# EXPERIMENTAL AND MOLECULAR DYNAMICS STUDIES IN COPPER ELECTRO CHEMICAL MECHANICAL POLISHING (Cu-ECMP)

By

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#### Abstract:

In recent times, copper electro chemical mechanical polishing (Cu ECMP) has received a great deal of interest from electrochemists and the semiconductor manufacturing industry. This attention is primarily due to its potential for yielding relatively defect-free surfaces with improved surface integrity compared to chemical mechanical planarization (CMP). In this work, Cu ECMP apparatus integrated with a sensing and data acquisition system was developed to polish  $\Phi$ 4 inch ( $\Phi$ 100 mm) blank Cu wafer surfaces to a finish of R<sub>a</sub>< 15nm, and continuously gather voltage and current signals during Cu ECMP process at a sampling rate of 100Hz. Experimental studies were carried out to understand the effects of anodic voltage, pH, and pad pressure on the material removal rate (MRR) and surface roughness ( $R_a$ ). Understanding the process from an atomistic standpoint helps us gain better control over the process and aids us in optimizing the key process output variables (KPOV). In order to gain a better understanding of the process, the molecular dynamic simulation (MDS) technique was adopted to develop a model to depict the real-time formation of copper (II) hexa-hydrate molecule  $Cu[(H_2O)_6]^{2+}$ , which is one of the key elements of the passivation layer formed over the Cu surface during ECMP. The behavior of the complex molecule under an electric force field was simulated to observe the process from a molecular perspective. From the trajectory of Cu<sup>2+</sup>, it was found that the velocity of copper ion increased with increase in applied voltage. Furthermore, the current carried by a single Cu<sup>2+</sup> ion was computed based on the applied voltage and velocity of the ion.

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## CHAPTER I

## RESEARCH OVERVIEW AND OBJECTIVE

Chapter I elucidates the need for an advanced integrated circuit planarization process. A review of the literature related to different planarization processes and a conceptual review of the mechanisms of the copper electrochemical mechanical planarization (Cu-ECMP) process are discussed. Lastly, the contribution of this research work to the academic and industrial knowledge is presented and the organization of the reminder of the thesis is summarized.

#### 1.1 Significance of the electro chemical mechanical polishing (ECMP) process

"The number of transistors on integrated circuits doubles approximately every 18 months"

- Gordon E. Moore, 1965

Over the past four decades, as Moore's law played out, the size of Si-based semiconductor components decreased and the device density increased by some 4 orders of magnitude. As future technologies are increasing device density, shrinking device dimensions and integrating novel device structures, semiconductor processing has been considered an alternative approach to effectively and efficiently accommodating this increase in transistor density. Debatably, planarization is one of the most critical, costly, and poorly understood processes in semiconductor processing.

A good planarization process is required to achieve a desirable depth of focus while patterning the wafer by lithography. In order to achieve the critical planarity required by future (sub 22nm) technology nodes, there ought to be new technology drivers in planarization, for both metal and dielectrics. After the wafer patterning, a barrier film made primarily of tungsten or tantalum is formed on the Si surface either by physical vapor deposition (PVD) or chemical vapor deposition (CVD) methods. This layer prevents the doping of Si by copper. Deposition of the barrier film is followed by formation of interconnects. Previously aluminum and tungsten metal interconnects were primarily used in IC chips because of their good electrical conductivity and low density. But as the feature sizes are being significantly reduced, better materials such as copper, which has higher electrical conductivity and superior resistance to electro migration, are needed to replace Al interconnects. Aluminum lines have been fabricated by the subtractive dry etch process but dry etching of copper is difficult. However, Cu interconnections have been incorporated successfully by means of electroplating and dual damascene techniques [1]. Copper exhibits the possibility of diffusing into Si and  $SiO_2$ , thereby degrading the device performance by introducing deep electronic levels into the Si band gap, eventually reducing the life time of the minority carrier. Moreover, the presence of Cu-Si precipitates in critical regions intensely affects the reverse leaking current of p-n junctions. This diffusion problem was solved by using a tantalum, tungsten, titanium or Ta alloy diffusion barrier layer [2, 3], which were successfully formed on the Si wafer surface by means of physical vapor deposition (PVD) and chemical vapor deposition (CVD) processes.

To date, chemical mechanical planarization (CMP) is being used to planarize the surface. CMP combines the effect of chemical (through slurry) and mechanical (through pad pressure) processes to achieve the desired planarity. Although CMP is an effective planarization technique, several quality issues including within-wafer non-uniformity (WIWNU), surface scratches, erosion, dishing, delamination, and peeling are known to appear during the polishing of soft metallic interconnects and low-k dielectrics. These glitches eventually lead to high operation and process maintenance costs [4 - 6]. The crucial problem for CMP arose when low-k dielectric materials were introduced in the mid-1990s by the microelectronics industry to address the need for preventing gate leakage and to reduce the heat dissipation. The only way to increase the material removal rate (MRR) in CMP is by increasing the pad pressure or down force. However, higher pad pressures might lead to delamination and other defects at the Cu interconnect and with low-k dielectric materials, which have a deleterious effect on wafer yields due to the bending of the pad, as low-k dielectric materials are porous and brittle.

Electro-chemical polishing (ECP), another polishing process, has the potential to eliminate some significant defects produced in CMP such as scratches, pitting, and dishing. ECP basically works on the principle of anodic leveling, which is associated with the removal of imperfections on the surface by the dissolution of metal ions in the presence of current, distributed evenly on the surface [7]. Though ECP was able to produce a defect-free surface, better control over the surface reactions and there-by the material removal rate could not be achieved. Hence a process which can produce a defect-free surface and also better control over the material removal rate is called for.

Many such investigations in CMP and ECP gave rise to a new revolutionary planarization process called the electro chemical mechanical process (ECMP). Sato *et al.* [8] proposed the ECMP process based on the electrochemical dissolution of Cu and reported that it could be used to form erosion-free and scratch-free damascene Cu interconnects at a pressure which is one-tenth of that used for conventional CMP technology. Semiconductor industries look for fast, efficient, and cost effective processes, and ECMP holds the key to achieving effective and efficient planarization of copperdeposited semiconductor wafers for manufacturing thin integrated chips (IC) for electronic products like cell phones, tablets, and laptops and for automobile and aerospace applications.

## 1.2 Mechanics of electro chemical mechanical polishing (ECMP) process

In the Cu ECMP process, the Cu wafer is made the anode and a metal such as aluminum, which can reduce hydrogen rather than undergoing oxidation when compared to copper, is made the cathode. Electrolytic corrosion ensues when the anode and the cathode (different metals) are immersed in a conductive electrolytic bath, i.e., because of the anodic electric potential, the electrons migrate from anode to cathode, and the copper at the wafer surface is oxidized, leading to the formation of  $Cu^{2+}$  ions. The oxidation reaction at the Cu anode is given by

 $Cu(s) \rightarrow Cu^+ + e^-$ 

$$Cu(s) \rightarrow Cu^{2+} + 2e^{-1}$$

In ECMP,  $Cu^{2+}$  ions are prominent than  $Cu^+$  ions because of the applied anodic potential, which promotes electron liberation from the anode. The  $Cu^{2+}$  combines with the oxidizer and other components such as corrosion inhibitor, a complexing agent and an electrolyte present in the slurry leading to the formation of a 0.1 -  $10^3$  nm thin film of a soft copper oxide complex that serves as a passivation layer. So the basic composition of the passivation layer should have metallic oxides and hydroxides as well as metallic complexes. Formation of the oxides present in the passivation layer on the anode (the Cu wafer in this case) can be summarized as follows [9]:  $2Cu^+ + H_2O \rightarrow Cu_2O + 2H^+$  $Cu^{2+} + 2OH^- \rightarrow Cu (OH)_2$  $Cu_2O + H_2O \rightarrow 2CuO + 2H^+ + 2e^-$ 

Formation of hydroxides and the copper complex with phosphoric acid slurry is given as follows [10, 11, 21, and 28]:

$$\operatorname{Cu}^{2+} + \operatorname{H}_3\operatorname{PO}_4 + \operatorname{H}_2\operatorname{O} \rightarrow [\operatorname{Cu}\operatorname{OH}]^+ [\operatorname{H}_2\operatorname{PO}_4]^- + 2\operatorname{H}^+ + 2\operatorname{e}^-$$

The passivation layer is then gently removed by the mechanical action of the polishing pad and the abrasive particles present in the electrolyte [12]. Since an electrolytic dissolution of material accompanies the chemical mechanical polishing action, a very low mechanical down force, as low as 1/10th of that applied for CMP, is sufficient to planarize the asperities [11]. Therefore, many defect-generating mechanisms in the CMP process are absent, which leads to improve wafer quality. Low down force is the most preferred process condition for polishing soft metals and low-k dielectric materials that are porous and soft. The ECMP polished surfaces tend to be devoid of defects like surface cracking, pits, and delamination.

In general, 4 electrochemical regions are observed during the ECMP process, based on an increase in anodic potential:

- The active region, where the applied anodic voltage is very low. The material removed in this region is primarily due to the mechanical action of the polishing pad. Some authors have also claimed that there can be no material removal at this region [11, 12]. If the process is operated in a low pH region, prominent corrosion marks could be smeared over the surface rather than being removed, as the pad pressure is low in ECMP. Also there is no improvement in surface quality, and the surface quality might deteriorate.
- 2) The passive region is characterized by an increase in corrosion and formation of the passivation layer. More pad pressure is required to remove the continuous formation passivation layer. This removal tends to induce more scratches, delamination, and peeling defects. Although in this region the MRR might increase, the surface quality tends to be poor due to the increase in pad pressure.
- 3) The transient region is the best region to work with in ECMP. The passivation layer becomes softer and less stable at this region. The electrolyte present in the slurry acts against the applied anodic potential, leading to an increase in Cu<sup>+</sup>/Cu<sup>2+</sup> stripping from the copper surface. These copper ions further react with the components of the slurry and form a soft passivation layer over the copper wafer which can be easily removed using low pad pressure.
- 4) In the trans-passive region, film formation and breakdown happens as a part of a dynamic equilibrium. So with the optimal pH (near neutral), we can achieve the good surface finish with high MRR.

Figure 1.1 provides a pictorial elucidation of the working principle of Cu ECMP. With application of anodic voltage, a complex oxidation layer called passivation layer is formed over the Cu wafer surface. The passivation layer is removed by the action of polishing with very low pad pressure.



Figure 1.1: Pictorial representation of the Cu ECMP process

## 1.3 Advantages of the electro chemical mechanical polishing (ECMP) process

The principal advantages of ECMP over CMP are speed and efficiency. The ability to planarize the copper oxides off the wafer in less time and with better surface quality and fewer defects makes ECMP more advantageous than CMP. Such abilities can be achieved only by reducing the pad pressure applied over the wafer to planarize the excess copper. In order to tackle the slow-signal relay problem, SiO<sub>2</sub> was replaced by low-k dielectric materials, as a barrier between the interconnects. However, low-k dielectric materials are 4 times softer than SiO<sub>2</sub>. CMP, which primarily relies on the mechanical force exerted on the wafer for removing material, cannot produce the desired level of planarization without destroying the underlying low-k dielectric isolations. Subsequently, ECMP has become the solution because it is able to remove the material quickly and efficiently with less force compared to CMP [13]. With CMP, precise control over the material removal and also determining the level of planarization is difficult. But in ECMP, by controlling the current density applied to the wafer, we can easily determine the material removal rate (MRR) and also the level of planarization.

After the electro-deposition process, a one-micrometer thick mound of copper is formed over the wafer surface, thicker in the middle than at the edges [14]. CMP tends to remove more material from the edge while removing the mound from the center, which consequently increases the resistance in the wires, rendering many chips useless [15]. However, as ECMP's principle mechanism of material removal is electrochemical, it removes material globally; and by just sweeping off the surface, it gives a clean and smoother surface than can be produced in CMP.

It has been officially claimed by Allied Materials that the cost of slurry could be cut down by nearly 30% by using ECMP compared to CMP. The cost of slurry is around \$15 - \$20 for the CMP process and can be reduced by \$4.50 - \$6 in the ECMP process, which can eventually eliminate much of the processing cost from one of the most expensive processes (i.e., CMP is second only to lithography) in the IC manufacturing process.

To summarize, the chief advantages of ECMP include efficient planarization of low-k-dielectric materials with low pad pressure; less polishing time compared to CMP; elimination of surface defects like scratches, delamination, and erosion as seen in CMP; and much easier end point detection effective controlling of applied current density and efficient slurry handling and waste disposal systems [31].

#### 1.4 Challenges in electrochemical mechanical polishing

Despite the great effort put into understanding the process characteristics, two conflicting theories are put forth to explain the formation of the passivation layer in the Cu ECMP process. The first theory proposes that the  $Cu^{2+/}Cu^+$  ions resulting from the oxidation process interact with the inhibitor present in the slurry to form a soft passivation layer on the wafer surface, which in turn is removed easily by the sweeping action of the polishing pad with minimal pad pressure [33]. The second theory is that a

passivation layer of some sort of polymer is deposited on the Cu surface and that this layer has a higher resistance than the exposed areas, and thus mechanical removal by the pad allows for localized Cu dissolution/removal [33]. Hence there is a lack of understanding of how the passivation layer is formed. Consequently end point detection in ECMP becomes difficult.

In order to better understand the process mechanics, methods which can help us to study the interaction of atoms in the fundamental level need to be developed. Such methods can help us observe the behavior of the  $Cu-H_2O$  complex ions formed during the process. Molecular dynamics simulation, one such method, can help us mimic the formation of copper complex molecules precisely in the way they are formed during the actual ECMP process. Studying Cu-ECMP from an atomistic standpoint requires attention, as it can aid us in determining the process mechanism accurately and thereby optimizing the MRR and surface quality in the ECMP process. The result might be process improvement and an ability to perform real time quality control of the process.

#### **1.5 Objective statement**

Copper electro chemical mechanical polishing (Cu- ECMP) is a relatively new polishing process which has the potential to replace the currently employed chemical mechanical polishing (CMP) in the semiconductor manufacturing industry in terms of process efficiency and the quality of surface produced. ECMP is known for its higher MRR and effective planarization at low down force. Moreover many of the defects generated by CMP, such as dishing, erosion, and delamination, can be reduced/completely eliminated. But the optimum parameters for planarization efficiency are yet to be identified. Dynamic changes in electrochemical regions in ECMP, i.e., active, passive, transient, and trans-passive regions, require better understanding. Also the passivation layer formation mechanism is least understood and requires more insight so that we can design the process so as to achieve better qualitative and quantitative yield. The objectives of this study are as follows:

- Improve the efficiency of the existing Cu-ECMP setup and integrate a data acquisition system with the ECMP setup to capture the time frequency patterns of V-I dynamics to identify the process mechanisms as well as to estimate the effects of various process parameters and thus optimize MRR and surface quality in ECMP.
- 2. Develop a molecular dynamics simulation model to study the formation of the cuprous hexa-hydrate complex molecule, i.e.,  $Cu[(H_2O)_6]^{2+}$ , which is one of the significant components of the passivation layer and slurry formed during the Cu-ECMP process. Also to observe the behavior of the complex molecule in the presence of electrode potential to better understand the process.

## CHAPTER II

#### BACKGROUND AND REVIEW OF LITERATURE

#### 2.1 Prior experimental investigations into ECMP process mechanisms

Characterization of ECMP process, including the effects of applied voltage, current density, down force and chemical composition has received some attention in the literature [2, 15-18, 26, 27]. In this chapter let us discuss some of the works researchers have done related to the Cu-ECMP process.

