where flow effects are suppressed.

Ackerson and Clark (1981) examined the possibility of using dynamic light scattering as a way to study the intensity correlation function for dense solutions of Brownian particles subjected to low shear rates. Both model calculations and experimental results were given and compared. An equation was given for the correlation time-dependent signal-to-background-noise factor. This term was directly related to the measurement of the intensity correlation function. An experimental setup was described which minimized flow effects on the signal-to-background-noise factor, but no experimental data was found using this setup. The experimental setup consisted of an arrangement where the scattered wave vector (the difference of incident and collected light) was perpendicular to the flow velocity vector. Instead of that arrangement, two other experimental setups for flow suppression were described and experiments were performed using those setups. Results from those experiments showed the modest effects of varying stop and pinhole diameters to suppress velocity effect.

In 1998 Hoppenbrouwers and van de Water used a different experimental setup to study dynamic light scattering in fluids subjected to shear flows. Their experimental setup incorporated two beams from either a helium-neon ( 632.8 nm wavelength) or an Argon-Ion ( 514.5 nm wavelength) laser that crossed within the flow volume of a Couette device. The scattered signal was collected by using a 1.0 mm diameter pinhole, a lens with a focal length of 100 mm , a multimode fiber optic cable, a photomultiplier, and a
real time digital correlator. They claimed that this setup was capable of uncoupling the effects of particle convection due to shear and particle diffusion due to Brownian effects within the correlation function. Doing this would allow the diffusion coefficient to be measured so that particle characterization could be done. They were successful in showing that they could characterize the diameter of $0.204 \mu \mathrm{~m}$ polystyrene spheres in a water-picoline mixture to within $15 \%$ of the known value.

## CHAPTER III

## THEORY

### 3.1 Introduction

The primary goal of this thesis is to show experimentally how two independent suppression techniques can be combined to produce usable data from a dense suspension of particles in a flowing setup. To better understand how these experimental suppression techniques work, one must have some general knowledge of the basic light scattering equations, the theory behind multiple scattering suppression, and the concepts that lead to the suppression of flow effects. The purpose of this chapter is merely to give the reader an overview and not a detailed description of the ideas described above. More in-depth coverage can be seen in the cited references and in Sundaresan (1999).

### 3.2 Basic Dynamic Light Scattering

A light scattering experiment consists of an incident beam, $\overline{\mathrm{k}}_{\mathrm{i}}$, entering a sample and being scattered in a new direction towards a detection device. The scattered wave vector, $\overrightarrow{\mathrm{k}}_{\mathrm{s}}$, defines this new direction. The magnitude of the two wave vectors is given by

$$
\begin{equation*}
\mathrm{k}_{\mathrm{i}, \mathrm{~s}}=\left|\stackrel{\rightharpoonup}{\mathrm{k}}_{\mathrm{i}, \mathrm{~s}}\right|=\frac{2 \mathrm{n} \pi}{\lambda_{\mathrm{i}, \mathrm{~s}}} \tag{3-1}
\end{equation*}
$$

where n is the index of refraction of the sample, $\lambda_{i, s}$ is the wavelength of the incident and the scattered beams, respectively. The angle between the line of transmission and the position of the detection device is the scattering angle, $\theta$. The scattering wave vector, $\overrightarrow{\mathrm{q}}$, is the difference between the incident wave vector and the scattered wave vector. Assuming elastic scattering and using the law of sines, the scattering wave vector is defined by the following equation and can be seen in Fig. 1 [Berne and Pecora (1976)]:

$$
\begin{equation*}
\overline{\mathrm{q}}^{2}=\left|\overrightarrow{\mathrm{k}}_{\mathrm{i}}-\overrightarrow{\mathrm{k}}_{\mathrm{s}}\right|^{2} \Rightarrow|\overrightarrow{\mathrm{q}}|=\frac{4 \pi \mathrm{n}}{\lambda_{\mathrm{i}}} \sin \frac{\theta}{2} \tag{3-2}
\end{equation*}
$$



Figure 1: The incident wave vector, $\overrightarrow{\mathrm{k}}_{\mathrm{i}}$, the scattered wave vector, $\overrightarrow{\mathrm{k}}_{\mathrm{s}}$, and the scattering angle, $\theta$, used to find the scattering wave vector, $\overrightarrow{\mathrm{q}}$.

Dynamic light scattering studies the fluctuations of the scattered intensity in a sample over time. For static fluid conditions, the fluctuations are due to Brownian motion of the particles. The motion of the particles changes the scattered electric field, which effects the intensity given by [Berne and Pecora (1976)]

$$
\begin{equation*}
I_{s}(t) \propto\left|E_{s}(t)\right|^{2} \tag{3-3}
\end{equation*}
$$

A measure of the normalized correlation function for the scattered intensity signal is defined by [Nobbmann (1991), (1997)]

$$
\begin{equation*}
\mathrm{g}^{2}(\stackrel{\rightharpoonup}{\mathrm{r}}, \overline{\mathrm{k}}, \tau)=\frac{\langle\mathrm{I}(\stackrel{\rightharpoonup}{\mathrm{r}}, \stackrel{\rightharpoonup}{\mathrm{k}}, \mathrm{t}) \mathrm{I}(\stackrel{\rightharpoonup}{\mathrm{r}}, \overrightarrow{\mathrm{k}}, \mathrm{t}+\tau)\rangle}{\langle\mathrm{I}(\stackrel{\mathrm{r}}{\mathrm{r}}, \overline{\mathrm{k}}, \mathrm{t})\rangle\langle\mathrm{I}(\stackrel{\mathrm{r}}{\mathrm{r}}, \overrightarrow{\mathrm{k}}, \mathrm{t})\rangle} \tag{3-4}
\end{equation*}
$$

where $\tau$ is the delay time and $<>$ is an ensemble average. A second way to define the scattered signal is with the normalized electric field correlation function [Nobbmann (1991), (1997)]

$$
\begin{equation*}
\mathrm{g}^{1}(\stackrel{\rightharpoonup}{r}, \overline{\mathrm{k}}, \tau)=\frac{\left\langle\mathrm{E}(\stackrel{\rightharpoonup}{\mathrm{r}}, \overrightarrow{\mathrm{k}}, \mathrm{t}) \mathrm{E}^{\bullet}(\overline{\mathrm{r}}, \stackrel{\rightharpoonup}{\mathrm{k}}, \mathrm{t}+\tau)\right\rangle}{\langle\mathrm{E}(\stackrel{\rightharpoonup}{\mathrm{r}}, \overline{\mathrm{k}}, \mathrm{t})\rangle^{2}} \tag{3-5}
\end{equation*}
$$

where $\mathrm{E}^{*}$ is the complex conjugate of the electric field. If Gaussian statistics dominate in the scattered light, the $g^{2}(\vec{r}, \vec{k}, \tau)$ function and the $g^{1}(\stackrel{\rightharpoonup}{r}, \vec{k}, \tau)$ function are related by the Siegert relation as follows:

$$
\begin{equation*}
g^{2}(\stackrel{\rightharpoonup}{\mathrm{r}}, \stackrel{\rightharpoonup}{\mathrm{k}}, \tau)=1+\gamma^{2} \mid \mathrm{g}^{1}(\stackrel{\rightharpoonup}{\mathrm{r}}, \stackrel{\mathrm{k}}{\mathrm{k}}, \tau)^{2} \tag{3-6}
\end{equation*}
$$

where $\gamma^{2}$ is the experimental signal-to-noise ratio.
For DLS experiments involving dilute monodisperse samples, the electric field correlation function is found to be [Nobbmann (1991), (1997)]

$$
\begin{equation*}
g^{1}(\stackrel{\rightharpoonup}{\mathrm{q}}, \tau)=\exp \left(-\mathrm{D}_{\mathrm{o}} \mathrm{q}^{2} \tau\right) \tag{3-7}
\end{equation*}
$$

where q is defined by Eq. (3-2), $\mathrm{D}_{\mathrm{o}}=\frac{\mathrm{k}_{\mathrm{B}} \mathrm{T}}{6 \pi \eta \mathrm{a}}$ (the Einstein diffusion constant), $\mathrm{k}_{\mathrm{B}}$ is the Boltzmann constant, $T$ is the absolute temperature of the sample, $\eta$ is the viscosity of the solvent, and a is the radius of the particles.

### 3.3 Multiple Scattering Suppression

The purpose of this thesis is to show how a dynamic light scattering technique can be applied to a dense suspension of flowing particles to determine size. When dealing with such a system, the first problem encountered is multiple scattering effects. These effects make interpretation of the collected data difficult. The second problem is flow effects, which can add unwanted fluctuations to the light scattering data. This section deals with the theory for a multiple scattering suppression technique. Section 3.4 will discuss the theory on which the technique for flow effect suppression is based.

The concept of using two single-mode fiber optic detectors and a single incident laser beam to suppress multiple scattering was developed by Meyer et al. (1997a,b). Theory matching the results from Meyer's experiments for this technique was also presented by Lock (1997a,b). For the system described by Meyer et al., the separation between the two single-mode detectors was not variable during an experiment Nobbmann et al. (1997) described a system that allowed the tilt angle between the two detectors to be adjusted so that a sweep of tilt angles could be mapped (shown clearly in a future figure). That same system was utilized in the research for this thesis.

This system allowed the detectors to remain focused on the same detection area while being separated by small angles (in mrad). According to the van Cittert-Zernike theorem for large angular separations, single scattering will produce larger intensity fluctuations from a laser beam focussed within a multiple scattering sample than the entire multiple scattering from the rest of the sample [Mandel and Wolf (1995)]. This
means that as the scattering volume is decreased, the single scattering detection area increases. Proof of this concept can be seen in Meyer et al. (1997a,b) and Nobbmann et al. (1997).

It is important to note that the area of useful separation between the two detectors was defined by the following two conditions. First, the cross-correlation of singly scattered light was not possible for too large of a separation between the detectors. Second, if the separation was not sufficient, multiple scattering effects dominated. For a cross-correlation function, mapping the signal-to-noise ratio versus the tilt angle showed all three areas. As mentioned above in Eq. (3-6), $\gamma^{2}$ was the signal-to-noise ratio. Through the use of simple approximations and by using the assumptions listed below, the square root of the signal-to-noise ratio was derived by [Nobbmann et al. (1997)] as a function of tilt angle between detectors to be:

$$
\begin{equation*}
Y(\phi)=A_{m} \frac{\exp \left[\frac{-q^{2} \phi^{2}}{8 \alpha_{m}}\right]}{\left[\alpha_{m}^{2}\left(\alpha_{m} \phi^{2} 4^{-1}+\delta_{m}\right)\right]^{1 / 2}}+B_{m} \frac{\exp \left[\frac{-q^{2} \phi^{2}}{4 \beta_{m}}\right]}{\beta_{m} \sin (\theta) \sqrt{\alpha_{m}}} \tag{3-8}
\end{equation*}
$$

where the ratio of $A_{m}$ to $B_{m}$ represents the ratio of multiple-to-single scattering, and the following assumptions were made during the derivation of this equation. The first set of assumptions made was that a small scattering volume created by a focused beam, an intermediate detection width, and a large sample volume, $\beta_{m} » \alpha_{m}>\delta_{m}$ respectively, were present in the system. Also assumed was that the tilt angle, $\phi$ (in mrad), was much smaller than the scattering angle, $\theta$ (in rad) $\phi \ll \theta$. Therefore, the $\sin (\phi) \cong \phi$ and the cos $(\phi) \cong 1$. Finally, it was assumed that $\beta_{m} \sin \left(\theta^{2}\right) » \delta_{m}$. Equation (3-8) is used in

Sundaresan (1999) to illustrate theoretical fits to data taken from three different concentrations of 0.107 mm PSL particles in a rectangular cell at a $90^{\circ}$ scattering angle.

As mentioned above on page 22, three areas can be seen in a signal-to-noise ratio versus tilt angle mapping (see Fig. 2). The first area is the 'peak' of the graph and corresponds to the tilt angle for which multiple scattering dominates. No separation between the two detectors gives the maximum of the 'peak' and corresponds to the highest signal-to-noise ratio. Both detectors are focused on the same detection area from within the multiple scattering threshold. The multiple scattering threshold is defined as a spherical volume around the incident beam in the sample where multiple scattering dominates. Data collected from this region generally undersize the diameters of the particles in the suspension. As the tilt angle increases, the signal-to-noise ratio decreases. At approximately $\pm 1.5 \mathrm{mrad}$, the 'peak' region turns into the 'shoulder' region for the mapping of Fig. 2. This region corresponds to a region where the bulk of the multiple scattering is suppressed. Evaluation of data collected from within this region generally give correct particle sizes. The final region begins where the signal-to-noise ratio becomes very low. In this region, the separation between the two detectors becomes too large and their signals are not correlated well. For Fig. 2, this region starts at approximately $\pm 5 \mathrm{mrad}$ and continues for the remainder of the mapping.

The experiments described in this thesis utilized a correlation software package to analyze the data collected for both the static fluid and flowing fluid setups (see Appendix I). Particle sizes were determined by combining Eqs. (3-6) and (3-7) with a two-cumulant expansion as follows [Nobbmann et al. (1997)]:

$$
\begin{equation*}
g^{2}(\tau)=1+\gamma^{2} \exp \left(-2 u \tau+2 w \tau^{2}\right) \tag{3-9}
\end{equation*}
$$

where $u$ and $w$ were two constants used for the cumulant fit. The first cumulant $u$ was used to determine the diameter of the particles within the suspension and was defined as

$$
\begin{equation*}
\mathrm{u}=\mathrm{D}_{\mathrm{o}} \mathrm{q}^{2} \tag{3-10}
\end{equation*}
$$

where $D_{0}$ is defined after Eq. (3-7) and $q$ is given in Eq. (3-2). The second cumulant was normalized by $w / u^{2}$ and was used to determine the exponentiality of the correlation function. For single scattering suspensions, the normalized second cumulant should ideally vanish, but is generally found to be less than 0.04 . For multiple scattering samples, a typical normalized second cumulant less than 0.10 indicates suppression of multiple scattering effect.


Figure 2: Signal-to-noise ratio versus tilt angle for a cross-correlation measurement taken at a scattering angle of $\theta=90^{\circ}$. A sample of $0.107 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.1330 percent by weight was used.

### 3.4 Flow Effect Suppression

As mentioned in Section 3.2, the purpose of this thesis is to show how two separate suppression techniques can be used together to determine particle size in a dense flowing suspension of particles. Section 3.3 gave a brief view of the theory behind multiple scattering suppression. The focus of this section will be to describe the theory behind flow effect suppression in systems where flow is present.

Dynamic light scattering experiments focus on fluctuating signals from particles undergoing Brownian motion. As previously stated in Eq. (3-3), this motion affects the electric field correlation, which in turn affects the intensity correlation. A well-defined solution exists for static fluid setups that utilize the fluctuations caused by Brownian motion to determine particle size. In a flowing system, a second source of fluctuations occurs due to the motion of the particles flowing through the scattering volume. The following derivation shows how flow effects influence the electric field correlation function. The unnormalized version of the electric correlation function is

$$
\begin{equation*}
\mathrm{G}^{1}(\tau)=\left\langle\mathrm{E}(\overline{\mathrm{q}}, \tau) \mathrm{E}^{*}(\overline{\mathrm{q}}, 0)\right\rangle \propto \mathrm{N}\left|\varepsilon_{\mathrm{oi}}\right|^{2}\left\langle\mathrm{e}^{\mathrm{iq}(\overline{\mathrm{q}}(\mathrm{f}(\tau)-\overline{\mathrm{F}}(0))}\right\rangle \tag{3-11}
\end{equation*}
$$

where N is the number of particles in the scattering volume, $\varepsilon_{0 i}$ is the amplitude of the incident wave, $\bar{r}$ is the spatial coordinate of a particle being correlated, and $<>$ is an ensemble average. Brownian motion and flow change the spatial coordinate of a viewed particle with time. Therefore, the following equation is used to define the motion of the particle through the scattering volume:

$$
\begin{equation*}
\overline{\mathrm{r}}(\tau)=\delta \overline{\mathrm{r}}(\tau)_{\mathrm{BM}}+\overline{\mathrm{v}}(\mathrm{r}) \cdot \tau+\overline{\mathrm{r}}(0) \tag{3-12}
\end{equation*}
$$

where $\delta \overline{\mathrm{r}}(\tau)_{\mathrm{BM}}$ is the change in the position of the particle due to Brownian motion, $\overline{\mathrm{v}}(\mathrm{r}) \cdot \tau$ is the change in position of the particle caused by the fluid velocity and $\mathrm{r}(0)$ is the initial position of the particle. Substitution of Eq. (3-12) into Eq. (3-11) gives

$$
\begin{equation*}
\mathrm{G}^{1}(\tau) \propto \mathrm{N}\left|\varepsilon_{\text {oi }}\right|^{2}\left\langle\mathrm{e}^{\mathrm{i} \bar{q} \bullet \delta \overline{\mathrm{~F}}(\tau))_{\mathrm{BM}}+(\mathrm{i} \bar{q} \bullet \stackrel{\mathrm{v}}{ }(\mathrm{r}))_{\tau}}\right\rangle \tag{3-13}
\end{equation*}
$$

As with Eq. (3-7), if single scattering of monodisperse particles dominates the system, then the normalized version of Eq. (3-13) is found to be:

$$
\begin{equation*}
g^{1}(\tau)=\exp \left[-D_{0} q^{2} \tau+[i \bar{q} \bullet \bar{v}(r)] \tau\right] \tag{3-14}
\end{equation*}
$$

where normalization occurred by dividing Eq. (3-13) by $N|\notin o i|^{2}$.
Equation (3-13) shows how the influence of flow interferes with the analysis of DLS experiments. To accurately size particles in a suspension with a mean flowing velocity, the influence of flow effects must be suppressed. The technique employed here to accomplish this involved the dot product seen in Eq. (3-14). To suppress flow effects, a setup was designed that allowed the scattering wave vector, $\overline{\mathrm{q}}$, and the flow vector, $\overline{\mathrm{v}}$, to be perpendicular. Discussion of this setup is given in Chapter $V$ of this thesis and in Sundaresan (1999). This geometric setup produced a dot product that was zero and therefore theoretically eliminated the flow effects. A similar setup was explained in Ackerson and Clark (1981) for a system where no mean flow was present.


Figure 3: A top view of the geometry that facilitates flow suppression by allowing the scattering wave vector, $\mathbf{q}$, and the fluid velocity vector, $\mathbf{v}$, to be perpendicular. The incident wave vector, $\mathbf{k}_{\mathbf{i}}$, and the scattered wave vector, $\mathbf{k}_{\mathbf{s}}$, are also shown.

## CHAPTER IV

## STATIC FLUID CONDITIONS: EXPERIMENTAL SETUP, ALIGNMENT, AND PROCEDURE

### 4.1 Introduction

The primary goal of this thesis is to demonstrate the feasibility of applying two unrelated suppression techniques together for particle sizing. Those two techniques were multiple scattering suppression and flow effect suppression. The technique described within this chapter involved a setup used for multiple scattering suppression in a nonflowing sample. Although the focus of this thesis deals with a flowing system, a good understanding of multiple scattering suppression was needed. As a means of accomplishing this, experiments were first performed on an existing static fluid setup. This not only allowed for means with which to gain experience with the equipment and data taking process, but also allowed the work previously done by Nobbmann et al. (1997) to be expanded. This chapter describes the static fluid suppression setup, the alignment procedure for that setup, the computer program used and its settings, a description of the samples, and the procedure used to collect data with the setup.

### 4.2 Experimental Setup

A schematic of the experimental setup used for particle sizing in static fluid samples can be seen in Fig. 4. The goniometer was broken up into three main sections: the light source arm, the sample stand, and the detection device arm. The first component necessary to conduct a dynamic light scattering experiment is the light source. In the static fluid experiments, a Helium-Neon ( $\mathrm{He}-\mathrm{Ne}$ ) laser ( 632.8 nm wavelength) was used to provide this light source and was attached to the stationary laser arm of the goniometer. The 20 mW laser was vertically polarized and had an original beam diameter of 0.68 mm . To ensure proper alignment of the laser, a holder capable of small vertical adjustments mounted it to the laser arm.

Also mounted on the laser arm were an adjustable attenuator and a lens holder. The approximate position of these components can be seen in Fig. 4. The attenuator allowed the amount of power coming from the laser to be controlled over a range from 0.08 to 9.54 mW . The range of attenuation was found by direct measurement using a power meter with a wand attachment (see Appendix I). Mounted on the laser arm between the attenuator and the sample stand was a lens holder. The lens holder was designed to allow small adjustments in all planes ( $x, y$, and $z$ ). This was valuable during the alignment process, which will be discussed in Section 4.3. A 12.7 mm diameter lens with a focal length of 37 mm was used throughout the experiments. Sundaresan's thesis (1999) contains the heights for and distances between all components used in both the non-flowing and flowing setups.


Figure 4: The experimental setup for static fluid samples.


Figure 5: Static fluid setup used for fiber multiple scattering suppression. Shown are the scattering angle, $\theta$, and the tilt angle, $\phi$.

The next essential part for a dynamic light scattering experiment is the sample (Fig. 4). For the static fluid setup, a sample stand was centered along the goniometer's axis of rotation. The stand consisted of a milled aluminum platform sized to hold a water bath. A rod attached the platform to the base of the goniometer allowing only vertical adjustments. A water bath was used in the static fluid setup to accomplish three goals. First, the water bath allowed greater freedom in aligning the laser in the sample. The larger diameter of the water bath meant that slight adjustment had less effect on how the incident beam moved in the sample. Second, the water bath helped to maintain the sample at a constant temperature. Third, the water bath container provided the mechanism used to hold the test tube vertically along the goniometer's axis of rotation. The container was formed in the shape of a small beaker from a section of glass tubing (see Appendix I). To enable proper vertical alignment of the test tube in the water bath, a lid and base were constructed from Teflon tubing. Each piece contained a test tube sized hole (approximately 10 mm in diameter) through the center. The base was designed to sit in the bottom of the container and was used to ensure that the test tube was held vertically. Fisher Scientific test tubes were used throughout the static fluid experiments (see Appendix I). The sample stand, water bath and its parts were all designed to ensure that the samples were held centered along the goniometer's axis of rotation.

The final segment necessary to complete the dynamic light scattering setup is the detection device. The detector housing (Fig. 6) was manufactured from aluminum and was mounted to a plate that in turn could be mounted to the detection device arm of the goniometer (Fig. 4). The detection device consisted of a beamsplitter and two singlemode fiber optic cables (see Appendix I); each connected to a photomultiplier tube.


Figure 6: Side view of the detector housing showing where the fiber mounts attach and the location of the beamsplitter.

The detector housing allowed for the mounting of the beamsplitter and the devices used to adjust the fiber optic cables. It was necessary to use a beamsplitter due to the physical constraints of the fiber optic mounts.

As discussed in the theory, each fiber optic detector must look at the same point within the sample, with an angular separation of only a few milliradians. This was physically impossible without the use of the beamsplitter. The nonpolarizing wavelength specific beamsplitter (see Appendix I) divided the scattered signal to each of the fiber optic detectors equally. One fiber optic detector saw direct transmission of the scattered signal and was called channel 0 , while the other fiber optic detector, channel 1 , was
mounted $90^{\circ}$ from the line of transmission on top of the detector housing (Figs. 5 and 6). Each fiber optic detector was mounted on a platform capable of small angle tilts as well as two direction translations. The mount for the top fiber was constructed out of two translation stages (see Appendix I) labeled X and Y in Figs. 7 and 8 rotated $90^{\circ}$ for translation in both the x - and y -directions. Those two stages were then mounted on an aluminum platform that was capable of tilting through small angles. Two different views of the fiber mount using micrometers are given in Figs. 7 and 8 with approximate dimensions given in Appendix I.

Micrometers were used to tilt and translate the top fiber so that its position could be recorded. Two micrometers were used for tilting the fiber. One micrometer was mounted at a front corner of the tilt plate, while a brass ball was placed under the other front corner to act as a pivot. The second micrometer was mounted at the rear corner of the tilt plate diagonally opposite to the front micrometer and along the same side as the pivot. This setup was mounted to the top of the detector housing by two compression springs and allowed controlled tilting of the top fiber. Two additional micrometers were used to control the translation of the fiber, one in the $x$-direction the other in the $y$ direction.

The back fiber mount was similar in its design, but did not utilize micrometers. Instead of micrometers, a series of set screws and springs were used to facilitate the back fiber's movement. Figures 9 and 10 display representations of the mount using set screws. The approximate dimensions of the back fiber mount are given in Appendix I. The lack of micrometers on the back fiber mount was acceptable because, once it was aligned, it remained fixed until realignment was needed for the next experiment. As


Figure 7: Front view of the top fiber mount. The X and Y denote the direction each translation stage moves, while BM and FM denote the back and front tilt micrometers, respectively.


Figure 8: Top view of the top fiber mount. BM, FM, X, and Y are repeated from Fig. 7.


Figure 9: Back view of the rear fiber mount.


Figure 10: Top view of the rear fiber mount.
described in Section 4.3, the alignment of the setup involved several components including the test cell. Changing the test tube altered the alignment each time. Therefore, the system was realigned after each experiment.

To conduct experiments at several scattering angles, the detector arm railing was mobile and could travel throughout a range of scattering angles from $0^{\circ}$ to $130^{\circ}$. No experiments were performed at angles below $30^{\circ}$ due to high intensity levels that could damage the photomultiplier tubes. However, a scattering angle of $0^{\circ}$ was necessary to properly align the detectors (see Section 4.3). As discussed in the theory, multiple scattering effects must be suppressed to correctly size particles. Although the majority of the suppression came from tilting the top fiber, some additional suppression was accomplished by the use of a polarizer. Single scattering of light is polarization preserving, while multiple scattered light generally does not maintain polarization. Therefore a polarizer that was oriented to transmit only vertically polarized light, i.e., polarized light from the laser, should suppress the multiply scattered signal left after the suppression caused by separating the detectors. The polarizer was mounted on the detection arm of the goniometer between the sample and the detector housing (Fig. 4).

The final pieces of equipment associated with the detection segment of the goniometer were the photomultiplier tubes (Fig. 4). The photomultiplier tubes were connected to the detector housing by single-mode fiber optic cables. Each fiber optic cable carried a scattered signal to one of the photomultiplier tubes. The photomultiplier tubes required two power supplies (see Appendix I). The power supplies were maintained at a constant 12 volts and 5 amps throughout the experiments. Each
photomultiplier converted the incoming signal into usable electronic pulses. These pulses were then read and processed by a commercial digital correlator card (ALV-5000).

### 4.3 Alignment

Like the work done by Phillies (1981), Mos et al. (1986), and the other authors cited in the literature review, the success of the experiments described in this thesis depended on the alignment of the setup. Through a trial and error process of data collecting, a defined alignment procedure was developed. Some observations made during this process will be discussed in Chapter VI. The alignment procedure began with the light source itself, the laser. As mentioned previously, the laser was attached to a holder capable of small vertical adjustments. The holder consisted of an aluminum plate mounted on top of another plate. The top plate contained four screws, one at each corner. The screws allowed the laser to be raised uniformly or tilted slightly. A level was placed on the laser and the screws were manipulated until the laser was leveled at the proper height. As a second check in the alignment process, a reference point was then placed on the wall at the spot where the laser contacted it. This spot marked the level of the laser prior to the start of the experiments. The dot was also valuable during the alignment of other components on the goniometer. The alignment of the laser was done once, but was checked prior to realignment for each new experiment and prior to any needed adjustments. Realignment of the laser was usually not necessary, but occasionally was needed due to relaxation of the set screws used to control its positioning.

With the laser aligned, the next step was to align each fiber optic cable. For the alignment of the fiber optic cables, all components were removed from the setup. Alignment of the cables was accomplished by first moving the detector arm to $0^{\circ}$, so that the laser and detector housing were directly opposite of each other. This allowed for each fiber optic cable to see the incident light source directly through the use of the nonpolarizing beamsplitter. The beamsplitter (see Appendix I) was designed to deliver $50 \pm$ $3 \%$ of the incident beam's power both to the fiber seeing direct transmission and to the fiber seeing reflection (see Fig. 6). Each fiber optic cable was then manipulated by use of the micrometers or the screws (Figs. 7-10) until a beam of light was visible from its unattached end. Upon detection of a visible beam, the next step was to maximize the intensity of the beam exiting each fiber by moving the micrometers for channel 1 and by moving the set screws for channel 0 . This was accomplished with the use of the previously mentioned power meter setup (see Appendix I). Holders, mounted on the optical table, were used to fix the position of both the wand of the power meter and the free end of the fiber optic cables during this step of the alignment procedure in order to ensure the consistency of the intensity measurements. With the intensity in each fiber maximized, the next step in the alignment process was to add all components to the goniometer.

The first components to be added to the goniometer were the water bath and the test tube. The water bath alone was added to the sample stand (Fig. 4) and positioned so that the beam passed through its center. Proper positioning was determined by both a visible check of the beam in the water bath as well as by checking for the strongest light from the fibers. The test tube was the next piece to be added. As described in the
previous section, the water bath was equipped with a lid and base designed to vertically hold the test sample along the axis of rotation of the goniometer. If needed, the bath and sample were then repositioned by hand to allow the highest intensity light possible to reach the fibers.

The attenuator and polarizer were the next components added to the goniometer. The polarizer was positioned on a stand between the detector housing and the sample (Fig. 4). Alignment of the polarizer was determined by which position allowed for the brightest intensity light from the fiber optics. The attenuator was mounted to the laser arm in front of the laser beam. Since the purpose of the attenuator was to decrease the power of the laser beam, a different alignment criterion was needed. Proper positioning of the attenuator was determined when the back reflection of the laser was directed back onto the edge of the laser opening. Because of possible damage to the laser, direct back reflection into the laser was avoided. The use of the back reflection showed that the attenuator was as perpendicular as possible without introducing laser instabilities. This alignment was desired to limit any stray reflections (caused by misalignment) from being sent to the detection fibers.

The final component positioned onto the goniometer was the lens. The purpose of the lens was to focus the incident beam within the sample. As described in the theory, a tightly focused incident beam was valuable, because it was necessary to reduce the multiple scattering area within the sample. Therefore, proper alignment of the lens was critical in obtaining reliable data. The process of aligning the lens involved both the dot on the wall and the use of the ALV-5000 correlation package. Prior to the alignment of the lens, the detection arm of the goniometer had to be moved to the appropriate
scattering angle for the upcoming experiment. Next, the lens holder was attached to the goniometer between the laser and the sample (Fig. 4). The first step in the alignment of the lens was to translate it both vertically and horizontally until the beam from the laser was directed through the center of the sample and on the dot on the wall. The lens was then adjusted along the length of the laser (x-direction) until the 'waist' of the beam was focused at the center of the sample. Due to diffraction effects, the focused beam could not be focused to a point with the lens, but instead was focused down to a condition where a 'waist' appeared. The 'waist' (Fig. 11) was considered the region where the smallest beam diameter occurred and was calculated by the following equation:

$$
\begin{equation*}
\mathrm{d} \cong \frac{4 \cdot \mathrm{f} \cdot \lambda}{\mathrm{D}}=\frac{4 \cdot(33 \mathrm{~mm}) \cdot(632.8 \mathrm{~nm})}{(0.68 \mathrm{~mm})}=39.1 \mu \mathrm{~m} \tag{4-1}
\end{equation*}
$$

where $d$ is the 'waist' diameter, $D$ is the incident beam diameter, $f$ is the focal length of the lens, and $\lambda$ is the wavelength of the incident beam.

In order to complete the alignment of the lens, the fiber optic cables had to be connected to the photomultiplier tubes and the ALV-5000 computer program had to be setup. A discussion of the computer program and its settings will follow in Section 4.4. With the computer parameters set and the computer in the count rate mode, the final alignment of the lens could be completed. The count rate mode was used because it only displayed the intensities seen by the two photomultipliers, channel 0 and channel 1 . The lens was adjusted vertically until the count rate of both channels was maximized. The maximized count rates showed that the detectors were looking at the center of the focused beam. After the maximization of the count rates, the intensity of the incident beam had to be attenuated to a workable range ( 50 to 250 kHz ).


Figure 11: Representation of the beam 'waist' caused by diffraction effects in the test tube, where $D$ is the incident beam diameter, $d$ is the waist diameter, and $f$ is the focal length of the lens. (Not shown to scale)

With all of the components added to the goniometer and aligned, the final adjustments could be made to channel 1. These adjustments were used to produce a signal-to-noise ratio near one and involved the top fiber. A high signal-to-noise ratio or $y$-intercept represented the level of useful input detected to the amount of background noise detected. It also implied that for a cross-correlation both fibers were looking at the same point in the sample from within the multiple scattering area. The ALV-5000 program was used in the run mode to determine when the top fiber was properly aligned. The program was set up to take multiple runs of 10 seconds per run. The short duration multiple runs were useful in giving immediate information on how a small adjustment of channel 1 affected the signal-to-noise ratio. Maximization of the signal-to-noise ratio was accomplished by adjusting the four micrometers of channel 1 (Figs. 7 and 8). The signal-to-noise ratio was generally maximized around a value of 0.900 . With the
alignment of channel 1 completed, the micrometer settings were recorded and the experiment could begin.

### 4.4 The ALV Computer Program

The correlation software package used in these experiments to read and analyze the data from the photomultiplier tubes was the ALV-5000 Multiple Tau Digital Correlator (see Appendix I). The software contains four main menus with several settings listed under each menu. Table 1 shows the parameters that were set at the beginning of each experiment. The rest of the settings were either kept at their default values or were not of interest to the research and therefore were not activated. Information on all parameters can be found in the ALV-5000 manual (1993) given in the reference section of this thesis.

Under the main menu, three subdirectories were useful in setting up the computer. The first subdirectory was the SampOpt subdirectory. This subdirectory allowed input of values that described the sample and setup. The first parameter was the wavelength of the laser used. For the static fluid experiments, the wavelength was 632.8 nm , while the wavelength used in the flowing fluid experiments was 532.5 nm . Other input parameters included refractive index of the sample, absolute temperature of the sample in $K$, and the type of solvent used in the sample. From these parameters, the ALV-5000 program calculated the viscosity of the sample. The other useful subdirectories under the main menu, were Angle and Multi. The subdirectory Angle allowed input of the scattering
angle, $\theta$, for each experiment. The subdirectory Multi was used to turn multiple runs on or off. This was valuable during the alignment of the setup (see Section 4.3).

Similar to the main menu, the control menu contained three subdirectories useful in the set up of the ALV-5000 program. The Scale subdirectory enabled the program to reduce the number of bits of information in the form of intensity that it was reading from each channel. This prevented intensity overflows in the two channels and allowed for an optimum performance of data fitting by the program. The Setup subdirectory contained several parameters that pertained to the type of experiment being performed. Duration was the amount of time each run would last. For the experiments performed in this thesis, a single correlation measurement was used, which had an initial sample time of 0.2 $\mu \mathrm{s}$. Both auto- and cross-correlation measurements were performed, depending on the type of experiment being performed. During autocorrelation measurements, it was necessary to select which channel was to be used, channel 0 or channel 1 . The final subdirectory of the control menu was the FileOpt subdirectory. This gave the user the option of saving information as an ASCII or a Binary data file. ASCII was useful in that it saved information in data sets containing the correlation data and its corresponding time. ASCII data was used to produce Figs. 22-25 in Section 6.5.1. Binary files were used the remainder of the time because they contained all of the pertinent information in a better disk space conserving manner.

Table 1: Input values for the ALV-5000 computer program.


### 4.5 Sample Preparation

The experiments performed using the static fluid setup dealt with particle sizing of nanometer spheres. Therefore, it was necessary to use special care when preparing the samples and not contaminate them. Duke Scientific supplied the core samples used in the experiments (see Appendix I). The core samples were polystyrene latex spheres at 10 percent solids by weight, which were then diluted down to a desired volume fraction for an upcoming experiment.

Preparation of the samples followed an established procedure to ensure consistency in the experiments. The first step was to wash a new test tube with deionized water to remove any dust. The deionized water came from an E-pure deionizer (see

Appendix I) and was stored in a sealed container. The storage of the water in the lab was done to allow the water to reach room temperature. With the test tube washed and dried, the next step was to weigh it and record the weight. This was done with an electronic lab scale (see Appendix I). The test tube was then filled with water and weighed again, with that weight also being recorded. When making these samples, a known volume fraction was desired. Therefore a calculation was made to determine the total weight of the test tube, the water and the sample necessary to produce that volume fraction. From that calculation, the required amount of the core sample by weight was then added using a pipette.

Hand calculations were used to compute the exact volume fraction using the weight of the test tube dry, the weight of the test tube and the water, the total weight of the test tube, water, and sample, and the percent solids of the core sample. These calculations are covered in Sundaresan (1999). The final step was to seal the test tube with Parafilm (see Appendix I) and label it with the volume fraction and the date prepared. The life expectancy of most samples was two to three weeks.

### 4.6 Experimental Procedure

The main goal of this thesis was to show how the two suppression techniques could be applied to a flowing multiple scattering sample for particle sizing. This chapter deals with the technique used to suppress the effects of multiple scattering. The previous sections described the procedure necessary to set up and align a non-flowing, dynamic light scattering experiment. This section describes the steps followed in carrying out the
tilt angle sweep experiments. The purpose of these experiments was to separate the fibers enough to map the speckle profile from multiple to single scattering.

As described in the theory, the method used to suppress multiple scattering involved two single-mode fiber optic cables that were separated by milliradians. Referring to the alignment procedure (Section 4.3), both fiber optics were adjusted to look at the focused beam from within the multiple scattering area. This resulted in a signal-to-noise ratio near unity and an intercept of the correlation function, displayed by the correlation package, near its theoretical limit of 2.0 [Nobbmann (1997)]. It was from this position where the first run was taken for a duration of 120 seconds. For each run, the initial delay time was 200 nanoseconds, and was increased automatically by the software to an optimum value. Data from each run was stored as an intensity correlation, $\mathrm{g}^{2}(\tau)$, and was converted by the software to the electric field correlation, $\mathrm{g}^{1}(\tau)$, by Eq. (37). This data was then saved on disk.

The top fiber was then tilted one or two divisions by the use of the back micrometer (Figs. 7 and 8). The movement of the top micrometer was calibrated by reversing a laser back through the fiber optic and noting the change in position of the light on a screen at a known length from the fiber's axis of rotation. From this calibration, it was determined that one division on the micrometer equaled 0.1125 mrad . After each tilt, the fiber was translated along the scattering direction until the intensity of channel 1 was maximized (Figs. 5 and 7). The ALV-5000 program was used in the multiple run mode for a duration of 10 seconds to determine when the intensity was maximized. For each scattering angle tested, this procedure was repeated in both tilt
angle directions until the signal-to-noise ratio became too small to give reliable data. The results from these experiments can be seen in Chapter VI.

## CHAPTER V

## FLOWING FLUID CONDITIONS: EXPERIMENTAL SETUP, ALIGNMENT, AND PROCEDURE

### 5.1 Introduction

Chapter IV explained the equipment used for multiple scattering suppression in a static fluid sample and its alignment. This chapter will focus on an experimental setup for flowing media capable of applying the same multiple scattering suppression technique. A technique for suppression of flow effects will also be discussed. The goniometer and flow circuit used in integrating the two suppression techniques were designed using the same principles as the static fluid setup. Details of the actual goniometer design are given in Sundaresan (1999). Therefore, this chapter will only give a brief description of the components used and any modifications that were made in switching from the static fluid setup. This will be followed by a discussion of the flow circuit and its design. Also discussed in this chapter will be the alignment of the setup and the experimental procedure used to collect data. The processes used to set up the ALV-5000 computer program and produce samples for the flowing setup were similar to the ones covered in Chapter IV. Any deviations from those processes will be mentioned.

### 5.2 Experimental Setup

Figure 12 gives a schematic diagram of the experimental setup used to perform particle sizing in a flowing fluid. Similar to the goniometer described in Chapter IV, this one can also be broken into three sections. The laser arm, the sample stand, and the detector arm were all modeled after the parts previously described in Chapter IV. A diagram of the dual detector setup used here can be seen in Figs. 5 and 6. However to accomplish both types of suppression, some of the equipment from the static fluid setup had to be replaced or modified. Since a detailed description of the goniometer design is presented in Sundaresan (1999), this section will give only a brief description of the equipment used and its general layout.


Figure 12: Experimental setup for simultaneous suppression of multiple scattering effects and of flow effects.

Like the static fluid setup, the first component of the setup was the light source. The light source was changed from the $\mathrm{He}-\mathrm{Ne}$ laser used in the static fluid setup to a Neodymium-Yttrium-Aluminum-Garnet (Nd: YAG) laser ( 532.5 nm wavelength). The 100 mW Nd : YAG laser was vertically polarized and had an original beam diameter of 0.70 mm . As in the static fluid setup, a holder mounted the laser to the goniometer arm. The holder maintained the horizontal alignment of the laser, while being capable of vertical adjustments. The same adjustable attenuator, lens, and lens holder from the static fluid setup were the final components mounted to the laser arm (see Fig. 12). Unlike the static fluid setup, the laser arm for the flowing setup was mobile and could travel over a range of $0^{\circ}$ to $48^{\circ}$. This range is marked as $\alpha_{\mathrm{L}}$ and can be seen (Fig: 12) to run from a position where the laser points in the direction perpendicular to the test cell $\left(\alpha_{L}=0^{\circ}\right)$ to where the sample stand interfered with the arm's movement ( $\alpha_{L \max }=48^{\circ}$ ).

For detection of the scattered signals, the same detector housing (see Fig. 6) was used as in the static fluid setup. This allowed the same technique of multiple scattering suppression to be utilized. The fiber optic cables and beamsplitter used in the static fluid setup were replaced with two new single-mode fiber optic cables specified for a 533 nm wavelength laser and a multi-band, nonpolarizing beamsplitter (see Appendix I) to accommodate the new laser. The new beamsplitter delivered $45 \pm 5 \%$ of the laser's power to both channels equally. The same polarizer was used here as in the static fluid setup, and it was mounted in the detector housing in front of the beamsplitter. Like the laser arm, the detector arm was mobile over a range of $0^{\circ}$ to $48^{\circ}$. This range started at the point where the detector arm was
perpendicular to the test cell ( $\alpha_{D}=0^{\circ}$ in Fig. 12) and ran to where the sample stand interfered with the arm's movement. This range is noted by $\alpha_{D}$ in Fig. 12. The photomultiplier tubes, the power supplies and the ALV-5000 correlation package described in Sections 4.2 and 4.3 were used in this setup as well.

The final section of the goniometer was the sample holder (Fig. 13). Like the static fluid setup, the sample holder was mounted on the axis of rotation of the goniometer. The sample holder for the flowing setup was a filter holder (see Appendix I) used to hold the rectangular test cell in place. The holder was mounted to a dovetail slide on top of the goniometer. The dovetail allowed the entire test cell to be moved in a direction that adjusted the depth of the laser inside the cell. The dovetail was equipped with a micrometer, so the position of the test cell could be marked and the depth of the laser within the test cell could be estimated (Fig. 13). The reason this depth was estimated will be discussed in Section 5.3. The slide portion of the dovetail carried the sample holder and contained a rotation stand made from Teflon. The rotation stand allowed the entire sample holder to be rotated through an angle of $\delta= \pm 60^{\circ}$ (Fig. 13). The usefulness of this feature will be discussed in Chapter VI, where results will be shown for experiments in which $\delta$ was varied through $\pm 10^{\circ}$.

The final component of the experimental setup was the flow system. The flow system was used to continuously circulate the fluids during the experiments. The flow system consisted of two holding tanks, a shuttle pump, the rectangular test cell, and Tygon tubing (see Appendix I). The holding tanks were fabricated from


Figure 13: Diagram of the test cell mounted on the rotation stand and the angle of rotation, $\delta$. A negative $\delta$ represented a rotation of the test cell toward the laser (counter clockwise), while a positive value meant a rotation toward the detector housing (clockwise). Also shown is the dovetail and micrometer used to move the test cell.

Plexiglas (see Appendix I). Both tanks were designed with large lids for easy filling and cleaning and had a capacity of 250 ml each (Fig. 14).

A dual channel shuttle pump was used to circulate the fluid throughout the system (see Appendix I). The shuttle pump was selected because it was capable of pumping at flow rates ranging from 1.25 to $25 \mathrm{ml} / \mathrm{min}$ without damaging the particles. This was a valuable aspect in the design of the flow system, because the theory used in determining the particle size was based on light scattering by spherical particles. The shuttle pump's design was similar to a peristaltic pump in that an oscillating mechanism squeezed the tubing to generate the flow.


Figure 14: Side view of the holding tanks, lid, and gasket used in the flow circuit.

However, the shuttle pump's design produced flow without the tubes being squeezed to occlusion, which significantly reduced the possibility of damaging the micron sized test particles. Results from experiments utilizing the flow system will be presented in Chapter VI for various flow rates. Table 2 shows experimentally determined values for the different flow rates used throughout the research and the corresponding velocities to those flow rates. The average velocities were calculated by dividing the flow rate by the cross sectional area of the test cell $\left(48 \mathrm{~mm}^{2}\right)$.

Table 2: Experimentally determined flow rates and their corresponding calculated average velocities.

| Pump <br> Settings | Calculated <br> Flow Rate <br> $(\mathrm{ml} / \mathrm{min})$ | Calculated <br> Average <br> Velocity <br> $(\mathrm{mm} / \mathrm{sec})$ |
| :---: | :---: | :---: |
| $0 \%$ Flow | 0.00 | 0.00 |
| $50 \%$ Flow | 9.46 | 3.28 |
| $75 \%$ Flow | 15.07 | 5.23 |
| $100 \%$ Flow | 22.40 | 7.78 |

The most vital component in the flow system was the rectangular test cell (see Appendix 1). The test cell used in the flowing experiments was made from quartz glass and had dimensions of $6 \mathrm{~mm} \times 8 \mathrm{~mm} \times 30.48 \mathrm{~cm}$. The length of the tube was necessary to produce a fully developed laminar flow profile and to reduce the possibility of entrance effects interfering with the experiments. The following simple calculations were performed to show that the flow system produced the situation described above. The Reynolds Number was first calculated with the following equation [Janna (1993)]:

$$
\begin{align*}
\operatorname{Re}=\frac{D_{\mathrm{h}} \varrho v}{\mu}= & \\
& \frac{\left(6.86 \times 10^{-3} \mathrm{~m}\right)\left(997.3 \mathrm{~kg} / \mathrm{m}^{3}\right)\left(7.78 \times 10^{-3} \mathrm{~m} / \mathrm{s}\right)}{\left(947.95 \times 10^{-6} \mathrm{~Pa} \cdot \mathrm{~s}\right)}=56 \tag{5-1}
\end{align*}
$$

where $\rho$ is the fluid density, $\mu$ is the dynamic viscosity of the fluid, $v$ is the fluid velocity, and $D_{h}=\frac{4 A}{P}$ is the hydraulic diameter with $A$ being the cross-sectional area and $P$ being the wetted perimeter. From the Reynolds Number, it was determined that the flow regime was laminar ( $\operatorname{Re}<2300$ ), which was desired for suppression of flow effects. With the Reynolds Number determined, calculation of the length necessary for fully developed flow was performed using this equation [Janna (1993)]:

$$
\begin{equation*}
L_{e}=(0.06)\left(D_{h}\right)(\mathrm{Re})=(0.06)(6.86 \mathrm{~mm})(56)=23 \mathrm{~mm} . \tag{5-2}
\end{equation*}
$$

Experiments using the flowing fluid system were performed near the center ( $\sim 150$ mm from the entrance) of the test cell, and therefore were conducted in a region of fully developed flow. Further calculations depicting the fully developed flow profile
can be seen in Appendix IV. Also contained in Appendix IV is a discussion of the time constants associated with fluid flow and Brownian motion of the particles.

### 5.3 Alignment

It has already been demonstrated in Chapter IV that the alignment of a dynamic light scattering device is critical. This setup is no different from any of the others that have been previously mentioned. The alignment process of the flow setup was very similar to the process described in Section 4.3.

The first component of the flowing setup to be aligned was the laser. To accomplish this, the laser arm (Fig. 12) was moved to $\alpha_{L}=0^{\circ}$ or the point were the laser arm was perpendicular to the test cell. To aid in the alignment of the laser and the fiber optics, a mirror was placed in the sample holder. With the mirror in place, the laser was adjusted until the back reflection was co-linear with the incident beam, which was determined by reflecting the beam almost back into the hole where the beam exits the laser. This positioning meant that the laser was aligned both horizontally and vertically. Once alignment was achieved, the laser was moved to the desired angle, $\alpha_{\mathrm{L}}$, for the upcoming experiment. Scattering angle calculations, corresponding to different laser and detector angles, are given in Section 5.4. The detector arm was then positioned at an angle that corresponded to $\alpha_{L}=\alpha_{D}$. As described in Chapter III, this allowed for the suppression of the flow effects by creating a scattering vector perpendicular to the flow vector. A second benefit of this positioning was apparent during the alignment of the multiple scattering
suppression equipment (i.e., the detectors). Laser arm and detector arm positions corresponding to $\alpha_{L}=\alpha_{D}$ matched the angles necessary for the incident beam to be reflected into the detector housing by the mirror. This positioning was equivalent to that described in Section 4.3, where the detector arm was moved to $\theta=0^{\circ}$ in order to directly view the incident beam. A similar alignment procedure as that described in Section 4.3 was followed for the detectors. With the laser and the fiber optic cables aligned, the mirror could be removed, and the following components could be added: the attenuator, the polarizer, the lens, and the test cell.

The first components added to the setup were the attenuator and the polarizer (Fig. 12.). The alignment procedure for these two pieces was identical to that described in Section 4.3. After the addition of the polarizer and attenuator, the fiber optic cables were connected to the PMT's, and the computer was set up following the procedure previously described in Section 4.4.

The next component added to the setup was the test cell. Initially the test cell was placed in the holder so that the 8 mm wall appeared to be visibly perpendicular to the incident laser beam. Because of the use of a rectangular test cell, direct specular reflections from the front and back walls caused intensity levels too high for the PMT detectors (see Fig. 15). The reflections also caused heterodyning that interfered with the collection of data. To alleviate these problems, the test cell was slightly rotated from its vertical orientation (approximately $5.0^{\circ}$ about an axis parallel to the cell's length) to direct the reflections away from the fiber optics. This method of reducing the reflections at the detector was only partially successful, and a second method had to be employed to suppress specular scattering.


Figure 15: Schematic showing the positioning of the electrical tape to block reflections.

The second method consisted of applying two small strips of electrical tape to the sample holder (see Fig. 15). One strip blocked light reflecting from the front wall, while the second strip blocked the reflection from the back wall. The two strips had to be placed on the holder in a manner that blocked only the reflections and did not interfere with the scattered signal. The count rate mode, of the ALV5000 software, was useful in applying the strips of tape to the holder by showing when a strip blocked out the scattered signal and not just the reflection. The last step in aligning the test cell was to move it so that the detection area was just inside the front wall. This was accomplished by moving the dovetail slide toward the laser and detectors with a micrometer while making autocorrelation readings with the ALV5000 program (see Fig. 12). Readings of the signal-to-noise ratio were taken every 10 divisions turned on the micrometer ( 5 divisions $=0.05 \mathrm{~mm}$ ). It was found that after approximately 100 divisions ( 1.0 mm ), the signal-to-noise ratio had risen to a
usable level near 0.9. All experimental results described in Chapter VI were performed at this approximate depth into the test cell.

With the test cell in place and the reflections blocked, the final component to be added to the setup was the lens. The lens was placed on the goniometer between the attenuator and the test cell. The lens was first adjusted vertically and horizontally until the count rate in both channels was maximized. The lens was then adjusted along the direction of the laser beam until the 'waist' of the laser beam was focused at the location where the detectors were looking. This alignment was determined also by maximizing the count rate of both channels. As in the static fluid setup, the ALV-5000 software was used during this process to determine the lens position that maximized the count rate. At this point, the last step in the alignment procedure was to repeat the steps outlined in Section 4.3 to maximize the signal-tonoise ratio.

### 5.4 Scattering Angle Calculations

Knowledge of the scattering angle is a necessary part in the success of any dynamic light scattering experiment in which the goal is to determine particle size. This is because Eq. (3-7) depends on the scattering angle to calculate particle size, and therefore an incorrect scattering angle would lead to an incorrect calculation of the particle size. For the flowing setup, the scattering angle had to be calculated before any runs could be made. The scattering angle is defined as the angle between the light directly transmitted through the sample volume and the direction from the sample volume to the detector and is given by the symbol $\theta$ in Fig. 16. The
determination of the scattering angle is generally a straightforward process, but for the flowing system, it was complicated by refraction of the incident beam in the sample. As seen in Fig. 16, the scattering angle is altered by refraction effects at the boundary between the air and the sample. The air to glass and the glass to sample boundaries are assumed to be parallel. Therefore, the net refraction effects from those two boundaries sums to zero. Due to these refractive effects, the scattering angle becomes a function of several variables.

The following calculations were used to determine the different scattering angles used in the flowing experiments. The results of these calculations are summarized in Table 3. Also all variables used in the following section refer to Fig.


Figure 16: Refraction effects on the scattering angle in the flow cell. (Not drawn to scale)
16. Calculation of the scattering angle was governed by simple trigonometric functions as well as Snell's Law, which states:

$$
\begin{equation*}
\mathrm{n}_{\text {solutuon }}\left[\sin \left(\beta_{1,2}\right)\right]=\mathrm{n}_{\text {air }}\left[\sin \left(\alpha_{\mathrm{L}, \mathrm{D}}\right)\right] . \tag{5-3}
\end{equation*}
$$

The scattering angle, $\theta$ was given as follows:

$$
\begin{equation*}
\theta=\pi-\beta_{1}-\beta_{2}=\pi-2 \beta \tag{5-4}
\end{equation*}
$$

where $\beta_{1}$ and $\beta_{2}$ were calculated by Snell's Law and were equal by the fact that $\alpha_{L}=$ $\alpha_{D}$, and the bisector between those angles was perpendicular to the test cell in the experiments performed for this thesis.

### 5.5 Experimental Procedure

This chapter has dealt with the experimental setup used to apply two independent suppression techniques to a dense flowing suspension for particle sizing. The previous sections have described the equipment used as well as the alignment procedure necessary to obtain reliable data from the setup. This section describes the experimental procedures used to carry out different tests with the setup. The first type of experiment performed with the flowing setup was performed on single scattering samples to show how rotation of the test cell would affect flow suppression. The second type of experiment was performed to demonstrate how effective the suppression of multiple scattering would be in a flowing media.

Table 3: Summary of scattering angle calculation for experiments where $\alpha_{L}=\alpha_{D}$ and $\beta_{1}=\beta_{2}$.

| $\alpha_{L}=\alpha_{D}$ |  | $\beta_{1}=\beta_{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (Degrees) |  |  |$\theta$ (Degrees)

### 5.5.1 Cell Rotation Effects

Preliminary experiments were performed with the flow setup to determine the sensitivity of the suppression technique to flow effects. These experiments were conducted on dilute samples, the results of which will be given in Chapter VI. From the alignment procedure (Section 5.3), the laser and detector arms were already
positioned where $\alpha_{L}=\alpha_{D}$ and the test cell was perpendicular to the angular bisector between the two arms. The suppression experiments involved taking autocorrelation measurements over a range of cell rotation angles from $\delta=-10^{\circ}$ to $\delta=+10^{\circ}$ (Fig. 12). The first experiments were done for a cell rotation of $\delta=0^{\circ}$. Autocorrelation measurements were taken for flow rates of $0 \%, 50 \%$, and $100 \%$ (see Table 2). The test cell was then rotated to $-5^{\circ},-10^{\circ},+5^{\circ}$, and $+10^{\circ}$ and the same autocorrelation measurements at all of the flow rates listed above were taken at each rotation angle. A negative $\delta$ represented a rotation of the flow cell toward the laser, while a positive $\delta$ represented a rotation toward the detector housing (Fig. 13). Each time the test cell was rotated, the electrical tape used to block reflections was adjusted according to the procedure given in Section 5.3. This process was repeated for $0.204 \mu \mathrm{~m}$ and $0.300 \mu \mathrm{~m}$ PSL particles.

### 5.5.2 Particle Sizing Experiments

The experiments performed using the flow setup were done to show how both multiple scattering effects and flow effects could be suppressed in order to determine particle size. The same steps from the procedure used on the static fluid setup (Sections 4.3 and 4.6 ) were repeated with the flowing setup to accomplish the suppression of the multiple scattering effects. Many of the steps already described (see Section 5.3) during the alignment of the flow setup helped to ensure that the flow effects were suppressed. Those steps included positioning the laser (incident wave vector) and detector (scattered wave vector) arms where $\alpha_{L}=\alpha_{D}$ and having the bisector between the two arms perpendicular to the test cell. Both of these facts
produced a scattering wave vector perpendicular to the flow vector. As described in Chapter III, suppression of flow effects was accomplished by the dot product of these two vectors equaling zero (Eq. (3-14)). Results from experiments with the flow setup on signal-to-noise ratio mapping and diameter mapping will be shown in Chapter VI for a variety of concentrations, particle sizes, and flow rates at a scattering angle of $112^{\circ}$.

## CHAPTER VI

## RESULTS AND DISCUSSION

### 6.1 Static Fluid Setup: Introduction

The purpose of this thesis is to show how two independent suppression techniques can be applied to a flowing system to determine particle size. The experimental setup described in Chapter IV was used to show how a single-beam, two detector system could be used to suppress multiple scattering in a static fluid sample. Several of the preliminary experiments performed using the static fluid setup did not yield the necessary data to determine particle size. However, these experiments were not failures, as they led to a discovery of nine characteristics of the static fluid setup. Four of these characteristics will be discussed in detail in Section 6.2 and the five other characteristics are discussed in Sundaresan (1999). The last 17 experiments performed using the static fluid setup did yield data usable to determine particle sizes. Three experiments were performed and repeated at a scattering angle of $90^{\circ}$ on samples in a square cell. The experiments included various volume fractions of $0.107 \mu \mathrm{~m}$ PSL particles and are discussed in Sundaresan (1999). Also described in Sundaresan (1999) is a match of those experiments to Eq. (3-8) from the theory. The last 11 static fluid setup experiments were performed
using a circular cell and a water bath. These experiments were performed at a variety of scattering angles on different volume fractions of $0.107 \mu \mathrm{~m}$ PSL particles (see Appendix I). Section 6.3 will review the results from experiments performed using the static fluid setup. Table 4 contains a summary of the last 17 static experiments performed, and Table 5 shows detailed results of the experimental data corresponding to the static fluid figures presented in this chapter. Both tables are given in Appendix II.

### 6.2 Static Fluid Setup Characteristics

As previously mentioned, the first static fluid experiments were performed to gain experience in employing the multiple scattering suppression technique. Through this experience, nine characteristics were discovered about the static fluid setup and are listed below. A discussion of the first four will follow the list, while the last five are covered by Sundaresan (1999).

1. Flaws in the beaker used to hold the water bath interfered with data collection.
2. To align the detectors properly, maximize the count rate in both channels at a scattering angle of $0^{\circ}$ first. Then move the detector arm to the desired scattering angle and maximize the signal-to-noise ratio by adjusting channel 1. These steps should be performed prior to each new experiment.
3. During the tilt angle sweep, the intensity of channel 1 should be maximized by translating the fiber after each tilt as described in Section 4.6.
4. The theory describing the suppression of multiple scattering is based on the idea that the detection area of each fiber overlaps the area through which the incident beam passes in the sample. To ensure this all components should be aligned horizontally.
5. Measurements at scattering angles less than $60^{\circ}$ were hard to align due to uncontrollable intensities detected due to direct transmission from the sample.
6. As explained in Section 4.3, the positioning and alignment of the lens is critical to the collection of usable data.
7. Also critical to the collection of usable data is the positioning of the detector housing.
8. The use of $0.107 \mu \mathrm{~m}$ PSL particles was found the most reliable diameter particles to work with.
9. Larger particles were thought to settle and therefore not behave as Brownian particles, which led to problems in the determination of their sizes.

The first characteristic that was found using the existing static fluid setup involved the water bath. It was originally contained in a 150 ml beaker that had been modified to hold the water bath. The top of the beaker had been cut off flat and a lid had been manufactured from Teflon to hold the test tube vertically at the center of the beaker. Several problems arose during the use of this beaker to hold the water bath, including interference of the incident beam by the labeling printed on the beaker's side. As previously mentioned in Chapter IV, component alignment was critical to the collection of reliable data. The writing on the beaker often interfered with that alignment, as special care was necessary to ensure that neither the incident nor the scattering light directions passed through the writing. A second similar problem stemmed from the lines in the beaker glass. These lines were a product of how the beaker was made and caused problems with the purity of the incident beam. To alleviate the problems, a piece of glass tubing (see Appendix I) was formed by the Oklahoma State University Glass Shop into a new container for the water bath. The new container was free of these problems and therefore helped to ease alignment and to provide data that was more reliable.

The second characteristic dealt with the alignment of the detectors. The preliminary experiments preformed with the static fluid setup led to the procedure
described in Section 4.3. The fiber optic detectors were initially aligned at a scattering angle of $0^{\circ}(\theta$ in Fig. 4) to start each experiment. This was done to ensure that the detection areas of both fibers were coincident with the incident beam. Next, the detection arm was moved to the desired scattering angle, and the remaining steps of the alignment procedure were completed. With the alignment procedure completed, the final step was to maximize the signal-to-noise ratio, which is also described in Section 4.3. This step was performed in order to ensure that both channels were focused on the same location within the sample.

The third characteristic dealt with the alignment of the top fiber throughout the tilt angle sweep experiments described in Section 4.6. From the alignment described in Section 4.3, it was assumed that both channels were focused on the same point at the center of incident beam. This point should correspond to a maximum count rate in both channels. Each time that the detector for channel 1 was tilted, it was no longer looking at the same point as the detector for channel 0 ; and therefore it was translated in the proper $x$-direction (Fig. 5) until its count rate was maximized. If the detector for channel 1 was tilted upward, a translation towards the sample was necessary, while a downward tilt corresponded to a translation in the opposite direction.

Characteristic 4 dealt with the horizontal alignment of the components on the goniometer. As mentioned in Section 3.3, an overlap between the detection area of the fiber optics and the focused beam in the sample was necessary to achieve a correlation function. This alignment was controlled at several locations on the flow system. In particular, the lens and laser holders were useful in accomplishing the desired horizontal alignment. The laser holder controlled the height of the laser to begin the experiment.

The lens holder overcame any refraction problems from the water bath or sample by adjusting the height of the incident beam before it entered the water bath.

The four characteristics described above along with the five discussed in Sundaresan (1999) led to the procedural steps described in Section 4.3. The knowledge gained by performing these preliminary experiments with the static fluid setup developed a foundation for later success. The following section will describe the results from Experiments 38-48.

### 6.3 Static Fluid Setup Experiments

Although the experiments performed on the static fluid setup were not the primary focus of this research project, those experiments provided useful information and helped in the design of the flowing setup. The following sections will discuss the two types of experiments performed on the static fluid samples. A discussion of the results from a scattering angle sweep experiment is given in Section 6.3.1. Results from tilt angle experiments are shown in Section 6.3.2. Knowledge gathered from those two experimental processes was then used during the design of the flowing setup.

### 6.3.1 Scattering Angle Sweep

A scattering angle sweep was performed on a single scattering (dilute) sample of $0.107 \mu \mathrm{~m}$ PSL particles (see Appendix I). A scattering angle sweep requires no adjustment of the top fiber once alignment has been accomplished (see Section 4.3). For
a scattering angle sweep experiment, autocorrelation measurements are taken (where in the data collected by each detector is self-correlated). If the sample is single scattering, the data should reveal the correct radius and the measured second cumulant should be near zero (see Section 3.3). Autocorrelation measurements utilizing both channels were taken for scattering angles ranging from $15^{\circ}$ to $120^{\circ}$.

Two trends were discovered from the scattering angle sweep experiments. The first noticeable trend was seen in the intensity versus scattering angle plot of Fig. 17. The intensity remained relatively constant between $120^{\circ}$ and $50^{\circ}$, seeing only a 10 kHz increase. As the scattering angle sweep decreased below $50^{\circ}$, the intensity in both channels increased by as much as $100+\mathrm{kHz}$. A comparison of the two extreme scattering angles, $15^{\circ}$ and $120^{\circ}$, showed that, at the smaller scattering angle, the intensities increased. The increase in intensities for channel 0 and channel 1 at $15^{\circ}$ was 4.40 and 4.81 times the value measured at $120^{\circ}$, respectively.

Figure 18 is a plot of the intensity from Fig. 17 times the sine of the scattering angle versus the scattering angle. This was done to compensate for the larger detection area seen at the lower scattering angles. If larger detection areas were the only cause, compensating with the $\sin (\theta)$ should produce a linear curve for intensity as a function of scattering angle, but Fig. 18 shows the same trend as seen in Fig. 17. The intensity increased as the scattering angle decreased.

A plot of the radius versus the scattering angle (Fig. 19) showed the second trend associated with the static fluid setup. The radius determined from the data decreased as the scattering angle decreased until a minimum was reached around $40^{\circ}$. Data points for the remaining five scattering angles seem to show an upward trend with an exception


Figure 17: Intensity versus scattering angle for a single scattering sample of $0.107 \mu \mathrm{~m}$ PSL particles in water. Autocorrelation measurements were taken using both channel 0 (circles) and channel 1 (inverted triangles). Data corresponds to Experiment 42.


Figure 18: Intensity times $\sin (\theta)$ versus scattering angle to determine effects of scattering angle for a single scattering sample of $0.107 \mu \mathrm{~m}$ PSL particles in water. Autocorrelation measurements were taken using both channel 0 (circles) and channel 1 (inverted triangles). Data corresponds to Experiment 42.


Figure 19: Radius versus scattering angle for a single scattering sample of $0.107 \mu \mathrm{~m}$ PSL particles in water. Autocorrelation measurements were taken using both channel 0 (circles) and channel 1 (inverted triangles). Data corresponds to Experiment 42.
occurring at $30^{\circ}$, which dropped back down. For the range of scattering angles from $105^{\circ}$ to $45^{\circ}$, the autocorrelation data was accurate in calculating a radius within the acceptance range (shown by the horizontal dashed lines in Fig. 19). Problems arose at the extreme values of the scattering angle range, especially at the smaller angles $\left(\theta<45^{\circ}\right)$. At the smaller angles, the diameter calculated from the data became erratic. This erratic behavior corresponded to the area where the intensity increased rapidly (Fig. 17). A second scattering angle sweep was performed on a different dilute sample and the data from that experiment (Exp. 45 in Appendix II) shows the same trends. The following discussion of the tilt angle sweep results will further illustrate the problems encountered at small scattering angles.

### 6.3.2 Tilt Angle Sweep

This section will discuss the result from the tilt angle sweep experiments. These experiments were conducted using the procedure described in Section 4.6. Crosscorrelation experiments were performed on samples of $0.107 \mu \mathrm{~m}$ PSL particles (see Appendix $I$ ) at volume fractions ranging from 0.1330 to 0.5025 percent by weight. The data shows that the best results occurred at scattering angles of $90^{\circ}$ and $120^{\circ}$, while scattering angles of $30^{\circ}, 45^{\circ}$, and $60^{\circ}$ gave poor results. These results agree with the ones given in the scattering angle sweep discussion.

Figure 20 shows the signal-to-noise ratio versus tilt angle sweep for the five different scattering angles listed above. The tilt angle sweep experiments were performed on samples of $0.107 \mu \mathrm{~m}$ PSL particles at volume fractions of approximately


Figure 20: Signal-to-noise ratio versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at volume fractions of approximately 0.3 percent by weight for various scattering angles. Data corresponds to Experiments 39, $40,43,47$, and 48.
0.30 percent by weight. The signal-to-noise ratio for all five scattering angles is shown to peak near 0.90 . All of the curves with the exception of $45^{\circ}$ show long 'shoulders', which should correspond to areas of single scattering. As will be shown in Fig. 21, this is not necessarily the case.

The second figure, Fig. 21, shows the radius as a function of tilt angle for the same five scattering angles. A comparison of Figs. 20 and 21 shows how multiple scattering effects are presumably more influential in forward scattering (scattering at small angles). These effects cause faster decay rates, which lead to apparently smaller diameters. For scattering angles of $30^{\circ}$ and $60^{\circ}$, the 'shoulder' region (Fig. 20) was well defined and should correspond to single scattering results. However, in Fig. 21, the radius map for $\theta=30^{\circ}$ was lower than the expected range, while the radius map for $\theta=$ $60^{\circ}$ barely reached the acceptance range. From all of the figures (17-21), the difficulties of dealing with a scattering angle around $45^{\circ}$ can be seen as no 'shoulder' is present in Fig. 20 and the radius never reaches the acceptable range in Fig. 21. As discussed in Section 2.3, Meyer et al. (1997a) covered several scattering angles, with $60^{\circ}$ being the smallest. Meyer reported no problems with taking data at any scattering angle.

The influences of concentration on the static fluid experiments performed with a water bath and circular test tube are limited. No more than two concentrations were used at one scattering angle (see Appendix II). However, some general trends can be seen in Figs. 22 and 23. Figure 22 shows the signal-to-noise ratio versus tilt angle for two different concentrations of $0.107 \mu \mathrm{~m}$ particles. The peaks for both concentrations are near the same point, but the more concentrated sample's signal-to-noise ratio drops much faster. This causes the peak region to be more pronounced. In addition, Fig. 22 shows


Figure 21: Radius versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at volume fractions of approximately 0.3 percent by weight various for scattering angles. Data corresponds to Experiments 39, 40, 43, 47, and 48.


Figure 22: Signal-to-noise ratio versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at a scattering angle of $90^{\circ}$ for two volume fractions. Data corresponds to Experiments 38 and 39.


Figure 23: Radius versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at a scattering angle of $90^{\circ}$ for two volume fractions. Data corresponds to Experiments 38 and 39.
that the shoulder region for the less concentrated sample extends 1.5 mrads farther in tilt angle than that of the more concentrated sample. The shoulder region begins at a higher signal-to-noise ratio for the less concentrated sample as well.

Figure 23 shows another trend that was connected to concentration level. The radius values given at the peak were much lower for the more concentrated sample, $\sim 31$ nm as compared to $\sim 46 \mathrm{~nm}$. This was due to the faster decay rates associated with higher concentrations. This fact meant that greater separation was necessary between the two detectors to achieve multiple scattering suppression. These trends were also seen at scattering angles of $60^{\circ}$ and $120^{\circ}$ (see Appendix II). Sundaresan (1999) contains further discussion on the effects of concentration.

### 6.4 Flowing Fluid Setup: Introduction

The goal of this thesis is to demonstrate that two independent suppression techniques can be used to determine particle size in a dense flowing system. The experimental setup depicted in Chapter V was used to suppress both multiple scattering effects and flow effects. Preliminary autocorrelation experiments were performed with the flowing setup to determine the sensitivity of measurements to a perpendicular alignment between the bisector (of the laser and detector arms) and the flow cell. These experiments were performed on $0.107 \mu \mathrm{~m}, 0.204 \mu \mathrm{~m}$, and $0.300 \mu \mathrm{~m}$ PSL particles (see Appendix I) at cell rotation angles ranging from $\delta=-10^{\circ}$ to $\delta=+10^{\circ}$ (see Fig. 13). The results from these experiments are described in Section 6.5.1. After the rotation
experiments, the test cell was returned to $0^{\circ}$ (perpendicular to the bisector) so that the tilt angle sweep experiments could be performed.