## 2.1.1 Effect of electrochemical actions on the process

Work done by Seo *et al.* [14] provides vital information about the characterization of the electrochemical mechanisms of the ECMP process. He studied the voltageactivated electrochemical reactions of Cu with different concentrations of HNO<sub>3</sub> based slurry. From the I-V characteristic curve obtained through concentrations of HNO<sub>3</sub> in electrolyte sweep voltammetry (LSV) and cyclic voltammetry (CV) methods, he tried to evaluate the electrochemical characteristics of Cu, such as the active, passive, transient, and trans-passive states. His work reveals that the I-V relationship depends strongly on the concentration of the electrolyte. Also the active, passive, transient and trans- passive regions has been clearly mapped out using I-V plots. One interesting feature of these plots is that the first V = IR drop indicates the beginning of a passive state and the second IR drop indicates the beginning of a trans-passive state, as shown in dotted boxes (A) and (B) of Figure 2.1. This drop indicates an abrupt increase in resistance at these regions and this phenomenon indicates that a passive oxide such as CuO or Cu<sub>2</sub>O was formed on the Cu surface.



**Figure 2.1:** I–V curve of  $HNO_3$  electrolyte with different concentrations. (A) Full I–V curve, (B) Detailed I–V curve between -0.5 And 0.5 V [12]

The DC voltage-controlled deposition and stripping mechanisms of Cu were investigated through the cyclic voltammetry (CV) curve. Noticeably different signs of deposition and stripping of Cu were observed. The cathodic current peaks correspond to the deposition of  $Cu^{2+}$ , whereas the anodic peaks represent the stripping of bulk and under-potentially deposited copper. Seo *et al.* [14] have concluded that the HNO<sub>3</sub>

electrolyte acts against voltage-activated dissolution of Cu by forming a passivation layer composed of oxides on the Cu surface. Verification of the stripping and deposition mechanism through SEM analyses has clearly depicted the surface of Cu at different electrochemical states, as shown in Figure 2.2.



**Figure 2.2:** SEM surface photographs at different operating voltage in 0.1 wt. % HNO<sub>3</sub> electrolyte. (A) Original surface of no bias, (B) active state: -1.0 V, (C) passive state: -0.25 V and (D) trans-passive state: 0.7 V [12].

Figure 2.2 (A) shows the original Cu surface with some scratches, and no bias was applied over the surface. In the active state, as shown in Figure 2.2 (B), the number of scratches was reduced but some bubble marks along with pitting type corrosion were observed. Figure 2.2 (C) depicts the passive state, where severe corrosion and a passivation layer composed of oxides were observed. In the trans-passive region, as shown in Figure 2.2 (D), a sponge-like non-uniform surface was observed, mainly because of the stripping of Cu as a result of electrochemical reactions on the Cu surface. Thus we can infer that by proper combinations of voltage activation and electrolyte, accurate control over the material removal rate and surface quality can be achieved in the ECMP process.

Han et al. [15] studied the current-voltage (I-V) curves to evaluate the effect of electrolytic concentrations of nitric acid (HNO<sub>3</sub>)<sup>-</sup> and potassium nitrate (KNO<sub>3</sub>)<sup>-</sup> based slurry on the electrochemical surface reaction of the Cu electrode. In order to study the electrochemical action in detail, they conducted the experiments in the absence of mechanical polishing. I-V curves were plotted as a function of slurry concentrations of 5%, 10%, and 20%. They observed that for both KNO<sub>3</sub> based slurries, as the concentration increased the trans-passive voltages decreased and the current density increased. Thus a passivation region is formed because of the formation of a continuous passive film composed of oxide and metallic hydroxide which remains on the Cu surface. As the voltage is further increased, the metallic ions get pulled off through the passive film, because of the strength of the potential. This same observation was recorded for  $HNO_3$  slurry, but the trans-passive state was steeper. From this work, we can conclude that the I-V relationship strongly depends on the concentration of the electrolyte. Han et al. [15] also investigated the DC voltage-controlled re-deposition and stripping mechanism of Cu from the cyclic voltammetry (CV) curve. They observed distinctly

different signs of deposition and stripping of Cu. The cathodic current peaks correspond to deposition of  $Cu^{2+}$ , whereas the anodic peaks represent the stripping of bulk and underpotentially deposited copper.



Figure 2.3: I-V curves as a function of electrolyte concentration of (a)  $KNO_3$  and (b)  $HNO_3$  for Cu [11]

The I-V characteristics study by Lee *et al.* [17] on ECMP of SKD11 steel suggests that the transient state is the most suitable for achieving optimal surface quality, and the trans-passive state is preferred for better polishing efficiency (i.e., MRR) [15, 24]. Recent studies in ECMP with  $H_3PO_4$  slurry by Chen *et al.* [12] strongly suggest that at low electric potential, material removal is predominantly due to the mechanical action of the pad whereas at higher potentials, the Cu removal rate is increased due to the effect of the electric potential enhanced by mechanical force.

### 2.1.2 Effect of slurry composition on the process

In addition to voltage activated material removal, the slurry plays a vital role in the formation of the passivation layer. Corrosion inhibitor, oxidizer, and a pH regulator are three main ingredients of the slurry. Together they govern the formation of the passivation layer. Usually BTA is added as a corrosion inhibitor,  $H_2O_2$  as an oxidizer, and organic salts like KOH for adjusting the pH in acidic ECMP slurry. Kristin *G et al.* [23] conducted planarization studies using phosphoric-acid-based slurry containing 1.0 M potassium phosphate salt concentration with a pH value of 2 and a benzotriazole (BTA) concentration of 0.001 M to check the planarization efficiency of patterned Cu structures. The authors claim that because the ECMP process exploits the combined effect of the anodic dissolution of a metal via an applied voltage and the delicate abrasion of the sample surface from a polishing pad, there is no need for an oxidizer and abrasive slurry particles, which are known to contribute to surface defects. They have focused their studies mainly towards the corrosion inhibition characteristics of BTA.

According to Lin *et al.* [24] and Finsgar *et al.* [25], the BTA inhibition layer can be formed by chemisorption of Cu and BTA in the form of Cu (BTAH)<sup>ad</sup> or by forming a polymeric BTA-Cu complex film over the Cu surface. In our case (acidic slurry), it is always BTAH, which is chemisorbed onto the surface. Though BTAH is not as strong a layer as polymeric BTA-Cu complex, it is able to protect the Cu surface from electrochemical dissolution. This layer is then easily removed by slight abrasion from a polishing pad, as it is lightly attached to the Cu surface. This mechanism has been elucidated in Figure 2.4.



Figure 2.4: Pictorial Representation of BTA removal mechanism [20]

Initially Lin *et al.* [24] and Finsgar *et al.* [25] compared the planarization capability of a custom-built ECMP tool with a rotating disk electrode-type (RDE) method using the same phosphoric-acid-based slurry with different concentrations of BTA as corrosion inhibitors. The following equation was used to predict the planarization capabilities of the ECMP electrolyte using electrochemical methods on a bench top ECMP tool, using blanket wafer fragments, and a Cu RDE:

$$\varepsilon = \frac{S_0 - S_f}{\lambda_{Avg}}$$

where,  $S_0$  and  $S_f$  are the pre- and post-process step heights, and  $\lambda_{Avg}$  is the average material thickness removed, calculated from the work piece mass before and after the process. Increasing values of  $\varepsilon$  imply an increase in feature scale planarization. In Figure

2.5 (A), the anodic polarization curves for oxidation of a Cu RDE with a rotation speed of 100 RPM have been shown.



Figure 2.5 (A): Anodic Polarization for oxidation of a Cu RDE with a rotation speed of 100 RPM.

Curves (a) and (b) are electrolytes of concentrated phosphoric acid (pH ~0) containing (a) cBTA = 0, (b) cBTA = 0.01 M. Curves (c) and (d) are electrolytes containing 1.0 M potassium phosphate salt at a pH value of 2 containing (c) cBTA = 0.0 M, (d) cBTA = 0.001 M., after Kristin G *et al.* [23]. From these curves, we can infer that in concentrated phosphoric acid, the presence of BTA, even at a concentration on an order of magnitude higher than that in a pH value of 2, has little effect on the polarization curve. Figure 2.5 (B) depicts the anodic polarization curves for oxidation of a Cu wafer on the ECMP tool operating without pad contact.



**Figure 2.5 (B):** Anodic polarization for oxidation of a Cu wafer on the ECMP tool operating without pad contact.

Electrolytes containing 1.0 M potassium phosphate salt at a pH value of 2 containing (a) cBTA = 0.0 M, and (b) cBTA = 0.001 M, after Kristin G *et al.* [23]. This electrochemical experiment demonstrates that in a specific electrical potential region, good passivation at low inhibitor concentrations and high MRR can be achieved; additionally with its simplified bath chemistry, this electrolytic slurry offers considerable advantages over previous electrolytes studied.

Kristin G *et al.* [23] have also conducted planarization experiments using three different pads, Suba<sup>TM</sup>, IC1000<sup>TM</sup>, and D100<sup>TM</sup> with an electrolyte containing 1.0 M potassium phosphate salt at a pH value of 2 with a BTA concentration of 0.001 M at 0.5 V. From this work we can infer that when no potential is applied to the surface of the wafer, no Cu is removed, even with pad contact. Similarly when 0.5 V is applied, if there

is no pad contact, there is no Cu removal. Therefore consistent with *Shattuck et al.* [26], it has been shown that both potential and pad contact are required for material removal in the ECMP process.

Tripathi et al. [27, 28] studied the influence of pH and a complexing agent (glycine) on the surface quality and Cu removal rate during ECMP by using a blanket copper wafer. One of the major issues with Cu ECMP electrolytes was that, although they provided excellent Cu removal rates and planarization efficiencies [27], the surface roughness was considerably higher than with the Cu CMP. Tripathi et al. [28] used Cu ECMP slurry composed of purified terephthalic acid (PTA) at various pH values. They demonstrated a few important results concerning surface quality. One is that the surface roughness increases following Cu ECMP relative to Cu CMP, decreases as the pH of ECMP is increased from 3 to 6, and decreases with the substitution of a less aggressive complexing agent (glycine) for oxalic acid. pH is a very critical variable in Cu ECMP because it determines in part whether Cu oxide forms or not. As the pH is raised, the Cu oxide that forms on the Cu surface becomes more stable [26-29]. Due to the importance of pH in the passive film formation mechanism, the effects of pH on Cu surface roughness during ECMP were also investigated. Figure 2.6 (A) illustrates the variation of pH with surface roughness in a glycine-containing ECMP slurry.

From Figure 2.6 (A), we observe that as the pH is raised into the range where a stable Cu oxide forms, the surface roughness value is drastically decreased



Figure 2.6 (A): RMS surface roughness vs pH for ECMP electrolytes containing 0.29 m HEDP, 0.13 m glycine, 1 mm PTA, and 3 wt. % colloidal silica. ECMP polishing is done at 3.5 V [26]

Figure 2.6 (B) shows the variation of Cu removal rate with the pH of the slurry. It can be observed that as the pH is raised from 3 to 6, the Cu removal rate declines from 620 to 400nm/min. Analysis of the passive film using Electrochemical Impedance Spectroscopy (EIS) suggests that the passive film formed at an applied anodic potential during Cu ECMP is thicker than that which forms during Cu CMP at the open circuit potential [26].


**Figure 2.6 (B):** Cu removal rate as a function of pH for ECMP electrolytes containing 0.29 m HEDP, 0.13 m glycine, 1 mm PTA, and 3 wt. % colloidal silica. ECMP polishing is done at 3.5 V, [26]

Kulyk *et al.* [5] studied the planarization of patterned PCB with copper interconnects with the ECMP process using an alkali based slurry containing benzotriazole (BTA) as a surfactant. They claim that the main purpose of BTA in their work is to overcome the dishing effect. They attempted to optimize the ECMP process by controlling some of the process parameters such as current, pad type, polishing time, pad rotation speed, and additives in the ECMP slurry. They used two Cu/PCB samples: one with trenches of 150  $\mu$ m ~a few mm wide and 35  $\mu$ m deep and one with trenches of 30 $\mu$ m ~a few mm wide and 30  $\mu$ m deep. Both are overfilled with copper by electro deposition.



Figure 2.7: Cross-Sectional (A) SEM image of Copper ON PCB filled in narrow trenches and (B) Optical Microscopic image of Copper layer filled in millimeter-scale trenches before ECMP process.

Studies by Oh *et al.* [33] indicate that the pH value of the electrolyte should be 14 or higher in order to get a higher polishing rate. This higher polishing rate is due to the formation of mechanically soft phases such as Cu (OH) <sub>2</sub>, HCuO<sup>2-</sup> and CuO<sub>2</sub><sup>-2</sup> under anodic potentials. But from the Pourbaix diagram of copper, Figure 2.8, a hard phase of Cu<sub>2</sub>O is generated mainly in the pH range of 6 - 14 at applied voltages of -0.4 - 0.3 V, as shown by Ng *et al.* [30]. In this work, the electrolyte is composed of a 20-30wt% solution of potassium hydroxide (KOH), a 0.5-3wt% of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>) solution, which acts as an oxidizer, and 0.005-0.02M of BTA as an additive for uniformity improvement. The pH was maintained around 6 by adding phosphoric acid to the slurry in order to avoid the hard Cu<sub>2</sub>O phase, and the anodic potential of -0.4 - 1 V was applied.



**Figure 2.8:** Pourbaix diagram of Copper in  $Cu^{2+}-H_2O$  system [30]

Some typical slurry compositions used in ECMP are the following:

(i) H<sub>3</sub>PO<sub>4</sub> 6 vol. %, ethylenediamine 2 vol. %, ammonium citrate 2 wt. %, BTA 0.3 wt. %, H<sub>2</sub>O<sub>2</sub> 1.5 vol. %, silica abrasives 0.15 wt. % and de-ionized water was used as the solvent [9],

- (ii) H<sub>3</sub>PO<sub>4</sub> 5.1 vol. %, ammonium citrate 2 wt. %, BTA 0.3 wt. %, polymeric inhibitor XP-1296 0.5 vol. %, PolyEthylenImine (PEI) 0.025 vol. % and rest deionized water [9],
- (iii) 1 M potassium phosphate, pure ortho-phosphoric acid, potassium chloride, BTA and de-ionized water [9],
- (iv) Ortho-phosphoric acid, potassium phosphate, BTA and de-ionized water [29].

#### 2.1.3 The effect of polishing pad, down force, and polishing speed on the process

Liu *et al.*'s [13] study of the effect of down force in ECMP of patterned Cu wafers suggests that a low down force significantly reduces the erosion of dense features, which is usually observed in CMP due to high down-force polishing. Pad structure is another main determinant of the planarization quality and efficiency in the ECMP process. Polishing pads are graded based mainly on their hardness and chemical inertness. Hard polyurethane pads are generally used in Cu-ECMP as they are found to offer high resistance to acidic and alkaline slurries and also a minimal dishing effect in the center of wider trenches is observed [8, 22].

Oh *et al.* [33] studied the influence of soft and hard polymer material pads on the planarized samples keeping the experimental conditions identical. Polyvinyl alcohol pads (soft pads) were found to preferentially remove copper on the edges of the trenches and also produce a dishing effect in the center of wider trenches. Oh *et al.* observed that the "dishing effect" originates from the initial non-uniformities in the copper layer obtained

after the electroplating process. When the copper layer is polished out all the way to the wafer level, the second effect starts to appear, namely the "corner over-removal effect." This effect is caused by the deformation of the soft pad on the copper-PCB board interface, which leads to a higher pressure in the pad-copper layer interaction and, as a result, a higher polishing rate in the interface region. Using a polyurethane pad (hard pad), such defects are reduced to a larger extent. Their observation clearly shows that uniform planarization was obtained on the copper layer and that pad structure is one of the main determinants of the planarization quality and efficiency in the ECMP process.

Liu *et al.*'s [13] study of the effect of down force in ECMP of patterned Cu wafers suggests that low down force significantly reduces the erosion of dense features, which is usually observed in CMP due to high down force polishing. ECMP can planarize porous low-k materials which would be damaged in a CMP polishing system because of the high load or down force. Lower down force serves as a main advantage of ECMP process. The copper left over from the formation of interconnects is removed from the low-k film without ripping the insulating layer.