As with the static setup, experiments were also performed through a tilt angle sweep to suppress multiple scattering effects. The purpose of these experiments was to show that both flow suppression and multiple scattering suppression could be used simultaneously to determine particle size. The following four parameters were investigated to determine their effects on particle sizing: flow velocity, particle size, concentration, and scattering angle. This thesis concentrates on varying the first three parameters at a scattering angle of $112^{\circ}$. The effects of varying the last parameter, scattering angle, is presented in Sundaresan (1999), where the effects of the first two parameters are studied for scattering angles of $122^{\circ}$ and $136^{\circ}$. Preliminary tilt angle sweep experiments were conducted on samples of $0.107 \mu \mathrm{~m}$ PSL at volume fractions of 0.066 percent by weight and 0.198 percent by weight for various flow rates. The sample of $0.107 \mu \mathrm{~m}$ particles at a volume fraction of 0.198 percent by weight as well as two additional particles sizes, $0.098 \mu \mathrm{~m}$ and $0.203 \mu \mathrm{~m}$ PSL particles (see Appendix I) were used to conduct the remaining tilt angle experiments. Experiments were performed at two flow rates for the $0.098 \mu \mathrm{~m}$ and $0.203 \mu \mathrm{~m}$ particles, while the $0.107 \mu \mathrm{~m}$ particles were test at three flow rates. For the $0.098 \mu \mathrm{~m}$ particles, volume fractions of 0.30 percent and 0.86 percent by weight were used, while a volume fraction of 0.20 percent by weight was used for the $0.203 \mu \mathrm{~m}$ particles. All of these results are discussed in Section 6.5.2.

### 6.5 Flowing Fluid Setup Experiments

The experiments performed with the flowing fluid setup demonstrate that the two independent suppression techniques can be used together to determine particle size in a dense flowing system of particles. The following two sections will describe the results of the two experiments used to prove that the suppression techniques work. Section 6.5.1 shows the degree of sensitivity of the setup to cell rotation, which affects the relationship between the bisector of the laser and detector arms to the flow direction. Results from tilt angle sweep experiments are shown in Section 6.5.2. It was the results from those experiments, which showed that the two suppression techniques could be applied together to a dense flowing system of particles in order to determine particle size.

### 6.5.1 Cell Rotation Effects

As described in Chapters III and V , the alignment of the flowing fluid setup was critical to the success of flow effect suppression. The geometric alignment described in Chapter V allowed the flow vector and the scattering wave vector to be positioned perpendicular to each other. The sensitivity of this alignment is examined in this section. To test this sensitivity, autocorrelation measurements were performed on single scattering (dilute) samples of $0.107 \mu \mathrm{~m}, 0.204 \mu \mathrm{~m}$, and $0.300 \mu \mathrm{~m}$ PSL particles. The first experiment was conducted on $0.107 \mu \mathrm{~m}$ PSL particles at a cell rotation angle of $\delta=0^{\circ}$ (see Fig. 13) for flow rates of $0 \%, 50 \%$ and $100 \%$ (see Table 2 for corresponding velocities). This section of the research was completed by five additional experiments
that followed the procedure outlined in Section 5.5 .1 for cell rotation angles of $\delta=-5^{\circ}$, $10^{\circ},+5^{\circ}$, and $+10^{\circ}$ and flow rates of $0 \%, 50 \%$ and $100 \%$. To interpret the effects of flow on the data taken during these experiments, the natural logarithm of the normalized field autocorrelation function $\left(g^{1}(\tau)\right)$ was plotted versus the delay time. To attain these values, the normalized field correlation function was found from the intensity field correlation data given by the ALV-5000 program using the following equation:

$$
\begin{equation*}
g^{1}(\tau)=\sqrt{\left(g^{2}(\tau)-1\right)} \tag{6-1}
\end{equation*}
$$

which is Eq. (3-6) solved for $g^{1}(\tau)$ assuming $\gamma^{2}(\theta)$ is equal to 1.0. Substituting Eq. (3-13) into Eq. (6-1) gives:

$$
\begin{equation*}
\exp \left(-D_{0} q^{2} \tau+(i \bar{q} \bullet \overline{\mathrm{v}}(\mathrm{r})) \tau\right)=\sqrt{\left(\mathrm{g}^{2}(\tau)-1\right)} \tag{6-2}
\end{equation*}
$$

In order to utilize this equation, the second term of the exponential on the left-hand side must be eliminated or suppressed. This is accomplished by forming the geometry described in Chapter V. If this term is suppressed, it can be seen that taking the natural logarithm and solving for the diffusion constant will lead to determination of the particle size by using the slope of the line as seen in Figs. 24-27. Therefore, the closer to linear that the plot of this equation is, the more reliable the data is and the better the determination of the particle size.

Figure 24 shows the natural logarithm of the normalized field autocorrelation function of channel 0 versus delay time using $0.107 \mu \mathrm{~m}$ PSL particles for no flow at five different cell rotation angles, $\delta$. To allow the graphs to be clearer, data from only channel 0 is given as a representation of the effect of $\delta$. However, a similar effect was found from the data collected by channel 1 . From this graph, it can be seen that the cell rotation has
very little effect on a static fluid case. Figure 25 shows the same information as Fig. 24, but for the case of $100 \%$ flow. The effect of flow can be seen to cause the plots represented by $\delta= \pm 10^{\circ}$ to curve. The effect of flow is more suppressed in the three smaller rotation angles, $\delta=0^{\circ}$ and $\delta= \pm 5^{\circ}$.

The effects of flow are even more pronounced when comparing Fig. 26 to Fig. 27, which are repeats of Figs. 24 and 25 except for the use of $0.204 \mu \mathrm{~m}$ PSL particles. Figure 26 is similar to Fig. 24 showing that cell rotation has little effect on static fluid conditions. Two facts are immediately apparent when viewing Fig. 27. The first fact is that as particle size is increased, rotating the cell away from a position perpendicular to the bisector had a greater effect on the influence of flow for the data taken. Secondly, rotation of the test cell toward the laser has a greater impact on the linearity of the data than rotation toward the detectors.

Experiments to determine the effects of test cell rotation on particle sizing of $0.300 \mu \mathrm{~m}$ particle samples were tried unsuccessfully. The intensity levels measured at scattering angles of $112^{\circ}, 136^{\circ}$, and $150^{\circ}$ were too low to produce useable correlations. This was determined to be due to the Rayleigh Gans form factor that is associated with $0.300 \mu \mathrm{~m}$ PSL particles. The Rayleigh Gans form factor is a measure of the amount of intensity scattered in a certain direction. Larger particles tend to be more forward scattering and the flow setup was built to take measurements from backscattering. Sundaresan (1999) contains calculations of the Rayleigh Gans form factors for the particles used in this research.

Three characteristics of the flowing fluid setup were found from the cell rotation experiments. First, the effect of flow can be suppressed if the bisector between the laser


Figure 24: Normalized field autocorrelation function for channel 0 versus delay time for a dilute sample of $0.107 \mu \mathrm{~m}$ PSL particles at a $0 \%$ flow rate. Data corresponds to Experiment 58.


Figure 25: Normalized field autocorrelation function for channel 0 versus delay time for a dilute sample of $0.107 \mu \mathrm{~m}$ PSL particles at a $100 \%$ flow rate. Data corresponds to Experiment 58.


Figure 26: Normalized field autocorrelation function for channel 0 versus delay time for a dilute sample of $0.204 \mu \mathrm{~m}$ PSL particles at a $0 \%$ flow rate. Data corresponds to Experiment 60.


Figure 27: Normalized field autocorrelation function for channel 0 versus delay time for a dilute sample of $0.204 \mu \mathrm{~m}$ PSL particles at a $100 \%$ flow rate. Data corresponds to Experiment 60.
and detector arms is kept perpendicular to the flow direction. Second, the plots show that, while the setup is sensitive to cell rotation, this sensitivity is minor for cell rotation angles of $-5^{\circ} \leq \delta \leq 5^{\circ}$. The final characteristic is that cell rotation toward the laser has a greater impact on data collection than rotation toward the detector. This characteristic can be attributed to the fact that a rotation toward the laser caused the detection area to be move deeper within the cell where higher velocities were present (see Appendix IV).

### 6.5.2 Tilt Angle Sweep

Tilt angle sweep experiments for the flowing system were similar to the ones described in Section 4.6 for the static fluid setup. These experiments resulted in a means of mapping out the 'peak' and 'shoulders' associated with multiple and single scattering, respectively. As with the static fluid setup, the preliminary experiments performed on the useful data to determine particle size. Experiments 69-92 (see Appendix III) can be categorized as preliminary experiments. From those experiments, two characteristics about the flowing fluid setup were discovered.

The first characteristic of the flowing system dealt with the procedure used to take data. For the static fluid setup, each time the tilt angle was altered during a tilt angle sweep experiment, the top detector was translated to maximize the intensity. Through the preliminary experiments, it was determined that translation of the top detector to maintain intensity was more effective than maximizing intensity. As can be seen in Fig. 28, a plot of radius versus tilt angle for both maximizing and maintaining intensities, the maintaining intensity plot reaches the acceptable radius range first and stays there for a


Figure 28: Radius versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.198 percent by weight and at a scattering angle of $112^{\circ}$. Data corresponds to Experiments 86 and 87 .
larger range of tilt angles. Due to possible statistical errors in the data collected to produce Fig. 28, the differences between maximizing and maintaining intensities seen in Fig. 28 cannot be the lone reason for making the decision to maintain intensities. Due to personal experience, the technique of maintaining intensities in the top detector was adopted for use in all flowing fluid experiments.

The second characteristic involved the alignment of the system. A major focus of this paper was the fact that alignment of a dynamic light scattering system is crucial to collecting reliable data. The preliminary flowing experiments were performed with the lens improperly focused. The 'waist' of the incident beam inside the sample (Chapter IV) was focused in a way that allowed only part of the beam to be in the detection volume. The misalignment was caused by the lens being too close to the sample. This positioning greatly shortened the amount of tilt angle covered by the shoulder. Figure 29 shows this effect by comparing two curves of $0.107 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.198 percent by weight. One curve shows the signal-to-noise ratio versus tilt angle for data collected before the lens was repositioned, while the second curve uses data taken after the lens was refocused. The second curve produced a tilt angle range four times that of the first curve.

With these changes in alignment and procedure made, the last nine experiments listed in Table 6 were performed at a scattering angle of $\theta=112^{\circ}$. As mentioned in Section 6.4, the purpose of these experiments was to study the effects of three parameters on determining particle size. These three parameters are as follows: flow velocity, particle size, and concentration. Plots of the signal-to-noise ratio and the radius versus


Figure 29: Signal-to-noise ratio versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.198 percent by weight and at a scattering angle of $112^{\circ}$. Data corresponds to Experiments 77 and 93.
tilt angle were created to illustrate the effect that each parameter had on particle measurements.

The flow velocity effect will be examined first. Three different particle sizes were used at various flow rates. The first particle size used in' a tilt angle sweep experiment was $0.107 \mu \mathrm{~m}$ PSL at a volume fraction of 0.198 percent by weight. Figure 30 shows a plot of the radius as a function of tilt angle for 3 flow rates: $0 \%, 50 \%$, and $100 \%$ (see Table 2 for corresponding velocities). From this plot, it can be seen that flow velocity affected the size computed for the multiple scattering area of the plot. As the velocity increased, the initial size given by the ALV-5000 program decreased. This also affected how much tilt was necessary to move into the single scattering area. These results were caused by the intermediate scattering wave vectors, $\mathbf{q}$, not being perpendicular to the velocity vector, $\mathbf{v}$. Due to this non-perpendicular alignment, the second term in the exponential of Eq. (3-14) was not suppressed and allowed the influence of the flow to affect the correct determination of particle size. In this way, the effect of velocity is similar to an increase in multiple scattering (compare Fig. 30. with Fig. 22 to see the effect of higher concentration).

In the peak region, the correlation function decays faster and therefore computations made from data within this region results in a smaller particle size. This effect can also be seen in Fig. 31, a plot of radius versus tilt angle for $0.203 \mu \mathrm{~m}$ PSL particles. Figure 31 shows a problem associated with the $0.203 \mu \mathrm{~m}$ PSL particles. The radii determined in the single scattering area by the ALV-5000 were slightly higher than the manufacturer's range of acceptance. At most, the determined radius was within 5 nm of the acceptance range. Therefore, it was decided that this error was acceptable.


Figure 30: Radius versus tilt angle for cross-correlation measurements of $0.107 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.198 percent by weight and at a scattering angle of $112^{\circ}$. Data corresponds to Experiments 93, 94, and 95.


Figure 31: Radius versus tilt angle for cross-correlation measurements of $0.203 \mu \mathrm{~m}$ PSL particles at a volume fraction of 0.200 percent by weight and at a scattering angle of $112^{\circ}$. Data corresponds to Experiments 111 and 112.

The effect caused by the second parameter, particle size, was examined next. Figure 32 shows the normalized radius versus tilt angle for both $0.107 \mu \mathrm{~m}$ and $0.203 \mu \mathrm{~m}$ PSL particles at volume fractions of 0.198 and 0.200 percent by weight, respectively. The comparisons shown in Figs. 32 and 33 were based on the volume fraction and not the optical thickness of the sample, which was different. Figure 33 compares the two particle sizes with respect to their signal-to-noise ratios. The data for Figs. 32 and 33 was collected at a flow rate of $100 \%$. For ease of comparison, the data collected from the ALV-5000 program for both particle sizes was normalized by dividing the data by the correct particle radius.

When comparing these two particle sizes, it can be seen that there is very little difference between the starting point of the two curves in Fig. 32. The peaks associated with both curves start at approximately the same value (0.65). Two major differences exist between the data used to produce the two curves in Fig. 32. The first difference is that, once the $0.107 \mu \mathrm{~m}$ PSL particles reach the acceptance range (the solid lines on Fig. 32) for their size, the radius determined by the data stays within the acceptance range for the remainder of the curve. The curve for the $0.203 \mu \mathrm{~m}$ particles does not stay within its acceptance range (the dashed lines on Fig. 32). The second difference is that the tilt angle range for the shoulder of the $0.107 \mu \mathrm{~m}$ particle curve is 4 mrads larger than the tilt angle range for the shoulder of the $0.203 \mu \mathrm{~m}$ particles. Figure 33 shows a greater effect of particle size. The curve for the $0.203 \mu \mathrm{~m}$ particles is much steeper and does not extend to as large of a tilt angle range as the curve for the $0.107 \mu \mathrm{~m}$ particles.

The final parameter studied was particle concentration. As with the static fluid setup, the data pertaining to this parameter was limited to Experiments 103-106. For a


Figure 32: Normalized radius versus tilt angle for cross-correlation measurements taken at a scattering angle of $112^{\circ}$ and flow rate of $100 \%$. The samples were both at a volume fraction of approximately 0.200 percent by weight. Data corresponds to Experiments 95 and 112.


Figure 33: Signal-to-noise ratio versus tilt angle for cross-correlation measurements taken at a scattering angle of $112^{\circ}$ and flow rate of $100 \%$. Data corresponds to Experiments 95 and 112.
particle size of $0.098 \mu \mathrm{~m}$, data from two volume fractions was available, 0.320 and 0.860 percent by weight. Similar results were found for the flowing system as for the static fluid data shown in Figs. 22 and 23. Figure 34 shows the signal-to-noise ratio versus tilt angle for the samples listed above at a flow rate of $100 \%$. As in Fig. 22, both curves in Fig. 34 peak at near the same signal-to-noise ratio, while the more concentrated sample shows a faster drop in its signal-to-noise ratio. This would give a narrower peak region if both shoulders were plotted. Additional effects of concentration can be seen in Fig. 35, where the radius is plotted as a function of tilt angle. The more concentrated sample gives a lower peak value for the radius of the particle, $\sim 31 \mathrm{~nm}$ as compared to ~ 33 nm .

This effect is not as great as that seen in Fig. 36, where the same effect of concentration is shown for a flow rate of $0 \%$. In Fig. 36, the radius given at the peak is $\sim 39 \mathrm{~nm}$ for the more concentrated sample as compared to $\sim 43 \mathrm{~nm}$ for the less concentrated sample. As in the static fluid sample, this difference in starting values is consistent with the less concentrated sample reaching the correct radius range at a smaller tilt angle than that of the more concentrated sample.

Results from the static fluid setup and flowing fluid setup are expanded upon in Sundaresan (1999). For the static fluid setup, the effects of concentration are examined for a sample of $0.107 \mu \mathrm{~m}$ PSL particles in a square test cell. For the flowing fluid setup, the effects of flow velocity, particle size, concentration, and cell rotation are covered in Sundaresan (1999) for scattering angles of $\theta=122^{\circ}$ and $\theta=135^{\circ}$.


Figure 34: Signal-to-noise ratio versus tilt angle for cross-correlation measurements taken from a sample of $0.098 \mu \mathrm{~m}$ PSL particles at a scattering angle of $112^{\circ}$ and flow rate of $100 \%$. Data corresponds to Experiments 104 and 106.


Figure 35: Radius versus tilt angle for cross-correlation measurements taken from a sample of $0.098 \mu \mathrm{~m}$ PSL particles at a scattering angle of $112^{\circ}$ and flow rate of $100 \%$. Data corresponds to Experiments 104 and 106.


Figure 36: Radius versus tilt angle for cross-correlation measurements taken from a sample of $0.098 \mu \mathrm{~m}$ PSL particles at a scattering angle of $112^{\circ}$ and flow rate of 0\%. Data corresponds to Experiments 103 and 105.

## CHAPTER VII

## CONCLUSIONS AND RECOMENDATIONS

### 7.1 Conclusions

The use of dynamic light scattering data to determine particle characteristics is applicable to several industrial operations. Industries are looking for a method to accurately measure particle size without interrupting their operations. These types of measurements are termed non-intrusive in-situ and allow measurements to occur during normal production or testing operations. The research described in this thesis focused on a light scattering system that was capable of determining particle size in dense flowing solutions.

The purpose of this thesis was to demonstrate the feasibility of integrating two independent suppression techniques to determine particle size. One technique dealt with the suppression of multiple scattering effects. The second technique was used to suppress flow effects. The experimental setup, alignment, and procedure used for suppression of multiple scattering effects were discussed first. The essential components necessary to suppress multiple scattering effects from the static fluid setup were then integrated into a setup designed to suppress flow effects. This combined system was then used to determine particle size. Using both the static fluid and the flowing fluid setups, several
parameters were studied to determine their effects on particle sizing. Those parameters include: scattering angle, sample concentration, particle size, flow rate, and alignment of the flowing setup test cell.

The first round of experiments involved the static fluid setup. The purpose of those experiments was twofold; the first purpose was to develop a familiarity with the suppression equipment and to develop the understanding of the suppression technique. The second purpose was to expand on the work done by Nobbmann et al. (1997). The static fluid experiments described within this thesis were conducted on $0.107 \mu \mathrm{~m}$ PSL particles at a variety of scattering angles ranging from $15^{\circ}$ to $120^{\circ}$ for volume fractions ranging from dilute to 0.5025 percent by weight. From the experiments performed on the static fluid setup, useful information about the characteristics of general dynamic light scattering and the data collected from it was discovered. Those characteristics include erratic behavior at small (forward) scattering angles and the system's sensitivity to misalignment. From these characteristics, a set procedure was adopted to deal with alignment of the system and for data collection. These characteristics as well as the others mentioned in Sections 6.2 and 6.3 led to the design of the flowing fluid setup.

The flowing fluid setup was designed to accomplish two forms of suppressions at once. First, the system utilized the same multiple scattering suppression procedure as that described in Chapter IV. Second, the system was aligned in a manner, which facilitated the suppression of flow effects. Various experiments were performed to demonstrate the suppression of flow effects and then the simultaneous suppression of multiple scattering and flow effects. The first set of experiments conducted with the flowing fluid setup demonstrated sensitivity to the system's alignment. The setup was
positioned with the scattering wave vector perpendicular to the velocity vector of the fluid. The test cell was then rotated about the axis centered through its height to show the effect of non-perpendicular alignment between those two vectors. These experiments were conducted for various flow rates on samples of $0.107 \mu \mathrm{~m}$ and $0.204 \mu \mathrm{~m}$ PSL particles in a dilute (single scattering) solution. It was determined from these experiments that misalignment between $\pm 5^{\circ}$ had little effect on the data collected. It was also determined that misalignment had a greater impact on the data used to determine the larger sized particles.

The second types of experiments performed with the flowing fluid setup were tilt angle mapping experiments (Sections 5.5 .2 and 6.5.2). For cases of $0 \%$ flow, these experiments were similar to those performed with the static fluid setup of Chapter IV. When flow was added, the tilt angle experiments demonstrated the system's ability to suppress multiple scattering within flowing suspensions. Experiments were performed at a scattering angle of $112^{\circ}$ on $0.098 \mu \mathrm{~m}, 0.107 \mu \mathrm{~m}$, and $0.203 \mu \mathrm{~m}$ PSL particles at various volume fractions and for various flow rates. Several trends were discovered during these experiments concerning the nature of particle sizing in a flowing fluid. The first trend discovered dealt with the effect of fluid velocity. For higher velocities, larger tilt angles were required in order to suppress multiple scattering effects. This effect was similar to the effect caused by an increase in concentration. This effect held true for the flowing fluid setup as well as the static fluid setup. A final trend associated with the flowing fluid setup was that larger particle sizes ( $\geq 0.2 \mu \mathrm{~m}$ ) appear to be more difficult to correctly size.

### 7.2 Recommendations

The theories utilized by the research described within this thesis are relatively new, and, to the knowledge of those involved in the experiments, have never been used together to determine particle characteristics. This thesis only discusses the use of the suppression techniques to determine one particle characteristic, particle size or diameter. Although several parameters (particle size, concentration, scattering angle, and flow velocity) were studied in order to evaluate their effects on the determination of particle size, many more parameters exist that were not researched. Parameters such as more scattering angles so that bigger particle sizes can be used, polarization, index of refraction, scattering/absorbing particles, higher velocity, etc., should be investigated. Also, more real world situations should be covered, such as solutions involving particle distributions instead of just monodisperse solutions. The current system must be modified if some of these recommendations are to be realized.

Some realistic minor modifications to the current setup are necessary to further the research. The first change would be to replace the current rear fiber mount (Figs. 9 and 10) with one similar in design to that of the top mount (Figs. 7 and 8). The ability to record the starting position as well as the position of alignment (through the micrometers on the top mount) would be of great benefit in the alignment process for the back mount.

The second change would be to modify the body of the goniometer (specifically the sample stand of Fig. 12) to allow the arms to achieve a wider range of angular travel. The current setup allows for each arm to travel to a maximum angle of $48^{\circ}$, which corresponds to a scattering angle of $112^{\circ}$. Due to the Rayleigh-Gans form factors
associated with larger particles [Sundaresan (1999)] this scattering angle is not large enough to produce useable intensity levels for particles larger than the $0.2 \mu \mathrm{~m}$ particles. A modification of the setup to allow detector and laser arm to move to angles of $\delta=70^{\circ}$ (corresponding to a scattering angle of $\theta=90^{\circ}$ from Table 3) would improve the chance of achieving useable intensities. Although this increase should be enough, I believe the goniometer should be redesigned to allow for enough movement from the arms to be positioned across from each other ( $\alpha_{\mathrm{L}}=\alpha_{\mathrm{D}}=90^{\circ}$ ). This positioning would allow the detectors direct view of the incident beam for alignment. This is similar to the alignment discussed in Section 4.3 for the static fluid setup. Doing so would allow for direct alignment of the detectors with respect to the incident beam.

A device utilizing a micrometer for precise cell rotations needs to be added to the current flowing fluid setup. The current system employs angles marked on the rotation disk, which carries the test cell holder. This method of positioning the test cell holder only allows for estimated rotation angles. The new system would allow exact cell rotation measurements to be conducted to determine a better-defined range of acceptable cell rotation angles.

The final recommendations for the redesign of the flowing fluid setup deal with the test cell holder. To use the current setup, the test cell is tilted approximately $5^{\circ}$ from vertical in the holder to redirect reflections away from the detector housing. A better design would be to allow the entire top section (Fig. 13) to be tilted. A micrometer could be added to the rear of the device to allow for measuring of how much of a tilt is necessary in order to keep direct reflections away from the detectors. This information
would also allow for the calculation of the actual depth of the incident laser beam into the cell.

A second recommendation for improving the current setup deals with the method employed to block reflections from reaching the detector. Strips of electrical tape are currently used to accomplish this task. The tape has a tendency to stretch during application and come loose from the holder during an experiment. This allows reflections to interfere with the measurements. A second problem with the use of electrical tape is the repeatability of replacing the tape on the cell holder in the same spot each time. An improvement would be the addition of slits to the test cell holder, which would be a more reliable method to block the reflections.

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## APPENDICES

## APPENDIX I

## Equipment List

1. 20 mW Helium Neon laser manufactured by Uniphase with a wavelength of 632.5 nm, Model No. 1135P.
2. $\quad 100 \mathrm{~mW}$ Neodymium-Yttrium-Silver laser manufactured by Adlas with a wavelength of 532.5 nm , Model No. DPY315II.
3. Optical Power Meter manufactured by Newport Inc., Model No. 840 with wand Model No. 818-ST.
4. Test tube manufactured by Fisher Scientific from borosilicate glass with dimensions of $10 \mathrm{~mm} \times 75 \mathrm{~mm}$, Catalog No. 14-961-25.
5. Fiber Optic Cables manufactured by Oz Optics LTD. Part No. LPC-02-532-4/125-P-0.7-3.2GR-30-1-3-3.
6. Translation stages, Model No. 426a, manufactured by Newport and equipped with SM-25 micrometers were used for the top fiber mount. Dimensions of the stages were $89 \mathrm{~mm} \times 89 \mathrm{~mm} \times 25.4 \mathrm{~mm}$ with a 50.8 mm diameter hole in the center. The aluminum tilt plates were produced by the OSU Chemistry/Physics Machine Shop and had dimensions of $127 \mathrm{~mm} \times 134 \mathrm{~mm} \times 10 \mathrm{~mm}$.
7. The back fiber mount was manufactured by the OSU Chemistry/Physics Machine Shop and had two main pieces, the mount plate and the piece that carried the set screws. The dimensions of the mount plate were $101.6 \mathrm{~mm} \times 101.6 \mathrm{~mm} \times 12.7$ mm with a 34.5 mm diameter hole in the center. The dimensions of the second piece were $101.6 \mathrm{~mm} \times 101.6 \mathrm{~mm} \times 22.9 \mathrm{~mm}$ with a 34.5 mm diameter hole in the center.
8. Two power supplies produced by Global Specialties, Model Nos. 1310 and 1302.
9. 632.5 nm wavelength specific beamsplitter from Newport, Model No. 05BC16NP.4, for the non-flowing setup.
10. ALV-5000 Multiple Tau Digital Correlator by ALV-Laser Vertriebsgesellschaft m.b.H Germany.
11. E-pure deionizer, Model No. D4641, manufactured by Barnstead and Thermolyne.
12. Electronic lab scale, Model No. 31205, by Sartorious.
13. Multi-band ( $400-700 \mathrm{~nm}$ ), nonpolarizing beamsplitter, Model No. 05FC16-PB.3, by Newport for the flowing setup.
14. Filter holder manufactured by Newport, Model No. FH-1.
15. Shuttle pump manufactured by Instech Labs, Model No. S20P.
16. $6 \mathrm{~mm} \times 8 \mathrm{~mm} 30.5 \mathrm{~cm}$ rectangular test cell with a 0.9 mm wall thickness, manufactured by Wilmad Glass from clear fused quartz, Catalog No. WQR-0608.
17. $1 / 8^{\prime \prime} \times 1 / 16^{\prime \prime}$ tubing manufactured by Tygon, S-50-HL, Class VI.
18. Water bath container formed from 64-stock glass tubing in Oklahoma State Glass Shop with dimensions of 6.35 cm outside diameter having a wall thickness of 2.4 mm and a height of 8.89 cm .
19. Photomultiplier tubes manufactured by Thorn EMI Electron Tubes Inc., Model No. EBA-805.
20. Holding tanks with dimensions of $215 \mathrm{~mm} \times 69.5 \mathrm{~mm} \times 49.5 \mathrm{~mm}$ were manufactured from Plexiglas by the Oklahoma State University Physics Machine Shop. The side walls were 12 mm thick, while the lid and base were 9 mm thick. The lid was sealed to the base by 21 screws and a rubber gasket.
21. Core samples of polystyrene latex (PSL) particles from Duke Scientific:
A. $\quad 0.107 \mu \mathrm{~m}$ diameter; 10 percent solids by weight; $5.6 \%$ Coefficient of Variation, Catalog No. 5010A-Lot No. 16456
B. $\quad 0.098 \mu \mathrm{~m}$ diameter; 10 percent solids by weight; $6.2 \%$ C.V., Catalog No. 5010A-Lot No. 20259
C. $0.203 \mu \mathrm{~m}$ diameter; 10 percent solids by weight; $2.1 \%$ C.V., Catalog No. 5020A-Lot No. 20500
D. $\quad 0.204 \mu \mathrm{~m} \pm 6 \mathrm{~nm}$ diameter; 1 percent solids by weight; Catalog No. 3200A-Lot No. 20613
E. $\quad 0.300 \mu \mathrm{~m} \pm 5 \mathrm{~nm}$ diameter; 1 percent solids by weight; Catalog No. 3300A-Lot No. 20286.
22. Parafilm laboratory film manufactured by American National Can was used to seal test tubes.

## APPENDIX II

## Static Fluid: Experimental Data

Table 4: Summary of the static fluid experiments discussed in Sundaresan (1999) and in this thesis.

| Experiment <br> Number | Scattering <br> Angle <br> $($ deg $)$ | Test Cell <br> Type | Volume <br> Fraction (\% <br> by weight) | Particle <br> Diameter <br> $(\mu \mathrm{m})$ |
| :---: | :---: | :---: | :---: | :---: |
| 32 | 90 | Square | 0.3239 | 0.107 |
| 33 | 90 | Square | 0.3239 | 0.107 |
| 34 | 90 | Square | 0.1536 | 0.107 |
| 35 | 90 | Square | 0.1536 | 0.107 |
| 36 | 90 | Square | 0.4285 | 0.107 |
| 37 | 90 | Square | 0.4285 | 0.107 |
| 38 | 90 | Circular | 0.1330 | 0.107 |
| 39 | 90 | Circular | 0.3201 | 0.107 |
| 40 | 30 | Circular | 0.3201 | 0.107 |
| 41 | 30 | Circular | 0.3201 | 0.107 |
| 42 | Sweep | Circular | Single | 0.107 |
| 43 | 60 | Circular | 0.3271 | 0.107 |
| 44 | 60 | Circular | 0.1545 | 0.107 |
| 45 | Sweep | Circular | Single | 0.107 |
| 46 | 120 | Circular | 0.5025 | 0.107 |
| 47 | 120 | Circular | 0.3795 | 0.107 |
| 48 | 45 | Circular | 0.3795 | 0.107 |

Note: Detailed data for Experiments 32-37 appear in Sundaresan (1999), and detailed data for Experiments 38-48 appear in Table 5.

Table 5: Detailed description of experiments 38-48 described in summary Table 4.

| Exp 38 | Circular Cell Front Tilt = $\mathbf{1 6 . 9 1} \mathbf{~ d i v}$. |  | $0.107 \mu \mathrm{~m}$ PSL |  | V.F. $=0.133 \%$ |  | $\theta=90 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div) | Side Translation (div) | Tilt Angle (mrad) | Intensit <br> Ch 0 | $(\mathrm{kHz})$ <br> Ch 1 | $Y$ Intercept | Decay Rate (1/msec) | Radius ( nm ) | Normalized Second Cumulant | Time (sec) |
| 16.68 | 11.84 | -7.425 | 221.485 | 167.417 | 0.0350 | 1.23 | 62.1 | -0.2900 | 300 |
| 16.70 | 11.92 | -7.200 | 222.044 | 165.552 | 0.0396 | 1.43 | 53.6 | -0.0019 | 300 |
| 16.72 | 12.00 | -6.975 | 222.858 | 164.413 | 0.0460 | 1.47 | 52.3 | 0.0280 | 300 |
| 16.74 | 12.10 | -6.750 | 223.251 | 160.046 | 0.0559 | 1.50 | 51.0 | 0.0091 | 300 |
| 16.76 | 12.19 | -6.525 | 223.600 | 160.302 | 0.0674 | 1.45 | 52.9 | 0.0560 | 300 |
| 16.78 | 12.29 | -6.300 | 224.735 | 162.080 | 0.0748 | 1.45 | 52.9 | 0.0061 | 300 |
| 16.80 | 12.41 | -6.075 | 227.312 | 165.040 | 0.0862 | 1.49 | 51.4 | 0.0170 | 300 |
| 16.82 | 12.48 | -5.850 | 215.699 | 170.612 | 0.0964 | 1.48 | 53.3 | -0.0061 | 300 |
| 16.84 | 12.59 | -5.625 | 216.424 | 168.351 | 0.1130 | 1.49 | 52.7 | 0.0330 | 120 |
| 16.86 | 12.67 | -5.400 | 217.390 | 165.495 | 0.1240 | 1.44 | 54.5 | 0.0170 | 300 |
| 16.88 | 12.77 | -5.175 | 217.082 | 162.023 | 0.1350 | 1.48 | 53.3 | 0.0320 | 120 |
| 16.90 | 12.93 | -4.950 | 219.400 | 159.414 | 0.1640 | 1.43 | 55.3 | -0.0210 | 120 |
| 16.94 | 13.15 | -4.500 | 219.450 | 170.152 | 0.2280 | 1.47 | 53.6 | 0.0150 | 120 |
| 16.96 | 13.27 | -4.275 | 219.451 | 171.163 | 0.2480 | 1.47 | 53.7 | 0.0200 | 120 |
| 16.98 | 13.35 | -4.050 | 220.289 | 166.459 | 0.2590 | 1.43 | 55.0 | 0.0064 | 120 |
| 17.02 | 13.49 | -3.600 | 219.408 | 168.211 | 0.2820 | 1.46 | 54.1 | 0.0170 | 120 |
| 17.06 | 13.70 | -3.150 | 219.829 | 163.177 | 0.3300 | 1.48 | 53.2 | 0.0300 | 120 |
| 17.10 | 13.90 | -2.700 | 213.072 | 170.496 | 0.3880 | 1.45 | 54.3 | 0.0016 | 120 |
| 17.14 | 14.14 | -2.250 | 222.806 | 171.898 | 0.4460 | 1.47 | 53.6 | 0.0120 | 120 |
| 17.18 | 14.34 | -1.800 | 215.328 | 165.455 | 0.4980 | 1.48 | 53.3 | 0.0370 | 120 |
| 17.20 | 14.40 | -1.575 | 222.773 | 174.563 | 0.5180 | 1.47 | 53.6 | 0.0200 | 120 |
| 17.22 | 14.42 | -1.350 | 214.396 | 166.929 | 0.5270 | 1.48 | 53.2 | 0.0290 | 120 |
| 17.24 | 14.47 | -1.125 | 222.837 | 163.309 | 0.5450 | 1.48 | 53.1 | 0.0240 | 120 |
| 17.26 | 14.61 | -0.900 | 221.006 | 167.366 | 0.6190 | 1.52 | 51.7 | 0.0270 | 120 |
| 17.27 | 14.68 | -0.788 | 223.399 | 170.716 | 0.6650 | 1.56 | 50.6 | 0.0400 | 120 |
| 17.28 | 14.73 | -0.675 | 229.111 | 174.544 | 0.7130 | 1.61 | 49.0 | 0.0490 | 120 |
| 17.29 | 14.79 | -0.563 | 223.647 | 171.153 | 0.7790 | 1.62 | 48.6 | 0.0630 | 120 |
| 17.30 | 14.84 | -0.450 | 231.247 | 178.654 | 0.8370 | 1.67 | 47.3 | 0.0750 | 120 |
| 17.31 | 14.90 | -0.338 | 223.532 | 169.284 | 0.8910 | 1.68 | 46.9 | 0.0720 | 120 |
| 17.32 | 14.96 | -0.225 | 225.221 | 169.094 | 0.9280 | 1.68 | 46.8 | 0.0810 | 120 |
| 17.33 | 15.00 | -0.113 | 238.175 | 178.087 | 0.9340 | 1.70 | 46.4 | 0.0720 | 120 |
| 17.34 | 15.03 | 0.000 | 223.847 | 207.015 | 0.9370 | 1.67 | 46.0 | 0.0880 | 120 |
| 17.35 | 15.10 | 0.113 | 223.897 | 211.382 | 0.8880 | 1.66 | 46.1 | 0.0740 | 120 |
| 17.36 | 15.16 | 0.225 | 225.348 | 214.338 | 0.8490 | 1.63 | 47.2 | 0.0720 | 120 |
| 17.37 | 15.22 | 0.338 | 224.310 | 214.517 | 0.7970 | 1.60 | 47.9 | 0.0620 | 120 |
| 17.38 | 15.29 | 0.450 | 224.986 | 216.072 | 0.7480 | 1.55 | 49.5 | 0.0560 | 120 |
| 17.39 | 15.33 | 0.563 | 213.781 | 195.927 | 0.6850 | 1.55 | 49.4 | 0.0550 | 120 |
| 17.40 | 15.42 | 0.675 | 229.206 | 203.753 | 0.6380 | 1.53 | 50.3 | 0.0460 | 120 |
| 17.42 | 15.48 | 0.900 | 220.696 | 194.896 | 0.6030 | 1.49 | 51.3 | 0.0240 | 120 |
| 17.44 | 15.57 | 1.125 | 217.499 | 186.720 | 0.5570 | 1.47 | 52.2 | 0.0210 | 120 |
| 17.46 | 15.68 | 1.350 | 219.748 | 187.649 | 0.5170 | 1.45 | 52.7 | 0.0047 | 120 |
| 17.48 | 15.79 | 1.575 | 219.727 | 189.240 | 0.4900 | 1.46 | 52.6 | 0.0330 | 120 |
| 17.50 | 15.90 | 1.800 | 210.688 | 191.171 | 0.4670 | 1.44 | 53.3 | 0.0280 | 120 |
| 17.52 | 15.97 | 2.025 | 223.077 | 194.545 | 0.4540 | 1.47 | 52.1 | 0.0520 | 120 |
| 17.54 | 16.12 | 2.250 | 223.289 | 188.619 | 0.4210 | 1.45 | 52.8 | 0.0170 | 120 |
| 17.56 | 16.21 | 2.475 | 224.628 | 185.381 | 0.3940 | 1.45 | 52.7 | -0.0120 | 120 |
| 17.60 | 16.41 | 2.925 | 227.158 | 185.976 | 0.3550 | 1.43 | 53.7 | 0.0150 | 120 |
| 17.64 | 16.59 | 3.375 | 228.177 | 192.643 | 0.2970 | 1.44 | 53.3 | 0.0270 | 120 |
| 17.68 | 16.74 | 3.825 | 230.986 | 188.780 | 0.2590 | 1.46 | 52.6 | 0.0260 | 120 |
| 17.72 | 16.94 | 4.275 | 229.023 | 183.630 | 0.2190 | 1.48 | 51.9 | 0.0200 | 120 |
| 17.76 | 17.14 | 4.725 | 230.248 | 186.504 | 0.1860 | 1.51 | 50.9 | 0.0490 | 120 |
| 17.78 | 17.26 | 4.950 | 231.853 | 188.501 | 0.1650 | 1.48 | 51.6 | 0.0630 | 120 |
| 17.80 | 17.37 | 5.175 | 230.712 | 187.176 | 0.1430 | 1.46 | 52.4 | 0.0250 | 120 |
| 17.82 | 17.52 | 5.400 | 232.443 | 181.634 | 0.1180 | 1.41 | 54.4 | 0.0190 | 120 |
| 17.84 | 17.69 | 5.625 | 232.644 | 174.714 | 0.0961 | 1.45 | 53.0 | 0.0230 | 120 |
| 17.86 | 17.75 | 5.850 | 232.546 | 181.706 | 0.0883 | 1.46 | 52.1 | -0.0360 | 120 |
| 17.88 | 17.86 | 6.075 | 231.698 | 185.053 | 0.0769 | 1.48 | 51.9 | 0.0780 | 120 |
| 17.90 | 17.92 | 6.300 | 231.402 | 193.843 | 0.0718 | 1.49 | 51.3 | 0.0260 | 120 |
| 17.92 | 17.99 | 6.525 | 228.680 | 204.385 | 0.0631 | 1.48 | 51.0 | -0.0110 | 300 |
| 17.94 | 18.37 | 6.750 | 225.741 | 226.926 | 0.0371 | 1.53 | 50.0 | 0.0400 | 300 |
| 17.96 | 18.38 | 6.975 | 223.530 | 237.697 | 0.0321 | 1.57 | 48.7 | -0.1100 | 300 |


| Exp 39 | Circular Cell Front Tilt $=\mathbf{1 6 . 9 2}$ div. |  | $0.107 \mu \mathrm{~m} \text { PSL }$ |  | $\text { V.F. }=0.3201 \%$ |  | $\theta=90 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div) | Side Translation (div) | Tilt Angle (mrad) | Intens ch 0 | CHz) | $Y$ - <br> Intercept | Decay Rate (1/msec) | Radius (nm) | Normalized Second Cumulant | Time (sec) |
| 16.78 | 12.84 | -6.3000 | 190.974 | 138.293 | 0.0320 | 2.61 | 47.6 | 0.1500 | 300 |
| 16.85 | 12.65 | -5.5125 | 191.773 | 136.145 | 0.0402 | 2.50 | 51.1 | 0.0110 | 300 |
| 16.88 | 12.80 | -5.1750 | 193.137 | 133.760 | 0.0499 | 2.35 | 56.9 | -0.0500 | 300 |
| 16.91 | 12.99 | -4.8375 | 194.167 | 135.453 | 0.0642 | 2.45 | 52.9 | -0.0530 | 300 |
| 16.94 | 13.12 | -4.5000 | 196.272 | 143.100 | 0.0854 | 2.45 | 53.0 | -0.0600 | 300 |
| 16.97 | 13.36 | -4.1625 | 195.062 | 141.545 | 0.1060 | 2.52 | 50.6 | 0.0220 | 120 |
| 17.00 | 13.48 | -3.8250 | 194.487 | 143.382 | 0.1170 | 2.50 | 51.1 | -0.0580 | 120 |
| 17.03 | 13.61 | -3.4875 | 194.066 | 143.984 | 0.1310 | 2.46 | 52.5 | -0.0430 | 120 |
| 17.05 | 13.69 | -3.2625 | 193.207 | 143.644 | 0.1360 | 2.46 | 52.7 | -0.0270 | 120 |
| 17.07 | 13.75 | -3.0375 | 191.789 | 143.692 | 0.1470 | 2.51 | 50.7 | -0.0200 | 120 |
| 17.09 | 13.85 | -2.8125 | 192.159 | 142.517 | 0.1590 | 2.48 | 51.9 | -0.0420 | 120 |
| 17.11 | 13.91 | -2.5875 | 191.361 | 140.449 | 0.1710 | 1.51 | 50.7 | 0.0170 | 120 |
| 17.13 | 14.00 | -2.3625 | 198.755 | 146.124 | 0.1830 | 1.48 | 51.8 | 0.0590 | 120 |
| 17.15 | 14.11 | -2.1375 | 191.943 | 140.886 | 0.1990 | 1.48 | 51.7 | -0.0320 | 120 |
| 17.17 | 14.21 | -1.9125 | 191.921 | 142.451 | 0.2150 | 1.50 | 51.1 | 0.0180 | 120 |
| 17.18 | 14.26 | -1.8000 | 192.468 | 142.909 | 0.2260 | 1.54 | 49.9 | 0.0570 | 120 |
| 17.19 | 14.31 | -1.6875 | 192.696 | 143.023 | 0.2340 | 1.46 | 52.6 | -0.0230 | 120 |
| 17.20 | 14.40 | -1.5750 | 193.111 | 144.261 | 0.2580 | 1.59 | 48.1 | 0.0620 | 120 |
| 17.21 | 14.45 | -1.4625 | 195.315 | 140.741 | 0.2760 | 1.60 | 48.0 | 0.0530 | 120 |
| 17.22 | 14.50 | -1.3500 | 196.486 | 144.521 | 0.3020 | 1.65 | 46.5 | 0.0530 | 120 |
| 17.23 | 14.54 | -1.2375 | 194.597 | 144.546 | 0.3220 | 1.66 | 46.3 | 0.0830 | 120 |
| 17.24 | 14.60 | -1.1250 | 199.877 | 144.530 | 0.3610 | 1.71 | 44.9 | 0.0870 | 120 |
| 17.25 | 14.65 | -1.0125 | 200.553 | 144.123 | 0.4040 | 1.84 | 41.7 | 0.1100 | 120 |
| 17.26 | 14.69 | -0.9000 | 201.878 | 144.548 | 0.4640 | 1.92 | 40.0 | 0.1200 | 120 |
| 17.27 | 14.74 | -0.7875 | 202.547 | 143.611 | 0.5240 | 1.99 | 38.5 | 0.1500 | 120 |
| 17.28 | 14.79 | -0.6750 | 202.694 | 143.524 | 0.6180 | 2.12 | 36.2 | 0.1500 | 120 |
| 17.29 | 14.83 | -0.5625 | 202.290 | 142.237 | 0.7020 | 2.25 | 34.0 | 0.1700 | 120 |
| 17.30 | 14.89 | -0.4500 | 201.857 | 142.968 | 0.8050 | 2.34 | 32.8 | 0.1700 | 120 |
| 17.31 | 14.93 | -0.3375 | 199.960 | 140.740 | 0.8820 | 2.39 | 32.1 | 0.1800 | 120 |
| 17.32 | 14.97 | -0.2250 | 200.994 | 144.334 | 0.9250 | 2.41 | 31.8 | 0.1900 | 120 |
| 17.33 | 15.00 | -0.1125 | 199.800 | 140.067 | 0.9280 | 2.43 | 31.5 | 0.1800 | 120 |
| 17.34 | 15.03 | 0.0000 | 197.733 | 137.942 | 0.9060 | 2.43 | 31.6 | 0.1800 | 120 |
| 17.35 | 15.13 | 0.1125 | 183.733 | 131.087 | 0.7710 | 2.31 | 33.2 | 0.1800 | 120 |
| 17.36 | 15.20 | 0.2250 | 189.266 | 135.814 | 0.6540 | 2.19 | 35.0 | 0.1700 | 120 |
| 17.37 | 15.26 | 0.3375 | 190.043 | 137.176 | 0.5620 | 2.05 | 37.3 | 0.1400 | 120 |
| 17.38 | 15.31 | 0.4500 | 201.283 | 144.362 | 0.4840 | 1.98 | 39.2 | 0.1400 | 120 |
| 17.39 | 15.34 | 0.5625 | 201.671 | 143.886 | 0.4670 | 1.88 | 40.8 | 0.1200 | 120 |
| 17.40 | 15.40 | 0.6750 | 201.879 | 143.758 | 0.4100 | 1.79 | 42.8 | 0.1100 | 120 |
| 17.41 | 15.46 | 0.7875 | 200.451 | 141.430 | 0.3660 | 1.72 | 44.7 | 0.0920 | 120 |
| 17.42 | 15.53 | 0.9000 | 200.537 | 138.728 | 0.3260 | 1.67 | 46.0 | 0.0690 | 120 |
| 17.43 | 15.56 | 1.0125 | 200.330 | 137.650 | 0.3070 | 1.63 | 47.2 | 0.0640 | 120 |
| 17.44 | 15.63 | 1.1250 | 199.679 | 135.287 | 0.2830 | 1.59 | 48.1 | 0.0720 | 120 |
| 17.45 | 15.72 | 1.2375 | 199.200 | 133.505 | 0.2600 | 1.52 | 50.6 | 0.0360 | 120 |
| 17.46 | 15.80 | 1.3500 | 197.722 | 131.054 | 0.2400 | 1.52 | 50.6 | 0.0430 | 120 |
| 17.48 | 15.90 | 1.5750 | 196.559 | 131.539 | 0.2290 | 1.51 | 50.7 | 0.0110 | 120 |
| 17.50 | 15.93 | 1.8000 | 196.285 | 133.635 | 0.2240 | 1.50 | 51.1 | 0.0320 | 120 |
| 17.52 | 16.07 | 2.0250 | 195.547 | 129.260 | 0.2020 | 1.48 | 51.8 | -0.0100 | 120 |
| 17.54 | 16.17 | 2.2500 | 193.781 | 128.697 | 0.1940 | 1.52 | 50.4 | 0.0500 | 120 |
| 17.56 | 16.24 | 2.4750 | 191.672 | 129.836 | 0.1820 | 1.50 | 51.2 | 0.0067 | 120 |
| 17.58 | 16.33 | 2.7000 | 191.585 | 131.905 | 0.1740 | 1.52 | 50.5 | -0.0160 | 120 |
| 17.60 | 16.40 | 2.9250 | 192.785 | 133.651 | 0.1600 | 1.49 | 51.3 | 0.0660 | 120 |
| 17.62 | 16.47 | 3.1500 | 191.442 | 133.943 | 0.1490 | 1.52 | 50.6 | 0.0390 | 120 |
| 17.64 | 16.60 | 3.3750 | 190.962 | 141.031 | 0.1310 | 1.46 | 52.7 | -0.0013 | 120 |
| 17.66 | 16.71 | 3.6000 | 189.416 | 145.103 | 0.1170 | 1.52 | 50.6 | -0.0360 | 180 |
| 17.68 | 16.84 | 3.8250 | 188.678 | 147.104 | 0.1010 | 1.49 | 51.3 | 0.0650 | 180 |
| 17.72 | 17.07 | 4.2750 | 194.181 | 131.148 | 0.0803 | 1.43 | 53.6 | -0.0027 | 180 |
| 17.76 | 17.28 | 4.7250 | 190.555 | 129.179 | 0.0630 | 1.49 | 51.5 | 0.0290 | 180 |
| 17.80 | 17.44 | 5.1750 | 184.284 | 135.242 | 0.0475 | 1.38 | 55.6 | 0.1100 | 180 |
| 17.84 | 17.65 | 5.6250 | 183.606 | 134.728 | 0.0342 | 1.51 | 50.8 | 0.0660 | 180 |
| 17.88 | 17.77 | 6.0750 | 182.395 | 133.023 | 0.0260 | 1.55 | 49.6 | 0.0910 | 180 |

Exp 40
Circular Cell
Front Tilt = $\mathbf{1 6 . 9 2}$ div.

| Rear Tilt <br> (div) | Side <br> Translation <br> (div) | Tilt Angle <br> (mrad) |
| :---: | :---: | :---: |
| 17.33 | 15.00 | 0.0000 |
| 17.32 | 14.97 | -0.1125 |
| 17.31 | 14.94 | -0.2250 |
| 17.30 | 14.90 | -0.3375 |
| 17.29 | 14.83 | -0.4500 |
| 17.28 | 14.79 | -0.5625 |
| 17.27 | 14.76 | -0.6750 |
| 17.26 | 14.70 | -0.7875 |
| 17.25 | 14.65 | -0.9000 |
| 17.24 | 14.61 | -1.0125 |
| 17.23 | 14.55 | -1.1250 |
| 17.22 | 14.49 | -1.2375 |
| 17.21 | 14.43 | -1.3500 |
| 17.20 | 14.38 | -1.4625 |
| 17.19 | 14.34 | -1.5750 |
| 17.18 | 14.28 | -1.6875 |
| 17.17 | 14.22 | -1.8000 |
| 17.16 | 14.18 | -1.9125 |
| 17.15 | 14.13 | -2.0250 |
| 17.14 | 14.08 | -2.1375 |
| 17.13 | 14.04 | -2.2500 |
| 17.12 | 13.99 | -2.3625 |
| 17.11 | 13.86 | -2.4750 |
| 17.10 | 13.80 | -2.5875 |
| 17.09 | 13.76 | -2.7000 |
| 17.08 | 13.72 | -2.8125 |
| 17.07 | 13.66 | -2.9250 |

$0.107 \mu \mathrm{~m}$ PSL $\quad$ V.F. $=0.3201 \% \quad \theta=30 \mathrm{deg}$
Rear Translation $=11.23$ div.

| Intensity (kHz) |  | Intercept | Decay Rate$(1 / \mathrm{msec})$ | Radius (nm) | Normalized Second Cumulant | $\begin{aligned} & \text { Time } \\ & \text { (sec) } \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ch 0 | Ch 1 |  |  |  |  |  |
| 216.628 | 160.471 | 0.9060 | 0.699 | 14.7 | 0.2800 | 120 |
| 214.462 | 161.398 | 0.9020 | 0.697 | 14.7 | 0.2800 | 120 |
| 214.538 | 167.347 | 0.8710 | 0.677 | 15.2 | 0.2800 | 120 |
| 215.131 | 166.873 | 0.8180 | 0.646 | 15.9 | 0.2800 | 120 |
| 216.336 | 161.059 | 0.7250 | 0.596 | 17.2 | 0.2900 | 120 |
| 212.607 | 165.244 | 0.6500 | 0.530 | 19.4 | 0.2800 | 120 |
| 214.100 | 169.244 | 0.5870 | 0.500 | 20.6 | 0.2800 | 120 |
| 211.671 | 168.166 | 0.5110 | 0.411 | 25.0 | 0.2700 | 120 |
| 213.418 | 169.385 | 0.4570 | 0.372 | 27.4 | 0.2600 | 120 |
| 216.881 | 168.940 | 0.4330 | 0.330 | 31.1 | 0.2300 | 120 |
| 218.751 | 170.740 | 0.4000 | 0.291 | 35.3 | 0.1900 | 120 |
| 219.905 | 168.390 | 0.3720 | 0.289 | 35.6 | 0.1800 | 120 |
| 221.660 | 166.660 | 0.3430 | 0.280 | 36.6 | 0.1600 | 120 |
| 218.276 | 167.876 | 0.3320 | 0.270 | 38.0 | 0.1600 | 120 |
| 218.327 | 169.400 | 0.3050 | 0.279 | 36.9 | 0.1200 | 120 |
| 220.253 | 167.796 | 0.2980 | 0.267 | 38.5 | 0.1500 | 120 |
| 221.263 | 168.496 | 0.2880 | 0.262 | 39.2 | 0.1500 | 120 |
| 209.687 | 180.986 | 0.2780 | 0.261 | 39.4 | 0.1400 | 120 |
| 207.231 | 171.480 | 0.2670 | 0.261 | 39.3 | 0.1500 | 120 |
| 245.728 | 168.922 | 0.2530 | 0.257 | 40.0 | 0.0910 | 120 |
| 235.476 | 183.883 | 0.2430 | 0.253 | 40.7 | 0.0740 | 120 |
| 211.871 | 207.831 | 0.2290 | 0.246 | 42.5 | 0.0840 | 120 |
| 236.293 | 183.982 | 0.2000 | 0.247 | 42.2 | 0.0310 | 240 |
| 238.336 | 180.935 | 0.1890 | 0.252 | 41.6 | 0.0920 | 300 |
| 238.444 | 181.848 | 0.1800 | 0.239 | 43.7 | 0.0480 | 300 |
| 237.080 | 183.789 | 0.1680 | 0.252 | 41.5 | 0.0850 | 300 |
| 236.960 | 184.653 | 0.1550 | 0.240 | 43.5 | 0.0640 | 300 |

Exp 41
Circular Cell
Front Tilt = $\mathbf{1 6 . 9 2}$ div. Rear Till
(div)

### 17.33

17.32
17.31
17.30
17.29
$17.28 \quad 14.85 \quad-0.5625$
$17.27 \quad 14.82 \quad-0.6750$
$\begin{array}{lll}17.26 & 14.79 & -0.7875 \\ 17.25 & 14.76 & 0.9000\end{array}$
$\begin{array}{lll}17.25 & 14.76 & -0.9000 \\ 17.24 & 14.71 & -1.0125\end{array}$

| 17.23 | 14.69 | -1.1250 |
| :--- | :--- | :--- |
| 17.22 | 14.62 | -1.2375 |

$\begin{array}{lll}17.21 & 14.54 & -1.3500 \\ 17.20 & 14.50 & -1.4625\end{array}$
$\begin{array}{lll}17.19 & 14.45 & -1.5750 \\ 17.18 & 14.38 & -1.6875\end{array}$
$\begin{array}{lll}17.17 & 14.31 & -1.8000\end{array}$
$\begin{array}{lll}17.16 & 14.26 & -1.9125 \\ 17.15 & 14.16 & -2.0250\end{array}$
$\begin{array}{lll}17.14 & 14.08 & -2.1375 \\ 17.13 & 14.02 & -2.2500\end{array}$
$\begin{array}{lll}17.12 & 13.98 & -2.25625 \\ 17.11 & 13.92 & -2.4750\end{array}$

| 17.11 | 13.92 | -2.4750 |
| :--- | :--- | :--- |
| 17.10 | 13.88 | -2.5875 |
| 17.09 | 13.84 | -2.7000 |

$\begin{array}{lll}17.09 & 13.84 & -2.7000 \\ 17.08 & 13.80 & -2.8125\end{array}$
$0.107 \mu \mathrm{~m}$ PSL
V.F. $=0.3201 \%$
$\theta=30 \mathrm{deg}$
Rear Translation $=11.23$ div.

| Intensity (kHz) |  | $Y$ Intercept |  | Radius (nm) | Normalized Second Cumulant | Time |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ch 0 | Ch 1 |  |  |  |  | (sec) |
| 198.197 | 177.236 | 0.9090 | 0.69 | 15.6 | 0.2800 | 120 |
| 189.354 | 169.819 | 0.0892 | 0.70 | 15.3 | 0.2800 | 120 |
| 189.943 | 170.720 | 0.0874 | 0.67 | 16.0 | 0.2800 | 120 |
| 188.337 | 170.884 | 0.8300 | 0.66 | 16.3 | 0.2900 | 120 |
| 187.077 | 170.576 | 0.7810 | 0.64 | 16.8 | 0.2900 | 120 |
| 186.659 | 107.722 | 0.7230 | 0.59 | 18.2 | 0.2900 | 120 |
| 186.497 | 170.030 | 0.6710 | 0.53 | 20.1 | 0.2900 | 120 |
| 185.173 | 168.427 | 0.6060 | 0.49 | 21.8 | 0.3000 | 120 |
| 183.106 | 165.743 | 0.5500 | 0.46 | 23.5 | 0.2800 | 120 |
| 185.146 | 173.731 | 0.4830 | 0.42 | 25.8 | 0.3000 | 120 |
| 176.123 | 152.719 | 0.4670 | 0.39 | 27.5 | 0.2900 | 120 |
| 174.775 | 155.426 | 0.4200 | 0.34 | 31.7 | 0.2600 | 120 |
| 176.500 | 159.048 | 0.3850 | 0.32 | 33.8 | 0.2400 | 120 |
| 170.662 | 151.994 | 0.3640 | 0.31 | 35.0 | 0.2200 | 120 |
| 175.844 | 160.166 | 0.3380 | 0.29 | 36.7 | 0.2100 | 120 |
| 175.429 | 161.412 | 0.3310 | 0.26 | 41.2 | 0.1700 | 120 |
| 180.032 | 160.396 | 0.3000 | 0.26 | 41.4 | 0.1600 | 120 |
| 172.730 | 164.351 | 0.2810 | 0.25 | 42.5 | 0.1400 | 120 |
| 168.626 | 163.146 | 0.2470 | 0.22 | 48.9 | 0.0460 | 120 |
| 173.109 | 157.109 | 0.2320 | 0.24 | 44.2 | 0.1100 | 120 |
| 177.866 | 152.325 | 0.2060 | 0.23 | 44.8 | 0.0890 | 120 |
| 173.438 | 152.232 | 0.2030 | 0.24 | 44.1 | 0.0990 | 120 |
| 169.144 | 159.735 | 0.1810 | 0.22 | 48.0 | 0.0180 | 120 |
| 167.207 | 165.030 | 0.1630 | 0.23 | 46.4 | 0.0910 | 120 |
| 167.867 | 163.191 | 0.1620 | 0.22 | 48.3 | 0.0810 | 120 |
| 170.571 | 161.822 | 0.1500 | 0.23 | 46.2 | 0.0830 | 300 |


| Exp 42 <br> Scattering <br> Angle (deg) | Circular Cell |  | $\begin{array}{r} 0.107 \\ \text { ion }=\text { Sing } \end{array}$ | m PSL <br> e Scatter | $\theta=$ Sweep | Tilt Angle (mrad) $=0.000$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Intensit | Hz) | $Y$ - | Decay Rate | Radius | Normalized Second | ime (sec) |
|  | Ch 0 | Ch 1 | Intercept | (1/msec) | (nm) | Cumulant | Troe |
| 120.00 | 31.540 | 0.000 | 0.883 | 1.9900 | 58.1 | 0.0150 | 60 |
| 115.00 | 31.785 | 0.000 | 0.886 | 1.9200 | 57.0 | 0.0170 | 60 |
| 110.00 | 31.905 | 0.000 | 0.889 | 1.8200 | 56.6 | -0.0059 | 60 |
| 105.00 | 31.165 | 0.000 | 0.890 | 1.7200 | 56.3 | 0.0100 | 60 |
| 100.00 | 30.084 | 0.000 | 0.879 | 1.6500 | 54.6 | 0.0230 | 60 |
| 95.00 | 31.011 | 0.000 | 0.884 | 1.5000 | 55.8 | -0.0047 | 60 |
| 90.00 | 31.629 | 0.000 | 0.890 | 1.3900 | 55.2 | -0.0047 | 60 |
| 85.00 | 31.939 | 0.000 | 0.895 | 1.2600 | 55.7 | 0.0047 | 60 |
| 80.00 | 32.865 | 0.000 | 0.897 | 1.1700 | 54.2 | 0.0096 | 60 |
| 75.00 | 34.536 | 0.000 | 0.908 | 1.0100 | 56.2 | -0.0036 | 60 |
| 70.00 | 35.323 | 0.000 | 0.907 | 0.9430 | 53.5 | 0.0180 | 60 |
| 65.00 | 36.797 | 0.000 | 0.912 | 0.8180 | 54.1 | 0.0310 | 60 |
| 60.00 | 39.536 | 0.000 | 0.918 | 0.7230 | 53.0 | 0.0450 | 60 |
| 55.00 | 41.602 | 0.000 | 0.903 | 0.6200 | 52.8 | 0.0360 | 60 |
| 50.00 | 44.747 | 0.000 | 0.923 | 0.5230 | 52.4 | 0.0470 | 60 |
| 45.00 | 48.165 | 0.000 | 0.930 | 0.4390 | 51.2 | 0.0850 | 60 |
| 40.00 | 51.805 | 0.000 | 0.933 | 0.3570 | 50.3 | 0.0350 | 60 |
| 35.00 | 59.273 | 0.000 | 0.936 | 0.2770 | 50.1 | 0.0860 | 60 |
| 30.00 | 67.845 | 0.000 | 0.904 | 0.2110 | 48.8 | 0.0990 | 60 |
| 25.00 | 77.111 | 0.000 | 0.916 | 0.1460 | 49.1 | 0.1000 | 60 |
| 20.00 | 111.859 | 0.000 | 0.937 | 0.0819 | 56.5 | 0.1400 | 60 |
| 15.00 | 138.699 | 0.000 | 0.900 | 0.0492 | 53.1 | 0.0099 | 60 |
| 120.00 | 0.000 | 19.260 | 0.942 | 1.9900 | 58.0 | 0.0170 | 60 |
| 115.00 | 0.000 | 19.325 | 0.943 | 1.9200 | 56.9 | 0.0080 | 60 |
| 110.00 | 0.000 | 19.286 | 0.940 | 1.8100 | 57.1 | 0.0046 | 60 |
| 105.00 | 0.000 | 18.196 | 0.932 | 1.7600 | 55.1 | 0.0190 | 60 |
| 100.00 | 0.000 | 18.266 | 0.933 | 1.6200 | 55.6 | -0.0079 | 60 |
| 95.00 | 0.000 | 17.914 | 0.936 | 1.4900 | 56.1 | 0.0060 | 60 |
| 90.00 | 0.000 | 19.447 | 0.936 | 1.4100 | 54.5 | 0.0260 | 60 |
| 85.00 | 0.000 | 19.730 | 0.940 | 1.2500 | 55.8 | 0.0130 | 60 |
| 80.00 | 0.000 | 20.104 | 0.942 | 1.1900 | 53.3 | 0.0440 | 60 |
| 75.00 | 0.000 | 20.884 | 0.943 | 1.0300 | 55.4 | 0.0240 | 60 |
| 70.00 | 0.000 | 21.416 | 0.953 | 0.9150 | 55.1 | 0.0030 | 60 |
| 65.00 | 0.000 | 22.375 | 0.941 | 0.8350 | 53.0 | 0.0220 | 60 |
| 60.00 | 0.000 | 24.252 | 0.945 | 0.7230 | 53.0 | 0.0180 | 60 |
| 55.00 | 0.000 | 25.911 | 0.948 | 0.6150 | 53.1 | 0.0054 | 60 |
| 50.00 | 0.000 | 28.239 | 0.939 | 0.5470 | 50.1 | 0.0790 | 60 |
| 45.00 | 0.000 | 31.444 | 0.948 | 0.4410 | 50.9 | 0.0720 | 60 |
| 40.00 | 0.000 | 33.966 | 0.935 | 0.3740 | 47.9 | 0.0880 | 60 |
| 35.00 | 0.000 | 39.407 | 0.969 | 0.2710 | 51.1 | 0.1100 | 60 |
| 30.00 | 0.000 | 45.108 | 0.947 | 0.2120 | 48.6 | 0.1200 | 60 |
| 25.00 | 0.000 | 53.391 | 0.9520 | 0.1390 | 51.6 | 0.1300 | 60 |
| 20.00 | 0.000 | 73.847 | 0.8850 | 0.0889 | 52.0 | 0.1900 | 60 |
| 15.00 | 0.000 | 92.572 | 0.9390 | 0.0523 | 50.0 | 0.1500 | 60 |


| Exp 43 | Circular Cell Front Till = 16.92 div. |  | $0.107 \mu \mathrm{~m}$ PSL <br> Rear Translation |  | $\text { V.F. }=0.3045 \%$$23 \mathrm{div} .$ |  | $\theta=60 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div) | Side Translation (div) | Tilt Angle (mrad) | Inten Ch 0 | kHz) Ch 1 | $Y$ - <br> Intercept | Decay Rate (1/msec) | Radius (nm) | Normalized Second Cumulant | Time (sec) |
| 17.04 | 13.51 | -3.2625 | 163.332 | 98.935 | 0.0480 | 0.946 | 41.2 | 0.1100 | 180 |
| 17.05 | 13.57 | -3.1500 | 163.879 | 98.558 | 0.0559 | 0.841 | 47.7 | 0.0200 | 180 |
| 17.06 | 13.62 | -3.0375 | 166.555 | 99.080 | 0.0635 | 0.794 | 50.1 | 0.0900 | 180 |
| 17.07 | 13.67 | -2.9250 | 167.512 | 99.986 | 0.0712 | 0.774 | 51.3 | -0.0160 | 180 |
| 17.08 | 13.71 | -2.8125 | 167.869 | 98.943 | 0.0746 | 0.808 | 49.2 | 0.0590 | 180 |
| 17.09 | 13.74 | -2.7000 | 170.751 | 98.133 | 0.0822 | 0.798 | 49.8 | 0.0550 | 300 |
| 17.10 | 13.77 | -2.5875 | 173.652 | 98.939 | 0.0850 | 0.821 | 48.3 | 0.0600 | 300 |
| 17.11 | 13.80 | -2.4750 | 173.115 | 95.832 | 0.0921 | 0.811 | 49.0 | 0.0190 | 300 |
| 17.12 | 13.82 | -2.3625 | 179.398 | 104.382 | 0.0928 | 0.779 | 51.7 | -0.0800 | 300 |
| 17.13 | 13.89 | -2.2500 | 182.377 | 105.658 | 0.1010 | 0.834 | 48.3 | 0.0520 | 300 |
| 17.14 | 13.94 | -2.1375 | 185.377 | 109.776 | 0.1050 | 0.773 | 52.1 | -0.0340 | 300 |
| 17.15 | 13.99 | -2.0250 | 173.833 | 102.341 | 0.1260 | 0.803 | 50.4 | 0.0800 | 300 |
| 17.16 | 14.05 | -1.9125 | 175.688 | 104.995 | 0.1320 | 0.816 | 49.6 | 0.0660 | 300 |
| 17.17 | 14.10 | -1.8000 | 179.041 | 103.804 | 0.1460 | 0.817 | 49.5 | 0.0400 | 300 |
| 17.18 | 14.17 | -1.6875 | 181.803 | 104.938 | 0.1670 | 0.819 | 49.4 | 0.0160 | 300 |
| 17.19 | 14.27 | -1.5750 | 176.887 | 105.533 | 0.1850 | 0.847 | 47.8 | 0.0550 | 300 |
| 17.20 | 14.34 | -1.4625 | 173.526 | 107.272 | 0.2000 | 0.886 | 45.7 | 0.0800 | 300 |
| 17.21 | 14.39 | -1.3500 | 176.078 | 108.545 | 0.2240 | 0.890 | 45.4 | 0.0790 | 120 |
| 17.22 | 14.44 | -1.2375 | 177.891 | 108.136 | 0.2470 | 0.963 | 42.0 | 0.1200 | 120 |
| 17.23 | 14.50 | -1.1250 | 179.997 | 107.705 | 0.2800 | 1.020 | 39.6 | 0.1500 | 120 |
| 17.24 | 14.57 | -1.0125 | 181.581 | 109.587 | 0.3310 | 1.140 | 35.6 | 0.1800 | 120 |
| 17.25 | 14.65 | -0.9000 | 182.177 | 107.734 | 0.3890 | 1.260 | 32.2 | 0.2200 | 120 |
| 17.26 | 14.68 | -0.7875 | 181.588 | 107.700 | 0.4320 | 1.340 | 30.2 | 0.2300 | 120 |
| 17.27 | 14.74 | -0.6750 | 171.099 | 102.415 | 0.5110 | 1.470 | 27.4 | 0.2300 | 120 |
| 17.28 | 14.78 | -0.5625 | 176.482 | 104.478 | 0.5930 | 1.610 | 25.1 | 0.2500 | 120 |
| 17.29 | 14.83 | -0.4500 | 178.808 | 107.210 | 0.6790 | 1.700 | 23.8 | 0.2500 | 120 |
| 17.30 | 14.87 | -0.3375 | 187.333 | 114.901 | 0.7430 | 1.820 | 22.3 | 0.2400 | 120 |
| 17.31 | 14.92 | -0.2250 | 184.195 | 115.343 | 0.8070 | 1.800 | 21.2 | 0.2500 | 120 |
| 17.32 | 14.95 | -0.1125 | 184.603 | 116.222 | 0.8540 | 1.860 | 20.6 | 0.1600 | 120 |
| 17.33 | 15.00 | 0.0000 | 174.626 | 110.801 | 0.9190 | 1.760 | 21.8 | 0.2500 | 120 |
| 17.34 | 15.06 | 0.1125 | 181.189 | 115.671 | 0.8730 | 1.720 | 22.3 | 0.2600 | 120 |
| 17.35 | 15.11 | 0.2250 | 173.986 | 112.051 | 0.8000 | 1.610 | 23.8 | 0.2500 | 120 |
| 17.36 | 15.18 | 0.3375 | 174.775 | 114.132 | 0.7040 | 1.500 | 25.5 | 0.2500 | 120 |
| 17.37 | 15.27 | 0.4500 | 174.529 | 115.138 | 0.5680 | 1.360 | 28.2 | 0.2300 | 120 |
| 17.38 | 15.35 | 0.5625 | 174.823 | 114.732 | 0.4780 | 1.200 | 31.9 | 0.2300 | 120 |
| 17.39 | 15.37 | 0.6750 | 179.472 | 118.442 | 0.4360 | 1.140 | 33.7 | 0.2100 | 120 |
| 17.40 | 15.40 | 0.7875 | 174.985 | 114.761 | 0.4010 | 1.080 | 35.4 | 0.1700 | 120 |
| 17.41 | 15.44 | 0.9000 | 173.497 | 113.580 | 0.3710 | 1.020 | 37.6 | 0.1600 | 120 |
| 17.42 | 15.49 | 1.0125 | 165.504 | 108.787 | 0.3170 | 0.977 | 39.3 | 0.1400 | 120 |
| 17.43 | 15.54 | 1.1250 | 165.466 | 114.179 | 0.2900 | 0.927 | 41.4 | 0.1400 | 120 |
| 17.44 | 15.57 | 1.2375 | 164.993 | 108.413 | 0.2790 | 0.905 | 42.4 | 0.1400 | 120 |
| 17.45 | 15.64 | 1.3500 | 163.863 | 108.294 | 0.2500 | 0.859 | 44.6 | 0.0780 | 120 |
| 17.46 | 15.69 | 1.4625 | 174.031 | 107.427 | 0.2330 | 0.858 | 44.7 | 0.0980 | 120 |
| 17.47 | 15.73 | 1.5750 | 165.064 | 109.015 | 0.2130 | 0.826 | 46.4 | 0.0540 | 120 |
| 17.48 | 15.78 | 1.6875 | 165.408 | 108.125 | 0.2040 | 0.814 | 47.1 | 0.0290 | 120 |
| 17.49 | 15.85 | 1.8000 | 165.135 | 108.259 | 0.1840 | 0.815 | 47.1 | 0.0630 | 300 |
| 17.50 | 15.90 | 1.9125 | 166.160 | 107.914 | 0.1720 | 0.816 | 47.0 | 0.0810 | 300 |
| 17.51 | 16.00 | 2.0250 | 167.723 | 105.975 | 0.1540 | 0.814 | 47.1 | 0.1000 | 300 |
| 17.52 | 16.05 | 2.1375 | 184.502 | 107.562 | 0.1240 | 0.804 | 47.7 | 0.0560 | 300 |
| 17.53 | 16.11 | 2.2500 | 185.472 | 105.696 | 0.1150 | 0.803 | 47.8 | 0.1100 | 300 |
| 17.54 | 16.21 | 2.3625 | 170.026 | 101.422 | 0.1080 | 0.806 | 47.6 | 0.0570 | 300 |
| 17.55 | 16.26 | 2.4750 | 173.602 | 104.330 | 0.0966 | 0.781 | 49.1 | 0.0360 | 300 |
| 17.56 | 16.32 | 2.5875 | 171.936 | 102.644 | 0.0869 | 0.788 | 48.7 | 0.0320 | 300 |
| 17.57 | 16.37 | 2.7000 | 173.918 | 101.464 | 0.0791 | 0.754 | 50.8 | -0.0054 | 300 |
| 17.58 | 16.42 | 2.8125 | 177.909 | 106.383 | 0.0731 | 0.756 | 50.7 | 0.0220 | 300 |
| 17.59 | 16.47 | 2.9250 | 182.680 | 107.160 | 0.0670 | 0.757 | 50.6 | -0.0230 | 300 |
| 17.60 | 16.52 | 3.0375 | 185.526 | 109.364 | 0.0625 | 0.752 | 51.0 | 0.0450 | 300 |
| 17.61 | 16.57 | 3.1500 | 188.553 | 111.130 | 0.0526 | 0.746 | 51.4 | 0.0028 | 300 |
| 17.62 | 16.60 | 3.2625 | 191.031 | 108.706 | 0.0437 | 0.760 | 50.4 | -0.0023 | 300 |
| 17.63 | 16.67 | 3.3750 | 184.565 | 109.328 | 0.0357 | 0.878 | 43.6 | 0.0440 | 300 |