Polishing speed is another important advantage of ECMP over CMP. Standard CMP systems cannot achieve very high speeds unless they increase the pad pressure on the wafer. But if the speed is increased, the high pad pressure cracks and peels away the low-k dielectric film [13]. However, decreasing the pad speed means an increase in the polishing time of the wafer. ECMP can achieve the same planarization in less time because the electrochemical mechanism is accompanied by mechanical polishing.

Figure 2.9 (A) shows the cracks observed in the dielectric layer in chemical mechanical polishing, after *Brown* [18] and Figure 2.9 (B) shows the ECMP with flattened wafer and no damage to the dielectric layer, after *Brown* [18].



Figure 2.9 (A): CMP process

Figure 2.9 (B): ECMP process

Economikos *et al.*, [19] compared CMP process with ECMP process and have shown that ECMP exhibits low dishing than CMP process.

# **2.2 Elements of the ECMP slurry**

The composition of the slurry is a significant factor in ECMP as it should be able to passivate the surface of the wafer [9, 10]. The chemical composition of the slurry should be appropriate so that topographical defects are not formed on the surface. The slurry should have some fixed components, which are essential for producing a good mirror finish in electrochemical mechanical polishing.

The essential components of the slurry are:

a. Acid based ele	ectrolyte	g. Abrasiv	e particulates
-------------------	-----------	------------	----------------

b. Chelating agent

h. Solvent

- c. Corrosion inhibitor
- d. Inorganic/organic salt
- e. pH adjusting agent
- f. Oxidizer

An acid-based electrolyte ensures electrical conductivity in the slurry. Phosphoric acid, acetic acid, nitric acid, per chloric acid, and citric acid are some of the suitable acidic electrolytes used in Cu ECMP. The slurry can contain 1-30% (Vol.) of electrolyte depending on the required conductivity for polishing [8].

Chelating agents or complexing agents form a complex compound with the copper ions and improve the material removal rate. Uniformity in anodic dissolution can also be achieved by these chelating agents. The formation of the passivation layer is also more efficient with the presence of chelating agents. Ethylenediamine (EDA), ammonium citrate, and acetic acid are some good chelating agents [8]. The concentration of these compounds can be around 0.1-15% by volume or weight of the total slurry composition

A corrosion inhibitor is another important aspect of ECMP slurry. It forms a passivation layer between the metal and the slurry, thereby reducing the chemical

interaction between them. Thus the oxidation or corrosion of the substrate is significantly reduced. The concentration of corrosion inhibitors can be between 0.001% and 5% by weight. Organic compounds like benzotriazole (BTA) and benzimidazole and polymeric inhibitors such as polyalkylaryl ether phosphate can be used as corrosion inhibitors [8].

Inorganic or organic salts enhance the formation of complex compounds in the slurry. The concentration of these salts can be between 0.1 - 15% by volume or weight of the total composition. Ammonium or potassium salts are commonly used in Cu ECMP slurry [8].

pH adjusting agents are used in the slurry to achieve and maintain a desired pH value. The permissible pH range is between 2 and 10 in the ECMP process. Usually 0.2 – 25% by volume of the slurry is composed of the pH adjusting agent. The concentration of the pH adjusting agent can be varied according to the composition of other chemicals in the slurry. Common examples include potassium hydroxide, ammonium hydroxide, and sodium hydroxide [8].

Oxidizers are used to enhance the formation of the passivation layer, which improves the material removal rate. In copper polishing, the copper ions from +1 state  $(Cu^+)$  are oxidized to +2 state  $(Cu^{+2})$ , which aids passivation. Hydrogen peroxide is a common oxidizing agent in ECMP slurry [8].

$$Cu(s) \rightarrow Cu^+ + e^-$$
  
 $Cu(s) \rightarrow Cu^{2+} + 2e^-$ 

Abrasive particles are used in the slurry to produce a good surface finish. They also enhance the material removal rate. When abrasives are used, ECMP can achieve a surface roughness equal to that of a conventional CMP process. The concentration of the abrasives can be from 0.001% to 5% by weight of the total composition. Silica, alumina, zirconia, titania and ceria are suitable examples of abrasive particles. Finally, solvent, generally deionized water is used as the remainder of the polishing composition [8].

### 2.3 Relationship between the KPIV and KPOV in ECMP process

Table 2.1 depicts the relationship between the Key Process Input Variables (KPIV) and Key Process Output Variables (KPOV) in ECMP process.

FACTORS	KEY PROCESS INPUT VARIABLES (KPIVS)/ PROCESS PARAMETERS	STATE VARIABLES	KEY PROCESS OUTPUT VARIABLES (KPOVS)/ PERFORMANCE/RESPONSE VARIABLES
Machine (Including the External Power Source)	Down Force, Platen speed, Wafer carrier speed, Input Voltage, Input current	Stress Distribution, Velocity Distribution, Passivation Layer formation	Endpoint thickness of the wafer, Output Voltage &Output Current
Polishing Pad	Stiffness, Topography, Conditioning, Pad Uniformity, Porosity	Contact area, Pad Pressure, Degradation and Temperature Distribution	Material Removal Rate(MRR)(nm/min)

TABLE 2.1 Relationships between KPIV and KPOV

Slurry	Acid based electrolyte, Chelating agent, Inorganic/Organic salt, Oxidizer, pH adjusting agent, Corrosion Inhibitor, Abrasive particulates, Solvent	pH drifts, Concentration, Size distribution, Dispersion, Agglomeration, Aggregation	Within Wafer Non-Uniformity (WIWNU), Wafer to Wafer Non- Uniformity, Within Die Non Uniformity (WIDNU), Erosion, Micro scratch & Pits
Wafer	Size, Curvature, Properties of coating, Coating thickness variation & Pattern Geometry	Direct contact, Hydro planning, Mixed contact	Surface finish, Roughness

## 2.4 MD Simulation models of electrochemical reactions

### 2.4.1 General principles in MD simulation

Molecular dynamics (MD) is a computer simulation method by which the physical movement of atoms and molecules and their interactions can be mapped with respect to time by integrating Newton's equations of motion. MD simulation acts as a bridge between macroscopic laboratory results and microscopic time and length scales. This simulation enables us to predict the bulk properties from the molecular interactions up to a certain degree of accuracy. Thus, physical processes can be explored and analyzed at the fundamental atomic and molecular levels. Every material, irrespective of its complexity, is composed of a large number of atoms ( $\sim 10^{23}$ ) even when the analysis is constrained to a nano-metric scale. Such a constraint requires an enormous computing resource in order to study the atomic interactions. MD simulations also act as a bridge

between theory and experiment. A theory can be tested by conducting a simulation using the same model. This model can also be tested by comparison with experimental results. We may also carry out simulations on the computer that are very difficult or impossible in the laboratory (for example, working at extremes of temperature or pressure).

Another significant issue involves the temporal scale in available MD simulations. In order to replicate atomic vibrations precisely, the maximum time step with which Newton's equations can be integrated numerically is essentially restricted to very small values of about 1 fs. Because of this constraint, many notable atomic motions require a long computational time for the simulations. Such motions are comparatively slower than the fast oscillations of bond lengths and bond angles that limit the integration time step. Generally, MD simulations are limited to molecular interactions and processes that happen in relatively short times, in other words "rare event" phenomena. Such infrequent processes are simulated quickly. Dissociation of a chemical bond, one such process, is completed on the pico-seconds (ps) time scale. However, many experimental events occur on time scales that are orders of magnitude longer than the dissociation time. This feature makes it difficult to observe such events in MD studies. Two approaches are generally followed to overcome the time scale restriction of MD simulations. In the first approach, coarse graining, the dynamic variables are divided into slow and fast degrees of freedom and averaging is performed on the fast degrees of freedom. In the second approach fast (and rare) trajectories between desired states are

computed explicitly, and their probabilities (relative to non-reactive trajectories) are calculated.

#### **2.4.2 Interatomic potentials**

The selection of appropriate interatomic potentials plays a vital role in determining the accuracy of the MD simulation results of any lattice structure model. The accuracy level of the trajectory calculations is also affected by the choice of the potential energy functions. In general, the total energy of the system is the summation of the potential energy and the kinetic energy. The computation of kinetic energy is simple but the computation of potential energy is complex, as it relies on the interatomic distances between the interacting atoms. The computational time required to complete the simulation is determined primarily by the complexity of the potential energy functions used in the model. The force acting on an atom is proportional to the negative slope of the potential curve, i.e., the negative first derivative of the potential function. Evaluation of the forces acting on the atoms is essential to calculate new positions of atoms. The functional form of the potential is chosen in a way that can depict the actual potential in the most realistic way for specific material under specific conditions. Two important steps are involved in designing the potential function for a model:

1. Selection of the analytical form for the potential. The analytical form could be the sum of pairwise terms, where the energy of the pair depends on the interatomic

distance, or a many-bodied form appropriate for a specific type of bonding involving bond distances, bond angles, and coordination number.

2. After the analytical form is decided, the specific parameters of the functions are determined. The set of parameters which best fits the data from *ab initio* calculations or experimental properties such as density and cohesive energy, or from a data set containing a combination of both theoretical and experimental observations, is found.

#### 2.4.3 Calculation of the trajectory and potential energy function

As discussed previously, classical dynamics involves simulating the atomic motion by integrating Newton's equations of motion. We now explore how the trajectory on an atom "*i*" is calculated based on integration of Newton's equation of motion. A trajectory depicts the position of an atom as a function of time. Let the position of the i<sup>th</sup> particle with respect to time be denoted as  $r_i$  (*t*). Let the mass of the atom be  $m_i$ . Then according to Newton's second law of motion,

$$F_i = m_i a_i = m_i \cdot \frac{dv_i}{dt} = m \cdot \frac{d^2 r_i}{dt^2}$$

where  $F_i$  is the total force acting on atom *i*, and  $v_i$  and  $a_i$  are the velocity and acceleration of atom *i* as a function of time. Force is related to acceleration and, in turn, to potential energy. Acceleration  $(a_i)$  and velocity  $(v_i)$  are further represented as,

$$a_i = \frac{dv_i}{dt}$$

The trajectory  $x_i$  (t) is obtained by integrating the initial positions and velocity  $(x_o \text{ and } v_o)$ , as shown in Equation 2.1.

$$x_i(t) = v \cdot t + x_o = a \cdot \frac{t^2}{2} + v_o t + x_o$$
(2.1)

If the force on  $i^{th}$  body due to  $j^{th}$  body is denoted as  $F_{ij}$ , then

$$F_i = \sum_{\substack{j=1\\j \neq 1}}^N F_{ij} \tag{2.2}$$

Apart from the trajectory, the accuracy of the simulation results depends primarily on the interatomic potentials used to model the lattice structure. Also the accuracy of the trajectories obtained from MD simulations is affected by the choice of the potential energy function. The total energy of the system is the sum of kinetic energy (KE) and potential energy (PE). The KE is simple to compute but the PE computation is complex since it depends on the positions of all interacting atoms. Therefore the computational time required depends on the complexity of the potential energy function. The force acting on an atom is proportional to the first derivative of the potential energy function. The largest part of a molecular dynamics simulation is the estimation of the forces required to calculate new positions of atoms. Two approaches are available for determining the potential energy of an interacting system. One approach is to perform *ab initio* calculations by solving Schrodinger's equation for the entire system at each integration step. The term "*ab initio*" is used to define a solution of the non-relativistic, time independent Schrodinger equation  $\hat{H}\psi = E\psi$ .  $\hat{H}$  is the Hamiltonian operator for the system, which is a function of the kinetic and potential energies of the particles of the system,  $\psi$  is the wave function, and E is the energy. The Hamiltonian operator for an N electron atom is shown in Equation 2.3 [44].

$$H = -\frac{\eta^2}{2m_n} \nabla_n^2 - \frac{\eta^2}{2m_e} \sum_{i=1}^N \nabla_i^2 - \sum_{i=1}^N \frac{Z_e^2}{r_i} + \sum_{i=1}^N \sum_{j=i+1}^N \frac{e^2}{r_{ij}}$$
(2.3)

where  $\frac{\eta^2}{2m_n} \nabla_n^2$  represents the kinetic energy of the nucleus that comprises the coordinates of the nucleus and derivatives with respect to the coordinates of the nucleus. It is also called the Zero-electron term, as it does not explain either the coordinates of the N electrons or their derivatives. The terms  $\frac{\eta^2}{2m_e}\sum_{i=1}^N \nabla_i^2$  and  $\sum_{i=1}^N \frac{Z_e^2}{r_i}$  correspond to the kinetic and potential energy of electrons respectively. The terms together are also called one-electron term. Finally,  $\sum_{i=1}^N \sum_{j=i+1}^N \frac{e^2}{r_{ij}}$  symbolizes the electron-electron repulsion and depends on the coordinates of both electrons; hence it is called two electron terms.

Even though theoretically this approach is best, it is not possible to perform such calculations for the entire system at every time step. For larger systems, comparatively simple and computationally efficient empirical potential functions are used, functions which take into account factors such as bond stretching, bending and torsions, and covalent bonding and long-range interactions such as Van der Waals and Coulombic forces. Rahman [95] used the Lennard-Jones potential and carried out perhaps the first simulation using a realistic potential in order to study the properties of liquid argon. The most common form of the Lennard-Jones potential is shown in Equation 2.4.

$$V_{LJ} = 4\epsilon \left[ \left(\frac{\sigma}{r}\right)^{12} - \left(\frac{\sigma}{r}\right)^6 \right]$$
(2.4)

where  $V_{LJ}$  is the Lennard-Jones potential,  $\epsilon$  is the depth of the potential well i.e., the region surrounding a local minimum of potential energy,  $\sigma$  is the finite distance at which the inter-particle potential is zero, and *r* is the distance between the particles. In this representation both the repulsive and attractive natures of the atomic interaction are well represented and because of its computational simplicity, it is used extensively though more accurate potentials are available. Rahman and Stillinger [96] also carried out a simulation of liquid water. Since then, several empirical potential functions suitable for different materials have been developed. The Lennard-Jones potential is an example of a simple pair potential. Some other examples of simple pair potentials are the Morse potential and the Born-Meyer potential. For complex systems, such as the one involving covalent bonding, generally-used potential functions are the Stillinger-Weber potential, the Tersoff potential, the Bolding-Anderson potential, the Brenner potential, and the embedded atom model (EAM) and modified embedded atom model (MEAM) potentials.

In MD simulation, the trajectories of the interacting particles are defined by virtue of their position and velocity as they change with time. The propagation of position and velocity with respect to finite time intervals can be determined by integrating a couple of first order differential equations. There are time integration algorithms, such as the Verlet algorithm, the leapfrog algorithm and the Beeman algorithm, which can aid us in improving the computational efficiency of the MD simulation. The selection of the algorithm is based on the conservation of energy, the least expensive time frame, and the integration time step. In general, these algorithms are based on the Taylor series expansion and truncated after some finite number of terms. The Taylor series development is shown in Equation 2.5.

$$x(t) = x_0 + v_0 t + a_0 \frac{t^2}{2} + a_0 \frac{t^3}{3!} + \cdots$$
 (2.5)

Once we know position x at time t, after a certain time passes,  $\Delta t$ , we can find  $x(t + \Delta t)$ , and from there we can get  $x(t + 2\Delta t)$ . In this way we can also find the change in velocity, v(t) to  $v(t + \Delta t)$ .

To summarize, every molecular dynamics simulation should have the following components as a part of the program:

- (i) A model that can depict the interactions of the constituents of the system (atomic, molecular and interaction with surfaces). It is often assumed that the molecular interactions are pair wise, examples of which are particles with fixed partial charges. These assumptions will reduce the computational time significantly and improve the efficiency of the simulation program.
- (ii) An integrator to depict the propagation of particle position and velocity from t to  $t + \Delta t$ . The integrator is a finite difference structure which moves trajectories discretely in time. In order to ensure stability of the integrator, i.e. the time step,  $\Delta t$  has to be precisely chosen so that there is no drift in the system's energy.
- (iii) A statistical ensemble, with a large number of systems each of which, when considered all at once, represents a possible state that the real system might be in and which are required to control thermodynamic quantities like temperature,

pressure or the number of particles in the system. Normally for any MD simulation, a micro canonical ensemble (all the systems are maintained in the same energy state within the total system) is preferred because the system's Hamiltonian potential without external potentials is a conserved quantity.

These components are essential to describe an MD simulation. With this technique, it is possible to obtain exact results with numerical precision.

### 2.4.4 MD simulation of metal-water interface

With recent advancements in technology, the number of experimental and theoretical studies directed towards analyzing the structure of the metal-water interface has significantly increased. The literature contains a wealth of information and summarizes findings related to the adsorption of water on different metal surfaces, including low coverage adsorption and mono, multilayer and thin water films [78-80, 83, 87, 91 and 93].

The controversy prevailing in the adsorption surface size in Cu-H<sub>2</sub>O interaction simulations motivated Nadler *et al.* [91] to investigate monomeric H<sub>2</sub>O adsorption in the low coverage regime to determine the size of the large super cell surface needed to obtain converged adsorption parameters. They indicate that the simulation of adsorption geometry of a single molecule can be compared to the adsorption geometry of H<sub>2</sub>O in condensed form over the Cu surface. Based on the X-ray scattering experiments done by Toney *et al.* [92] to study the water density profile of an Ag(111) surface, we learn that the first solvent layer is denser compared to the bulk layer not only for charged surfaces but also when the system is uncharged. Furthermore, based on their work, Schnur *et al.* [93] concluded that the structure and density growth depend on the interaction strength between the surface metal and the adsorbed water molecule. The authors reported a theoretical study of a water-Cu (111) interface based on quantum mechanical calculations performed by means of density functional theory (DFT). First, they calculated the binding energy and structure of water monomers deposited on the metal. Then, they performed first-principle MD simulations on the Born-Oppenheimer surface at two different temperatures to obtain trajectories in order to estimate the relevant interface properties such as the electrochemical potential, the planar distribution functions g(z), the angular probability distribution functions  $p(\varphi,z)$  and  $p(\theta,z)$ , the water layer structure, and the vibrational spectra . The adsorption energy is calculated by the formula

$$E_{ads} = (E_{tot} - E_{Cu} - n. E_{H_20})/n$$
(2.6)

where  $E_{tot}$  is the energy of the total system,  $E_{Cu}$  is the energy of the slab alone,  $E_{H_2O}$  is the energy of one molecule of H<sub>2</sub>O, and *n* is the number of water molecules adsorbed. The molecules are assumed to be adsorbed on one side of the slab. The negative adsorption energies are equivalent to exothermic adsorption. A *p* (3×3) three-layer supercell model is used for the MD simulation of the water layer. Forty H<sub>2</sub>O molecules resulting in a 23.6 A° are utilized for the model. The Lorentz-Berthelot mixing rules are applied to get the Lennard-Jones parameters for the Cu-H<sub>2</sub>O interactions. These simulations are done using the molecular mechanics driver FIST that is part of the CP2K package. To get a clear understanding of the orientation of hydrogen and oxygen with the water molecule, the angle probability distribution function on P( $\Phi$ , Z) and P( $\Theta$ , Z) are analyzed and the authors conclude that the H atoms are located roughly at the same height and, therefore, the molecular plane is tilted away from the slab surface by 30°. Figure 2.10 illustrates the angle distribution function P ( $\phi$ , Z) and P ( $\theta$ , Z) for the Cu27 (H2O) 40 model



Figure 2.10: Angle distribution functions for Cu<sub>27</sub> (H<sub>2</sub>O) 40 model [45]

The authors also conclude that the adsorption energies, Eads, and the tilt angle  $\alpha$  of water monomers on the Cu (111) surface are sensitive measures and vary strongly with the coverage. Figure 2.11 shows the top view and side view of a distorted pentagonal water structure, which is observed on the upper surface. In Figure 2.11 (a), the super-cell is indicated in yellow. Only the outermost Cu layer is shown in Figure 2.11 (b).



Figure 2.11: Distorted pentagonal water structure Top (a) and side (b) views [45]

Figure 2.12 further depicts the trajectories of the O- atoms corresponding to the first peak closest to the upper slab surface. The lines might be discontinuous as the molecules move further away from the surface. The gray shades indicate the mean location of the corresponding Cu surface atoms



Figure 2.12: Trajectories of the O- atoms corresponding to the first peak closest to the upper slab surface. [45]

The analysis Schnur *et al.* [93] conducted based on  $\Theta$  and  $\Phi$  angles reveals that in the first solvent layer, the molecules are mainly arranged in an O- down orientation with their molecular plane tilted away from the metal surface, which is similar to the adsorption geometry of the H<sub>2</sub>O monomer at low coverage. The preferred site for adsorption is on top of the Cu surface atoms. In the second solvent layer the molecular planes are nearly parallel to the z-axis and one or two OH-bonds point toward the surface. Finally the authors conclude that the presence of the Cu surface alters the H-bond network, which is similar to the condition observed while increasing the temperature. They suggest that the increase in energy is mainly caused by the weakening of the Hbond network and simultaneous strengthening of the OH bonds.

The accuracy of the MD simulation technique relies mainly on how well the dynamics of the metallic phase and metal-electrolyte interactions are represented. Another important work in the simulation of the interaction between Cu (110)-H<sub>2</sub>O was done by Izvekov *et al.* [83], who claim that it was the first fully *ab initio* simulation of a metal-electrolyte interface. Ab initio MD simulations are very efficient in depicting the interfacial structure. They also provide insight into the electronic structure of each species in the system. To investigate whether in a metal-electrolyte contact, the interface structure is governed by water-water and ion-water interactions rather than direct interaction between the water and the metal, the authors have carried out an *ab initio* Car-Parrinello MD simulation of the Cu (110)-water interface. The calculations were done using a plane-wave-based pseudo potential formalism with a generalized gradient

correction to the local density approximation of the exchange-correlation energy as proposed by Becke and Lee, Yang, and Parr (BLYP) [83]. This approach has been shown to provide a reasonable description of hydrogen bonding in water. The system was represented by a tetragonal super cell having dimensions  $0.76683 \times 1.08433 \times 1.8690$ nm and containing a seven-layer slab of metal atoms with H<sub>2</sub>O molecules filled into the nonmetal region. The exposed face of the slab was the Cu (110) surface of the FCC structure, with 9 atoms per layer, so that the total number of metal atoms per cell was 63. The liquid phase was represented by 24 H<sub>2</sub>O molecules. The simulation was carried out using the code "CPMD Version 3.0f." Analysis of the wall-oxygen distribution function, together with the individual trajectories of the H<sub>2</sub>O molecules, showed that the interface had a bilayer structure. Through their simulation, Izvekov et al. [83] have showed that the H<sub>2</sub>O molecules in the first layer are chemisorbed on the top metal sites of the topmost metal layer, and that these metal sites are partially occupied. The chemisorbed  $H_2O$ molecules in the over layer are significantly inclined relative to the normal metal surface and appear to form a square ice-like network through hydrogen bonding. Also the authors have addressed the way the water interface affect the crystal and surface electronic states in the metal slab. Figure 2.13 depicts the wall-water angle distribution function g  $(\cos(\theta), \cos(x))$  for H<sub>2</sub>O molecules in the contact layer, as well as the rest of the H<sub>2</sub>O molecules. The view of the distribution function for the molecules in the contact layer along the  $\cos(\theta)$  axis, is shown as an insert in the figure.



Figure 2.13: The Wall–Water angle distribution function (a) in the contact layer, and (b) the rest of the  $H_2O$  molecules. [83]

The work done by Natalia *et al.*[98] and Gonzalez-Lafont *et al.* [99] provides significant insight into the atomic interaction between the Cu<sup>2+</sup> ion and water molecules. Their work addresses two important aspects of Cu<sup>2+</sup> $-H_2O$  interaction: The first aspect is whether non-additive terms sufficient for describing solvation of four bodies (e.g. Ion – (H<sub>2</sub>O)<sub>3</sub>) and higher order terms should also be considered. In general, ab initio calculation of a set of trimers is the preferred method for estimation of three-body

interactions i.e., the difference between the interaction energy of the trimer and the sum of the interaction energy of the dimers therein. During non-pairwise additive electrostatic calculation, the remainder of the interaction energy is largely a pairwise additive, especially in the intermediate to longer range. At shorter distances, the intrinsic nonadditivity is represented by the exchange terms in the interaction energy or through intermolecular perturbation theory.

The second aspect is that the average ion water binding energy is markedly lower than the corresponding interaction energy of the ion water dimer. This imbalance of energy has to be corrected by the introduction of a three-body or higher order correction. If it is not corrected, the attraction of the water molecules to the metallic ion will be unreasonably high. This high attractive force can be avoided if an effective potential is introduced to represent the average ion-water interaction energy in a typical solvation condition, instead of a conventional ab initio two-body potential. Gonzalez-Lafont *et al.* [99] introduced an effective two-body potential model function, to the *ab initio* interaction energies of the ion – water clusters while the non-pairwise additive polarization is considered, and a comparative study between conventional and effective ab initio potentials is done. The effects of the counter poise (CP) correction on the ab initio gas surfaces used in the fitting procedure are also examined. Next, two and (two + three body) model potentials for a simulation of Cu<sup>2+</sup>-(H<sub>2</sub>O) <sub>20</sub> clusters at T = 298 K are tested. The performance and validity of these models are also discussed. Natalia *et al.*  [98] described the  $Cu^{2+}-H_2O$  pair-additive interaction energies by a potential function similar to that done by Gonzalez-Lafont *et al.* [51]

$$\begin{aligned} V_2 &= c_1 \left[ \frac{2q_0}{r_a} + 2q_H \left( \frac{1}{r_b} + \frac{1}{r_c} \right) \right] + c_2 \left[ \frac{1}{r_a^2} - \frac{1}{2} \left( \frac{1}{r_b^2} + \frac{1}{r_c^2} \right) \right] + c_3 \left[ \frac{1}{r_a^3} - \frac{1}{2} \left( \frac{1}{r_b^3} + \frac{1}{r_c^3} \right) \right] + \frac{c_4}{r_a^4} \\ &+ c_5 \left( \frac{1}{r_b^4} + \frac{1}{r_c^4} \right) + \frac{c_6}{r_a^5} + c_7 \left( \frac{1}{r_b^5} + \frac{1}{r_c^5} \right) - \frac{c_8}{r_a^6} - c_9 \left( \frac{1}{r_b^6} + \frac{1}{r_c^6} \right) + \frac{c_{10}}{r_a^{12}} \\ &+ c_{11} \left( \frac{1}{r_b^{12}} + \frac{1}{r_c^{12}} \right) \end{aligned}$$

where, atomic charges  $q_0 = -0.684$  &  $q_H = +0.342$ 

 $r_a, r_b$  and  $r_c\,$  are the Cu - 0 and Cu - H interatomic distances.

Coefficients  $c_i$  (i =1,2...,11) were determined by fitting the two body  $V_2$  function to the ab initio-calculated interaction energies of the dimer  $Cu^{2+}-H_2O$  for 5 planar and 6 nonpolar orientations of the water molecule around the Cu(II) ion

Gonzalez-Lafont *et al.* [99] studied the energetic and structural properties of  $Cu^{2+}$ - $(H_2O)_{20}$  clusters in Monte Carlo simulations at 298K. The same properties are analyzed for a diluted aqueous solution of  $Cu^{2+}$  by MC simulation of a  $Cu^{2+}$ - $(H_2O)_{200}$  complex. All simulations are carried out with an NVT ensemble using a Metropolis algorithm and cubic boundary conditions with minimum cutoff. The authors claim that if good energetic results are required, then repulsive exchange non-additivity should also be included. In addition, they suggest that the effective pair potential proposed in this work might be

improved by increasing the number of configurations used in the fit with non-octahedral structure.

In this work, water molecules have been assumed to have a rigid structure and the interaction between O-H is not included. Hence the exact depiction of  $Cu^{2+}$  ions in water under electric potential cannot be fully visualized.

In a similar work, Lybrand and Kollman [100], developed potential functions for water-water and water-ion interactions that include terms to model nonadditive or a variety of body effects. They used the potentials to estimate solvation enthalpies and coordination numbers for several ionic species:  $Na^+$ ,  $K^+$ ,  $Mg^{++}$  and Cl<sup>-</sup> They used Coulombic or electrostatic interactions to account for the repulsions for O-O and H-H and the Morse-potential function for the O-H interaction. They indicate that many body potentials enable us to model water-solute interactions (specifically water-ion interactions) more realistically than any set of potential functions developed in the past. One important point they did not consider is inclusion of intramolecular degrees of freedom (i.e., bond stretching and angle bending) for the water molecules to correct for the inadequacies of rigid water molecule models.

Natalia *et al.*'s [98] paper discusses the effects of three-body and, in a few cases, four-body and higher-order terms for several  $Cu^{2+}-(H_2O)_n$  (n up to 8) clusters proposed as models of the first hydration shell of the  $Cu^{2+}$  ion. Their many-body results show that the three-body terms are non-negligible and lead to strong repulsive contributions to the total

pairwise additive interaction energy of the  $Cu^{2+}$ -(H<sub>2</sub>O)<sub>n</sub> system. This work is one of the pioneering systematic studies of three-body and higher-order effects on the interaction between a transition metal ion and water. The authors emphasized the use of intermolecular potentials, which explicitly include the three-body terms, at least those of the type ion-water-water. Also they report from their experimental and computational observations that the water-water-water trimer contributions to the total interaction energy of the metal clusters are small relative to those of the metal-water-water trimers, and consideration of water-water-water three-body contributions will make the simulations of the ionic solutions far slower. Moreover, the effects of non-additivity discussed in this paper are limited to the first hydration shell, and the authors claim that some other divalent metal solutions may require inclusion of metal-water-water trimers with one molecule in each of the first two hydration shells to accurately run the ionic solution simulation.

The effect of voltage on the shape and trajectory of the complex ion-water molecules has received the least attention in the literature, but holds the key to understanding any electrochemical process subjected to an electric potential. Such studies will enable us to predict the ionic mass transfer and material removal rates in electrochemical-mechanical planarization process and will help us optimize the process in a better qualitative way.

## CHAPTER III

## EXPERIMENTAL STUDY OF THE Cu-ECMP PROCESS

The prime objective of this experimental study is to identify the different electrochemical phases as stated by several authors [12, 13, 14, and 15] and observe the influence of those electrochemical phases on the ECMP process. Through this study, we will be able to identify the best region to work with in ECMP so that it will be possible to optimize the KPOV, which is outside the scope of this work. This study focuses on understanding the ECMP process by exploring the various electrochemical phases.

# **3.1 Machine Construction**

A face-up ECMP apparatus (as shown in Figure 3.1) was designed and fabricated for polishing  $\Phi 100$  mm blank Cu wafers. Two mini vertical milling machines (Model 44991 from Central Machineries) were arranged such that one milling machine housed the wafer carrier and the other provided rotary motion for the wafer carrier by means of a drive pulley and a V-belt. The axes of the wafer and pad holder tables were offset so that the complete wafers are polished as they rotate relative to each other. An aluminum bar was used as a cathode, considering its electro negativity relative to Cu. The rotational speeds of the polishing pad and wafer carrier were controlled using the knobs in the first and second milling machines, respectively. The working table was moved with the cross feed handle wheel and the longitudinal feed handle wheel to adjust the position of the wafer carrier with respect to the polishing pad. The pad was offset so that the complete wafer is polished as wafer and pad holders rotate relative to each other. A Shenzhen Mastech power supply, having a range of 0.1-12 V, was used to power the wafer (anode) and the cathode. A wooden base was used to insulate the first mini milling machine from the positively charged anode in the wafer carrier assembly. The apparatus was placed in a fume hood to exhaust chemical fumes formed during the process.