Exp 44
Circular Cell $\quad 0.107 \mu \mathrm{mPSL} \quad$ V.F. $=0.1545 \% \quad \theta=60$ deg
Front Tilt $=16.92$ div. Rear Translation $=11.23$ div.

| Rear Tilt (div) | Side Translation (div) | Tilt Angle (mrad) | Intensity (kHz) |  | $Y$ - <br> Intercept | Decay <br> Rate <br> ( $1 / \mathrm{msec}$ ) | Radius ( m ) | Normalized Second Cumulant | Time (sec) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Ch0 | Ch 1 |  |  |  |  |  |
| 16.89 | 12.72 | -4.9500 | 188.501 | 113.884 | 0.0262 | 0.638 | 61.7 | 0.096 | 300 |
| 16.91 | 12.34 | -4.7250 | 189.840 | 115.051 | 0.0407 | 0.731 | 53.9 | 0.0750 | 240 |
| 16.93 | 12.93 | -4.5000 | 188.748 | 114.986 | 0.0513 | 0.374 | 53.3 | -0.0870 | 240 |
| 16.95 | 13.04 | -4.2750 | 189.572 | 114.984 | 0.0613 | 0.776 | 50.7 | 0.1000 | 240 |
| 16.97 | 13.12 | -4.0500 | 188.930 | 113.088 | 0.0685 | 0.761 | 51.8 | 0.0150 | 240 |
| 16.99 | 13.22 | -3.8250 | 188.623 | 113.767 | 0.0799 | 0.788 | 50.0 | 0.0180 | 180 |
| 17.01 | 13.28 | -3.6000 | 187.220 | 111.350 | 0.0903 | 0.780 | 50.5 | 0.0013 | 180 |
| 17.03 | 13.37 | -3.3750 | 187.280 | 113.600 | 0.1170 | 0.709 | 55.5 | 0.0580 | 180 |
| 17.05 | 13.47 | -3.1500 | 187.792 | 113.325 | 0.1310 | 0.765 | 51.5 | 0.0420 | 120 |
| 17.07 | 13.56 | -29250 | 185.859 | 114.551 | 0.1640 | 0.776 | 50.7 | 0.0028 | 120 |
| 17.09 | 13.64 | -27000 | 184.452 | 113.456 | 0.1870 | 0.772 | 51.0 | 0.0480 | 120 |
| 17.11 | 13.73 | -2.4750 | 183.871 | 112.671 | 0.2070 | 0.776 | 50.7 | 0.0580 | 120 |
| 17.13 | 13.80 | -2.2500 | 184.022 | 110.686 | 0.2300 | 0.788 | 50.0 | 0.0760 | 120 |
| 17.15 | 13.92 | -2.0250 | 183.123 | 113.562 | 0.2690 | 0.773 | 50.9 | 0.0077 | 120 |
| 17.16 | 13.97 | -1.9125 | 182.749 | 112.708 | 0.2750 | 0.757 | 52.0 | 0.0084 | 120 |
| 17.17 | 14.02 | -1.8000 | 182.572 | 111.684 | 0.3040 | 0.771 | 51.1 | 0.0470 | 120 |
| 17.18 | 14.12 | -1.6875 | 193.628 | 113.771 | 0.3500 | 0.772 | 51.0 | 0.0420 | 120 |
| 17.19 | 14.15 | -1.5750 | 191.148 | 113.770 | 0.3680 | 0.782 | 50.3 | 0.0770 | 120 |
| 17.20 | 14.19 | -1.4625 | 190.044 | 112.460 | 0.3770 | 0.753 | 52.3 | -0.0025 | 120 |
| 17.21 | 14.23 | -1.3500 | 190.830 | 111.003 | 0.3950 | 0.793 | 49.6 | 0.0770 | 120 |
| 17.22 | 14.36 | -1.2375 | 185.740 | 113.178 | 0.4510 | 0.804 | 49.0 | 0.0440 | 120 |
| 17.23 | 14.42 | -1.1250 | 184.952 | 113.051 | 0.4790 | 0.801 | 49.2 | 0.0840 | 120 |
| 17.24 | 14.50 | -1.0125 | 183.584 | 114.410 | 0.5250 | 0.843 | 46.7 | 0.0780 | 120 |
| 17.25 | 14.57 | -0.9000 | 182540 | 114.476 | 0.5950 | 0.866 | 45.5 | 0.1100 | 120 |
| 17.26 | 14.62 | -0.7875 | 181.238 | 113.594 | 0.6320 | 0.884 | 44.5 | 0.1100 | 120 |
| 17.27 | 14.67 | -0.6750 | 180.832 | 111.593 | 0.6680 | 0.918 | 42.9 | 0.1400 | 120 |
| 17.28 | 14.74 | -0.5625 | 180.647 | 111.328 | 0.7490 | 0.972 | 40.5 | 0.1600 | 120 |
| 17.29 | 14.77 | -0.4500 | 180.025 | 109.751 | 0.7890 | 1.010 | 39.2 | 0.1700 | 120 |
| 17.30 | 14.83 | -0.3375 | 180.376 | 109.879 | 0.8420 | 1.040 | 37.7 | 0.1800 | 120 |
| 17.31 | 14.91 | -0.2250 | 179.764 | 110.894 | 0.9190 | 1.080 | 36.5 | 0.1900 | 120 |
| 17.32 | 14.96 | -0.1125 | 182306 | 112601 | 0.9310 | 1.100 | 35.9 | 0.1800 | 120 |
| 17.33 | 15.00 | 0.0000 | 186.309 | 116.340 | 0.9270 | 1.100 | 35.9 | 0.1800 | 120 |


| Exp 45 | Circular Cell |  | $0.107 \mu \mathrm{~m}$ PSL |  | $\theta=$ Sweep | Tilt Angle $(\mathrm{mrad})=0.000$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Scattering Angle (deg) | Intens <br> Ch 0 | (kHz) Ch 1 | Y- <br> Intercept | Decay Rate (1/msec) | Radius (nm) | Normalized Second Cumulant | Time (sec) |
| 120.00 | 61.675 | 0.000 | 0.958 | 2.170 | 55.3 | 0.00440 | 60 |
| 115.00 | 60.878 | 0.000 | 0.950 | 2.080 | 54.6 | 0.01400 | 60 |
| 110.00 | 60.021 | 0.000 | 0.955 | 1.980 | 54.0 | 0.01200 | 60 |
| 105.00 | 58.945 | 0.000 | 0.952 | 1.850 | 54.3 | 0.02400 | 60 |
| 100.00 | 58.808 | 0.000 | 0.952 | 1.750 | 53.5 | 0.03000 | 60 |
| 95.00 | 58.630 | 0.000 | 0.941 | 1.640 | 52.9 | 0.00830 | 60 |
| 90.00 | 59.109 | 0.000 | 0.942 | 1.480 | 53.9 | -0.00052 | 60 |
| 85.00 | 60.722 | 0.000 | 0.938 | 1.400 | 52.1 | 0.02000 | 60 |
| 80.00 | 58.110 | 0.000 | 0.940 | 1.270 | 52.1 | 0.01400 | 60 |
| 75.00 | 65.162 | 0.000 | 0.948 | 1.130 | 52.4 | 0.05200 | 60 |
| 70.00 | 66.242 | 0.000 | 0.942 | 0.999 | 52.6 | 0.00650 | 60 |
| 65.00 | 70.873 | 0.000 | 0.939 | 0.911 | 50.6 | 0.03100 | 60 |
| 60.00 | 75.175 | 0.000 | 0.937 | 0.787 | 50.8 | 0.01100 | 120 |
| 55.00 | 79.033 | 0.000 | 0.938 | 0.676 | 50.4 | 0.04500 | 120 |
| 50.00 | 88.918 | 0.000 | 0.949 | 0.549 | 52.0 | 0.11000 | 120 |
| 45.00 | 90.126 | 0.000 | 0.938 | 0.481 | 48.6 | 0.06600 | 120 |
| 40.00 | 102.658 | 0.000 | 0.946 | 0.380 | 49.1 | 0.12000 | 120 |
| 35.00 | 128.749 | 0.000 | 0.973 | 0.250 | 57.8 | 0.21000 | 120 |
| 30.00 | 118.902 | 0.000 | 0.944 | 0.199 | 53.7 | 0.17000 | 120 |
| 120.00 | 0.000 | 42.037 | 0.962 | 2.200 | 54.6 | 0.02900 | 60 |
| 115.00 | 0.000 | 40.941 | 0.958 | 2.080 | 54.5 | 0.01100 | 60 |
| 110.00 | 0.000 | 40.053 | 0.965 | 1.990 | 54.0 | 0.04100 | 60 |
| 105.00 | 0.000 | 39.832 | 0.963 | 1.880 | 53.4 | 0.02800 | 60 |
| 100.00 | 0.000 | 39.184 | 0.970 | 1.740 | 53.9 | 0.00840 | 60 |
| 95.00 | 0.000 | 39.111 | 0.957 | 1.640 | 53.0 | 0.01700 | 60 |
| 90.00 | 0.000 | 39.295 | 0.959 | 1.480 | 54.0 | -0.01800 | 60 |
| 85.00 | 0.000 | 40.480 | 0.953 | 1.370 | 53.1 | 0.00460 | 60 |
| 80.00 | 0.000 | 41.395 | 0.951 | 1.280 | 51.7 | 0.01500 | 60 |
| 75.00 | 0.000 | 43.381 | 0.947 | 1.150 | 51.3 | 0.02600 | 60 |
| 70.00 | 0.000 | 44.198 | 0.956 | 1.020 | 51.3 | 0.02500 | 60 |
| 65.00 | 0.000 | 47.204 | 0.961 | 0.920 | 50.1 | 0.01400 | 60 |
| 60.00 | 0.000 | 50.774 | 0.971 | 0.804 | 49.7 | 0.05200 | 120 |
| 55.00 | 0.000 | 54.505 | 0.947 | 0.684 | 49.8 | 0.06800 | 120 |
| 50.00 | 0.000 | 58.292 | 0.961 | 0.590 | 48.3 | 0.06600 | 120 |
| 45.00 | 0.000 | 63.990 | 0.961 | 0.483 | 48.4 | 0.08500 | 120 |
| 40.00 | 0.000 | 71.897 | 0.952 | 0.394 | 47.5 | 0.12000 | 120 |
| 35.00 | 0.000 | 87.921 | 0.956 | 0.291 | 49.6 | 0.09900 | 120 |
| 30.00 | 0.000 | 99.764 | 0.966 | 0.206 | 52.0 | 0.18000 | 120 |

## Exp 46

Circular Cell
$0.107 \mu \mathrm{~m}$ PSL V.F. $=0.50255 \%$
$\theta=120 \mathrm{deg}$
Front Tilt $=16.92$ div. $\quad$ Rear Translation $=11.72$ div.

| Rear Tilt <br> (div) | Side <br> Translation <br> (div) | Tilt Angle <br> (mrad) | Intensity (kHz) | Y- <br> (ntercept | Decay <br> Rate <br> $(1 / \mathrm{msec})$ | Radius <br> (nm) | Normalized <br> Second <br> Cumulant | Time <br> (sec) |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 17.35 | 15.04 | 0.0000 | 156.721 | 100.844 | 0.9010 | 5.16 | 22.9 | 0.210 | 120 |
| 17.34 | 14.99 | -0.1125 | 157.468 | 101.650 | 0.8840 | 5.14 | 23.0 | 0.210 | 120 |
| 17.33 | 14.95 | -0.2250 | 161.346 | 104.309 | 0.7920 | 4.96 | 23.8 | 0.200 | 120 |
| 17.32 | 14.90 | -0.3375 | 161.785 | 104.385 | 0.6630 | 4.90 | 24.1 | 0.210 | 120 |
| 17.31 | 14.86 | -0.4500 | 162.250 | 105.025 | 0.5220 | 4.74 | 24.9 | 0.210 | 120 |
| 17.30 | 14.80 | -0.5625 | 162.265 | 106.094 | 0.3530 | 4.44 | 26.6 | 0.200 | 120 |
| 17.29 | 14.74 | -0.6750 | 160.324 | 103.418 | 0.2270 | 4.00 | 29.5 | 0.190 | 120 |
| 17.28 | 14.67 | -0.7875 | 160.816 | 103.462 | 0.1490 | 3.52 | 33.6 | 0.180 | 120 |
| 17.27 | 14.61 | -0.9000 | 157.853 | 102.419 | 0.1010 | 3.22 | 36.7 | 0.180 | 120 |
| 17.26 | 14.56 | -1.0125 | 157.184 | 102.815 | 0.0811 | 3.04 | 38.9 | 0.180 | 120 |
| 17.25 | 14.51 | -1.1250 | 159.558 | 103.692 | 0.0685 | 2.68 | 44.1 | 0.074 | 300 |
| 17.24 | 14.46 | -1.2375 | 163.422 | 102.924 | 0.0564 | 2.58 | 45.8 | 0.011 | 300 |
| 17.23 | 14.42 | -1.3500 | 161.738 | 103.037 | 0.0505 | 2.48 | 47.7 | 0.100 | 300 |
| 17.22 | 14.35 | -1.4625 | 160.432 | 104.361 | 0.0424 | 2.35 | 50.2 | 0.032 | 420 |


| Exp 47 | Circular Cell Front Tilt = 16.92 div. |  | $0.107 \mu \mathrm{~m}$ PSL |  | V.F. $=0.3795 \%$ |  | $\theta=120 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Rear Tr | siation | 23 div |  |  |  |  |
| Rear <br> Tilt (div) | Side Translation (div) | Tilt Angle (mrad) | Intens <br> Ch 0 | $\mathrm{kHz})$ Ch 1 | Y-Intercept | $\begin{gathered} \text { Decay } \\ \text { Rate } \\ (1 / \mathrm{msec}) \end{gathered}$ | Radius (nm) | Normalized Second Cumulant | Time (sec) |
| 17.06 | 13.53 | -3.2625 | 92.551 | 58.181 | 0.0426 | 2.07 | 57.0 | -0.12000 | 420 |
| 17.08 | 13.66 | -3.0375 | 92.767 | 59.478 | 0.0520 | 2.32 | 51.0 | 0.02300 | 420 |
| 17.10 | 13.74 | -2.8125 | 92.895 | 60.895 | 0.0608 | 2.12 | 55.8 | -0.01400 | 420 |
| 17.12 | 13.82 | -2.5875 | 92.613 | 60.892 | 0.0660 | 2.26 | 52.4 | 0.06200 | 420 |
| 17.13 | 13.85 | -2.4750 | 92.699 | 61.088 | 0.0677 | 2.29 | 51.7 | 0.00970 | 420 |
| 17.14 | 13.93 | -2.3625 | 92.154 | 61.162 | 0.0685 | 2.32 | 51.0 | 0.03200 | 420 |
| 17.15 | 13.97 | -2.2500 | 92.921 | 61.525 | 0.0751 | 2.30 | 51.5 | 0.04800 | 420 |
| 17.16 | 14.04 | -2.1375 | 93.191 | 61.250 | 0.0769 | 2.24 | 52.7 | 0.04800 | 420 |
| 17.17 | 14.07 | -2.0250 | 93.590 | 61.347 | 0.0822 | 2.13 | 55.4 | 0.01200 | 420 |
| 17.18 | 14.11 | -1.9125 | 94.406 | 61.500 | 0.0949 | 2.17 | 54.5 | -0.02100 | 420 |
| 17.19 | 14.16 | -1.8000 | 94.249 | 61.475 | 0.0903 | 2.28 | 51.8 | 0.02500 | 300 |
| 17.20 | 14.20 | -1.6875 | 94.028 | 61.550 | 0.0975 | 2.23 | 52.9 | 0.01500 | 300 |
| 17.21 | 14.25 | -1.5750 | 94.512 | 61.697 | 0.1070 | 2.26 | 52.2 | 0.03000 | 300 |
| 17.22 | 14.33 | -1.4625 | 94.881 | 62.582 | 0.1200 | 2.33 | 50.8 | 0.06200 | 300 |
| 17.23 | 14.38 | -1.3500 | 94.411 | 62.643 | 0.1290 | 2.27 | 52.0 | 0.03300 | 300 |
| 17.24 | 14.45 | -1.2375 | 94.250 | 63.053 | 0.1430 | 2.39 | 49.4 | 0.04900 | 300 |
| 17.25 | 14.53 | -1.1250 | 94.830 | 63.270 | 0.1780 | 2.44 | 48.5 | 0.03600 | 300 |
| 17.26 | 14.59 | -1.0125 | 94.416 | 62.259 | 0.1980 | 2.54 | 46.5 | 0.07300 | 120 |
| 17.27 | 14.65 | -0.9000 | 94.599 | 62.215 | 0.2410 | 2.66 | 44.5 | 0.07300 | 120 |
| 17.28 | 14.71 | -0.7875 | 94.073 | 61.186 | 0.2980 | 2.82 | 41.9 | 0.14000 | 120 |
| 17.29 | 14.75 | -0.6750 | 95.199 | 61.752 | 0.3710 | 2.96 | 40.0 | 0.12000 | 120 |
| 17.30 | 14.80 | -0.5625 | 94.710 | 60.309 | 0.4610 | 3.12 | 37.8 | 0.16000 | 120 |
| 17.31 | 14.84 | -0.4500 | 95.231 | 60.953 | 0.5750 | 3.32 | 35.6 | 0.15000 | 120 |
| 17.32 | 14.88 | -0.3375 | 95.902 | 60.904 | 0.7060 | 3.40 | 34.8 | 0.16000 | 120 |
| 17.33 | 14.94 | -0.2250 | 94.763 | 90.100 | 0.8250 | 3.51 | 33.7 | 0.17000 | 120 |
| 17.34 | 14.98 | -0.1125 | 94.991 | 59.076 | 0.8830 | 3.58 | 33.0 | 0.17000 | 120 |
| 17.35 | 15.04 | 0.0000 | 93.645 | 58.386 | 0.9220 | 3.57 | 33.1 | 0.17000 | 120 |
| 17.36 | 15.09 | 0.1125 | 92.084 | 56.589 | 0.8430 | 3.55 | 33.3 | 0.16000 | 120 |
| 17.37 | 15.14 | 0.2250 | 92.289 | 56.873 | 0.7700 | 3.46 | 34.1 | 0.17000 | 120 |
| 17.38 | 15.18 | 0.3375 | 92.049 | 56.356 | 0.6270 | 3.36 | 35.1 | 0.15000 | 120 |
| 17.39 | 15.23 | 0.4500 | 92.411 | 56.531 | 0.5130 | 3.25 | 36.4 | 0.16000 | 120 |
| 17.40 | 15.27 | 0.5625 | 92.090 | 56.687 | 0.4130 | 3.07 | 38.5 | 0.14000 | 120 |
| 17.41 | 15.30 | 0.6750 | 92.341 | 56.695 | 0.3390 | 2.91 | 40.6 | 0.13000 | 120 |
| 17.42 | 15.36 | 0.7875 | 92.258 | 56.808 | 0.2610 | 2.72 | 43.5 | 0.10000 | 120 |
| 17.43 | 15.41 | 0.9000 | 91.953 | 56.814 | 0.2110 | 2.62 | 45.1 | 0.08100 | 120 |
| 17.44 | 15.47 | 1.0125 | 92.092 | 57.287 | 0.1730 | 2.46 | 48.0 | 0.08000 | 120 |
| 17.45 | 15.54 | 1.1250 | 92.496 | 58.246 | 0.1440 | 2.33 | 50.6 | 0.04200 | 300 |
| 17.46 | 15.64 | 1.2375 | 92.543 | 58.369 | 0.1210 | 2.24 | 52.7 | 0.03000 | 300 |
| 17.47 | 15.70 | 1.3500 | 97.946 | 59.616 | 0.1140 | 2.21 | 53.6 | 0.00460 | 300 |
| 17.48 | 15.77 | 1.4625 | 97.291 | 59.404 | 0.1040 | 2.17 | 54.5 | -0.02000 | 300 |
| 17.49 | 15.86 | 1.5750 | 96.951 | 57.960 | 0.0985 | 2.18 | 54.3 | 0.01200 | 300 |
| 17.50 | 15.93 | 1.6875 | 95.474 | 56.937 | 0.0904 | 2.20 | 53.7 | -0.00026 | 300 |
| 17.51 | 16.00 | 1.8000 | 93.921 | 56.528 | 0.0848 | 2.24 | 52.7 | 0.05100 | 300 |
| 17.52 | 16.05 | 1.9125 | 94.095 | 56.753 | 0.0823 | 2.28 | 51.8 | 0.03200 | 300 |
| 17.53 | 16.08 | 2.0250 | 95.189 | 56.818 | 0.0801 | 2.21 | 53.6 | 0.01000 | 300 |
| 17.55 | 16.17 | 2.2500 | 94.894 | 57.078 | 0.0723 | 2.18 | 54.3 | -0.02600 | 300 |
| 17.57 | 16.27 | 2.4750 | 95.537 | 54.335 | 0.0635 | 2.14 | 55.3 | -0.00015 | 300 |
| 17.59 | 16.33 | 2.7000 | 97.177 | 54.765 | 0.0613 | 2.34 | 50.6 | 0.06400 | 300 |
| 17.61 | 16.43 | 2.9250 | 97.702 | 54.228 | 0.0509 | 2.25 | 52.4 | 0.05200 | 300 |
| 17.63 | 16.53 | 3.1500 | 98.370 | 55.510 | 0.0428 | 2.23 | 53.1 | 0.00460 | 300 |
| 17.65 | 16.67 | 3.3750 | 97.348 | 55.062 | 0.0350 | 1.96 | 60.4 | -0.01700 | 300 |

Circular Cell
$0.107 \mu \mathrm{~m}$ PSL V.F. $=0.3795 \%$
$\theta=45 \mathrm{deg}$
Front Tilt $=16.92$ div. Rear Translation $=11.23$ div .

| Rear <br> Tilt | Side <br> Translation <br> (div) | Tilt Angle <br> (div) |  | Intensity $(\mathrm{kHz})$ | Y- <br> Inad) | Decay <br> Rate | Radius <br> (nm $)$ | Normalized <br> Second <br> Cumulant | Time <br> (sec) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 17.19 | 14.29 | -1.8000 | 107.519 | Ch 1 | 66.934 | 0.0205 | 0.499 | 47.5 | 0.170 |
| 17.20 | 14.33 | -1.6875 | 106.929 | 66.240 | 0.0302 | 0.494 | 48.0 | 0.038 | 300 |
| 17.21 | 14.37 | -1.5750 | 117.776 | 71.795 | 0.0369 | 0.553 | 42.8 | 0.093 | 300 |
| 17.22 | 14.42 | -1.4625 | 117.407 | 71.182 | 0.0501 | 0.578 | 41.0 | 0.220 | 300 |
| 17.23 | 14.49 | -1.3500 | 117.730 | 71.963 | 0.0797 | 0.600 | 39.5 | 0.140 | 300 |
| 17.24 | 14.56 | -1.2375 | 116.807 | 71.892 | 0.1150 | 0.673 | 35.2 | 0.230 | 120 |
| 17.25 | 14.61 | -1.1250 | 117.496 | 71.988 | 0.1530 | 0.774 | 30.6 | 0.240 | 120 |
| 17.26 | 14.65 | -1.0125 | 117.397 | 71.905 | 0.1870 | 0.843 | 28.1 | 0.260 | 120 |
| 17.27 | 14.70 | -0.9000 | 117.545 | 72.305 | 0.2520 | 0.908 | 26.1 | 0.260 | 120 |
| 17.28 | 14.77 | -0.7875 | 117.539 | 72.107 | 0.3550 | 1.030 | 23.1 | 0.280 | 120 |
| 17.29 | 14.80 | -0.6750 | 118.275 | 72.059 | 0.4230 | 1.130 | 20.9 | 0.270 | 120 |
| 17.30 | 14.83 | -0.5625 | 116.934 | 70.877 | 0.5090 | 1.250 | 19.0 | 0.280 | 120 |
| 17.31 | 14.89 | -0.4500 | 95.834 | 58.459 | 0.6290 | 1.460 | 15.9 | 0.270 | 120 |
| 17.32 | 14.93 | -0.3375 | 95.202 | 57.811 | 0.7340 | 1.530 | 15.0 | 0.270 | 120 |
| 17.33 | 14.97 | -0.2250 | 97.229 | 59.374 | 0.8170 | 1.600 | 14.4 | 0.270 | 120 |
| 17.34 | 15.01 | -0.1125 | 95.749 | 58.322 | 0.8650 | 1.630 | 14.1 | 0.270 | 120 |
| 17.35 | 15.04 | 0.0000 | 93.631 | 57.130 | 0.9060 | 1.650 | 14.0 | 0.270 | 120 |
| 17.36 | 15.11 | 0.1125 | 114.548 | 83.647 | 0.8110 | 1.530 | 15.5 | 0.270 | 120 |
| 17.38 | 15.20 | 0.3375 | 113.055 | 82.046 | 0.6310 | 1.370 | 17.3 | 0.270 | 120 |
| 17.40 | 15.33 | 0.5625 | 112.781 | 83.176 | 0.3890 | 1.060 | 22.3 | 0.270 | 120 |
| 17.42 | 15.43 | 0.7875 | 111.094 | 82.259 | 0.2470 | 0.903 | 26.2 | 0.260 | 120 |
| 17.44 | 15.52 | 1.0125 | 111.310 | 81.246 | 0.1590 | 0.712 | 33.3 | 0.230 | 120 |
| 17.46 | 15.66 | 1.2375 | 99.378 | 79.271 | 0.0797 | 0.603 | 39.5 | 0.140 | 120 |
| 17.47 | 15.73 | 1.3500 | 107.423 | 74.826 | 0.0505 | 0.554 | 42.8 | 0.130 | 120 |
| 17.48 | 15.77 | 1.4625 | 108.206 | 75.337 | 0.0370 | 0.508 | 46.7 | 0.029 | 120 |
| 17.49 | 15.87 | 1.5750 | 109.300 | 75.985 | 0.0153 | 0.781 | 30.4 | 0.370 | 120 |

## APPENDIX III

## Flowing Fluid: Experimental Data

Table 6: Summary of the flowing fluid experiments conducted at a scattering angle of $\theta$ $=112^{\circ}$. N/A means temperature was not available.

| Experment Number | Volume Fraction (\% by weight) | Particle Diameter ( $\mu \mathrm{m}$ ) | $\begin{aligned} & \text { Flow } \\ & \text { Rate (\%) } \end{aligned}$ | Rotation Angle (degrees) | Temperature (K) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 58 | Single Scattering | 0.107 | Sweep | Sweep | N/A |
| 60 | Single Scattering | 0.107 | Sweep | Sweep | N/A |
| 69 | 0.066 | 0.107 | 25 | 0 | N/A |
| 70 | 0.066 | 0.107 | 25 | 0 | N/A |
| 71 | 0.066 | 0.107 | 50 | 0 | N/A |
| 72 | 0.066 | 0.107 | 50 | 0 | N/A |
| 73 | 0.066 | 0.107 | 75 | 0 | 296 |
| 74 | 0.066 | 0.107 | 75 | 0 | 296 |
| 75 | 0.066 | 0.107 | 100 | 0 | 296 |
| 76 | 0.066 | 0.107 | 100 | 0 | 296 |
| 77 | 0.198 | 0.107 | 0 | 0 | 295 |
| 78 | 0.198 | 0.107 | 0 | 0 | 295 |
| 79 | 0.198 | 0.107 | 25 | 0 | N/A |
| 80 | 0.198 | 0.107 | 25 | 0 | N/A |
| 81 | 0.198 | 0.107 | 75 | 0 | N/A |
| 82 | 0.198 | 0.107 | 75 | 0 | N/A |
| 83 | 0.198 | 0.107 | 100 | 0 | 295 |
| 84 | 0.198 | 0.107 | 100 | 0 | 295 |
| 86 | 0.198 | 0.107 | 25 | 0 | 295 |
| 87 | 0.198 | 0.107 | 25 | 0 | 295 |
| 93 | 0.198 | 0.107 | 0 | 0 | 295 |
| 94 | 0.198 | 0.107 | 50 | 0 | 295 |
| 95 | 0.198 | 0.107 | 100 | 0 | 295 |
| 103 | 0.320 | 0.098 | 0 | 0 | 296 |
| 104 | 0.320 | 0.098 | 100 | 0 | 296 |
| 105 | 0.860 | 0.098 | 0 | 0 | 296 |
| 106 | 0.860 | 0.098 | 100 | 0 | 296 |
| 111 | 0.200 | 0.203 | 0 | 0 | 296 |
| 112 | 0.200 | 0.203 | 100 | 0 | 296 |

Note: The experiments not listed in this summary table did not contain any pertinent information to the subjects discussed in this thesis.

Table 7: Detailed description of experiments described in summary Table 6.