Figure 3.1: Cu ECMP Apparatus

## **3.1.1** Modification in the machine design

A few modifications were made to the existing ECMP setup in order to accommodate the data acquisition system and ensure that the sensor data is free of errors.

(i) The wafer carrier provided a 4 inch diameter and 0.25 inch depth in order to house the copper wafer. The retaining ring had not been properly designed, initially, so the wafer exhibited some movement during the process. This movement was caused by the simultaneous rotation of the wafer carrier and polishing pad. This movement hindered the current supply to the wafer and thereby affected the electrochemical reaction at the surface of the Cu wafer. To stop this wafer movement during the process and to ensure continuous flow of current, a retaining ring was designed, machined out of aluminum, and placed around the anode so that the entire surface of the Cu wafer was in full contact with the positive potential, i.e. the anode. The wafer carrier after being fitted with the conditioning ring is shown in Figure 3.2.



Figure 3.2: Wafer carrier with anodic conditioning ring

(ii) To stop the vibrations of the wafer relative to the polishing pad, the pad carrier was secured to a 1-inch-diameter collet and then attached to the spindle of the milling machine housing the wafer carrier. This arrangement ensured a constant and uniform contact between the polishing pad and Cu wafer, so that the material removal would be even and no defects produced with respect to the polishing pad over the surface of the Cu wafer.

### **3.1.2 Installation of the sensor system**

The voltage and current signals from the setup were collected by means of an NI DAQmx data acquisition driver. The basic circuitry used for signal acquisition from the ECMP process is shown in Figure 3.3. Three nodes  $-a_0$ ,  $a_1$  and  $a_2$ , of the DAQ driver were assigned to the process current (I), the output voltage from process (V<sub>1</sub>), and the input anodic potential (V), respectively. The DAQ was linked to LabVIEW software so that the analog outputs could be digitized and continuous signals could be recorded from the process. These signals were then discretized with a sampling rate of 100Hz. A DC power supply was used to control the anodic potential (V). As shown in Figure 3.3, a 10 $\Omega$  resistor was connected across the system to measure the process current (I).



Figure 3.3: Sensor integrated ECMP Circuitry

### 3.2 Wafer preparation and cleaning

Copper wafers 4 inches (100 mm) in diameter and 0.05 inch thick were used in this study. Copper has the tendency to react with the oxygen present in the surrounding air and the surface becomes oxidized, which prevents the underlying copper from being oxidized. For our study we needed pure copper, so the surfaces were initially cleaned with distilled water and acetone to remove the oxidation products. Some of the wafers undergo deep oxidation in air and bluish/green marks are left on the surface. These marks were removed by fine abrasive paper of grit size 15.3  $\mu$ m. The abrasive paper was rubbed smoothly, with very light pressure, so as not to scratch the surface. Generally, processes like CMP and ECMP are final finishing processes and hence the material removal rate is

generally lower compared to grinding and lapping processes. The cleaned wafers were then subjected to a lapping procedure to reduce the initial wafer surface roughness to  $\approx$ 20-30 nm. The wafers were lapped in a bench-mounted LAPMASTER<sup>®</sup> 12 lapping machine (see Figure 3.4). The wafers were lapped using 1200 grit SiC sheets and then Imperial cloth polishing sheets loaded with fine grit alumina particles. Finally, the wafers were cleaned with distilled water before they were used in the experiments.



Figure 3.4: LAPMASTER<sup>®</sup> 12 lapping machine

After conducting the ECMP experiments, the copper wafers were cleaned immediately using de-ionized water to remove the slurry debris from the surface. Then the wafers were subjected to ultra-sonic cleaning to ensure that they were completely devoid of oxidation layer residues, which generally are not removed by acetone cleaning. Figure 3.5 shows the ultra-sonic cleaner used for wafer cleaning.



Figure 3.5: Ultra-sonic cleaner

## 3.3 Measurement of MRR and surface characterization

The weight of the Cu blanket wafers was measured in micrograms ( $\mu$ g) using a Sartorius digital weighing machine (see Figure 3.6) and then converted to grams (g). Then the material removal rate (MRR) was calculated using the following formula:

MRR

= Initial Wt of the Cu wafer (before polishing) – Final Wt of the Cu wafer (after polishing) Polishing Time (in min)

The MRR was then converted to a appropriate unit i.e., nm/min

The surface roughness  $(R_a)$  of the copper wafers was obtained using MicroXAM, an optical laser interference microscope from ADE Phase Shift Technologies.  $R_a$  was measured in nm.



Figure 3.6: Sartorius Digital Weighing Machine



Figure 3.6: MICROXAM, Optical laser interference microscope (From ADE phase shift technologies)
### **3.4 Experimental results and discussions**

This experimental study focuses on understanding the process mechanism, under different experimental conditions. A phosphoric-acid-based slurry with a composition of 6% (vol.) of phosphoric acid, 2% (vol.) of ethylene glycol, 0.3% (wt.) of benzo tri-azole (BTA), 4% (wt.) of potassium hydroxide (KOH), 0.45 % (vol.) of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>), 0.15% (wt.) 0.05 and  $\mu$ m de-agglomerated gamma alumina particles and Buehler Chemo-met<sup>®</sup> polishing pads were used for all experiments. All experiments had an iteration time of 10 minutes. The pad speed was maintained at 325 RPM and wafer carrier speed was maintained at 22 RPM.

Figure 3.8 depicts the Pourbaix diagram of copper ECMP, which has 4 regions based on the anodic dissolution of Cu and electrochemical reactions happening at the Cu surface. In order to optimize the MRR and R<sub>a</sub>, many experiments were carried out based on the different electrochemical regions of the Pourbaix diagram. The phosphoric-acid based-slurry was used for all experiments, as Munnangi [34] found that it provided a good surface roughness and MRR. Hard polyurethane Buehler Chemo-met polishing pads were used for all experiments, as they were found to have the desired mechanical and chemical stability.



Figure 3.8: Pourbaix diagram of Copper-water system, after Ng et al.,

From the Pourbaix diagram, we can observe that at acidic pH (<6) and voltages above 0.5 V, copper undergoes severe corrosion because the anodic dissolution is very high and no sacrificial surface layer is allowed to form on the copper. As the pH is increased i.e., greater than 6, the solubility of the metal interface limits the mass transfer of the metal ions into the solution. A passivation layer is formed when the concentrated ions accumulate at the interface. This layer, which now acts as a sacrificial layer, is removed by pad during ECMP.

# **3.4.1 Influence of pH in the Process**

For this study, the pH of the slurry was adjusted by varying the proportion of potassium hydroxide in the slurry. Figure 3.9 and Figure 3.10, depict the variations in MRR at different voltages and 2 different pH conditions, one acidic (pH = 2.5) and the other near neutral (pH = 6.5).



Figure 3.9: Variation of MRR with pH at different anodic voltages

рН	Observation	<b>Reason for the observation</b>
2.5	<ul> <li>MRR is normally high.</li> <li>MRR gradually increases with voltage.</li> </ul>	<ul> <li>Anodic dissolution rate is very high due to rapid oxidation of the surface.</li> <li>The dissolution rate is further accelerated with increase in voltage as there is no formation of passivation layer. Cu<sup>2+</sup> ions are stripped out of the</li> </ul>

Table 3.1: Scientific explanation for influence of pH on MF
---

		surface in a rate, that doesn't provide sufficient time for the formation of passivation layer.
6.5	<ul> <li>MRR is comparatively lower than that observed at pH = 2.</li> <li>MRR increases with voltage.</li> </ul>	<ul> <li>The solubility of the metal interface limits the mass transfer of the metal ions into the solution.</li> <li>A passivation layer is formed when the concentrated ions accumulate at the interface, which is then removed by the action of polishing pad.</li> </ul>



Figure 3.10: Variation of R<sub>a</sub> with pH at different anodic voltages

рН	Observation	<b>Reason for the observation</b>
2.5	<ul> <li>Severe pitting is observed over the surface.</li> <li>R<sub>a</sub> is very high and increases with voltage.</li> </ul>	• The electrochemical dissolution rate of the Cu anode is high, which leads to stripping of Cu <sup>2+</sup> ions from the surface.

	• Surface roughness values are		
6.5	low.	• Anodic dissolution is limited by the	
	• Occurrence of pitting is	formation of passive layer.	
	reduced and is completely	• Passivation layer formation protects the	
	absent at voltages above 4V.	surface from severe corrosion and gets	
	• R <sub>a</sub> decreases with increase in	planarized by the polishing pad.	
	voltage.		

#### **3.4.2 Effect of anodic voltage**

For this study, the pH was maintained around 6.5 (near neutrality) and the other conditions were maintained as described earlier. Figure 3.11 shows the effect of applied voltage on MRR and surface quality, from a process standpoint.

As shown in Figure 3.11 (a), the MRR tends to be low in the low potential region (1-2 V). This region corresponds to the active region, which is generally characterized by low MRR. In the active region, the dissolution rate of Cu<sup>2+</sup> ions from the surface is low, as the applied voltage is low; in other words, the oxidation rate is less. With very few Cu<sup>2+</sup> ions available in the system to react with slurry components, the rate of formation of the passivation layer tends to be low. The MRR in this region is primarily dependent on pad pressure (which for ECMP is low) rather than electrochemical dissolution. Hence the MRR tends to be low in the active region. With further increase in anodic voltage (3–4V), the MRR is observed to increase slightly because the process has entered the passivation region. In this region, the dissolution rate of Cu<sup>2+</sup> increases, which promotes the growth of the passivation layer. With more oxidation over the copper wafer

surface, the MRR tends to improve at this region even with a lower pad pressure. Figure 3.11 (a) summarizes the effect of anodic voltage on material removal rate (MRR).



Figure 3.11 (a): Variation of MRR (nm/min) with anodic voltage (V)

At 4-5V the process enters the transient region (as shown in Figure 3.11 (a)) where the MRR increases to a larger extent, indicating that the electrolyte concentration is acting against the applied voltage leading to stripping of the Cu<sup>2+</sup> ions from the surface of the wafer. Several authors have claimed that the transient region is the best region to work with if high MRR is desired [add the references]. Increasing the voltage above 5V positions the process in the trans-passive electrochemical region, where the film

formation and breakdown happen as part of dynamic equilibrium. Hence the MRR tends to increase linearly with the increase in the anodic potential.



Figure 3.11 (b) summarizes the effect of anodic voltage on surface roughness (R<sub>a</sub>).

Figure 3.11 (b): Variation of Surface roughness (R<sub>a</sub>) with anodic voltage (V)

As shown in Figure 3.11 (b), the  $R_a$  value tends to be high and does not change much in the active region (1–2 V). The surface tends to be rough ( $R_a \sim 25$ nm) due to severe pitting and corrosion accompanying ohmic oxidation. The removal of asperities is uneven in this region, and as the planarization happens only through the action of the polishing pad, because of the low dissolution rate, the surface quality is deteriorated. With an increase in voltage (2-3 V), the process enters the passivation region, where the R<sub>a</sub> value is comparatively low. The passivation promotes higher material removal from asperities compared to pits. Consequently the surface quality improves gradually in this passivation phase (see Figure 3.11 (b)). For applied voltage in the 3–4 V range, the passivation film becomes softer and less stable, leading to a transient phase. Here, the surface quality is much improved (see Figure 3.11(b)). For voltages above 4 V, the ECMP process follows a trans-passive electrochemical transfer, where the film formation and breakdown happen as part of dynamic equilibrium. Here surface quality deteriorates with increasing voltage because of frequent breakdowns of the film and removal of surface asperities.

#### **3.4.3 Surface texture analysis**

Figure 3.12 shows the images of copper wafer surface captured at different electrochemical regions through laser interferometer, MicroXAM. Through surface analysis, we can physically observe the effects of active, passive, transient and transpassive regions on the Cu wafer surface. In the active region, we can observe that the surface tends to be rough ( $R_a \sim 25$ nm) due to severe pitting and corrosion. Here the dissolution rate is poor and the passivation region is not being effectively formed and removed. We can also observe smeared marks on the surface, which is the result of polishing pad smearing the poorly formed passivation region during the process. In the passive region, we can clearly see the existence of corrosion marks and parts of the

passivation layer. In the transient and trans-passive regions, we can see a clear improvement in the surface. The surface roughness values are also very low ( $R_a \sim 13$ -15nm), with most of the scratch marks, corrosion marks and pits removed. Some deep scratches still remain on the surface as ECMP process is a nano manufacturing process; it is incapable of removing deep scratch marks produced during the initial stages of slicing from the ingot.



Figure 3.12: Surface Texture at (a)  $1 \vee (R_a = 25.209 \text{ nm})$ , (b)  $3 \vee (R_a = 16.511 \text{ nm})$ , (c)  $4 \vee (R_a = 15.822 \text{ nm})$ , and (d)  $5 \vee (R_a = 13.709 \text{ nm})$ 



Figure 3.13: Initial condition of Copper wafer (Left) After Cu-ECMP (Right)

Figure 3.13 shows the actual and micrograph images of random regions of the wafer. The image to the left shows the initial surface of the wafer with lot of scratches. The  $R_a$  was brought down from 60-70 nm to 25-30 nm using Lapping process. The image to the right shows the surface of the copper wafer after ECMP process with a run time of 10 min and at an anodic potential of 4.5 V

#### 3.4.4 Characterization of electrochemical regions from I-V curves

The observed variation of MRR, and the surface morphology and roughness with the applied voltage may be traced back to the electrochemical phase prevailing during the specified process conditions. The trends of the I-V characteristics (Figure 3.14) can be used to identify the process conditions, and hence track MRR, Ra and other quality and performance variables of ECMP process in real-time. In contrast with prevailing approaches to delineate I-V characteristics, we not only capture the first moment, averaged trend, but also capture time-frequency patterns of the current signals at various applied voltage conditions. As shown in Figure 4.4, the first moment trend indicates that for low applied voltage of < 1 V; the average current of < 0.005 A was noted, corresponding to active dissolution. The average current exhibits a slight increase as the voltage is increased to 2 V. Upon further increase in anodic voltage to 3 V, the current shows a marked increase (average current of 0.08 A at 3 V), and increases linearly until 5 V. Upon further increasing trend becomes somewhat sharper, suggesting a transition from a transient to trans-passive electrochemical phase.



Figure 3.14: I-V curves of H<sub>3</sub>PO<sub>4</sub> electrolyte under different voltage conditions

#### CHAPTER IV

#### MOLECULAR DYNAMICS STUDY ON Cu-ECMP PROCESS

#### 4.1 Description of the molecular system

Here we consider the effects of the applied voltage on the shape and the trajectory of an isolated copper (II) hexa hydrate  $Cu[(H_2O)_6]^{2+}$  complex molecular system. The potentials were adapted based on the previous work done by Natalia *et al.*, [98]. In this work, the intermolecular and intramolecular potentials depicting the water-water interaction and O-H interaction have been considered. Also harmonic potential function to depict the stretching and bending of the water molecules have also been considered. The molecule is simulated at room temperature condition (298K). From the mobility and velocity of  $Cu^{2+}$ , the current carried by the  $Cu^{2+}$  can be computed. With further extension of this model i.e., by determining the current carried by H<sup>+</sup> and PO<sub>4</sub><sup>3-</sup> we can calculate the total current of the system. If MRR can be expressed as a function of total current, then the MRR can be predicted at different voltage conditions. Furthermore from I-V characteristic curve the transition of the electrochemical phase can be elucidated.

#### **4.2 Design of pairwise interatomic potentials**

Pair potentials are appropriate for metals such as copper, aluminum and silver for atomistic simulation. Figure 4.1 shows the structure of copper (II) hydrate molecule used for determining the global minima position.



Figure 4.1: Structure of Copper (II) Hydrate molecule

The potential between  $Cu^{2+}$ -H<sub>2</sub>O (V<sub>2</sub>) and potential within H<sub>2</sub>O molecule (V<sub>H<sub>2</sub>O</sub>) were modeled separately and then combined to obtain the total potential energy of the system.