Exp 60
Flow Cell
$\theta=112 \mathrm{deg}$

| Cell <br> Rotation | Intensity (kHz) |  | $Y \text { - }$ | Decay Rate | Radius | Normalized Second | Flow <br> Rates |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Angle (deg) | Ch 0 | Ch 1 |  | (1/ms) |  | Cumulant | (\%) |
| 0.00 | 30.727 | 0.000 | 0.9210 | 1.52 | 103 | 0.0320 | 0 |
| 0.00 | 0.000 | 30.458 | 0.9470 | 1.47 | 106 | 0.0350 | 0 |
| 0.00 | 30.724 | 0.000 | 0.9270 | 1.53 | 103 | 0.0220 | 50 |
| 0.00 | 0.000 | 30.413 | 0.9530 | 1.50 | 104 | 0.0150 | 50 |
| 0.00 | 31.272 | 0.000 | 0.9180 | 1.52 | 103 | 0.0460 | 100 |
| 0.00 | 0.000 | 31.717 | 0.9570 | 1.52 | 103 | 0.0230 | 100 |
| -5.00 | 29.934 | 0.000 | 0.8940 | 1.50 | 104 | 0.0130 | 0 |
| -5.00 | 0.000 | 29.716 | 0.9420 | 1.48 | 106 | 0.0200 | 0 |
| -5.00 | 30.444 | 0.000 | 0.8980 | 1.52 | 103 | -0.0600 | 50 |
| -5.00 | 0.000 | 29.136 | 0.9430 | 1.52 | 103 | -0.0230 | 50 |
| -5.00 | 30.051 | 0.000 | 0.8970 | 1.68 | 92.3 | -0.0460 | 100 |
| -5.00 | 0.000 | 28.735 | 0.9430 | 1.65 | 94.9 | 0.0220 | 100 |
| -10.00 | 24.397 | 0.000 | 0.8570 | 1.44 | 108 | 0.0260 | 0 |
| -10.00 | 0.000 | 25.618 | 0.9360 | 1.46 | 107 | 0.0140 | 0 |
| -10.00 | 27.348 | 0.000 | 0.9500 | 3.45 | 45.4 | 0.2100 | 50 |
| -10.00 | 0.000 | 25.833 | 0.9440 | 1.70 | 92.0 | -0.0400 | 50 |
| -10.00 | 27.128 | 0.000 | 0.9690 | 6.44 | 24.3 | 0.3300 | 100 |
| -10.00 | 0.000 | 25.803 | 0.9360 | 1.89 | 82.9 | -0.5900 | 100 |
| 5.00 | 29.538 | 0.000 | 0.8710 | 1.44 | 109 | 0.0200 | 0 |
| 5.00 | 0.000 | 26.857 | 0.9360 | 1.46 | 107 | 0.0450 | 0 |
| 5.00 | 29.107 | 0.000 | 0.8750 | 1.48 | 106 | -0.0033 | 50 |
| 5.00 | 0.000 | 27.213 | 0.9360 | 1.49 | 105 | 0.0280 | 50 |
| 5.00 | 29.082 | 0.000 | 0.8700 | 1.67 | 93.5 | 0.0930 | 100 |
| 5.00 | 0.000 | 26.325 | 0.9380 | 1.58 | 99.1 | 0.0530 | 100 |
| 10.00 | 29.134 | 0.000 | 0.8690 | 1.41 | 111 | 0.0290 | 0 |
| 10.00 | 0.000 | 24.449 | 0.9330 | 1.44 | 109 | 0.0360 | 0 |
| 10.00 | 28.701 | 0.000 | 0.8780 | 1.65 | 95.0 | 0.0910 | 50 |
| 10.00 | 0.000 | 24.627 | 0.9380 | 1.80 | 86.8 | 0.1200 | 50 |
| 10.00 | 28.747 | 0.000 | 0.8640 | 1.76 | 89.1 | 0.0530 | 100 |
| 10.00 | 0.000 | 24.523 | 0.9280 | 1.70 | 91.8 | 0.0720 | 100 |


| Exp 69 | Flow Cell <br> Front Tilt $=18.01$ div. |  | $\begin{gathered} 0.107 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.066 \% \\ \text { Rear Translation }=12.23 \text { div. } \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=25 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt | Side | Tilt Angle | Inte | kHz) | $Y$ - | Decay Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) |  | Cumulant | (sec.) |
| 15.99 | 11.63 | 0.0000 | 96.383 | 97.181 | 0.8130 | 3.23 | 48.6 | 0.06300 | 120 |
| 16.00 | 11.71 | 0.1125 | 93.988 | 97.796 | 0.8050 | 3.20 | 49.0 | 0.05900 | 120 |
| 16.01 | 11.75 | 0.2250 | 93.341 | 96.512 | 0.7800 | 3.19 | 49.2 | 0.05200 | 120 |
| 16.02 | 11.78 | 0.3375 | 93.682 | 99.646 | 0.7620 | 3.14 | 50.1 | 0.03000 | 120 |
| 16.03 | 11.80 | 0.4500 | 93.520 | 99.253 | 0.7260 | 3.11 | 50.5 | 0.04000 | 120 |
| 16.04 | 11.84 | 0.5625 | 92.885 | 97.347 | 0.6760 | 3.12 | 50.4 | 0.03300 | 120 |
| 16.05 | 11.87 | 0.6750 | 92.482 | 99.978 | 0.6220 | 3.09 | 50.8 | 0.03300 | 120 |
| 16.06 | 11.92 | 0.7875 | 91.579 | 96.562 | 0.5440 | 3.07 | 51.2 | 0.02600 | 120 |
| 16.07 | 11.96 | 0.9000 | 92.280 | 97.020 | 0.4770 | 3.03 | 51.9 | 0.01700 | 120 |
| 16.08 | 11.96 | 1.0125 | 92.401 | 107.209 | 0.4700 | 3.02 | 52.0 | 0.00970 | 120 |
| 16.09 | 12.00 | 1.1250 | 92.117 | 105.475 | 0.3990 | 3.00 | 52.3 | -0.00140 | 120 |
| 16.10 | 12.03 | 1.2375 | 92.327 | 107.345 | 0.3430 | 3.01 | 52.1 | 0.02000 | 120 |
| 16.11 | 12.08 | 1.3500 | 92.233 | 102.482 | 0.2780 | 2.97 | 52.8 | -0.00079 | 180 |
| 16.12 | 12.12 | 1.4625 | 92.314 | 102.029 | 0.2240 | 2.90 | 54.1 | -0.04700 | 180 |
| 16.13 | 12.15 | 1.5750 | 91.943 | 102.598 | 0.1810 | 2.92 | 53.8 | -0.01800 | 300 |
| 16.14 | 12.19 | 1.6875 | 91.896 | 98.486 | 0.1420 | 2.96 | 53.0 | -0.02200 | 300 |
| 16.15 | 12.21 | 1.8000 | 91.903 | 100.762 | 0.1190 | 2.94 | 53.5 | 0.02600 | 300 |
| 16.16 | 12.25 | 1.9125 | 91.972 | 101.419 | 0.0924 | 2.84 | 54.4 | 0.00820 | 300 |
| 16.17 | 12.28 | 2.0250 | 91.750 | 101.364 | 0.0690 | 2.99 | 52.5 | -0.00017 | 420 |
| 16.18 | 12.32 | 2.1375 | 90.057 | 96.931 | 0.0497 | 2.95 | 53.3 | 0.02400 | 420 |
| 16.19 | 12.35 | 2.2500 | 89.846 | 98.880 | 0.0334 | 2.87 | 54.8 | -0.03000 | 420 |
| 16.20 | 12.39 | 2.3625 | 89.836 | 98.270 | 0.0199 | 3.25 | 48.4 | -0.01200 | 420 |

Exp 70
Repeat of Exp 69 at random points to verify repeatability

| 15.99 | 11.63 | 0.0000 | 89.4580 | 92.030 | 0.8030 | 3.26 | 48.2 | 0.06600 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.02 | 11.77 | 0.3375 | 92.9670 | 98.257 | 0.7520 | 3.16 | 49.7 | 0.05200 | 120 |
| 16.05 | 11.85 | 0.6750 | 87.8050 | 99.354 | 0.6460 | 3.09 | 50.8 | 0.02900 | 120 |
| 16.08 | 11.94 | 1.0125 | 89.5980 | 102.565 | 0.4730 | 3.06 | 51.4 | 0.02500 | 120 |
| 16.11 | 12.07 | 1.3500 | 89.2930 | 103.931 | 0.2900 | 2.96 | 53.1 | -0.00280 | 180 |
| 16.14 | 12.17 | 1.6875 | 89.2890 | 105.570 | 0.1630 | 2.90 | 54.2 | 0.01000 | 300 |
| 16.17 | 12.27 | 2.0250 | 89.2790 | 104.906 | 0.0710 | 3.00 | 52.4 | 0.06900 | 420 |
| 16.20 | 12.39 | 2.3625 | 88.4870 | 101.755 | 0.0201 | 3.37 | 46.6 | 0.21000 | 420 |

Exp 71
Flow Cell
Front Tilt $=\mathbf{1 8 . 0 1}$ div.
$\begin{array}{cc}\text { Rear Tilt } & \begin{array}{c}\text { Side } \\ \text { (div.) }\end{array} \\ \begin{array}{c}\text { Translation } \\ \text { (div.) }\end{array} & \begin{array}{c}\text { Tilt Angle } \\ \text { (mrad) }\end{array}\end{array}$

| Rear Tilt |  |
| :---: | :---: |
| (div.) | Stanslation <br> (div.) | | Tilt Angle |
| :---: |
| (mrad) |

$$
\begin{gathered}
0.107 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.066 \% \\
\text { Rear Translation }=12.23 \text { div. }
\end{gathered}
$$

| Intensity (kHz) | Y- | Decay <br> Rate |  |
| :---: | :---: | :---: | :---: |
| Ch 0 | Ch 1 | Intercept |  |
| $(1 / \mathrm{ms})$ |  |  |  |


| 100.176 | 101.437 |
| :---: | :---: |
| 97.312 | 104.902 |
| 96.200 | 101.725 |

16.01
16.02
$\begin{array}{lll}16.03 & 11.83 & 0.4500\end{array}$
$16.04 \quad 11.86 \quad 0.5625$
$\begin{array}{lll}16.05 & 11.90 & 0.6750 \\ 10.06 & 11.93 & 0.7875\end{array}$
16.06
16.07
$\begin{array}{lll}16.08 & 12.00 & 1.0125\end{array}$
$16.09 \quad 12.04 \quad 1.1250$
$16.10 \quad 12.08 \quad 1.2375$
$16.11 \quad 12.11 \quad 1.3500$
$\begin{array}{lll}16.12 & 12.15 & 1.4625 \\ 16.13 & 12.18 & 1.5750\end{array}$
$\begin{array}{lll}16.13 & 12.18 & 1.5750 \\ 16.14 & 12.21 & 1.6875\end{array}$
$16.15 \quad 12.25 \quad 1.8000$
$\begin{array}{lll}16.16 & 12.28 & 1.9125\end{array}$
$16.17 \quad 12.33 \quad 2.0250$
$\begin{array}{lll}16.18 & 12.36 & 2.1375\end{array}$

## Repeat of Exp 71 at random points to verify repeatability

| 16.00 | 11.70 | 0.1125 | 92.3340 | 101.821 | 0.7950 | 3.34 | 47.1 | 0.0620 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.03 | 11.82 | 0.4500 | 92.0590 | 100.387 | 0.7000 | 3.26 | 48.2 | 0.0550 | 120 |
| 16.06 | 11.93 | 0.7875 | 91.0770 | 93.234 | 0.5300 | 3.14 | 50.0 | 0.0360 | 120 |
| 16.09 | 12.03 | 1.1250 | 91.9700 | 100.024 | 0.3600 | 3.06 | 51.3 | 0.0031 | 120 |
| 16.12 | 12.14 | 1.4625 | 91.6440 | 98.922 | 0.2080 | 2.97 | 53.0 | 0.0210 | 180 |
| 16.15 | 12.24 | 1.8000 | 91.6940 | 99.132 | 0.1090 | 3.04 | 51.7 | 0.0410 | 300 |
| 16.18 | 12.35 | 2.1375 | 91.6730 | 100.160 | 0.0424 | 3.11 | 50.5 | 0.0240 | 420 |
| 16.19 | 12.38 | 2.2500 | 91.3870 | 97.592 | 0.0308 | 3.20 | 49.0 | 0.0830 | 420 |


| Exp 73 | Flow Cell <br> Front Tilt $=18.01$ div. |  | $\begin{aligned} & 0.107 \mu \mathrm{~m} \text { PSL } \\ & \text { Rear Transla } \end{aligned}$ |  | V.F. $=0.066 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) | Side Translation (div.) | Tilt Angle (mrad) | $\begin{aligned} & \text { Inten } \\ & \text { Ch } 0 \end{aligned}$ | $\begin{aligned} & \mathrm{kHz}) \\ & \mathrm{Ch} 1 \end{aligned}$ | Intercept | Decay Rate <br> (1/ms) | Radius (nm) | Normalized Second Cumulant | Time (sec.) |
| 16.00 | 11.70 | 0.0000 | 93.124 | 101.554 | 0.7710 | 3.43 | 45.9 | 0.08500 | 120 |
| 16.01 | 11.76 | 0.1125 | 91.844 | 100.323 | 0.7700 | 3.43 | 45.8 | 0.08200 | 120 |
| 16.02 | 11.79 | 0.2250 | 91.407 | 99.447 | 0.7310 | 3.41 | 46.1 | 0.07700 | 120 |
| 16.03 | 11.83 | 0.3375 | 91.130 | 99.162 | 0.6800 | 3.37 | 46.6 | 0.06300 | 120 |
| 16.04 | 11.87 | 0.4500 | 91.321 | 96.756 | 0.6660 | 3.32 | 47.4 | 0.06200 | 120 |
| 16.05 | 11.89 | 0.5625 | 91.089 | 99.185 | 0.5950 | 3.28 | 47.9 | 0.03900 | 120 |
| 16.06 | 11.92 | 0.6750 | 91.035 | 99.624 | 0.5580 | 3.23 | 48.7 | 0.04000 | 120 |
| 16.07 | 11.96 | 0.7875 | 91.265 | 100.647 | 0.4970 | 3.20 | 49.1 | 0.05000 | 120 |
| 16.08 | 12.00 | 0.9000 | 91.128 | 100.360 | 0.4340 | 3.13 | 50.3 | 0.01000 | 120 |
| 16.09 | 12.03 | 1.0125 | 91.129 | 101.234 | 0.3660 | 3.11 | 50.6 | 0.00590 | 120 |
| 16.10 | 12.08 | 1.1250 | 90.642 | 98.037 | 0.2960 | 3.09 | 50.8 | 0.00860 | 180 |
| 16.11 | 12.11 | 1.2375 | 90.999 | 98.877 | 0.2600 | 3.01 | 52.2 | -0.02300 | 180 |
| 16.12 | 12.15 | 1.3500 | 90.879 | 100.588 | 0.2160 | 3.07 | 51.3 | 0.01500 | 300 |
| 16.13 | 12.18 | 1.4625 | 90.703 | 101.921 | 0.1710 | 3.07 | 51.1 | 0.02400 | 300 |
| 16.14 | 12.21 | 1.5750 | 90.706 | 99.225 | 0.1470 | 3.01 | 52.1 | -0.01300 | 300 |
| 16.15 | 12.25 | 1.6875 | 90.892 | 100.228 | 0.1130 | 3.01 | 52.1 | 0.01100 | 300 |
| 16.16 | 12.28 | 1.8000 | 90.779 | 99.507 | 0.0927 | 2.95 | 53.2 | -0.00083 | 420 |
| 16.17 | 12.32 | 1.9125 | 90.720 | 99.269 | 0.0652 | 3.00 | 52.3 | -0.00280 | 420 |
| 16.18 | 12.35 | 2.0250 | 90.647 | 100.953 | 0.0492 | 3.20 | 49.2 | 0.01400 | 420 |

Exp $74 \quad$ Repeat of Exp 73 at random points to verify repeatability

| 16.03 | 11.82 | 0.3375 | 89.2920 | 97.468 | 0.7030 | 3.32 | 47.3 | 0.0700 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.08 | 12.01 | 0.9000 | 90.1530 | 98.759 | 0.4620 | 3.07 | 51.2 | 0.0130 | 120 |
| 16.12 | 12.16 | 1.3500 | 90.1000 | 97.756 | 0.2480 | 3.05 | 51.5 | 0.0130 | 240 |
| 16.15 | 12.25 | 1.6875 | 90.0900 | 100.291 | 0.1430 | 3.04 | 51.6 | 0.0300 | 300 |


| Exp 75 | Flow Cell <br> Front Tilt $=18.01$ div. |  | $\begin{gathered} 0.107 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.066 \% \\ \text { Rear Translation }=12.23 \text { div. } \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=50 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) | Side | Tilt Angle | Inte | kHz) | Y- | Decay <br> Rate | Radius | Normalized Second | Time |
|  | Translation (div.) | (mrad) | Ch 0 | Ch 1 | Intercept | (1/ms) |  | Cumulant |  |
| 16.00 | 11.70 | 0.0000 | 89.663 | 99.434 | 0.8020 | 3.68 | 42.7 | 0.1200 | 120 |
| 16.01 | 11.73 | 0.1125 | 88.699 | 103.900 | 0.7930 | 3.65 | 43.0 | 0.1100 | 120 |
| 16.02 | 11.73 | 0.2250 | 89.237 | 99.858 | 0.7860 | 3.61 | 43.6 | 0.0990 | 120 |
| 16.03 | 11.83 | 0.3375 | 92.450 | 100.135 | 0.6990 | 3.53 | 44.5 | 0.0860 | 120 |
| 16.04 | 11.87 | 0.4500 | 92.564 | 99.513 | 0.6500 | 3.47 | 45.3 | 0.0770 | 120 |
| 16.05 | 11.90 | 0.5625 | 91.827 | 97.968 | 0.6000 | 3.41 | 46.1 | 0.0680 | 120 |
| 16.06 | 11.94 | 0.6750 | 92.397 | 99.752 | 0.5450 | 3.34 | 47.0 | 0.0500 | 120 |
| 16.07 | 11.98 | 0.7875 | 92.118 | 98.654 | 0.4820 | 3.31 | 47.4 | 0.0530 | 120 |
| 16.08 | 12.01 | 0.9000 | 92.103 | 99.432 | 0.4240 | 3.27 | 48.0 | 0.0310 | 120 |
| 16.09 | 12.04 | 1.0125 | 92.176 | 99.682 | 0.3710 | 3.24 | 48.5 | 0.0330 | 120 |
| 16.10 | 12.08 | 1.1250 | 91.921 | 99.960 | 0.3250 | 3.17 | 49.6 | 0.0083 | 180 |
| 16.11 | 12.12 | 1.2375 | 91.632 | 98.774 | 0.2680 | 3.12 | 50.4 | -0.0150 | 180 |
| 16.12 | 12.15 | 1.3500 | 91.316 | 99.962 | 0.2310 | 3.13 | 50.3 | 0.0180 | 180 |
| 16.13 | 12.18 | 1.4625 | 91.599 | 99.285 | 0.1900 | 3.12 | 50.4 | 0.0320 | 180 |
| 16.14 | 12.22 | 1.5750 | 91.067 | 100.822 | 0.1550 | 3.14 | 50.1 | -0.0150 | 180 |
| 16.15 | 12.25 | 1.6875 | 91.298 | 100.465 | 0.1300 | 2.98 | 52.7 | -0.0230 | 300 |
| 16.16 | 12.29 | 1.8000 | 91.189 | 99.810 | 0.0981 | 3.01 | 52.2 | -0.0410 | 300 |
| 16.17 | 12.32 | 1.9125 | 91.025 | 100.110 | 0.0798 | 3.07 | 51.2 | 0.0350 | 300 |
| 16.18 | 12.36 | 2.0250 | 90.903 | 98.882 | 0.0588 | 3.12 | 50.3 | 0.0660 | 300 |
| 16.19 | 12.39 | 2.1375 | 90.701 | 100.552 | 0.0427 | 3.17 | 49.6 | 0.1300 | 420 |

Exp 76
Repeat of Exp 75 at random points to verify repeatability

| 16.00 | 11.70 | 0.0000 | 88.5150 | 97.700 | 0.7990 | 3.61 | 43.5 | 0.1100 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.04 | 11.87 | 0.4500 | 88.3630 | 97.485 | 0.6520 | 3.42 | 45.9 | 0.0780 | 120 |
| 16.08 | 12.02 | 0.9000 | 87.9990 | 97.614 | 0.4290 | 3.25 | 48.3 | 0.0280 | 120 |
| 16.12 | 12.16 | 1.3500 | 88.3700 | 97.715 | 0.2300 | 3.17 | 49.6 | 0.0480 | 180 |
| 16.16 | 12.29 | 1.8000 | 88.3180 | 98.765 | 0.1030 | 3.06 | 51.4 | 0.0077 | 300 |


| Exp 77 | Fow Cell <br> Front Tilt $=18.02 \mathrm{dv}$. |  | $\begin{gathered} 0.107 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.198 \% \\ \text { Rear Translation }=123 \text { div. } \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Fow Rate }=0 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) | Side |  |  | (k-z) |  |  |  | Normalized |  |
|  |  |  | Cho | Ch 1 | Interce | (1/ms) |  | Curnul | (sec.) |
| 16.00 | 11.70 | 0.0000 | 93.279 | 101.081 | 0.7580 | 3.58 | 44.1 | 0.0990 | 120 |
| 16.01 | 11.77 | 0.1125 | 92.919 | 99.043 | 0.6900 | 3.53 | 44.5 | 0.0960 | 120 |
| 16.02 | 11.80 | 0.2250 | 92588 | 99.955 | 0.6610 | 3.49 | 45.0 | 0.0900 | 120 |
| 16.03 | 11.84 | 0.3375 | 92148 | 99.209 | 0.5970 | 3.45 | 45.5 | 0.0920 | 120 |
| 16.04 | 11.87 | 0.4500 | 93.21 | 100.120 | 0.5470 | 3.41 | 46.0 | 0.1000 | 120 |
| 16.05 | 11.91 | 0.5625 | 93.086 | 98.237 | 0.4820 | 3.33 | 47.1 | 0.0760 | 120 |
| 16.06 | 11.94 | 0.6750 | 93.046 | 99.058 | 0.4340 | 3.29 | 47.8 | 0.0660 | 120 |
| 16.07 | 11.98 | 0.7875 | 92.619 | 96.482 | 0.3700 | 3.20 | 49.1 | 0.0590 | 120 |
| 16.08 | 1201 | 0.9000 | 93.179 | 98.033 | 0.3190 | 3.19 | 49.3 | 0.0670 | 120 |
| 16.09 | 1204 | 1.0125 | 92.218 | 99.302 | 0.2770 | 3.17 | 49.6 | 0.0570 | 180 |
| 16.10 | 1207 | 1.1250 | 91.785 | 102089 | 0.2480 | 3.16 | 49.7 | 0.0570 | 180 |
| 16.11 | 1212 | 1.2375 | 91.826 | 99.498 | 0.1930 | 3.12 | 50.4 | 0.0260 | 240 |
| 16.12 | 1216 | 1.3500 | 91.736 | 98.252 | 0.1540 | 3.04 | 51.7 | 0.0019 | 240 |
| 16.13 | 1219 | 1.4625 | 91.526 | 97.775 | 0.1280 | 295 | 53.3 | 0.0093 | 240 |
| 16.14 | 1223 | 1.5750 | 91.594 | 97.802 | 0.1090 | 3.01 | 521 | 0.0340 | 300 |
| 16.15 | 1226 | 1.6875 | 91.715 | 98.085 | 0.0897 | 3.03 | 51.9 | 0.0630 | 300 |
| 16.16 | 1229 | 1.8000 | 91.644 | 98.426 | 0.0727 | 289 | 54.4 | -0.0350 | 420 |
| 16.17 | 1232 | 1.9125 | 91.608 | 99.210 | 0.0609 | 3.03 | 51.8 | 0.0650 | 420 |
| 16.18 | 1236 | 20250 | 91.491 | 98.358 | 0.0460 | 3.08 | 51.0 | 0.0350 | 420 |
| 16.19 | 1239 | 21375 | 91.792 | 98.753 | 0.0345 | 3.16 | 49.7 | 0.1400 | 420 |

Exp $78 \quad$ Repeat of Exp 77 at random points to verify repeatability

| 16.00 | 11.70 | 0.0000 | 93.902 | 102306 | 0.7510 | 3.69 | 42.5 | 0.1100 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.04 | 11.88 | 0.4500 | 91.895 | 98.764 | 0.5210 | 3.42 | 45.9 | 0.0860 | 120 |
| 16.08 | 1202 | 0.9000 | 91.828 | 98.367 | 0.2840 | 3.19 | 49.2 | 0.0400 | 120 |
| 16.12 | 1217 | 1.3500 | 90.762 | 96.612 | 0.1280 | 298 | 527 | 0.0480 | 240 |
| 16.16 | 1230 | 1.8000 | 92806 | 97.575 | 0.0511 | 282 | 55.7 | -0.0460 | 420 |

Exp 79
Flow Cell
Front Tilt $=18.02$ div.

| Rear | Side | Tilt |
| :---: | :---: | :---: |
| Tilt | Translation | Angle |
| (div.) | (div.) | (mrad) |

$16.00 \quad 1170 \quad 0.000$
$16.01 \quad 11.79 \quad 0.1125$
$16.02 \quad 11.83 \quad 0.2250$
$16.03 \quad 11.86 \quad 0.3375$
$16.04 \quad 11.89 \quad 0.4500$
$16.05 \quad 11.92 \quad 0.5625$
$16.06 \quad 11.96 \quad 0.6750$
$16.07 \quad 11.99 \quad 0.7875$
$16.08 \quad 12.03 \quad 0.9000$ $16.09 \quad 12.06 \quad 1.0125$
$16.10 \quad 12.09 \quad 1.1250$
$\begin{array}{lll}16.10 & 12.14 & 1.2375 \\ 16.12 & 12.18 & 1.3500\end{array}$
$16.13 \quad 12.21 \quad 1.4625$
$16.14 \quad 12.24 \quad 1.5750$
$\begin{array}{lll}16.15 & 12.28 & 1.6875 \\ 16.16 & 12.31 & 1.8000\end{array}$
$\begin{array}{lll}16.17 & 12.35 & 1.9125 \\ 16.18 & 12.38 & 2.0250\end{array}$
$0.107 \mu \mathrm{~m}$ PSL V.F. $=0.198 \%$
Rear Translation $=12.23 \mathrm{div}$.
$\theta=112 \mathrm{deg}$
Flow Rate $=25 \%$
Normalized Time

| Intensity (kHz) | Y- | Decay <br> Rate |
| :--- | :---: | :---: |
| Ch 0 | Ch 1 | Intercept |
| $(1 / \mathrm{ms})$ |  |  |

$\begin{array}{llll}89.259 & 100.451 & 0.7230 & 4.08\end{array}$
$88.263 \quad 97.518 \quad 0.6320 \quad 3.86$
$\begin{array}{llll}89.072 & 99.527 & 0.5860 & 3.86\end{array}$
$\begin{array}{llll}89.169 & 100.247 & 0.5360 & 3.76\end{array}$
$\begin{array}{llll}89.289 & 99.850 & 0.4790 & 3.67 \\ 89.264 & 99.619 & 0.4260 & 3.60\end{array}$
$89.281 \quad 99.460 \quad 0.3730 \quad 3.45$
$89.213 \quad 99.458 \quad 0.3170 \quad 3.42$
$89.205 \quad 100.792 \quad 0.2650$
$\begin{array}{llll}89.825 & 101.524 & 0.2230 & 3.28 \\ 89.146 & 100.947 & 0.1820 & 3.26\end{array}$
$\begin{array}{cccc}89.146 & 100.947 & 0.1820 & 3.26 \\ 89.100 & 99.267 & 0.1390 & 3.03\end{array}$
$89.100 \quad 99.267 \quad 0.1390$
$\begin{array}{llll}88.766 & 99.083 & 0.1130 & 3.10\end{array}$
$\begin{array}{lll}88.871 & 97.644 & 0.0884\end{array}$
$88.754 \quad 99.489 \quad 0.0743 \quad 3.08$
$88.682 \quad 99.253 \quad 0.0500 \quad 2.88$
$88.742 \quad 99.879 \quad 0.0423 \quad 2.99$
$\begin{array}{llll}88.538 & 98.904 & 0.0314 & 3.09 \\ 88.695 & 98.120 & 0.0245 & 3.07\end{array}$

Repeat of Exp 79 at random points to verify repeatability
Exp 80

| 16.00 | 11.70 | 0.0000 | 90.0330 | 100.214 | 0.7400 | 4.01 | 39.2 | 0.0920 | 120 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 16.04 | 11.89 | 0.4500 | 90.0360 | 99.939 | 0.4960 | 3.64 | 43.1 | 0.0760 | 180 |
| 16.08 | 12.03 | 0.9000 | 90.0860 | 100.169 | 0.2840 | 3.26 | 48.2 | 0.0290 | 300 |
| 16.12 | 12.17 | 1.3500 | 89.9820 | 99.566 | 0.1380 | 3.06 | 51.3 | -0.0074 | 300 |
| 16.16 | 12.30 | 1.8000 | 89.8250 | 100.401 | 0.0584 | 3.01 | 52.2 | -0.0140 | 420 |


| Exp 81 | Flow Cell |  | $0.107 \mu \mathrm{mPSL}$ |  | V.F. $=0.198 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) |  |  | Inte | (kHz) |  |  |  |  |  |
|  |  |  | Ch 0 | Ch 1 |  | (1/ms) |  |  | (sec.) |
| 16.00 | 11.70 | 0.0000 | 90.254 | 100.828 | 0.7400 | 5.27 | 29.8 | 0.1800 | 120 |
| 16.01 | 11.80 | 0.1125 | 89.927 | 99.676 | 0.6500 | 4.91 | 32.0 | 0.1600 | 20 |
| 16.02 | 11.83 | 0.2250 | 89.950 | 100.234 | 0.6000 | 4.73 | 33.2 | 0.1500 | 120 |
| 16.03 | 11.86 | 0.3375 | 89.705 | 99.486 | 0.5490 | 4.60 | 34.1 | 0.1400 | 120 |
| 16.04 | 11.89 | 0.4500 | 89.842 | 100.372 | 0.5120 | 4.43 | 35.5 | 0.1400 | 120 |
| 16.05 | 11.92 | 0.5625 | 89.751 | 99.451 | 0.4500 | 4.29 | 36.7 | 0.1200 | 120 |
| 16.06 | 11.95 | 0.6750 | 89.593 | 100.052 | 0.4010 | 4.14 | 38.0 | 0.1100 | 120 |
| 16.07 | 11.99 | 0.7875 | 89.855 | 100.073 | 0.3490 | 3.97 | 39.6 | 0.0980 | 180 |
| 16.08 | 12.02 | 0.9000 | 89.768 | 101.210 | 0.2960 | 3.83 | 41.1 | 0.0970 | 180 |
| 16.09 | 12.07 | 1.0125 | 89.866 | 99.841 | 0.2390 | 3.6 | 43. | 0.0560 | 180 |
| 16.10 | 12.10 | 1.1250 | 89.613 | 100.061 | 0.2020 | 3.48 | 45.1 | 0.0390 | 80 |
| 16.11 | 12.13 | 1.2375 | 89.729 | 99.997 | 0.1660 | 3.44 | 45.6 | 0.0270 | 240 |
| 16.12 | 12.17 | 1.3500 | 89.635 | 100.618 | 0.1370 | 3.43 | 45.8 | 0.0530 | 240 |
| 16.13 | 12.20 | 1.4625 | 89.908 | 101.293 | 0.1150 | 3.20 | 49.0 | -0.0015 | 240 |
| 16.14 | 12.24 | 1.5750 | 89.610 | 100.140 | 0.0865 | 3.18 | 49.4 | -0.0043 | 300 |
| 16.15 | 12.27 | 1.6875 | 89.540 | 100.289 | 0.0711 | 3.17 | 49.5 | 0.0031 | 300 |
| 16.16 | 12.30 | 1.8000 | 89.644 | 100.708 | 0.0573 | 3.11 | 50.4 | 0.0042 | 360 |
| 16.17 | 12.34 | 1.9125 | 89.461 | 101.209 | 0.0420 | 2.89 | 54.3 | -0.2000 | 360 |
| 16.18 | 12.37 | 2.0250 | 88.489 | 98.303 | 0.0332 | 2.88 | 54.6 | -0.1200 | 420 |

## Exp $82 \quad$ Repeat of $\operatorname{Exp} 81$ at random points to verity repeatability

| 16.00 | 11.70 | 0.0000 | 87.8220 | 98.322 | 0.7380 | 5.37 | 29.3 | 0.1800 | 120 |
| :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| 16.04 | 11.88 | 0.4500 | 88.0750 | 101.124 | 0.5170 | 4.42 | 35.6 | 0.1200 | 120 |
| 16.08 | 12.01 | 0.9000 | 87.9660 | 101.056 | 0.3030 | 3.80 | 41.4 | 0.0710 | 240 |
| 16.12 | 12.17 | 1.3500 | 87.9360 | 98.051 | 0.1310 | 3.41 | 46.1 | 0.0440 | 240 |
| 16.16 | 12.31 | 1.8000 | 87.6700 | 100.784 | 0.0548 | 3.07 | 51.2 | -0.0083 | 420 |


| Exp 83 | Flow Cell <br> Front Tilt $=18.02$ div. |  | $\begin{aligned} & 0.107 \mu \mathrm{~m} \text { PSL V.F. }=0.198 \% \\ & \text { Rear Translation }=12.23 \text { div. } \end{aligned}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=100 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) | Side Translation (div.) | Tilt Angle (mrad) | Ch 0 | (kHz) Ch 1 | $Y$ - <br> Intercept | $\begin{aligned} & \text { Decay } \\ & \text { Rate } \end{aligned}$ (1/ms) | Radius (nm) | Normalized Second Cumulant | Time (sec.) |
| 16.00 | 11.70 | 0.0000 | 88.209 | 98.791 | 0.7350 | 6.50 | 24.2 | 0.2200 | 120 |
| 16.01 | 11.73 | 0.1125 | 88.213 | 100.974 | 0.7160 | 6.26 | 25.1 | 0.2200 | 120 |
| 16.02 | 11.75 | 0.2250 | 87.670 | 100.926 | 0.6910 | 6.11 | 25.7 | 0.2000 | 120 |
| 16.03 | 11.77 | 0.3375 | 88.009 | 101.683 | 0.6490 | 5.86 | 26.8 | 0.2000 | 120 |
| 16.04 | 11.88 | 0.4500 | 88.181 | 99.089 | 0.5120 | 5.30 | 29.7 | 0.1800 | 120 |
| 16.05 | 11.92 | 0.5625 | 87.859 | 99.409 | 0.4560 | 5.08 | 30.9 | 0.1800 | 180 |
| 16.06 | 11.96 | 0.6750 | 87.894 | 98.002 | 0.3950 | 4.74 | 33.1 | 0.1500 | 180 |
| 16.07 | 11.99 | 0.7875 | 87.973 | 98.076 | 0.3390 | 4.50 | 34.9 | 0.1400 | 180 |
| 16.08 | 12.03 | 0.9000 | 87.912 | 98.650 | 0.2910 | 4.37 | 36.0 | 0.1400 | 240 |
| 16.09 | 12.07 | 1.0125 | 87.934 | 97.922 | 0.2380 | 4.16 | 37.7 | 0.1300 | 240 |
| 16.10 | 12.10 | 1.1250 | 88.002 | 98.215 | 0.1970 | 3.98 | 39.4 | 0.1100 | 300 |
| 16.11 | 12.13 | 1.2375 | 87.938 | 99.646 | 0.1680 | 3.83 | 41.1 | 0.0760 | 300 |
| 16.12 | 12.17 | 1.3500 | 87.929 | 99.529 | 0.1390 | 3.72 | 42.2 | 0.0970 | 300 |
| 16.13 | 12.20 | 1.4625 | 88.026 | 98.887 | 0.1080 | 3.63 | 43.3 | 0.0670 | 420 |
| 16.14 | 12.24 | 1.5750 | 87.897 | 99.110 | 0.0880 | 3.54 | 44.4 | 0.0610 | 420 |
| 16.15 | 12.27 | 1.6875 | 88.245 | 99.329 | 0.0740 | 3.45 | 45.5 | 0.0780 | 840 |
| 16.16 | 12.31 | 1.8000 | 88.337 | 98.656 | 0.0555 | 3.27 | 48.1 | 0.0000 | 840 |
| 16.17 | 12.34 | 1.9125 | 88.324 | 100.398 | 0.0427 | 3.22 | 48.8 | 0.0069 | 840 |
| 16.18 | 12.37 | 2.0250 | 88.237 | 99.675 | 0.0309 | 3.16 | 49.7 | -0.0400 | 840 |
| 16.19 | 12.41 | 2.1375 | 88.221 | 99.270 | 0.0225 | 3.0 | 51.6 | -0.0440 | 60 |

Exp 84
Repeat of Exp 83 at random points to verify repeatability

| 16.00 | 11.70 | 0.0000 | 87.5940 | 98.642 | 0.7390 | 6.40 | 24.6 | 0.2200 | 120 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 16.04 | 11.89 | 0.4500 | 87.7010 | 99.272 | 0.4960 | 5.20 | 30.2 | 0.1800 | 120 |
| 16.08 | 12.03 | 0.9000 | 87.6290 | 99.547 | 0.2890 | 4.25 | 37.0 | 0.1300 | 240 |
| 16.12 | 12.17 | 1.3500 | 87.6970 | 100.147 | 0.1360 | 3.60 | 43.6 | 0.0880 | 300 |
| 16.16 | 12.31 | 1.8000 | 88.0050 | 99.835 | 0.0547 | 3.35 | 46.9 | 0.0720 | 480 |

Exp 86
Flow Cell
Front Tilt $=18.12$ div.

| Rear <br> Tilt <br> (div.) | Side <br> Translation <br> (div.) | Tilt Angle <br> (mrad) |
| :---: | :---: | :---: |
| 16.08 | 11.87 | 0.000 |
| 16.09 | 11.93 | 0.113 |
| 16.10 | 11.95 | 0.225 |
| 16.11 | 12.02 | 0.338 |
| 16.12 | 12.06 | 0.450 |
| 16.13 | 12.10 | 0.563 |
| 16.14 | 12.13 | 0.675 |
| 16.15 | 12.15 | 0.788 |
| 16.16 | 12.18 | 0.900 |
| 16.17 | 12.21 | 1.013 |
| 16.18 | 12.27 | 1.125 |
| 16.19 | 12.32 | 1.238 |
| 16.20 | 12.36 | 1.350 |
| 16.21 | 12.38 | 1.463 |
| 16.22 | 12.40 | 1.575 |
| 16.23 | 12.45 | 1.688 |
| 16.24 | 12.47 | 1.800 |

$0.107 \mu \mathrm{~m}$ PSL V.F. $=0.198 \%$
Rear Translation $=12.27$ div.

| Intensity ( kHz ) |  | Y-Intercept | Decay Rate (1/ms) | Radius (nm) |
| :---: | :---: | :---: | :---: | :---: |
| Ch 0 | Ch 1 |  |  |  |
| 81.284 | 96.266 | 0.8330 | 3.52 | 44.7 |
| 83.322 | 98.066 | 0.7440 | 3.48 | 45.1 |
| 83.826 | 95.619 | 0.7360 | 3.45 | 45.6 |
| 82.570 | 98.248 | 0.6390 | 3.41 | 44.7 |
| 83.791 | 99.288 | 0.5700 | 3.39 | 45.0 |
| 84.677 | 98.219 | 0.5060 | 3.35 | 45.5 |
| 84.704 | 98.358 | 0.4580 | 3.25 | 46.8 |
| 83.815 | 98.852 | 0.4190 | 3.29 | 46.3 |
| 83.102 | 100.827 | 0.3610 | 3.19 | 47.7 |
| 82.335 | 101.650 | 0.3130 | 3.23 | 47.3 |
| 82.917 | 101.798 | 0.2570 | 3.19 | 47.7 |
| 84.435 | 99.543 | 0.2010 | 3.17 | 48.0 |
| 85.324 | 99.610 | 0.1630 | 3.12 | 48.8 |
| 84.268 | 101.233 | 0.1400 | 3.04 | 50.0 |
| 82.914 | 99.888 | 0.1210 | 3.09 | 49.3 |
| 84.277 | 98.400 | 0.0901 | 2.98 | 51.1 |
| 83.138 | 99.673 | 0.0773 | 3.12 | 48.8 |

0.107 mm PSL V.F. $=0.198 \%$

Rear Translation $=12.27$ div.

Exp 87
Flow Cell
Front Tilt $=\mathbf{1 8 . 1 2} \mathbf{~ d i v}$.

| Rear <br> Tilt <br> (div.) | Side <br> Translation <br> (div.) | Tilt Angle <br> (mrad) |
| :---: | :---: | :---: |
| 16.08 | 11.87 | 0.0000 |
| 16.09 | 11.96 | 0.1125 |
| 16.10 | 11.99 | 0.2250 |
| 16.11 | 12.04 | 0.3375 |
| 16.12 | 12.07 | 0.4500 |
| 16.13 | 12.12 | 0.5625 |
| 16.14 | 12.17 | 0.6750 |
| 16.15 | 12.21 | 0.7875 |
| 16.16 | 12.25 | 0.9000 |
| 16.17 | 12.28 | 1.0125 |
| 16.18 | 12.31 | 1.1250 |
| 16.19 | 12.35 | 1.2375 |
| 16.20 | 12.40 | 1.3500 |
| 16.21 | 12.44 | 1.4625 |
| 16.22 | 12.48 | 1.5750 |
| 16.23 | 12.51 | 1.6875 |
| 16.24 | 12.54 | 1.8000 |
| 16.25 | 12.58 | 1.9125 |


| Exp 93 | Flow Cell <br> Front Tilt $=18.12$ div. |  | $0.107 \mu \mathrm{~m}$ PSL $\quad$ V.F. $=0.198 \%$Rear Translation $=12.27$ div. |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=0 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) |  |  |  | kHz) |  |  |  | Normalized Second |  |
|  |  |  | Cho | Ch 1 |  | (1/ms) |  | Cumulant |  |
| 16.04 | 11.80 | 0.0000 | 114.759 | 146.581 | 0.9090 | 3.31 | 46.7 | 0.05500 | 20 |
| 16.05 | 11.83 | 0.1125 | 114.547 | 146.792 | 0.9030 | 3.3 | 46.7 | 0.05000 | 20 |
| 16.06 | 11.89 | 0.2250 | 114.407 | 143.958 | 0.8850 | 3.28 | 47.1 | 0.06100 | 20 |
| 16.08 | 11.96 | 0.4500 | 114.857 | 146.804 | 0.8530 | 3.21 | 48.2 | 0.04700 | 120 |
| 16.10 | 12.04 | 0.6750 | 114.702 | 145.092 | 0.7960 | 3.28 | 48.6 | 0.03900 | 120 |
| 16.12 | 12.12 | 0.9000 | 114.304 | 147.589 | 0.7480 | 3.16 | 49.0 | 0.04200 | 120 |
| 16.14 | 12.20 | 1.1250 | 113.113 | 150.027 | 0.7020 | 3.09 | 50.2 | 0.03400 | 120 |
| 16.16 | 12.27 | 1.3500 | 114.656 | 147.925 | 0.6700 | 3.03 | 51.0 | 0.00950 | 120 |
| 16.18 | 12.35 | 1.5750 | 115.055 | 146.796 | 0.6350 | 3.02 | 51.3 | 0.01700 | 120 |
| 16.20 | 12.43 | 1.8000 | 114.946 | 145.883 | 0.6000 | 2.99 | 51.7 | 0.00730 | 120 |
| 16.22 | 12.50 | 2.0250 | 114.599 | 146.715 | 0.5750 | 2.97 | 52.2 | 0.01000 | 120 |
| 16.24 | 12.58 | 2.2500 | 115.581 | 148.146 | 0.5500 | 2.98 | 51.9 | 0.00790 | 120 |
| 16.26 | 12.65 | 2.4750 | 114.713 | 144.648 | 0.5290 | 2.95 | 52.4 | 0.01700 | 120 |
| 16.28 | 12.73 | 2.7000 | 114.489 | 145.750 | 0.5000 | 2.98 | 52.0 | 0.02800 | 120 |
| 16.30 | 12.79 | 2.9250 | 115.295 | 143.078 | 0.4790 | 2.96 | 52.3 | 0.01200 | 120 |
| 16.32 | 12.87 | 3.1500 | 117.121 | 147.693 | 0.4570 | 2.94 | 52.6 | 0.00045 | 120 |
| 16.34 | 12.95 | 3.3750 | 115.383 | 142.492 | 0.4300 | 2.96 | 52.3 | 0.01800 | 120 |
| 16.36 | 13.02 | 3.6000 | 115.714 | 146.768 | 0.4110 | 2.96 | 52.3 | 0.02600 | 120 |
| 16.38 | 13.11 | 3.8250 | 115.067 | 143.877 | 0.3850 | 2.97 | 52.0 | 0.01700 | 120 |
| 16.40 | 13.17 | 4.0500 | 115.522 | 145.983 | 0.3700 | 2.95 | 52.5 | 0.01500 | 120 |
| 16.43 | 13.29 | 4.3875 | 116.021 | 145.791 | 0.3400 | 2.85 | 54.3 | -0.01400 | 120 |
| 16.45 | 13.36 | 4.6125 | 115.724 | 145.424 | 0.3220 | 2.93 | 52.9 | 0.00290 | 120 |
| 16.47 | 13.43 | 4.8375 | 115.474 | 145.591 | 0.3090 | 2.94 | 52.7 | -0.01000 | 120 |
| 16.49 | 13.50 | 5.0625 | 115.511 | 145.857 | 0.2910 | 2.90 | 53.4 | -0.01400 | 120 |
| 16.51 | 13.58 | 5.2875 | 115.404 | 146.035 | 0.2760 | 2.92 | 53.0 | 0.00400 | 120 |
| 16.53 | 13.65 | 5.5125 | 115.569 | 144.355 | 0.2610 | 3.03 | 51.1 | 0.04000 | 120 |
| 16.55 | 13.72 | 5.7375 | 115.548 | 143.562 | 0.2490 | 2.98 | 52.0 | 0.02000 | 120 |
| 16.57 | 13.80 | 5.9625 | 115.792 | 142.496 | 0.2340 | 2.97 | 52.0 | 0.00830 | 120 |
| 16.59 | 13.87 | 6.1875 | 115.670 | 142.910 | 0.2240 | 2.89 | 53.3 | -0.02100 | 120 |
| 16.61 | 13.93 | 6.4125 | 115.403 | 144.216 | 0.2110 | 2.99 | 51.8 | -0.00045 | 120 |
| 16.6 | 14.01 | 6.6375 | 114.855 | 142.493 | 0.1960 | 2.91 | 52.2 | 0.02100 | 120 |
| 16.65 | 14.08 | 6.8625 | 114.614 | 141.048 | 0.1890 | 2.93 | 52.8 | 0.03300 | 120 |
| 16.67 | 14.14 | 7.0875 | 114.917 | 146.214 | 0.1790 | 2.84 | 54.5 | 0.02400 | 120 |
| 16.69 | 14.23 | 7.3125 | 119.281 | 138.562 | 0.1630 | 2.95 | 52.5 | 0.02300 | 180 |
| 16.71 | 14.30 | 7.5375 | 115.330 | 143.105 | 0.1560 | 2.91 | 53.2 | 0.00530 | 180 |
| 16.74 | 14.43 | 7.8750 | 114.748 | 144.387 | 0.1410 | 2.96 | 52.3 | -0.00230 | 180 |
| 16.77 | 14.54 | 8.2125 | 115.129 | 144.895 | 0.1290 | 2.93 | 52.8 | -0.04100 | 180 |
| 16.80 | 14.66 | 8.5500 | 115.475 | 145.015 | 0.1180 | 2.86 | 54.2 | -0.03700 | 240 |
| 16.83 | 14.78 | 8.8875 | 115.424 | 144.625 | 0.1050 | 2.91 | 53.2 | -0.01800 | 24 |
| 16.86 | 14.88 | 9.2250 | 115.398 | 145.536 | 0.0974 | 3.02 | 51.3 | -0.02500 | 30 |
| 16.89 | 14.96 | 9.5625 | 115.926 | 142.215 | 0.0918 | 2.87 | 54.0 | -0.06000 | 30 |


| Exp 94 | Fow Cell <br> Front Tilt $=18.12$ div. |  | $\begin{gathered} 0.107 \mu \mathrm{mPSL} \quad \text { V.F. }=0.198 \% \\ \text { Rear Translation }=1227 \mathrm{div} . \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=50 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Tilt Angle | Inten | - z ) |  | Decay | Radius | Normalized Second |  |
|  | (div.) |  | Cho | Ch1 | Interce | (1/ms) |  | Curmulant |  |
| 16.04 | 11.80 | 0.0000 | 118.276 | 146.859 | 0.9190 | 3.76 | 41.1 | 0.11000 | 120 |
| 16.05 | 11.85 | 0.1125 | 118.372 | 146.384 | 0.9140 | 3.76 | 41.1 | 0.10000 | 20 |
| 16.06 | 11.88 | 0.2250 | 118.288 | 146.323 | 0.9020 | 3.74 | 41.4 | 0.09800 | 20 |
| 16.08 | 11.96 | 0.4500 | 118.156 | 146.137 | 0.8640 | 3.65 | 42.4 | 0.08800 | 120 |
| 16.10 | 1203 | 0.6750 | 118.231 | 146.293 | 0.8220 | 3.45 | 44.8 | 0.06500 | 120 |
| 16.12 | 12.12 | 0.9000 | 118.013 | 145.154 | 0.7560 | 3.34 | 46.3 | 0.05800 | 120 |
| 16.14 | 1220 | 1.1250 | 118.288 | 144.817 | 0.7050 | 3.27 | 47.3 | 0.04900 | 120 |
| 16.16 | 12.27 | 1.3500 | 118.531 | 145.269 | 0.6670 | 3.12 | 49.6 | 0.02600 | 120 |
| 16.18 | 1234 | 1.5550 | 118.846 | 146.896 | 0.6280 | 3.10 | 49.9 | 0.03200 | 120 |
| 16.20 | 1242 | 1.8000 | 118.184 | 147.666 | 0.5930 | 3.08 | 50.2 | 0.03100 | 120 |
| 16.22 | 12.49 | 20250 | 118.282 | 147.242 | 0.5660 | 3.01 | 51.5 | -0.00032 | 120 |
| 16.24 | 1256 | 2.2500 | 118.108 | 147.175 | 0.5380 | 3.01 | 51.5 | 0.01500 | 120 |
| 16.26 | 12.64 | 24750 | 117.877 | 147.271 | 0.5050 | 2.99 | 51.7 | 0.00500 | 120 |
| 16.28 | 1272 | 27000 | 117.96 | 146.829 | 0.4720 | 298 | 52.0 | 0.00480 | 120 |
| 16.30 | 1280 | 29250 | 116.584 | 145.471 | 0.4400 | 3.00 | 51.6 | 0.00960 | 120 |
| 16.32 | 1288 | 3.1500 | 117.875 | 147.171 | 0.4100 | 297 | 52.0 | -0.00370 | 120 |
| 16.34 | 1297 | 3.3750 | 122.194 | 147.62 | 0.3780 | 2.91 | 53.2 | -0.03400 | 120 |
| 16.36 | 13.03 | 3.6000 | 111.456 | 145.233 | 0.3580 | 292 | 53.0 | -0.01200 | 120 |
| 16.38 | 13.12 | 3.8250 | 118.218 | 147.386 | 0.3270 | 298 | 520 | 0.00210 | 120 |
| 16.40 | 13.19 | 4.0500 | 117.507 | 148.234 | 0.3070 | 2.97 | 52.1 | -0.00270 | 120 |
| 16.42 | 13.27 | 4.2750 | 117.283 | 145.419 | 0.2810 | 295 | 525 | -0.00340 | 120 |
| 16.44 | 13.35 | 4.5000 | 118.815 | 144.088 | 0.2590 | 3.01 | 51.4 | 0.00170 | 120 |
| 16.46 | 13.40 | 4.7250 | 117.515 | 147.537 | 0.2470 | 3.01 | 51.4 | 0.01600 | 120 |
| 16.48 | 13.47 | 4.9500 | 117.230 | 147.173 | 0.2270 | 2.91 | 53.2 | -0.01900 | 120 |
| 16.50 | 13.56 | 5.1750 | 117.280 | 144.919 | 0.2090 | 203 | 51.1 | 0.02200 | 120 |
| 16.52 | 13.62 | 5.4000 | 117.511 | 145.890 | 0.1940 | 291 | 53.1 | -0.01700 | 120 |
| 16.54 | 13.70 | 5.6250 | 118.785 | 138.490 | 0.1790 | 292 | 52.9 | -0.02700 | 120 |
| 16.56 | 13.77 | 5.8500 | 117.440 | 145.850 | 0.1660 | 284 | 54.4 | -0.01000 | 120 |
| 16.58 | 13.84 | 6.0750 | 117.355 | 145.516 | 0.1540 | 2.94 | 52.6 | 0.00580 | 120 |
| 16.60 | 13.90 | 6.3000 | 117.386 | 147.207 | 0.1440 | 287 | 53.9 | -0.04500 | 120 |
| 16.62 | 13.98 | 6.5250 | 117.397 | 145.690 | 0.1310 | 2.96 | 523 | -0.00540 | 180 |
| 16.64 | 14.06 | 6.7500 | 117.347 | 145.074 | 0.1170 | 3.02 | 51.2 | 0.04100 | 180 |
| 16.66 | 14.13 | 6.9750 | 117.529 | 142.435 | 0.1100 | 2.96 | 523 | 0.01800 | 180 |
| 16.70 | 14.27 | 7.4250 | 117.191 | 146.628 | 0.0936 | 287 | 53.9 | -0.01800 | 180 |


| Exp 95 | Flow Cell <br> Front Tilt = 18.12 div . |  | $0.107 \mu \mathrm{~m} \text { PSL }$ <br> Rear Transl |  | V.F. $=0.198 \%$ |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=100 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rea | Side | Tilt Angle | Inten | (kHz) | $Y$ - | Decay Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) |  | Cumulant | .) |
| 16.04 | 11.80 | 0.000 | 117.281 | 142.977 | 0.9100 | 4.52 | 34.2 | 0.1900 | 120 |
| 16.06 | 11.88 | 0.225 | 116.510 | 144.543 | 0.8980 | 4.46 | 34.7 | 0.1900 | 120 |
| 16.08 | 11.95 | 0.450 | 116.650 | 143.331 | 0.8640 | 4.30 | 36.0 | 0.1800 | 120 |
| 16.10 | 12.03 | 0.675 | 116.766 | 141.372 | 0.8180 | 4.04 | 38.3 | 0.1600 | 120 |
| 16.12 | 12.11 | 0.900 | 117.078 | 142.159 | 0.7630 | 3.76 | 41.1 | 0.1200 | 120 |
| 16.14 | 12.19 | 1.125 | 117.276 | 139.888 | 0.7110 | 3.56 | 43.4 | 0.0920 | 120 |
| 16.16 | 12.26 | 1.350 | 117.809 | 143.627 | 0.6780 | 3.41 | 45.4 | 0.0710 | 120 |
| 16.18 | 12.34 | 1.575 | 119.677 | 144.017 | 0.6440 | 3.27 | 47.3 | 0.0440 | 120 |
| 16.20 | 12.41 | 1.800 | 118.019 | 141.195 | 0.6100 | 3.19 | 48.5 | 0.0330 | 120 |
| 16.22 | 12.48 | 2.025 | 118.653 | 141.415 | 0.5810 | 3.11 | 49.7 | 0.0200 | 120 |
| 16.24 | 12.55 | 2.250 | 118.481 | 143.786 | 0.5560 | 3.09 | 50.0 | 0.0067 | 120 |
| 16.26 | 12.63 | 2.475 | 118.880 | 142.615 | 0.5270 | 3.08 | 50.3 | -0.0048 | 120 |
| 16.28 | 12.71 | 2.700 | 118.937 | 143.273 | 0.5010 | 3.03 | 51.0 | 0.0170 | 120 |
| 16.30 | 12.79 | 2.925 | 112.174 | 143.265 | 0.4700 | 3.02 | 51.2 | 0.0050 | 120 |
| 16.32 | 12.87 | 3.150 | 112.174 | 142.709 | 0.4350 | 3.02 | 51.2 | 0.0190 | 120 |
| 16.34 | 12.96 | 3.375 | 112.683 | 142.093 | 0.4070 | 3.01 | 51.5 | 0.0110 | 120 |
| 16.36 | 13.03 | 3.600 | 116.956 | 141.895 | 0.3890 | 3.01 | 51.4 | 0.0160 | 120 |
| 16.38 | 13.10 | 3.825 | 119.963. | 145.791 | 0.3680 | 3.01 | 51.3 | 0.0430 | 120 |
| 16.40 | 13.17 | 4.050 | 116.862 | 143.643 | 0.3450 | 2.96 | 52.3 | -0.0017 | 120 |
| 16.42 | 13.25 | 4.275 | 121.701 | 143.690 | 0.3190 | 3.00 | 51.6 | -0.0085 | 120 |
| 16.44 | 13.33 | 4.500 | 116.509 | 142.787 | 0.3060 | 2.96 | 52.3 | 0.0180 | 120 |
| 16.46 | 13.40 | 4.725 | 116.272 | 141.820 | 0.2860 | 2.92 | 53.0 | 0.0019 | 120 |
| 16.48 | 13.47 | 4.950 | 116.242 | 141.507 | 0.2690 | 2.98 | 51.9 | 0.0088 | 120 |
| 16.50 | 13.54 | 5.175 | 116.175 | 142.155 | 0.2540 | 2.92 | 53.0 | -0.0180 | 120 |
| 16.52 | 13.61 | 5.400 | 115.664 | 142.633 | 0.2360 | 2.99 | 51.7 | -0.0048 | 120 |
| 16.54 | 13.68 | 5.625 | 115.568 | 143.095 | 0.2240 | 2.95 | 52.4 | -0.0110 | 120 |
| 16.56 | 13.76 | 5.850 | 114.865 | 146.472 | 0.2110 | 2.86 | 54.1 | -0.0350 | 120 |
| 16.58 | 13.84 | 6.075 | 115.400 | 141.599 | 0.1960 | 2.99 | 51.7 | -0.0200 | 120 |
| 16.60 | 13.91 | 6.300 | 115.190 | 140.385 | 0.1830 | 2.95 | 52.4 | 0.0083 | 120 |
| 16.62 | 13.97 | 6.525 | 115.302 | 142.617 | 0.1720 | 2.88 | 53.7 | -0.0700 | 120 |
| 16.64 | 14.05 | 6.750 | 115.138 | 141.599 | 0.1600 | 2.92 | 53.0 | 0.0120 | 120 |
| 16.66 | 14.12 | 6.975 | 110.402 | 142.851 | 0.1510 | 2.88 | 53.8 | -0.0290 | 120 |
| 16.70 | 14.26 | 7.425 | 115.026 | 144.266 | 0.1330 | 2.83 | 54.7 | -0.0350 | 180 |
| 16.74 | 14.42 | 7.875 | 114.915 | 142.421 | 0.1130 | 2.97 | 52.1 | 0.0400 | 240 |
| 16.78 | 14.59 | 8.325 | 115.453 | 142.348 | 0.0944 | 2.89 | 53.6 | -0.0750 | 240 |
| 16.82 | 14.74 | 8.775 | 114.782 | 144.606 | 0.0823 | 2.83 | 54.8 | 0.0440 | 240 |


| Exp 103 | Flow Cell |  | $0.098 \mu \mathrm{~m}$ PSL |  | V.F. $=0.320 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Front Tilt $=18.09$ div. |  | Rear Translation $=12.29$ div. |  |  |  | Flow Rate $=0 \%$ |  |  |
| Rear Tilt | Side | Tilt Angle | Intens | ( $\mathrm{k}-\mathrm{l} \mathrm{z}$ ) | Y- | Decay Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) |  | Cumulant | (sec.) |
| 16.04 | 11.77 | 0.000 | 118.401 | 156.753 | 0.8930 | 3.64 | 43.2 | 0.0680 | 120 |
| 16.06 | 11.85 | 0.225 | 118.383 | 156.349 | 0.8510 | 3.60 | 43.6 | 0.0670 | 120 |
| 16.08 | 11.92 | 0.450 | 117.927 | 155.190 | 0.7990 | 3.51 | 44.7 | 0.0470 | 120 |
| 16.10 | 12.00 | 0.675 | 118.071 | 153.624 | 0.7460 | 3.45 | 45.6 | 0.0500 | 120 |
| 16.12 | 12.07 | 0.900 | 118.435 | 156.151 | 0.6960 | 3.38 | 46.5 | 0.0340 | 120 |
| 16.14 | 12.14 | 1.125 | 118.627 | 155.112 | 0.6490 | 3.36 | 46.8 | 0.0460 | 120 |
| 16.16 | 12.23 | 1.350 | 118.097 | 157.115 | 0.5990 | 3.31 | 47.5 | 0.0120 | 120 |
| 16.18 | 12.31 | 1.575 | 117.946 | 153.308 | 0.5560 | 3.28 | 47.9 | 0.0170 | 120 |
| 16.20 | 12.38 | 1.800 | 117.923 | 155.572 | 0.5240 | 3.26 | 48.1 | 0.0200 | 120 |
| 16.22 | 12.46 | 2.025 | 117.740 | 156.777 | 0.4890 | 3.21 | 49.0 | 0.0140 | 120 |
| 16.24 | 12.53 | 2.250 | 116.983 | 155.247 | 0.4600 | 3.22 | 48.7 | 0.0170 | 120 |
| 16.27 | 12.65 | 2.588 | 116.824 | 151.083 | 0.4070 | 3.20 | 49.2 | -0.0016 | 120 |
| 16.30 | 12.76 | 2.925 | 116.781 | 153.858 | 0.3650 | 3.20 | 49.0 | 0.0260 | 120 |
| 16.33 | 12.89 | 3.263 | 116.949 | 156.675 | 0.3210 | 3.17 | 49.5 | 0.0230 | 120 |
| 16.36 | 12.99 | 3.600 | 118.126 | 156.177 | 0.2840 | 3.23 | 48.6 | 0.0420 | 120 |
| 16.39 | 13.10 | 3.938 | 117.121 | 158.184 | 0.2480 | 3.20 | 49.1 | 0.0270 | 180 |
| 16.42 | 13.22 | 4.275 | 116.644 | 152.989 | 0.2130 | 3.23 | 48.6 | 0.0320 | 240 |
| 16.45 | 13.32 | 4.613 | 116.662 | 153.841 | 0.1860 | 3.14 | 50.1 | 0.0110 | 240 |
| 16.48 | 13.41 | 4.950 | 117.515 | 154.081 | 0.1660 | 3.11 | 50.5 | 0.0220 | 240 |
| 16.51 | 13.52 | 5.288 | 118.482 | 152.235 | 0.1440 | 3.20 | 49.5 | -0.0017 | 240 |
| 16.54 | 13.61 | 5.625 | 117.994 | 150.291 | 0.1280 | 3.21 | 48.9 | 0.0083 | 240 |
| 16.57 | 13.72 | 5.963 | 117.688 | 152.806 | 0.1120 | 3.16 | 49.8 | 0.0540 | 240 |
| 16.60 | 13.83 | 6.300 | 116.946 | 151.540 | 0.0961 | 3.20 | 49.1 | 0.0001 | 300 |
| 16.64 | 14.00 | 6.750 | 116.537 | 149.403 | 0.0762 | 3.16 | 49.7 | 0.0460 | 300 |
| 16.68 | 14.16 | 7.200 | 118.634 | 149.587 | 0.0616 | 3.22 | 48.8 | 0.0150 | 420 |
| 16.72 | 14.32 | 7.650 | 118.826 | 149.439 | 0.0509 | 3.34 | 47.1 | 0.0310 | 420 |


| Exp 104 | Flow CellFront Tilt $=18.09 \mathrm{div}$. |  | $\begin{gathered} 0.098 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.320 \% \\ \text { Rear Translation }=12.29 \text { div. } \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Fow Rate }=100 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Rear Tilt (div.) | Side Translation (div.) | Tilt Angle (mrad) | Intensity | (kHzz) Ch 1 | Y- <br> Intercept | Decay <br> Rate <br> (1/ms) | Radius ( $n \mathrm{~m}$ ) | Normalized Second Cumulant | Time <br> (sec.) |
| 16.04 | 11.77 | 0.000 | 119.364 | 158.848 | 0.8980 | 4.72 | 33.3 | 0.1700 | 20 |
| 16.06 | 11.85 | 0.225 | 119.345 | 158.146 | 0.8610 | 4.49 | 35.0 | 0.1600 | 120 |
| 16.08 | 11.92 | 0.450 | 119.414 | 158.276 | 0.8080 | 4.27 | 36.8 | 0.1500 | 120 |
| 16.10 | 12.01 | 0.675 | 119.036 | 156.152 | 0.7330 | 4.03 | 39.0 | 0.1200 | 12 |
| 16.12 | 12.09 | 0.900 | 118.922 | 157.472 | 0.6780 | 3.77 | 41.7 | 0.0820 | 120 |
| 16.14 | 12.16 | 1.125 | 118.785 | 159.031 | 0.6290 | 3.62 | 43.4 | 0.0700 | 120 |
| 16.16 | 1224 | 1.350 | 118.099 | 159.778 | 0.5820 | 3.43 | 45.7 | 0.0400 | 120 |
| 16.18 | 12.32 | 1.575 | 119.146 | 158.552 | 0.5460 | 3.37 | 46.6 | 0.0340 | 120 |
| 16.20 | 12.39 | 1.800 | 119.216 | 158.321 | 0.5070 | 3.35 | 46.9 | 0.0170 | 120 |
| 16.22 | 12.46 | 2.025 | 119.070 | 157.154 | 0.4750 | 3.32 | 47.3 | 0.0210 | 120 |
| 16.24 | 12.53 | 2.250 | 119.295 | 158.094 | 0.4450 | 3.27 | 48.0 | 0.0210 | 120 |
| 16.27 | 12.65 | 2.588 | 119.630 | 157.957 | 0.3950 | 3.22 | 48.8 | 0.0380 | 120 |
| 16.30 | 12.77 | 2.925 | 119.302 | 157.063 | 0.3400 | 3.21 | 49.0 | 0.0120 | 20 |
| 16.33 | 12.90 | 3.263 | 119.450 | 157.597 | 0.2930 | 3.22 | 48.8 | 0.0130 | 120 |
| 16.36 | 13.02 | 3.600 | 119.664 | 159.625 | 0.2560 | 3.26 | 48.2 | 0.0270 | 180 |
| 16.42 | 13.25 | 4.275 | 125.478 | 158.664 | 0.2120 | 3.11 | 50.6 | 0.0530 | 240 |
| 16.48 | 13.47 | 4.950 | 125.405 | 157.113 | 0.1530 | 3.02 | 52.0 | 0.0300 | 300 |
| 16.50 | 13.54 | 5.175 | 126.026 | 154.316 | 0.1400 | 3.14 | 50.1 | 0.0200 | 300 |
| 16.54 | 13.62 | 5.625 | 122.759 | 156.508 | 0.1240 | 3.16 | 49.7 | 0.0270 | 300 |
| 16.58 | 13.80 | 6.075 | 124.684 | 159.036 | 0.0983 | 3.14 | 50.0 | 0.0019 | 42 |


| Exp 105 | Flow Cell |  | $0.098 \mu \mathrm{mPSL}$ |  | V.F. $=0.860 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Front Tilt $=18.09$ div. |  | Rear Translation $=12.29$ div. |  |  |  | Flow Rate = 0\% |  |  |
| Rear Tilt | Side | Tilt Angle | Inten | (kHz) | Y- | Decay <br> Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) |  | Cumulant | (sec.) |
| 16.04 | 11.77 | 0.000 | 139.072 | 184.903 | 0.8900 | 4.00 | 39.3 | 0.0950 | 120 |
| 16.06 | 11.84 | 0.225 | 139.630 | 185.198 | 0.8480 | 3.92 | 40.0 | 0.0960 | 120 |
| 16.08 | 11.92 | 0.450 | 141.678 | 182.834 | 0.7770 | 3.83 | 41.0 | 0.0810 | 120 |
| 16.10 | 12.00 | 0.675 | 139.999 | 184.336 | 0.7080 | 3.69 | 42.5 | 0.0620 | 120 |
| 16.12 | 12.08 | 0.900 | 140.559 | 185.916 | 0.6400 | 3.63 | 43.3 | 0.0740 | 120 |
| 16.14 | 12.17 | 1.125 | 140.595 | 185.113 | 0.5670 | 3.53 | 44.5 | 0.0660 | 120 |
| 16.16 | 12.24 | 1.350 | 140.492 | 184.097 | 0.5180 | 3.47 | 45.2 | 0.0560 | 120 |
| 16.18 | 12.32 | 1.575 | 141.202 | 185.131 | 0.4760 | 3.39 | 46.3 | 0.0370 | 120 |
| 16.20 | 12.40 | 1.800 | 141.167 | 184.500 | 0.4340 | 3.33 | 47.1 | 0.0170 | 120 |
| 16.22 | 12.47 | 2.025 | 141.377 | 184.714 | 0.4000 | 3.40 | 46.2 | 0.0330 | 120 |
| 16.24 | 12.54 | 2.250 | 141.730 | 186.144 | 0.3740 | 3.37 | 46.6 | 0.0390 | 120 |
| 16.26 | 12.62 | 2.475 | 142.055 | 182.742 | 0.3380 | 3.38 | 46.5 | 0.0480 | 120 |
| 16.28 | 12.69 | 2.700 | 142.279 | 186.109 | 0.3130 | 3.34 | 47.1 | 0.0380 | 180 |
| 16.30 | 12.78 | 2.925 | 143.151 | 184.721 | 0.2820 | 3.32 | 47.3 | 0.0190 | 180 |
| 16.33 | 12.89 | 3.262 | 142.853 | 185.888 | 0.2450 | 3.37 | 46.6 | 0.0380 | 180 |
| 16.36 | 13.01 | 3.600 | 143.350 | 185.861 | 0.2130 | 3.26 | 48.2 | 0.0067 | 240 |
| 16.39 | 13.13 | 3.938 | 143.212 | 184.384 | 0.1810 | 3.31 | 47.4 | 0.0130 | 240 |
| 16.42 | 13.24 | 4.275 | 143.471 | 184.890 | 0.1580 | 3.32 | 47.4 | 0.0110 | 240 |
| 16.45 | 13.34 | 4.613 | 143.698 | 184.999 | 0.1360 | 3.25 | 48.3 | -0.0190 | 240 |
| 16.48 | 13.46 | 4.950 | 143.880 | 183.444 | 0.1150 | 3.33 | 47.2 | 0.0340 | 240 |
| 16.51 | 13.56 | 5.288 | 144.194 | 183.940 | 0.1020 | 3.28 | 47.9 | 0.0180 | 240 |
| 16.54 | 13.67 | 5.625 | 144.687 | 185.494 | 0.0872 | 3.12 | 49.4 | -0.0290 | 240 |
| 16.57 | 13.78 | 5.963 | 145.198 | 185.027 | 0.0757 | 3.22 | 48.8 | 0.0250 | 300 |
| 16.60 | 13.89 | 6.300 | 144.593 | 185.923 | 0.0638 | 3.34 | 47.1 | 0.0380 | 300 |
| 16.64 | 14.04 | 6.750 | 144.644 | 182.030 | 0.0524 | 3.28 | 48.0 | 0.0170 | 300 |


| Exp 106 | Flow Cell |  | $0.098 \mu \mathrm{mPSL}$ |  | V.F. $=0.860 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Front Tilt $=18.09$ div. |  | Rear Translation $=12.29$ div. |  |  |  | Flow Rate = 100\% |  |  |
| Rear Tilt | Side <br> Translation | Tilt Angle | Inten | kHz) | Y- | Decay Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) | (nm) | Cumulant | (sec.) |
| 16.04 | 11.77 | 0.000 | 144.089 | 187.425 | 0.8990 | 5.01 | 31.4 | 0.1600 | 120 |
| 16.06 | 11.84 | 0.225 | 144.054 | 188.714 | 0.8540 | 4.76 | 33.0 | 0.1500 | 120 |
| 16.08 | 11.92 | 0.450 | 144.911 | 188.127 | 0.7790 | 4.49 | 35.0 | 0.1300 | 120 |
| 16.10 | 12.00 | 0.675 | 143.095 | 186.813 | 0.7030 | 4.25 | 37.0 | 0.1100 | 120 |
| 16.12 | 12.09 | 0.900 | 144.634 | 190.635 | 0.6280 | 4.01 | 39.2 | 0.0940 | 120 |
| 16.14 | 12.18 | 1.125 | 148.697 | 195.057 | 0.5630 | 3.74 | 42.0 | 0.0570 | 120 |
| 16.16 | 12.24 | 1.350 | 142.419 | 189.699 | 0.5190 | 3.68 | 42.7 | 0.0500 | 120 |
| 16.18 | 12.32 | 1.575 | 142.037 | 189.171 | 0.4680 | 3.55 | 44.3 | 0.0460 | 120 |
| 16.20 | 12.40 | 1.800 | 142.565 | 182.164 | 0.4310 | 3.49 | 45.1 | 0.0420 | 120 |
| 16.22 | 12.51 | 2.025 | 141.959 | 186.951 | 0.3770 | 3.45 | 45.6 | 0.0300 | 120 |
| 16.24 | 12.58 | 2.250 | 145.271 | 193.799 | 0.3500 | 3.32 | 47.3 | 0.0170 | 120 |
| 16.26 | 12.65 | 2.475 | 142.937 | 193.405 | 0.3140 | 3.37 | 46.7 | 0.0190 | 120 |
| 16.28 | 12.68 | 2.700 | 142.535 | 186.617 | 0.3020 | 3.35 | 46.9 | 0.0270 | 120 |
| 16.30 | 12.77 | 2.925 | 141.983 | 187.093 | 0.2710 | 3.30 | 47.5 | 0.0230 | 120 |
| 16.33 | 12.90 | 3.262 | 142.598 | 184.138 | 0.2330 | 3.31 | 47.4 | 0.0260 | 120 |
| 16.36 | 13.02 | 3.600 | 142.318 | 186.151 | 0.2000 | 3.17 | 49.5 | -0.0095 | 120 |
| 16.39 | 13.13 | 3.938 | 142.205 | 186.407 | 0.1690 | 3.25 | 48.4 | 0.0140 | 120 |
| 16.42 | 13.25 | 4.275 | 144.341 | 178.609 | 0.1430 | 3.20 | 49.1 | -0.0086 | 180 |
| 16.45 | 13.36 | 4.613 | 142.165 | 185.701 | 0.1230 | 3.24 | 48.5 | 0.0390 | 180 |
| 16.48 | 13.45 | 4.950 | 142.244 | 188.625 | 0.1070 | 3.25 | 48.3 | 0.0087 | 240 |
| 16.51 | 13.56 | 5.288 | 142.177 | 186.865 | 0.0916 | 3.19 | 49.3 | -0.0430 | 240 |
| 16.54 | 13.67 | 5.625 | 142.667 | 184.853 | 0.0771 | 3.24 | 48.5 | 0.0300 | 300 |
| 16.57 | 13.77 | 5.963 | 142.577 | 185.137 | 0.0658 | 3.14 | 50.1 | -0.0380 | 300 |
| 16.60 | 13.88 | 6.300 | 142.615 | 185.132 | 0.0579 | 3.10 | 50.6 | 0.0430 | 300 |
| 16.64 | 14.00 | 6.750 | 142.905 | 185.970 | 0.0495 | 3.30 | 47.7 | 0.0710 | 300 |