The pairwise additive interaction energies between  $Cu^{2+}$ -H<sub>2</sub>O are mathematically modeled using a multi-polar expansion, as shown in Equation 4.1. This expansion accounts for both repulsive and attractive interactions between  $Cu^{2+}$  and H<sub>2</sub>O [98, 99].

$$V_{2} = c_{1} \left[ \frac{2q_{0}}{r_{1}} + 2q_{H} \left( \frac{1}{r_{2}} + \frac{1}{r_{3}} \right) \right] + c_{2} \left[ \frac{1}{r_{1}^{2}} - \frac{1}{2} \left( \frac{1}{r_{2}^{2}} + \frac{1}{r_{3}^{2}} \right) \right] + c_{3} \left[ \frac{1}{r_{1}^{3}} - \frac{1}{2} \left( \frac{1}{r_{2}^{3}} + \frac{1}{r_{3}^{3}} \right) \right] + \frac{c_{4}}{r_{1}^{4}} + c_{5} \left( \frac{1}{r_{2}^{4}} + \frac{1}{r_{3}^{4}} \right) + \frac{c_{6}}{r_{1}^{5}} + c_{7} \left( \frac{1}{r_{2}^{5}} + \frac{1}{r_{3}^{5}} \right) - \frac{c_{8}}{r_{1}^{6}} - c_{9} \left( \frac{1}{r_{2}^{6}} + \frac{1}{r_{3}^{6}} \right) + \frac{c_{10}}{r_{1}^{12}} + c_{11} \left( \frac{1}{r_{2}^{12}} + \frac{1}{r_{3}^{12}} \right)$$
(4.1)

where,  $q_0$  is the atomic charge on oxygen = -0.6836

 $q_H$  is the atomic charge on hydrogen = + 0.3415  $r_1$ ,  $r_2$  and  $r_3$  are the  $Cu^{2+}$ –O,  $Cu^{2+}$ –H<sub>1</sub> and  $Cu^{2+}$ –H<sub>2</sub> interatomic distances Coefficients  $c_i$  (i = 1, 2, ..., 11) were determined by fitting the two-body V<sub>2</sub> function to *ab initio* calculated interaction energies of the dimer  $Cu^{2+}-H_2O$ for 5 planar and 6 nonpolar orientations of the water molecule around Cu(II) ion [98, 99].

We can observe that from Figure 4.2 that the global minima position is at 3.8635 bohrs<sup>1</sup> where the energy is -0.7396 hartree<sup>2</sup>.

 $<sup>\</sup>frac{1}{1}$  bohr radius = 5.29177 × 10<sup>-11</sup>m  $\frac{1}{1}$  hartree = 2625.5 kJ/mol



**Figure 4.2:** Potential Energy ( $V_2$ ) vs Cu<sup>2+</sup>-O ( $r_1$ ) distance.

Morse potential was used to model the pair potential between  $O-H_1(r_4)$  and  $O-H_2(r_5)$ . A simple harmonic function of bond angle  $h(\theta)$  was used to depict the bending of the  $H_1-O-H_2$  angle ( $\theta = 104.5^{\circ}$ ).

The mathematical expressions used for the Morse function and simple harmonic function of bond angle are given in Equations 4.2, 4.3 and 4.4.

$$V_{H_20} = f(r_4) + f(r_5) + h(\theta)$$

where,

$$f(r_4) = D[1 - \exp(-\alpha (r_4 - r_{eq}))]^2$$
(4.2)

$$f(r_5) = D[1 - \exp(-\alpha (r_5 - r_{eq}))]^2$$
(4.3)

$$h(\theta) = k[\theta - \theta_{eq}]^2 \tag{4.4}$$

with,

D = Well depth (defined relative to the dissociated atoms) = 463 kJ/mol

$$\alpha = \left[\frac{2\pi^2 \mu \gamma_0^2}{D}\right]^{\frac{1}{2}} = 1.145272 \text{ bohr}^{-1}$$

 $\widetilde{\gamma_b} = \text{Bending Frequency for } H_2 0 = 1595 \ \text{cm}^{-1}$ 

$$\therefore \gamma_{b} = \widetilde{\gamma_{b}}C = (1595 \text{ cm}^{-1})(2.997 \times 10^{10} \text{ cm s}^{-1})$$
$$= 4.780215 \times 10^{13} \text{ s}^{-1} \times \left(\frac{2.41888 \times 10^{-17} \text{ s}}{\text{Time Unit}}\right)$$
$$= 0.0015628 \text{ time unit}^{-1}$$

$$\mu = \text{Reduced mass of H}_2 O = \frac{m_H m_O}{(m_H + m_O)} = 1.57 \times 10^{-27} \text{kg}$$

 $r_{eq}$  = Equilibrium distance at stable molecular condition = 1.8085 bohr

$$k = 4\pi^2 \gamma_b{}^2 \mu r_{eq}^2 = 0.2975315$$
 hartree

 $\theta_{eq}$  = Equilibrium bond angle at stable molecular condition =  $\,1.823$  radians

Figure 4.3 and Figure 4.4 depict the normal shape of the Morse potential and simple harmonic function of bond angle, respectively.



Figure 4.4: Simple harmonic function curve

Morse potential and bond angle potential were then modeled using MATLAB and the following curves were obtained with variation of length of  $O-H_1$  ( $r_4$ ) and theta ( $\theta$ ).



Figure 4.5: Morse potential curve for  $[Cu(H_2O)]^{2+}$  system in equilibrium position



Figure 4.6: Bond angle potential curve for  $[Cu(H_2O)]^{2+}$  system in equilibrium position

At this point, the total potential energy of the system ( $V_{TOT}$ ) is computed from the summation of  $V_2$  and  $V_{H_2O}$  potential over 'n' water molecules.

$$V_{\text{TOT}} = \sum_{i=1}^{N} (V_2) + \sum_{i=1}^{2N} (V_{\text{H}_2 \text{O}})$$
(4.5)

#### 4.3 Atomic force computation

The force on acting on an atom is computed as the negative first order derivative of the potential function V i.e., the negative slope of the potential energy curve.

The partial differential equations of  $V_2$  with respect to  $r_1$ ,  $r_2$  and  $r_3$  are computed as shown below,

$$\frac{\partial V_2}{\partial r_1} = -\frac{2C_1q_0}{r_1^2} - \frac{C_2}{r_1^3} - \frac{C_3}{r_1^4} - \frac{C_4}{r_1^5} - \frac{C_6}{r_1^6} + \frac{C_8}{r_1^7} - \frac{C_{10}}{r_1^{13}}$$
$$\frac{\partial V_2}{\partial r_2} = -\frac{2C_1q_H}{r_2^2} + \frac{0.5C_2}{r_2^3} + \frac{1.5C_3}{2r_2^4} - \frac{4C_5}{r_2^5} - \frac{5C_7}{r_2^6} + \frac{6C_9}{r_2^7} - \frac{12C_{11}}{r_2^{13}}$$
$$\frac{\partial V_2}{\partial r_3} = -\frac{2C_1q_H}{r_3^2} + \frac{0.5C_2}{r_3^3} + \frac{1.5C_3}{2r_3^4} - \frac{4C_5}{r_3^5} - \frac{5C_7}{r_3^6} + \frac{6C_9}{r_3^7} - \frac{12C_{11}}{r_3^{13}}$$

The partial differential equations of  $V_{\rm H_2O}$  with respect to  $r_4$  ,  $r_5$  and  $r_6$  are computed as shown below,

$$\frac{\partial V_{H_2O}}{\partial r_4} = \frac{\partial f(r_4)}{\partial r_4} + \left[ \left( \frac{\partial h(\theta)}{\partial \theta} \right) \cdot \left( \frac{\partial \theta}{\partial r_4} \right) \right]$$
$$\frac{\partial V_{H_2O}}{\partial r_5} = \frac{\partial f(r_5)}{\partial r_5} + \left[ \left( \frac{\partial h(\theta)}{\partial \theta} \right) \cdot \left( \frac{\partial \theta}{\partial r_5} \right) \right]$$

$$\frac{\partial V_{H_2O}}{\partial r_6} = \left[ \left( \frac{\partial h(\theta)}{\partial \theta} \right) \cdot \left( \frac{\partial \theta}{\partial r_6} \right) \right]$$

The distance between the Cartesian coordinates of Particle 1  $[X_1 Y_1 Z_1]$  and Particle  $\alpha [X_{\alpha} Y_{\alpha} Z_{\alpha}]$  i.e.,  $r_{1\alpha}$  is calculated as shown in Eq. 4.6.

$$r_{1\alpha}^{2} = [(X_{1} - X_{\alpha})^{2} + (Y_{1} - Y_{\alpha})^{2} + (Z_{1} - Z_{\alpha})^{2}]$$
(4.6)

Differentiating (Eq. 4.7) partially with respect to  $r_{1\alpha}$  we get,

$$2 r_{1\alpha} \left( \frac{\partial r_{1\alpha}}{\partial X_1} \right) = 2(X_1 - X_\alpha)$$

$$\left( \frac{\partial r_{1\alpha}}{\partial X_1} \right) = \frac{(X_1 - X_\alpha)}{r_{1\alpha}}$$

$$(4.7)$$

$$\therefore \left( \frac{\partial r_{1\alpha}}{\partial X_1} \right) = \left( \frac{X_1 - X_\alpha}{[(X_1 - X_\alpha)^2 + (Y_1 - Y_\alpha)^2 + (Z_1 - Z_\alpha)^2]^{\frac{1}{2}}} \right)$$

Likewise, we can find other partial derivatives of  $r_1$ ,  $r_2$ ,  $r_3$  with respect to Y and Z directions.

The force along X, Y and Z axis was then determined using 12 chain rule

derivatives. They are  $\begin{cases} \begin{pmatrix} \frac{\partial V_2}{\partial X_i} \\ \begin{pmatrix} \frac{\partial V_2}{\partial Y_i} \\ \begin{pmatrix} \frac{\partial V_2}{\partial Z_i} \end{pmatrix} \end{cases}$  where  $1 \le i \le 4$ , so totally there will be 4 derivatives for

the force calculation, along every axis.

Force along X axis i.e., 
$$F_{X} = \frac{\partial V}{\partial X_{i}}$$
  

$$\frac{\partial V}{\partial X_{1}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial X_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial X_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial X_{1}}\right)$$

$$\frac{\partial V}{\partial X_{2}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial X_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial X_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial X_{2}}\right)$$

$$\frac{\partial V}{\partial X_{3}} = \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial X_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial X_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial X_{3}}\right)$$

$$\frac{\partial V}{\partial X_{4}} = \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial X_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial X_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial X_{4}}\right)$$

Force along Y axis i.e., 
$$F_{Y} = \frac{\partial V}{\partial Y_{1}}$$
  

$$\frac{\partial V}{\partial Y_{1}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial Y_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial Y_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial Y_{1}}\right)$$

$$\frac{\partial V}{\partial Y_{2}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial Y_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial Y_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial Y_{2}}\right)$$

$$\frac{\partial V}{\partial Y_{3}} = \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial Y_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial Y_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial Y_{3}}\right)$$

$$\frac{\partial V}{\partial Y_{4}} = \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial Y_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial Y_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial Y_{4}}\right)$$

Force along Z axis i.e., 
$$F_{Z} = \frac{\partial V}{\partial Y_{i}}$$
  

$$\frac{\partial V}{\partial Z_{1}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial Z_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial Z_{1}}\right) + \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial Z_{1}}\right)$$

$$\frac{\partial V}{\partial Z_{2}} = \left(\frac{\partial V_{2}}{\partial r_{1}} \cdot \frac{\partial r_{1}}{\partial Z_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial Z_{2}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial Z_{2}}\right)$$

$$\frac{\partial V}{\partial Z_{3}} = \left(\frac{\partial V_{2}}{\partial r_{2}} \cdot \frac{\partial r_{2}}{\partial Z_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{4}} \cdot \frac{\partial r_{4}}{\partial Z_{3}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial Z_{3}}\right)$$

$$\frac{\partial V}{\partial Z_{4}} = \left(\frac{\partial V_{2}}{\partial r_{3}} \cdot \frac{\partial r_{3}}{\partial Z_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{5}} \cdot \frac{\partial r_{5}}{\partial Z_{4}}\right) + \left(\frac{\partial V_{H_{2}O}}{\partial r_{6}} \cdot \frac{\partial r_{6}}{\partial Z_{4}}\right)$$

For simplicity, let us represent the first order derivatives of the potential energy with respect to all atoms in the X, Y and Z direction as Eq. 4.8

$$\begin{pmatrix}
\frac{\partial V}{\partial X_{i}} \\
\frac{\partial V}{\partial Y_{i}} \\
\frac{\partial V}{\partial Z_{i}}
\end{pmatrix}$$
(4.8)

In order to verify the correctness and accuracy of the values of derivatives, a 7-point slope method was adopted as shown in Figure 4.7.



Figure 4.7: Graphical depiction of 7-point slope method

Figure 4.7 depicts a graphical representation of a 7-point slope method. The X axis represents the change in the X coordinates of the atoms. The Y axis contains the corresponding Potential energy values calculated using the  $V_{TOT}$  function (Eq. 4.5) based upon different positions of the atoms. The method is based on the computation of S<sub>1</sub>, S<sub>2</sub> and S<sub>3</sub> can be found from Equations 4.9, 4.10 and 4.11, where h = 0.0005.

$$S_1 = \frac{(Y_1 - Y_{-1})}{h} \tag{4.9}$$

$$S_2 = \frac{(Y_2 - Y_{-2})}{h} \tag{4.10}$$

$$S_3 = \frac{(Y_3 - Y_{-3})}{h} \tag{4.11}$$

If the red curve shown in Figure 4.7 can be accurately fitted with a polynomial expansion in X of order 6 or less, then it can be rigorously shown that

$$\frac{\partial Y}{\partial x}\Big|_{X=X_0} = 0.75S_1 - 0.15S_2 + 0.01666S_3 \tag{4.12}$$

The results obtained from the comparison between the values computed from the derivatives shown in Eq. 8 and the values of the slope computed using 7-point slope method is presented in Table 1. Two water molecules have been considered for this test. So there are 7 atoms in this test model, which are designated as shown below,

 $1 - Cu^{2+}$  ion, 2 - O of the first water molecule,  $3 - H_1$  of the first water molecule,

 $4 - H_2$  of the first water molecule, 5 - O of the second water molecule,  $6 - H_1$  of the second water molecule,  $7 - H_2$  of the second water molecule.

Derivative of V with	Values computed from	Values computed from 7-point	
respect to atoms	derivatives – Eq. 8	slope method – Eq.12	
$\partial V / \partial X_1$	3.74647773971345E-10	3.74648789071506E-10	
$\partial V / \partial Y_1$	-4.55040393490804E-06	-4.55040393228684E-06	
$\partial V / \partial Z_1$	8.75514604295787E-09	8.75514604244237E-06	
$\partial V / \partial X_2$	2.26679800811792E-06	2.26679800728777E-06	
$\partial V / \partial Y_2$	4.46613209182386E-06	4.46613208940395E-06	
$\partial V / \partial Z_2$	-2.15011102692851E-08	-2.22631485868115E-06	
$\partial V / \partial X_3$	-2.21996213183842E-06	-2.21996213073261E-06	
$\partial V / \partial Y_3$	8.20151222379824E-09	8.20151179567898E-09	
$\partial V / \partial Z_3$	-3.19980359269939E-06	-3.19980359114016E-06	
$\partial V / \partial X_4$	6.15654086177955E-09	6.15654086482174E-06	
$\partial V / \partial Y_4$	1.59399220685803E-06	1.59399220618633E-06	
$\partial V / \partial Z_4$	1.35498613388521E-06	1.35498613379016E-06	
$\partial V / \partial X_5$	-6.52325226371036E-09	-6.57444815102026E-07*	
$\partial V / \partial Y_5$	-6.73517451912826E-07	-6.73517451390600E-07	
$\partial V / \partial Z_5$	6.41019405924549E-09	6.41019713447353E-09	
$\partial V / \partial X_6$	2.45635508271192E-06	2.45635508017717E-06	
$\partial V / \partial Y_6$	-1.97733906912618E-08	-1.97733906255392E-06	
$\partial V / \partial Z_6$	-1.21798071070880E-06	-1.21798070945804E-06	
$\partial V / \partial X_7$	-3.89824017398553E-06	-3.89824017265228E-06	
$\partial V / \partial Y_7$	8.32173805208994E-10	1.93648081453384E-06	
$\partial V / \partial Z_7$	1.94442893592877E-06	1.94442893514135E-06	

**Table 4.1:** Comparison table between  $Q_i$  values and 7-point slope method. All values aregiven in units of hartree bohr<sup>-1</sup>.