| Exp 111 | Flow Cell |  | $0.203 \mu \mathrm{mPSL}$ |  | V.F. $=0.20 \%$ |  | $\theta=112 \mathrm{deg}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Front Tilt $=18.10$ div. |  | Rear Translation $=12.31 \mathrm{div}$. |  |  |  | Flow Rate $=0 \%$ |  |  |
| Rear Tilt (div.) | Side <br> Translation | Tilt Angle | Inten | (kHz) | Y- | Decay Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) | nm) | Cumulant | (sec.) |
| 16.03 | 11.71 | 0.000 | 99.520 | 150.610 | 0.8990 | 1.68 | 93.8 | 0.0720 | 120 |
| 16.05 | 11.78 | 0.225 | 99.450 | 151.261 | 0.8600 | 1.67 | 93.9 | 0.0610 | 120 |
| 16.07 | 11.86 | 0.450 | 99.765 | 150.594 | 0.8020 | 1.63 | 96.6 | 0.0590 | 120 |
| 16.09 | 11.93 | 0.675 | 94.898 | 148.120 | 0.7350 | 1.61 | 97.7 | 0.0490 | 120 |
| 16.11 | 12.00 | 0.900 | 99.518 | 150.435 | 0.6490 | 1.58 | 99.7 | 0.0450 | 120 |
| 16.13 | 12.09 | 1.125 | 99.545 | 152.204 | 0.5870 | 1.56 | 101 | 0.0380 | 120 |
| 16.15 | 12.18 | 1.350 | 100.047 | 149.361 | 0.5230 | 1.56 | 101 | 0.0400 | 120 |
| 16.17 | 12.25 | 1.575 | 99.651 | 150.768 | 0.4880 | 1.52 | 103 | 0.0470 | 120 |
| 16.19 | 12.32 | 1.800 | 99.400 | 150.485 | 0.4450 | 1.52 | 104 | 0.0360 | 120 |
| 16.21 | 12.40 | 2.025 | 99.636 | 150.346 | 0.3980 | 1.52 | 103 | 0.0260 | 180 |
| 16.23 | 12.47 | 2.250 | 100.341 | 150.654 | 0.3670 | 1.52 | 104 | 0.0480 | 180 |
| 16.25 | 12.55 | 2.475 | 96.922 | 149.996 | 0.3300 | 1.52 | 103 | 0.0550 | 180 |
| 16.27 | 12.63 | 2.700 | 101.073 | 151.431 | 0.3010 | 1.48 | 106 | 0.0110 | 240 |
| 16.29 | 12.71 | 2.925 | 99.788 | 152.083 | 0.2680 | 1.51 | 104 | 0.0270 | 240 |
| 16.31 | 12.80 | 3.150 | 90.568 | 153.792 | 0.2370 | 1.51 | 104 | -0.0170 | 120 |
| 16.33 | 12.89 | 3.375 | 99.076 | 141.698 | 0.2140 | 1.51 | 104 | 0.0320 | 120 |
| 16.35 | 12.96 | 3.600 | 98.911 | 148.326 | 0.1960 | 1.52 | 103 | 0.0360 | 240 |
| 16.37 | 13.02 | 3.825 | 99.285 | 150.411 | 0.1800 | 1.55 | 101 | 0.0740 | 120 |
| 16.39 | 13.11 | 4.050 | 98.588 | 148.786 | 0.1580 | 1.56 | 101 | 0.0190 | 240 |
| 16.41 | 13.19 | 4.275 | 98.953 | 149.803 | 0.1420 | 1.53 | 102 | 0.0520 | 120 |
| 16.43 | 13.26 | 4.500 | 99.229 | 148.234 | 0.1320 | 1.47 | 107 | 0.0520 | 240 |
| 16.45 | 13.32 | 4.725 | 98.608 | 148.459 | 0.1160 | 1.59 | 98.6 | 0.0820 | 180 |
| 16.47 | 13.39 | 4.950 | 98.446 | 147.094 | 0.1020 | 1.51 | 104 | 0.0400 | 300 |
| 16.49 | 13.46 | 5.175 | 98.500 | 148.067 | 0.0933 | 1.45 | 108 | -0.0270 | 300 |


| Exp 112 | Flow Cell <br> Front Tilt $=18.10 \mathrm{div}$. |  | $\begin{gathered} 0.203 \mu \mathrm{~m} \text { PSL } \quad \text { V.F. }=0.20 \% \\ \text { Rear Translation }=12.31 \mathrm{div} . \end{gathered}$ |  |  |  | $\begin{gathered} \theta=112 \mathrm{deg} \\ \text { Flow Rate }=100 \% \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |  |  |  |
| Rear Tilt (div.) | Side <br> Translation | Tilt Angle | Inten | (kHz) | Y- | Decay <br> Rate | Radius | Normalized Second | Time |
|  | (div.) |  | Ch 0 | Ch 1 | Intercept | (1/ms) | (nm) | Cumulant | (sec.) |
| 16.03 | 11.72 | 0.000 | 96.776 | 146.167 | 0.8830 | 2.34 | 67.0 | 0.2000 | 120 |
| 16.05 | 11.81 | 0.225 | 98.672 | 144.693 | 0.8320 | 2.26 | 69.6 | 0.1900 | 120 |
| 16.07 | 11.89 | 0.450 | 99.143 | 146.527 | 0.7620 | 2.08 | 75.6 | 0.1600 | 120 |
| 16.09 | 11.97 | 0.675 | 98.867 | 147.319 | 0.6900 | 1.95 | 80.8 | 0.1400 | 120 |
| 16.11 | 12.05 | 0.900 | 98.522 | 145.951 | 0.6130 | 1.80 | 87.2 | 0.0890 | 120 |
| 16.13 | 12.13 | 1.125 | 98.351 | 146.407 | 0.5500 | 1.71 | 92.1 | 0.0650 | 120 |
| 16.15 | 12.20 | 1.350 | 101.118 | 147.986 | 0.4960 | 1.62 | 96.7 | 0.0330 | 120 |
| 16.17 | 12.28 | 1.575 | 98.508 | 145.524 | 0.4520 | 1.58 | 99.6 | 0.0380 | 120 |
| 16.19 | 12.35 | 1.800 | 93.937 | 142.503 | 0.4200 | 1.54 | 102 | 0.0140 | 120 |
| 16.21 | 12.42 | 2.025 | 98.620 | 148.365 | 0.3860 | 1.53 | 103 | 0.0015 | 120 |
| 16.23 | 12.49 | 2.250 | 98.761 | 147.941 | 0.3530 | 1.51 | 104 | -0.0052 | 120 |
| 16.25 | 12.58 | 2.475 | 98.954 | 146.634 | 0.3160 | 1.49 | 105 | 0.0270 | 180 |
| 16.27 | 12.67 | 2.700 | 98.551 | 145.583 | 0.2790 | 1.47 | 107 | -0.0150 | 120 |
| 16.29 | 12.74 | 2.925 | 98.332 | 146.225 | 0.2540 | 1.48 | 106 | -0.0012 | 240 |
| 16.31 | 12.81 | 3.150 | 98.014 | 148.159 | 0.2300 | 1.50 | 105 | 0.0099 | 240 |
| 16.33 | 12.87 | 3.375 | 95.551 | 154.264 | 0.2120 | 1.52 | 103 | 0.0480 | 120 |
| 16.35 | 12.96 | 3.600 | 98.223 | 149.720 | 0.1920 | 1.50 | 104 | 0.0230 | 240 |
| 16.37 | 13.04 | 3.825 | 98.316 | 149.800 | 0.1710 | 1.47 | 107 | 0.0180 | 240 |
| 16.39 | 13.11 | 4.050 | 98.002 | 147.499 | 0.1590 | 1.48 | 106 | 0.0400 | 120 |
| 16.41 | 13.20 | 4.275 | 98.118 | 147.093 | 0.1380 | 1.47 | 107 | -0.0051 | 240 |
| 16.44 | 13.32 | 4.613 | 97.756 | 144.985 | 0.1170 | 1.55 | 102 | -0.0035 | 180 |
| 16.48 | 13.44 | 5.062 | 97.910 | 148.611 | 0.0982 | 1.49 | 105 | 0.0085 | 180 |

## Appendix IV

## Flow Profile Discussion

The purpose of a dynamic light scattering experiment is to determine the diffusion constant associated with the Brownian motion of the particles within a test sample. Section 3.3 showed how the diffusion constant was then used to calculate particle size. When conducting experiments in a flowing fluid, factors other than Brownian motion can influence the motion of the particles, and therefore influence the data used to determine particle size. These factors are a product of the flow profile as it passes through the detection area. As mentioned in Section 5.2, fully developed laminar flow is necessary to conduct dynamic light scattering experiments in flowing samples. The calculations given in Section 5.2 prove that the flow associated with the experiments described in this thesis was laminar. Those calculations also prove that the possibility of entrance effects influencing the flow in the detection area was minimal due to the length of the test cell. Although the flow used in the experiments covered by this thesis were laminar, several questions remained about the flow profile. Therefore, this appendix will concentrate on the fully developed flow profile found in the detection area. The following equations were used to determine the flow profile [White (1991)]:

$$
\begin{align*}
v_{x}(y, z)= & \left(\frac{16 c^{2}}{\mu \pi^{3}}\right)\left(\frac{-d \hat{p}}{d x}\right) \sum_{j=1,3,5, \ldots}^{\infty}(-1)^{\frac{(j-1)}{2}}\left[1-\frac{\cosh [j \pi z /(2 c)]}{\cosh [j \pi b /(2 c)]}\right] \\
& {\left[\frac{\cos [i \pi y /(2 c)]}{j^{3}}\right] } \tag{A-IV-1}
\end{align*}
$$

where $v_{x}(y, z)$ is the $x$-direction velocity as a function of position in the $y-z$ plane, 2 c is the width of the test cell, 2 b is the height of the test cell, $\mu$ is the viscosity of the sample, and $\frac{-d \hat{p}}{d x}$ is the pressure gradient of the system. For the flow system associated with this thesis, the only unknown of Eq. (A-IV-1) is the pressure gradient, and it was solved for with the following equation [White (1991)]:

$$
\begin{equation*}
\frac{-\mathrm{d} \hat{\mathrm{p}}}{\mathrm{dx}}=\frac{\mathrm{Q}}{\left(\frac{4 \mathrm{bc}^{3}}{3 \mu}\right)\left[1-\frac{192 \mathrm{c}}{\mathrm{~b} \pi^{5}} \sum_{j=1,3,5, \ldots}^{\infty} \frac{\tanh [\mathrm{j} \pi \mathrm{c} /(2 \mathrm{a})]}{j^{5}}\right]} \tag{A-IV-2}
\end{equation*}
$$

where Q is the overall flow rate (see Table 2 in Section 5.2). These equations were used to determine the flow profile of the system. Mathcad 7.0, a commercial mathematical software package, was used to plot the profile as follows.

Once the flow profile is determined, the next step was to use the information from the time constants associated with the flow system. The first time constant calculated was the diffusion time constant. This time constant relates the time it takes a particle to move under Brownian motion through a unit length scale determined by the scattered wave vector. The single scattering time constant is defined by [Nobbmann (1997)]

$$
\begin{equation*}
\tau_{\text {diffusion }}=\frac{1}{\mathrm{D}_{\mathrm{o}} \mathrm{k}^{2}} \tag{A-IV-3}
\end{equation*}
$$

where $D_{0}$ is the diffusion constant, and $k$ is the scattering wave vector. The second time constant to be calculated was the transit time constant. This time constant relates the time it takes a particle to travel through the focused incident laser beam and is defined by

$$
\begin{equation*}
\tau_{\text {transit }}=\frac{L}{v} \tag{A-IV-4}
\end{equation*}
$$

where $\mathrm{L}=\frac{\mathrm{b}}{\cos \left(\alpha_{L}\right)}$ and is the length the particle travels through the focused laser beam, and $v$ is the centerline velocity of the flow. The centerline velocity was calculated using Eq. (A-IV-1). The final time constant calculated was the Doppler time constant. The Doppler time constant measured how misalignments affected the calculation of particle diameter. The Doppler time constant was defined by

$$
\begin{equation*}
\tau_{\text {Doppler }}=\frac{1}{\mathrm{k} \cos (90-\delta) \mathrm{v}} \tag{A-IV-5}
\end{equation*}
$$

where $\delta=5^{\circ}$ and is the cell rotation angle (see Fig. 13 in Section 5.2).

Knowledge about these three time constants are valuable in determining the likelihood of success in collecting data using DLS techniques. To make DLS measurements, a particle must diffuse while in the detection area. If the time it takes a particle to diffuse (i.e., the diffusion time constant) is larger than the transit time constant; no usable data will be collected.

The following pages contain the calculation of the flow profile and the time constants describe above. The first two pages of calculations describe the flow profile determined by the conditions associated with the flow setup used in this research. From the profile, it can be seen that the flow is fully developed and the maximum velocity associated with the setup is $0.0162 \mathrm{~m} / \mathrm{sec}$ at the centerline. The collection of data was conducted at an estimated depth of 1 mm and the velocity calculated there was found to be $0.00926 \mathrm{~m} / \mathrm{sec}$.

The second set of calculations presented is for the three time constants. The diffusion time constant was calculated for $0.107 \mu \mathrm{~m}$ PSL particles and was found to be 0.0019 seconds. The transit time constant associated with a centerline velocity was found to be 0.0022 seconds. From these calculations $\tau_{\text {diffusion }}<\tau_{\text {transit }}$; and therefore Brownian motion diffusion occurs in the detection area. The Doppler time constant for a misalignment of $\delta=5^{\circ}$ was calculated to be 0.045 msec . The effects of this misalignment were discussed in Section 6.5.1.

The same calculations were performed using the velocity at 1 mm from the wall and similar results were found. The difference between the diffusion time constant ( $\tau_{\text {diffusion }}=0.0019 \mathrm{sec}$ ) and the transit time constant $\left(\tau_{\text {transit }}=0.0038 \mathrm{sec}\right)$ increased by a factor of two. This was due to the lower velocities seen near the wall.

Time constants were then calculated for $0.204 \mu \mathrm{~m}$ PSL particles at a distance of 1 mm from the wall. For the $0.204 \mu \mathrm{~m}$ particles, the diffusion time constant $\left(\tau_{\text {diffusion }}=\right.$ 0.0036 sec ) was found to be nearly twice that for $0.107 \mu \mathrm{~m}$ particles. However, the diffusion time constant was still less than the transit time constant ( $\tau_{\text {transit }}=0.0038 \mathrm{sec}$ ) seen at 1 mm from the wall.

A final set of calculations was performed using a flow rate that was 10 times the value seen in the actual experiments. The flow profile was calculated first. From this calculation, it was determined that the velocity was proportional to the flow rate and therefore, increased as well by a factor of 10 . From the time constants associated with this higher flow rate, it was determined that $\tau_{\text {diffusion }}(0.0019 \mathrm{sec})>\tau_{\text {transit }}(0.00038 \mathrm{sec})$; and therefore no usable data could be collected. The calculations were performed using
the velocity 1 mm from the wall. Table 8 gives a summary of all of the calculations performed on the flow profile and time constants.

Table 8: Summary of results from flow profile and time constant calculations

| Flow Rate <br> Q <br> $(\mathrm{ml} / \mathrm{min})$ | Measure <br> ment <br> Location <br> $(\mathrm{mm})$ | Particle <br> Diameter <br> $(\mathrm{nm})$ | Velocity <br> $(\mathrm{mm} / \mathrm{sec})$ | $\boldsymbol{\tau}_{\text {Diflusion }}$ <br> $(\mathrm{sec})$ | $\boldsymbol{\tau}_{\text {Transit }}$ <br> $(\mathrm{sec})$ | $\tau_{\text {Doppler }}$ <br> $(\mathrm{sec})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 22.40 | 3 | 0.107 | 16.2 | 0.0019 | 0.0022 | 0.000045 |
| 22.40 | 1 | 0.107 | 9.26 | 0.0019 | 0.0038 | 0.000079 |
| 22.40 | 1 | 0.204 | 9.26 | 0.0036 | 0.0038 | 0.000079 |
| 224.03 | 1 | 0.107 | 92.6 | 0.0019 | 0.00038 | 0.000008 |

Flow profile calculations for the conditions and test cell used during experiments.

$$
\begin{array}{rl}
\mathrm{c}:=0.003 \mathrm{~m} \quad \mathrm{~b}:=0.004 \mathrm{~m} & \mathrm{Q} 1:=22.4025 \frac{\mathrm{~mL}}{\mathrm{~min}} \quad \mu:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
\text { convertq }:=\frac{0.000001}{60} & \mathrm{Q}:=\mathrm{Q} 1 \cdot \text { convertq } \quad \mathrm{Q}=3.734 \cdot 10^{-7} \frac{\mathrm{~m}^{3}}{\mathrm{sec}}
\end{array}
$$

$$
\mathrm{j}:=1,3 . .11
$$

Velocity at wall Velocity at 1 mm from wall Centerline velocity

$$
v(0.003,0)=0 \quad \frac{\mathrm{~m}}{\mathrm{sec}} \quad v(0.002,0)=9.2618 \cdot 10^{-3} \quad \frac{\mathrm{~m}}{\mathrm{sec}} \quad v(0,0)=0.0162 \quad \frac{\mathrm{~m}}{\mathrm{sec}}
$$

The values given below depict the position in the $y$ - and $z$-directions at which the velocity calculations were performed to create the following velocity table and flow profile graph. Integer designation was used on the graph, rather than the actual locations in meters.
$y^{T}=$

m
$z^{T}=$
$\begin{array}{llllllllllllll}-0.004 & -0.003 & -0.002 & -0.002 & -0.001 & 0 & 0.001 & 0.002 & 0.002 & 0.003 & 0.004\end{array}$
m

$$
\begin{aligned}
& \operatorname{dpdx}:=\frac{-Q}{\frac{4 \cdot b \cdot c^{3}}{3 \cdot \mu} \cdot\left[1-\frac{192 \cdot c}{\pi^{5} \cdot b} \square_{j}^{j^{5}}\right]} \\
& \mathrm{dpdx}=-4.5402 \quad \frac{\mathrm{~Pa}}{\mathrm{~m}} \\
& v(y, z):=\left(\frac{16 \cdot c^{2}}{\mu \cdot \pi^{3}}\right) \cdot(-d p d x) \cdot\left[\square_{j}(-1)^{\frac{j-1}{2}} \cdot\left[1-\frac{\cosh \left(\frac{j \cdot \pi \cdot z}{2 \cdot c}\right)}{\cosh \left(\frac{j \cdot \pi \cdot b}{2 \cdot c}\right)}\right) \cdot\left(\frac{\cos \left(\frac{j \cdot \pi \cdot y}{2 \cdot c}\right)}{j^{3}}\right)\right] \\
& \text { ny }:=10 \quad j:=0 \text {.. ny } n z:=10 \quad k:=0 . . n z \\
& y_{j}:=-c+2 \cdot c \cdot \frac{j}{n y} \quad z_{k}:=-b+2 \cdot b \cdot \frac{k}{n z} \quad V_{j, k}:=v\left(y_{j}, z_{k}\right)
\end{aligned}
$$

Velocities ( $\mathrm{m} / \mathrm{sec}$ ) associated with the positions given by y and z

$\mathrm{V}^{\mathrm{T}}=$| 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | 0.003 | 0.0048 | 0.006 | 0.0066 | 0.0068 | 0.0066 | 0.006 | 0.0048 | 0.003 | 0 |
| 0 | 0.0046 | 0.0077 | 0.0098 | 0.011 | 0.0114 | 0.011 | 0.0098 | 0.0077 | 0.0046 | 0 |
| 0 | 0.0055 | 0.0094 | 0.0121 | 0.0137 | 0.0142 | 0.0137 | 0.0121 | 0.0094 | 0.0055 | 0 |
| 0 | 0.006 | 0.0103 | 0.0133 | 0.0151 | 0.0157 | 0.0151 | 0.0133 | 0.0103 | 0.006 | 0 |
| 0 | 0.0061 | 0.0106 | 0.0137 | 0.0156 | 0.0162 | 0.0156 | 0.0137 | 0.0106 | 0.0061 | 0 |
| 0 | 0.006 | 0.0103 | 0.0133 | 0.0151 | 0.0157 | 0.0151 | 0.0133 | 0.0103 | 0.006 | 0 |
| 0 | 0.0055 | 0.0094 | 0.0121 | 0.0137 | 0.0142 | 0.0137 | 0.0121 | 0.0094 | 0.0055 | 0 |
| 0 | 0.0046 | 0.0077 | 0.0098 | 0.011 | 0.0114 | 0.011 | 0.0098 | 0.0077 | 0.0046 | 0 |
| 0 | 0.003 | 0.0048 | 0.006 | 0.0066 | 0.0068 | 0.0066 | 0.006 | 0.0048 | 0.003 | 0 |
| 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |



## Time constants of $0.107 \mu \mathrm{~m}$ PSL particles

Time constants associated with the Brownian motion of $0.107 \mu \mathrm{~m}$ particles

$$
\begin{aligned}
& \mathrm{a}:=0.107 \cdot 10^{-6} \mathrm{~m} \\
& \mathrm{~Kb}:=1.380658 \cdot 10^{-23} \frac{\mathrm{~J}}{\mathrm{~K}} \\
& \eta:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
& \mathrm{~T}:=296 \mathrm{~K} \\
& \mathrm{n}:=1.33 \\
& \lambda:=532.5 \cdot 10^{-9} \quad \mathrm{~m} \\
& \text { Do }:=\frac{\mathrm{Kb} \cdot \mathrm{~T}}{6 \cdot \pi \cdot \eta \cdot \mathrm{a}} \\
& \text { Do }=2.138 \cdot 10^{-12} \quad \frac{\mathrm{~m}^{2}}{\mathrm{~s}} \\
& k:=\frac{2 \cdot \pi \cdot n}{\lambda} \\
& \mathrm{k}=1.569 \cdot 10^{7} \\
& m^{-1} \\
& \text { चdiffusion }:=\frac{1}{\mathrm{Do} \cdot \mathrm{k}^{2}} \\
& \text { } \tau \text { diffusion }=0.0019 \quad \mathrm{sec}
\end{aligned}
$$

Time constant associated with the travel of the particle through the focused laser beam.

Focused beam diameter in sample $=b$ and is defined in Sundaresan (1999). The velocity used is the center line velocity found by the profile Eqs. (A-IV-1) and (A-IV-2).

$$
\mathrm{b}:=\frac{1}{\sqrt{18 \cdot 10^{8}}} \quad \mathrm{~b}=2.357 \cdot 10^{-5} \mathrm{~m} \quad \alpha \mathrm{~L}:=.837758040957 \mathrm{rad} \quad \mathrm{v}:=.0162 \frac{\mathrm{~m}}{\mathrm{~s}}
$$

Length traveled in the laser beam by particle $=L$

$$
\begin{aligned}
& \mathrm{L}:=\frac{\mathrm{b}}{\cos (\alpha \mathrm{~L})} \quad \mathrm{L}=3.523 \cdot 10^{-5} \mathrm{~m} \\
& \text { transit }:=\frac{\mathrm{L}}{\mathrm{v}} \quad \text { ttransit }=0.002174 \mathrm{sec}
\end{aligned}
$$

Time constant associated with Doppler beating if the bisector between the angles created by the laser arm and the detector arm is not perpendicular to the flow vector. For this calculation, it is assumed that the cell rotation is 5 degrees from perpendicular ( 85 degrees).
rdoppler $:=\frac{1}{\mathrm{k} \cdot \cos (1.4835298642) \cdot \mathrm{v}}$
tdoppler $=0.000045 \mathrm{sec}$

## Time constants of $0.107 \mu \mathrm{~m}$ PSL particles

Time constant associated with Brownian motion of the particles

$$
\begin{array}{cc}
\mathrm{a}:=0.107 \cdot 10^{-6} \mathrm{~m} & \mathrm{~Kb}:=1.380658 \cdot 10^{-23} \frac{\mathrm{~J}}{\mathrm{~K}} \quad \mathrm{n}:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
\mathrm{~T}:=296 \mathrm{~K} & \lambda:=532.5 \cdot 10^{-9} \quad \mathrm{~m} \\
\text { Do }:=\frac{\mathrm{Kb} \cdot \mathrm{~T}}{6 \cdot \pi \cdot \eta \cdot \mathrm{a}} \quad \text { Do }=2.138 \cdot 10^{-12} \frac{\mathrm{~m}^{2}}{\mathrm{~s}} \quad \mathrm{k}:=\frac{2 \cdot \pi \cdot \mathrm{n}}{\lambda} \quad \mathrm{k}=1.569 \cdot 10^{7} \quad \mathrm{~m}^{-1} \\
\text { 亿diffusion }:=\frac{1}{\mathrm{Do} \cdot \mathrm{k}^{2}} \quad \text { 亿diffusion }=0.0019 \quad \mathrm{sec}
\end{array}
$$

Time constant associated with the travel of the particle through the focused laser beam.

Focused beam diameter in sample $=\mathrm{b}$ and is defined in Sundaresan (1999). The velocity used is the value found by the flow profile Eqs. (A-IV-1) and (A-IV-2) at 1 mm from the wall.

$$
b:=\frac{1}{\sqrt{18 \cdot 10^{8}}} \quad b=2.357 \cdot 10^{-5} \mathrm{~m} \quad \quad \alpha L:=.837758040957 \mathrm{rad} \quad v:=9.2618 \cdot 10^{-3} \frac{\mathrm{~m}}{\mathrm{~s}}
$$

Length traveled in the laser beam by particle $=\mathrm{L}$

$$
\begin{aligned}
& \mathrm{L}:=\frac{\mathrm{b}}{\cos (\alpha \mathrm{~L})} \quad \mathrm{L}=3.523 \cdot 10^{-5} \mathrm{~m} \\
& \text { transit }:=\frac{\mathrm{L}}{\mathrm{v}} \quad \text { ttransit }=0.003803 \mathrm{sec}
\end{aligned}
$$

Time constant associated with Doppler beating if the bisector between the angles created by the laser arm and the detector arm is not perpendicular to the flow vector. For this calculation, it is assumed that the cell rotation is 5 degrees from perpendicular ( 85 degrees).
$\tau$ doppler $:=\frac{1}{\mathrm{k} \cdot \cos (1.4835298642) \cdot \mathrm{v}}$

$$
\text { Tdoppler }=0.000079 \mathrm{sec}
$$

## Time constants associated $0.204 \mu \mathrm{~m}$ PSL particles

Time constant for Brownian motion

$$
\begin{aligned}
& \mathrm{a}:=0.204 \cdot 10^{-6} \mathrm{~m} \\
& \mathrm{~Kb}:=1.380658 \cdot 10^{-23} \frac{\mathrm{~J}}{\mathrm{~K}} \\
& \eta:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
& T:=296 \mathrm{~K} \\
& \mathrm{n}:=1.33 \\
& \lambda:=532.5 \cdot 10^{-9} \quad \mathrm{~m} \\
& \text { Do }:=\frac{\mathrm{Kb} \cdot \mathrm{~T}}{6 \cdot \pi \cdot \eta \cdot \mathrm{a}} \quad \mathrm{Do}=1.121 \cdot 10^{-12} \quad \frac{\mathrm{~m}^{2}}{\mathrm{~s}} \quad \mathrm{k}:=\frac{2 \cdot \pi \cdot \mathrm{n}}{\lambda} \quad \mathrm{k}=1.569 \cdot 10^{7} \quad \mathrm{~m}^{-1} \\
& \text { rdiffusion : }=\frac{1}{\text { Do } \cdot \mathrm{k}^{2}} \quad \text { } \quad \text { diffusion }=0.003622 \mathrm{sec}
\end{aligned}
$$

Time constant associated with the travel of the particle through the focused laser beam.

Focused beam diameter in sample $=\mathrm{b}$ and is defined in Sundaresan (1999). The velocity used is the value found by the flow profile Eqs. (A-IV-1) and (A-IV-2) at 1 mm from the wall.

$$
b:=\frac{1}{\sqrt{18 \cdot 10^{8}}} \quad b=2.357 \cdot 10^{-5} \quad \mathrm{~m} \quad \alpha L:=837758040957 \mathrm{rad} \quad v:=9.2618 \cdot 10^{-3} \frac{\mathrm{~m}}{\mathrm{~s}}
$$

Length traveled in the laser beam by particle $=\mathrm{L}$

$$
\begin{aligned}
& \mathrm{L}:=\frac{\mathrm{b}}{\cos (\alpha \mathrm{~L})} \quad \mathrm{L}=3.523 \cdot 10^{-5} \mathrm{~m} \\
& \text { transit }:=\frac{\mathrm{L}}{\mathrm{v}} \quad \text { ttransit }=0.003803 \mathrm{sec}
\end{aligned}
$$

Time constant associated with Doppler beating if the bisector between the angles created by the laser arm and the detector arm is not perpendicular to the flow vector. For this calculation, it is assumed that the cell rotation is 5 degrees from perpendicular ( 85 degrees).
tdoppler $:=\frac{1}{\mathrm{k} \cdot \cos (1.4835298642) \cdot \mathrm{v}}$
tdoppler $=0.000079 \mathrm{sec}$

Flow profile calculations for flow 10 times the value found in the actual test cell during experiments.

$$
\begin{array}{rl}
\mathrm{c}:=0.003 \mathrm{~m} \quad \mathrm{~b}:=0.004 \mathrm{~m} & \mathrm{Q} 1:=224.025 \frac{\mathrm{~mL}}{\mathrm{~min}} \quad \mu:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
\text { convertq }:=\frac{0.000001}{60} & \mathrm{Q}:=\mathrm{Q} 1 \cdot \text { convertq } \quad \mathrm{Q}=3.734 \cdot 10^{-6} \frac{\mathrm{~m}^{3}}{\mathrm{sec}}
\end{array}
$$

$$
\mathrm{j}:=1,3 . .11
$$

$$
\begin{aligned}
& \text { dpdx }:=\frac{-Q}{\frac{4 \cdot b \cdot c^{3}}{3 \cdot \mu} \cdot\left[1-\frac{192 \cdot c}{\pi^{5} \cdot b} \sum_{j} \frac{\tanh \left(\frac{j \cdot \pi \cdot b}{2 \cdot c}\right)}{j^{5}}\right]} \quad \operatorname{dpdx}=-45.4015 \quad \frac{P a}{m} \\
& v(y, z):=\left(\frac{16 \cdot c^{2}}{\mu \cdot \pi^{3}}\right) \cdot(-\operatorname{dpdx}) \cdot\left[\square_{j}(-1)^{\frac{j-1}{2}} \cdot\left[1-\frac{\cosh \left(\frac{j \cdot \pi \cdot z}{2 \cdot c}\right)}{\cosh \left(\frac{j \cdot \pi \cdot b}{2 \cdot c}\right)}\right) \cdot\left(\frac{\cos \left(\frac{j \cdot \pi \cdot y}{2 \cdot c}\right)}{j^{3}}\right)\right]
\end{aligned}
$$

$$
\text { ny }:=10 \quad \text { j }:=0 . . \text { ny } \quad n z:=10 \quad k:=0 . . n z
$$

$$
y_{j}:=-c+2 \cdot c \cdot \frac{j}{n y} \quad z_{k}:=-b+2 \cdot b \cdot \frac{k}{n z} \quad V_{j, k}:=v\left(y_{j}, z_{k}\right)
$$

Velocity at wall Velocity at 1 mm from wall Centerline velocity

$$
v(0.003,0)=0 \quad \frac{\mathrm{~m}}{\sec } \quad v(0.002,0)=0.0926 \quad \frac{\mathrm{~m}}{\mathrm{sec}} \quad v(0,0)=0.1615 \quad \frac{\mathrm{~m}}{\mathrm{sec}}
$$

The values given below depict the position in the $y$ - and $z$-directions at which the velocity calculations were performed to create the following velocity table and flow profile graph. Integer designation was used on the graph, rather than actual locations in meters.

$\mathrm{y}^{\mathrm{T}}=$|  | -0.003 | -0.002 | -0.002 | -0.001 | -0.001 | 0 | 0.001 | 0.001 | 0.002 | 0.002 | 0.003 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | m



Velocities ( $\mathrm{m} / \mathrm{sec}$ ) associated with the positions given by $y$ and $z$

$\mathrm{V}^{\mathrm{T}}=$| 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| 0 | 0.0295 | 0.048 | 0.0598 | 0.0664 | 0.0684 | 0.0664 | 0.0598 | 0.048 | 0.0295 | 0 |
| 0 | 0.0455 | 0.077 | 0.0979 | 0.1099 | 0.1136 | 0.1099 | 0.0979 | 0.077 | 0.0455 | 0 |
| 0 | 0.0547 | 0.0942 | 0.1211 | 0.1367 | 0.1416 | 0.1367 | 0.1211 | 0.0942 | 0.0547 | 0 |
| 0 | 0.0595 | 0.1033 | 0.1335 | 0.1511 | 0.1568 | 0.1511 | 0.1335 | 0.1033 | 0.0595 | 0 |
| 0 | 0.061 | 0.1061 | 0.1374 | 0.1556 | 0.1615 | 0.1556 | 0.1374 | 0.1061 | 0.061 | 0 |
| 0 | 0.0595 | 0.1033 | 0.1335 | 0.1511 | 0.1568 | 0.1511 | 0.1335 | 0.1033 | 0.0595 | 0 |
| 0 | 0.0547 | 0.0942 | 0.1211 | 0.1367 | 0.1416 | 0.1367 | 0.1211 | 0.0942 | 0.0547 | 0 |
| 0 | 0.0455 | 0.077 | 0.0979 | 0.1099 | 0.1136 | 0.1099 | 0.0979 | 0.077 | 0.0455 | 0 |
| 0 | 0.0295 | 0.048 | 0.0598 | 0.0664 | 0.0684 | 0.0664 | 0.0598 | 0.048 | 0.0295 | 0 |
| 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |



## Time constants of $0.107 \mu \mathrm{~m}$ PSL particles for 10 times the actual flow rate

Time constant associated with Brownian motion of the particles

$$
\begin{aligned}
& \mathrm{a}:=0.107 \cdot 10^{-6} \mathrm{~m} \\
& \mathrm{~Kb}:=1.380658 \cdot 10^{-23} \frac{\mathrm{~J}}{\mathrm{~K}} \\
& \eta:=947.95 \cdot 10^{-6} \mathrm{~Pa} \cdot \mathrm{sec} \\
& \text { T:=296K } \\
& \text { n : }=1.33 \\
& \lambda:=532.5 \cdot 10^{-9} \quad \mathrm{~m} \\
& \text { Do }:=\frac{\mathrm{Kb} \cdot \mathrm{~T}}{6 \cdot \pi \cdot \eta \cdot \mathrm{a}} \quad \mathrm{Do}=2.138 \cdot 10^{-12} \quad \frac{\mathrm{~m}^{2}}{\mathrm{~s}} \quad \mathrm{k}:=\frac{2 \cdot \pi \cdot \mathrm{n}}{\lambda} \quad \mathrm{k}=1.569 \cdot 10^{7} \mathrm{~m}^{-1} \\
& \text { đdiffusion }:=\frac{1}{\text { Do } \cdot \mathrm{k}^{2}} \\
& \text { } \tau \text { diffusion }=0.0019 \quad \mathrm{sec}
\end{aligned}
$$

Time constant associated with the travel of the particle through the focused laser beam.

Focused beam diameter in sample $=b$ and is defined in Sundaresan (1999). The velocity used is the value found by the flow profile Eqs. (A-IV-1) and (A-IV-2) at 1 mm from the wall.
$b:=\frac{1}{\sqrt{18 \cdot 10^{8}}}$
$\mathrm{b}=2.357 \cdot 10^{-5} \mathrm{~m}$
$\alpha \mathrm{L}:=.837758040957 \mathrm{rad}$
$v:=9.2618 \cdot 10^{-2} \frac{\mathrm{~m}}{\mathrm{~s}}$

Length traveled in the laser beam by particle $=\mathrm{L}$

$$
L:=\frac{b}{\cos (\alpha L)} \quad L=3.523 \cdot 10^{-5}
$$

ttransit $:=\frac{\mathrm{L}}{\mathrm{v}} \quad$ ttransit $=0.00038 \quad \mathrm{sec}$

Time constant associated with Doppler beating if the bisector between the angles created by the laser arm and the detector arm is not perpendicular to the flow vector. For this calculation, it is assumed that the cell rotation is 5 degrees from perpendicular ( 85 degrees).
$\tau$ doppler $:=\frac{1}{\mathrm{k} \cdot \cos (1.4835298642) \cdot \mathrm{v}}$ $\tau$ doppler $=0.000008 \mathrm{sec}$

## VITA

## Ryan Matthew Cambern

Candidate for the Degree of
Master of Science

## Thesis: MULTIPLE SCATTERING SUPPRESSION CROSS-CORRELATION OF A FLOWING FLUID TO DETERMINE PARTICLE SIZE

Major Field: Mechanical Engineering
Biographical:
Personal Data: Born in Beaver, Oklahoma, January 10, 1974. The son of Charles Cambern and Joe and Lou Ann Perkins.

Education: Graduated from Union High School, Tulsa Oklahoma in May 1992; received Bachelor of Science degree in Mechanical Engineering from Oklahoma State University, Stillwater, Oklahoma in May 1997. Completed the requirements for the Master of Science Degree at Oklahoma State University in July 1999.