\*7 point slope method is based on the difference between the potential, so cannot be always accurate From Table 1 we can observe that the accuracy is of the order of 9 significant digits for derivatives whose magnitude is  $\geq 10^{-6}$  hartree bohr<sup>-1</sup>. For derivatives whose magnitude is  $\leq 10^{-9}$  hartree bohr<sup>-1</sup>, Eq.12 is not sufficiently accurate to evaluate  $\partial V / \partial X$ .

### 4.4 Initial spherical polar coordinates assignment

Initially, the coordinates of the atoms were assigned arbitrarily to validate whether the model complies with the Law of conservation of Energy.

The initial position of  $Cu^{2+}$  ion was taken to be origin [0 0 0]. Next, the Cu–O distance for each of the N oxygen atoms was taken to be the equilibrium  $Cu^{2+}$ –O distance obtained from Figure 4.1. The rotation angle,  $\theta$  and  $\emptyset$ , for vector  $\vec{R}$  were assigned randomly. Figure 4.8 illustrates the process.



Figure 4.8: Spherical polar coordinates arrangement

From Figure 4.8, the Cartesian coordinates of oxygen are

$$X_{O} = R_{CuO} \cdot \sin\theta \cdot \cos\emptyset$$
$$Y_{O} = R_{CuO} \cdot \sin\theta \cdot \sin\emptyset$$
$$Z_{O} = R \cdot \cos\theta$$

We can then cover all possible angles by using,

$$0 \le \theta \le \pi$$
 and  
 $0 \le \emptyset \le 2\pi$ .

In order to initially position the atoms, we use a random number  $0 \le \zeta \le 1$  and select the azimuthal and rotational angle in the following manner,

$$\theta_i = \zeta_i \cdot \pi; \quad \phi_i = \zeta \cdot 2\pi \quad (i = 1, 2, 3..., N)$$



Figure 4.9: Spherical polar coordinates initialization based on angles

The hydrogen atoms attached to each oxygen are initially taken to be in the (Y, Z) plane so that we have the coordinates of the O, H<sub>1</sub> and H<sub>2</sub> as,

O coordinates = 
$$\begin{pmatrix} 0 \\ R_{cuo} \sin\theta \\ R_{cuo} \cos\theta \end{pmatrix}$$
, H<sub>1</sub> coordinates =  $\begin{pmatrix} 0 \\ Y_3 \\ Z_3 \end{pmatrix}$  and H<sub>2</sub> coordinates =  $\begin{pmatrix} 0 \\ Y_4 \\ Z_4 \end{pmatrix}$ 

From Figure 4.9, using the law of cosines, we have,

$$(R_{CuH})^{2} = (R_{CuO})^{2} + (R_{OH})^{2} - 2 \cdot R_{CuO} \cdot R_{OH} \cdot Cos\beta$$
$$(R_{OH})^{2} = (R_{CuO})^{2} + (R_{CuH})^{2} - 2 \cdot R_{CuO} \cdot R_{CuH} \cdot Cos\gamma$$
$$Cos\gamma = \left[\frac{(R_{CuH})^{2} + (R_{CuO})^{2} - (R_{OH})^{2}}{2 \cdot R_{CuO} \cdot R_{CuH}}\right]$$
$$\therefore \gamma = \cos^{-1}\left[\frac{(R_{CuH})^{2} + (R_{CuO})^{2} - (R_{OH})^{2}}{2 \cdot R_{CuO} \cdot R_{CuH}}\right]$$

From Figure 4.8, we can also determine the values of  $\theta_{H_{N1}}$  as  $\theta - \gamma$  and  $\theta_{H_{N2}}$  as  $\theta + \gamma$ , which gives

Coordinates of 
$$H_1 = \begin{pmatrix} 0 \\ R_{CuH} \cdot \sin\theta_{H_{N_1}} \\ R_{CuH} \cdot \cos\theta_{H_{N_1}} \end{pmatrix}$$
 and coordinates of  $H_2 = \begin{pmatrix} 0 \\ R_{CuH} \cdot \sin\theta_{H_{N_2}} \\ R_{CuH} \cdot \cos\theta_{H_{N_2}} \end{pmatrix}$ 

After assigning the initial coordinates of O,  $H_1$  and  $H_2$ , we need to rotate the  $H_2O$  molecule about the Z axis through a rotational angle  $\emptyset'$  selected when O atoms were positioned.

In this case,  $\emptyset' = \left(\frac{\pi}{2} - \emptyset\right)$ . The coordinates of O, H<sub>1</sub> and H<sub>2</sub> are then rotated with the aid of a rotational matrix  $R_Z(\emptyset')$ , where

$$R_Z(\phi') = \begin{bmatrix} \cos \phi' & \sin \phi' & 0 \\ -\sin \phi' & \cos \phi' & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

For instance, let us examine how the rotational operation is performed for the  $O_N$  coordinates.

Figure 4.10 shows the structure of copper (II) hexa hydrate complex molecule used for the energy conservation verification test. The magenta sphere denotes the copper (II) ion, red spheres denote the oxygen atoms and green spheres denote the hydrogen atoms.



Figure 4.10: Copper (II) hexa hydrate complex molecule with randomly assigned positions based on angles and bond length

## 4.5 Initial momentum assignment and kinetic energy calculation

To test the derivatives and integration accuracy, we initially assign thermal energy at 298K to each of the atoms.

Let E be the initial energy assigned to the system, M be the mass of the atom, v is the velocity, R is the gas constant =  $8.314 J mol^{-1}K^{-1}$  and T is the system temperature = 298 K. In a thermal system, the average kinetic energy  $\langle T \rangle$ 

$$\langle \mathbf{T} \rangle = \langle \frac{1}{2} \mathbf{M} \mathbf{v}^2 \rangle = \frac{\mathbf{M}}{2} \langle \mathbf{v}^2 \rangle = \frac{3}{2} \cdot \mathbf{R} \mathbf{T}$$

The average square velocity of the system (v) can be calculated based on the positive square root of the sum of square of velocity components along X, Y and Z directions i.e.,  $v_X$ ,  $v_Y$  and  $v_Z$  as shown in Eq. 4.12 (a)

$$\langle v^{2} \rangle = \langle \frac{3RT}{M} \rangle \ m^{2} s^{-2}$$

$$\langle v^{2} \rangle^{\frac{1}{2}} = \left[ \frac{3RT}{M} \right]^{\frac{1}{2}} m/s$$

$$\langle v^{2} \rangle^{\frac{1}{2}} = \langle v_{X}^{2} + v_{Y}^{2} + v_{Z}^{2} \rangle^{\frac{1}{2}} = \langle 3v_{X}^{2} \rangle^{\frac{1}{2}} = \sqrt{3} \langle v_{X}^{2} \rangle^{\frac{1}{2}}$$

$$(4.21 a)$$

where, we assume an equipartitioning of kinetic energy among the three translation degrees of freedom.

$$\langle v_X \rangle^{\frac{1}{2}} = \frac{1}{\sqrt{3}} \langle v^2 \rangle^{\frac{1}{2}} = \frac{1}{\sqrt{3}} \left[ \frac{3RT}{M} \right]^{\frac{1}{2}} = \pm \left[ \frac{RT}{M} \right]^{\frac{1}{2}} m/s$$
 (4.13)

The signs of  $\left[\frac{RT}{M}\right]^{\frac{1}{2}}$  are chosen randomly.

The momentum components along X, Y and Z directions i.e.,  $P_X$ ,  $P_Y$  and  $P_Z$  are calculated based on the mass of the atom and corresponding momenta component.

$$\langle P_X \rangle = m_X \langle v_X^2 \rangle^{\frac{1}{2}}$$

The units of velocity  $(m s^{-1})$  are then converted into bohrs/time unit by the following conversion calculation.

$$\langle v_{X}^{2} \rangle^{\frac{1}{2}} \text{ m s}^{-1} \times \left[\frac{100 \text{ cm}}{\text{m}}\right] \times \left[\frac{10^{8} \text{ A}^{\circ}}{\text{ cm}}\right] \times \left[\frac{1 \text{ bohr}}{0.5429171 \text{ A}^{\circ}}\right] \times \left[\frac{1 \text{ bohr}}{0.5429171 \text{ A}^{\circ}}\right] \times \left[\frac{2.41888 \times 10^{-17}}{1 \text{ Time unit}}\right] = \langle v_{X}^{2} \rangle^{\frac{1}{2}} \text{ bohr}/\text{ time unit}$$

The mass of the atoms are converted into atomic mass unit (amu) as shown below for oxygen atom.

$$\mathbf{m}_{0} = \frac{16.0 \ g \ mol^{-1}}{6.022 \times 10^{23} \ mol^{-1}} \cdot \left[\frac{1 \ amu}{9.10939 \times 10^{-28} \ g}\right] = 1 \ amu$$

From the above calculation procedure, the magnitude of the momentum along X, Y and Z direction will be equal for individual atoms whereas the signs would be randomly designated i.e.,  $|P_X| = |P_Y| = |P_Z|$ 

The kinetic energy is then calculated by the summation of momentum values and mass of atoms over 'n' water molecules.

Kinetic Energy, 
$$\langle T \rangle = \sum_{i=1}^{n} \frac{P_i^2}{2m_i}$$
 (4.14)

where,  $P_i^2 = P_{X_i}^2 + P_{Y_i}^2 + P_{Z_i}^2$ 

 $m_i = mass of Cu^{2+} m_{Cu} (or) mass of Oxygen m_0 (or) mass of Hydrogen m_H$ 

 $\therefore$  The total energy of the system (E<sub>TOT</sub>) will be the sum of potential energy and kinetic energy.

$$E_{TOT at time(t)=0} = V_{TOT} + \langle T \rangle$$

#### 4.6 Calculation of trajectory

In order to test the integration accuracy of the model, we need to numerically integrate the coupled first-order differential equations of coordinates and momentum.

From Eq. 4.8 we have the first order derivatives of potential with respect to every atom. Equations 4.15, 4.16 and 4.17 depict the first-order time derivatives used for assigning the momentum to the atoms in the new position.

$$\frac{\partial P_{X_i}}{\partial t} = -\left(\frac{\partial V}{\partial X_i}\right) = \dot{P_{X_i}} \tag{4.15}$$

$$\frac{\partial P_{Y_i}}{\partial t} = -\left(\frac{\partial V}{\partial X_i}\right) = \dot{P_{Y_i}}$$
(4.16)

$$\frac{\partial P_{Z_i}}{\partial t} = -\left(\frac{\partial V}{\partial X_i}\right) = \dot{P}_{Z_i} \tag{4.17}$$

For convenience, let us denote  $\vec{P}_{X_i}$ ,  $\vec{P}_{Y_i}$  and  $\vec{P}_{Z_i}$  together as  $P_i$ . The time derivative for the coordinates of the atoms can be determined from Equations 4.18, 4.19 and 4.20

$$\frac{\partial X_i}{\partial t} = \frac{P_{X_i}}{m_i} = \dot{X}_i \tag{4.18}$$

$$\frac{\partial Y_i}{\partial t} = \frac{P_{Y_i}}{m_i} = \dot{Y}_i \tag{4.19}$$

$$\frac{\partial Z_i}{\partial t} = \frac{P_{Z_i}}{m_i} = \dot{Z}_i \tag{4.20}$$

Let us denote Equation 4.18, 4.19 and 4.20 i.e.,  $\dot{X}_i$ ,  $\dot{Y}_i$  and  $\dot{Z}_i$  as  $Q_i$ .

 $P_i$  and  $Q_i$  together constitute the coupled, first-order differential equations known as the Hamilton's equations of motion. They constitute the working equation for most trajectory calculations. In order to calculate the trajectory, Hamilton's equations of motion have to be numerically integrated from an initial state in the configuration space to some final state that lies in the configuration space.

The general form of Hamilton's equations is,

$$\frac{\mathrm{d}Y}{\mathrm{d}t} = \mathrm{f}[\mathrm{y}_{1,}\mathrm{y}_{2,\ldots,}\mathrm{y}_{6\mathrm{N}},\mathrm{t}]$$

where the elements of Y are the coordinates and their conjugate momenta:

$$Y_i = Q_i$$
 (i = 1, 2, ...., 3N)  
 $Y_i = P_i$  (i = 3N + 1, 3N + 2, ...., 6N)

Generally, the values of all elements of Y will be known at the initial condition when time  $t = t_0$ . In our simulation model we are using 4<sup>th</sup> order Runge-Kutta method. In this method the value of  $Y_i$  at  $t = t_0 + h$  (where h is the step size) is taken as

$$Y_{i}(t_{0} + h) = Y_{i}(t_{0}) + \frac{h[a_{1}k_{1}^{i} + a_{2}k_{2}^{i} + a_{3}k_{3}^{i} + a_{4}k_{4}^{i}]}{(a_{1} + a_{2} + a_{3} + a_{4})},$$

where,

$$k_1^i = f_i[y_1^0, y_2^0, \dots, y_{6N}^0, t_0],$$

$$\begin{split} k_2^i &= f_i \left[ Y_1^0 + \frac{k_1^1}{2}, \ Y_2^0 + \frac{k_1^2}{2}, \dots, Y_{6N}^0 + \frac{k_1^{6N}}{2}, t_0 + \frac{h}{2} \right], \\ k_3^i &= f_i \left[ Y_1^0 + \frac{k_2^i}{2}, \ Y_2^0 + \frac{k_2^2}{2}, \dots, Y_{6N}^0 + \frac{k_2^{6N}}{2}, t_0 + \frac{h}{2} \right], \\ k_4^i &= f_i \big[ Y_1^0 + k_3^1, \ Y_2^0 + k_3^2, \dots, Y_{6N}^0 + k_3^{6N}, t_0 + h \big], \end{split}$$

where the  $Y_1^0$  denotes  $Y_i(t_0)$ , the value of  $Y_i$  at  $t = t_0$ 

The Runge-Kutta procedure has several advantages. First, it is "self-starting" in that it is unnecessary to know the values of the elements of Y at times prior to  $t = t_0$ . Second, the local error in a given integration step is of the order  $h^5$ . The method therefore provides good accuracy. Finally, the method is stable and easy to program. The two major disadvantages of the this method are

- (1) It provides no estimate of the accuracy being achieved during the integration, and
- (2) 24N derivatives must be evaluated in each integration step, where N is the number of atoms in the model. Consequently, the method would appear to be very demanding of computer time.

In our model, there are 6 water molecules and one  $Cu^{2+}$  ion, so in total 19 atoms. Hence, there will be 57 first-order differential equations defining temporal variation of the coordinates and 57 first-order differential equations used for defining time dependence of the conjugate momentum of the atoms. The Y<sub>i</sub> would have totally 114 elements, which have to be numerically integrated over time, t. In our simulation, 1 atomic time unit =  $2.41888 \times 10^{-17}$  s, which is determined by our choice of energy, mass, and distance units to be hartrees, atomic mass units, and bohrs, respectively.

## 4.7 Conservation of energy and back integration

The Law of conservation energy, this states that "the total amount of energy in an isolated system remains constant over time". To verify whether the atomic model is in agreement with this law, the model was numerically integrated after providing initial momentum at room temperature as described previously. From Figure 4.11 and Figure 4.12, we can observe that the total energy ( $E_{TOT}$ ) and all momentum components are constant over the entire integration time length, as required by the conservation laws. Deviations are on the order of  $10^{-13}$  units.






Figure 4.12: Conservation of Linear Momenta along X, Y and Z direction

The integration accuracy is then further tested accessed by back integration of the trajectory. In this procedure, after n integration steps, all the  $P_i$  are replaced by -  $P_i$  in Y. The trajectory is then "back integrated" n steps. The extent of agreement of the final values of the elements of Y with their corresponding initial values provides a measure of integration accuracy. In general, if an agreement to at least three significant digits is obtained, the integration will be sufficiently accurate.

First the trajectory between Cu-O (R<sub>Cu-O</sub>) is tracked. Figure 4.13 shows the trajectory obtained as a result of forward integration with a time step of  $\Delta t = 1$  time unit between an initial time, t<sub>0</sub> = 0 and final time, t<sub>final</sub> = 600. Figure 4.14 shows the trajectory obtained as a result of backward integration with  $\Delta t = 1$  between an initial time, t<sub>0</sub> = 600 and final time, t<sub>final</sub> = 0.

We can observe from Figure 4.13 and Figure 4.14 that the trajectories perfectly overlap over each other. Next, the trajectory of  $O-H_1$  is tracked, as shown in Figure 4.15 and Figure 4.16



Figure 4.13: Forward integrated trajectory of Cu-O



Figure 4.14: Backward integrated trajectory of Cu-O



Figure 4.15: Forward integrated trajectory of  $O-H_1$ 



Figure 4.16: Backward integrated trajectory of  $O-H_1$ 

Even in this case, we can observe from Figure 4.15 and Figure 4.16 that the trajectories perfectly overlap. Clearly, the integrator is sufficiently accurate.

#### 4.8 Energy assignment based on vibration modes of water molecule

With three atoms in a nonlinear arrangement, H<sub>2</sub>O has three vibrational modes, according to the formula 3N - 6 (*i.e.*, 3(3) - 6 = 3). Their measured frequencies expressed in wave numbers are **3825**. 3 cm<sup>-1</sup> in the symmetric stretching mode  $v_S$  (as shown in Figure 4.17), **1653**. 9 cm<sup>-1</sup> in the bending mode  $v_B$  (as shown in Figure 4.18) and **3935**. 6 cm<sup>-1</sup> in the asymmetric stretching mode  $v_A$  (as shown in Figure 4.19)



Figure 4.17: Symmetric Stretching mode

Figure 4.18: Bending mode



Figure 4.19: Asymmetric Stretching mode

The frequencies are converted into zero-point energy (E) values in hartree using the following conversion factor,

$$E = (\omega) \left[ \frac{1 \text{ Hartree}}{2.19475 \times 10^5} \right] \left( \frac{1}{2} \right),$$

where  $\omega$  is the wave number in  $cm^{-1}$ .

The zero-point energy values in hartree are shown below:

$$E(v_S) = 0.0087147$$
 hartree  
 $E(v_B) = 0.0037679$  hartree  
 $E(v_A) = 0.0089659$  hartree

The zero-point energies in the equilibrium configuration will all be in the form of kinetic energy of  $H_2O$  molecule. Hence the total zero-point energy value will be equal to the total kinetic energy when  $H_2O$  is in its equilibrium configuration.

In the symmetric stretching mode (shown in Figure 4.17), the hydrogen atoms are each displaced in the negative Y direction by an amount  $S_1$ , which, for moment we will assume is our normal mode vibrational coordinate. To maintain the center of mass fixed, the oxygen atom must be displaced upward by an amount  $\Delta y_o$  during this phase of the vibrational motion. The requirement that the center of mass remain fixed means we must have

$$-\mathbf{m}_{\mathrm{H}}\mathbf{S}_{1} - \mathbf{m}_{\mathrm{H}}\mathbf{S}_{1} + \mathbf{m}_{\mathrm{O}}\,\Delta\mathbf{y}_{\mathrm{O}} = \mathbf{0}$$

This gives an oxygen atom displacement in Symmetric mode.

$$\Delta y_{o} = \frac{2m_{H}}{m_{O}} S_{1}$$

All other displacements are zero in this mode. These displacements are symmetric with respect to a  $180^{\circ}$  rotation about the Y-axis in that such a rotation does not change the magnitude of signs of the displacements. This makes it clear that this type of displacement symmetrically stretches the O-H bonds. Consequently, S<sub>1</sub> is termed as symmetric stretching coordinate.

In the bending mode (shown in Figure 4.18), the oxygen atom remains stationary while the hydrogen atoms H<sub>1</sub> and H<sub>2</sub> are displaced along the X-axis by amount  $\Delta x_2 = -S_2$  and  $\Delta x_3 = S_2$ . All other displacements in this mode are zero. A 180<sup>o</sup> rotation about

the symmetric (Y-axis) produces no change in the displacements. Consequently, this mode is also symmetric. The predominate action of the vibrational coordinate is to bend the molecule.  $S_2$  is, therefore, called a bending coordinate.

In the asymmetric bending mode (shown in Figure 4.19), the displacements are

$$\Delta x_2 = -S_3 \sin \theta$$
 and  $\Delta y_2 = S_3 \cos \theta$ ,  $\Delta x_3 = -S_3 \sin \theta$  and  $\Delta y_3 = -S_3 \cos \theta$ 

To maintain the center of mass fixed, we must have

$$m_H \Delta x_2 + m_H \Delta x_3 + m_O \Delta x_o = -2S_3 m_H \sin \theta + m_O \Delta x_o = 0$$

Therefore, the oxygen-atom displacement must be  $\Delta x_o = \frac{2m_H}{m_O}S_3 \sin \theta$ . All other displacements are zero. If we rotate the molecule by  $180^{\circ}$  about the Y-axis, all the displacements change directions. Therefore, this vibrational coordinate is asymmetric. It also either stretches or compresses the O–H bonds. We can, therefore, reasonably refer to S<sub>3</sub> as an asymmetric coordinate.

In our case, the three vibrational coordinates are independent and each exhibits harmonic motion, we can set the vibrational coordinates such that the potential and kinetic energies can be written in the form

Potential energy, 
$$V(S_1, S_2, S_3) = \sum_{i=1}^{3} a_i Q_i^2$$
 (4.21)

Kinetic energy, 
$$T = \sum_{i=3}^{3} b_i v_i^2$$
 (4.22)

where  $v_i = \frac{\partial Q_i}{\partial t}$  and the  $a_i$  and  $b_i$  are constants. The vibrational coordinates  $S_1$ ,  $S_2$  and  $S_3$  are called "normal mode coordinates" and the three vibrational modes are called "normal modes".

In order to make sure that their motion will be harmonic, the three vibrational coordinate  $S_1$ ,  $S_2$  and  $S_3$ , must satisfy Equations 4.21, 4.22. The total displacement for all atoms may be obtained by summing the displacement for the three modes.

The result is

$$\Delta x_{o} = \frac{2m_{H}}{m_{o}} S_{3} \sin \theta, \ \Delta y_{o} = \frac{2m_{H}}{m_{o}} S_{1}, \ \Delta z_{o} = 0$$
(4.23 a)

$$\Delta x_{H_1} = -S_2 - S_3 \sin \theta, \ \Delta y_{H_1} = -S_2 + S_3 \cos \theta, \ \Delta z_{H_1} = 0$$
(4.23 b)

$$\Delta x_{H_2} = S_2 - S_3 \sin \theta, \ \Delta y_{H_2} = -S_1 + S_3 \cos \theta, \ \Delta z_{H_2} = 0$$
(4.23 c)

In Cartesian coordinates, the kinetic energy is given by Eq. 4.24.

$$T = \sum_{i=3}^{3} 0.5 m_i \left[ \left( \frac{\partial \Delta x_i}{\partial t} \right)^2 + \left( \frac{\partial \Delta y_i}{\partial t} \right)^2 + \left( \frac{\partial \Delta z_i}{\partial t} \right)^2 \right]$$
(4.24)

If we take the time derivatives of Equations 4.23 a, b and c, making sure that  $\theta$  is fixed at the equilibrium angle for H<sub>2</sub>O, square each result and the sum of derivatives for the three atoms multiplied by the appropriate mass factors, the result is

$$T = \left[m_{H} + \frac{2m_{H}^{2}}{m_{0}}\right] \left(\frac{\partial S_{1}}{\partial t}\right)^{2} + m_{H} \left(\frac{\partial S_{2}}{\partial t}\right)^{2} + \left[m_{H} + \frac{2m_{H}^{2}\sin\theta}{m_{0}}\right] \left(\frac{\partial S_{3}}{\partial t}\right)^{2} = \left[m_{H} + \frac{2m_{H}^{2}}{m_{0}}\right] v_{S_{1}}^{2} + m_{H} v_{S_{2}}^{2} + \left[m_{H} + \frac{2m_{H}^{2}\sin\theta}{m_{0}}\right] v_{S_{3}}^{2}$$

$$(4.25)$$

When calculated, both LHS and RHS of Eq.4.25 were equal to 0.0214 Hartree, verifying that the normal mode coordinates were accurate and the H<sub>2</sub>O exhibits harmonic motion.

Figure 4.20 (a) shows the vibration of an O–H bond containing all 3 vibrational modes. Figure 4.20 (b) shows the FFT plot of the time series plot shown in Figure 4.20 (a). It can be clearly observed that the vibration has three prominent frequencies which correspond to the 3 modes of normal modes of vibration.



Figure 4.20 (a): Vibration of the O-H bond about the Equilibrium distance (1.8085 Bohrs)



Figure 4.20 (b): FFT of vibration of the O-H bond

# 4.9 Octahedral structural arrangement of [Cu (H<sub>2</sub>O)<sub>6</sub>]<sup>2+</sup> complex molecule

Before determining the coordinates of all 6 water molecules, the center of mass is determined using the formulation depicted in Eq. 4.8

$$m_0 X - 2m_H \left\{ \cos\left(\frac{\theta}{2}\right) R_{0-H} - X \right\} = 0$$
(4.26)

where,  $m_0 = mass of oxygen atom = 29165.7836 amu$ 

 $m_{\rm H}$  = mass of hydrogen atom = 1837.4041 amu.

 $\theta$  = bond angle of water H - 0 - H = 104.5<sup>o</sup> = 1.8239 radians

X = Center of mass in bohrs

 $R_{O-H}$  = Euilibrium distance between O - H in H<sub>2</sub>O molecule = 1.8085 Bohr

From Eq. 4.26 the center of mass of water is at 0.1239 bohrs. Hence, for all normal mode energy calculations, the oxygen atom is placed at distance of 0.1239 bohrs above  $Cu^{2+}$  assuming that  $Cu^{2+}$  is exactly located at the center of mass position. Now the system containing just one molecule of water i.e.,  $[Cu(H_2O)]^{2+}$  assigned with the energies calculated from normal mode coordinates is integrated until several thousand structures and their conjugate momenta values are obtained and saved.

The position of all O atoms about  $Cu^{2+}$  i.e.,  $R_{Cu-O}$  is then determined based on the density of H<sub>2</sub>O model as 5.506 bohrs. In order to introduce a minimal randomness in  $R_{Cu-O}$ , Eq. 4.27 was used.

$$R_{Cu-0} = 5.506 \pm 0.02(5.506)\xi \tag{4.27}$$

where,  $\xi$  is a random number such that  $-1 \le \xi \le +1$ 

The coordinates of Cu<sup>2+</sup> were fixed at [0 0 0]. Six random coordinates and their conjugate momenta values are chosen from the previously saved structures. The Z coordinates of the O, H<sub>1</sub> and H<sub>2</sub> atoms of the first water molecule are fixed as 5.506  $\pm$  0.02(5.506) $\xi$ , and X coordinates are made 0 as the molecule is placed in the Z-Y plane. Now the second structure is taken and rotated 180<sup>o</sup> about the Y axis using the Rotation matrix R<sub>Y</sub>. Its conjugate momenta values are also rotated 180<sup>o</sup> about the Y axis. The Z coordinates of the atoms in the second water molecule are added with a factor of 5.506  $\pm$  0.02(5.506) $\xi$ . The third structure is chosen, and is rotated 90<sup>o</sup> about the Y axis using the rotation

matrix  $R_Z$ . Now the third water molecule is moved towards its final position by adding the X and Y coordinates with a factor of  $\cos 45^{\circ} [5.506 \pm 0.02(5.506)\xi]$ . Its conjugate momenta values are subjected to the same treatment as the coordinates. Same procedure is adopted for fixing the fourth, fifth and sixth water molecules, but in the final step the coordinates are fixed by the following method,

For the fourth water molecule,

$$(Q_X) = (Q_X) - \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$$
  
 $(Q_Y) = (Q_Y) - \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$ 

For the fifth water molecule,

$$(Q_X) = (Q_X) - \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$$
$$(Q_Y) = (Q_X) + \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$$

For the sixth water molecule,

$$(Q_X) = (Q_X) + \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$$
  
 $(Q_Y) = (Q_X) - \cos 45^\circ [5.506 \pm 0.02(5.506)\xi].$ 

where,  $(Q_X)$  and  $(Q_Y)$  represents all the X and Y coordinates of O, H<sub>1</sub> and H<sub>2</sub> of fourth, fifth and sixth water molecule.

All the momenta values are subjected to the same treatment as their corresponding coordinates. The final octahedral arrangement of the copper (II) hexa-hydrate complex molecule is shown in Figure 4.21.



Figure 4.21: Octahedral structure of copper (II) hexa hydrate complex molecule

#### **4.9.1 Damped Trajectory Calculation**

After setting up the octahedral structure of the copper (II) hexa hydrate molecule, the average distance between  $Cu^{2+}-O$  at the global minima position had to be determined. We know that nature will always place the molecules in the lowest energy position. So by performing damped trajectory calculation, we can determine the average distance between  $Cu^{2+}$  and O, which would depict the approximate positions of the atoms, at lowest energy possible.



Figure 4.22: Marbel-Bowl experiment

In order to elucidate the concept of damped trajectory, consider an open bowl into which a marble 'm' is rolled, as shown in Figure 4.22. When m is in its initial position 1, it stores potential energy (PE), which eventually gets converted into kinetic energy (KE) until it reaches position 2. The marble loses its kinetic energy because of friction the exerted by the marble, in other words the momenta value becomes 0 at position 2. Again the marble rolls from position 2 to 3, loses some KE due to friction and rolls to position 3 where the momenta again become 0. This rolling action happens until the marble has reached the lowest energy position. Law of conservation of energy is the key principle behind this experiment.

To determine the average distance of  $Cu^{2+}-O$  a similar concept was formulated. The [Cu (H<sub>2</sub>O) <sub>6</sub>]<sup>2+</sup> system was set up in its initial conditions and was integrated through a particular time interval and the kinetic energy was monitored. Initially, the kinetic energy was increasing (PE was getting converted into KE) and the simulation was stopped at the point where the KE started to decrease. The momenta values at this position were zeroed out and again the system was integrated through some time interval until the KE started decreasing again. This procedure was repeated until the KE value converged and remained constant. The position at this point was stored and from this set of coordinates all the  $Cu^{2+}$ –O distances were calculated, in toto 6 distances corresponding to 6 O atoms in the system. The average  $Cu^{2+}$ –O was calculated out to be 6.5484 Bohrs, which was updated in the coordinates assigning code and the new initial coordinates of the atoms were defined.

#### 4.10 Water-water interaction potential

Having modeled octahedral arrangement of the complex molecule, next, the waterwater interaction potentials were added based on the Coulombic forces of interaction between the pairwise partial charges.

#### 4.10.1 O-O interactions (V<sub>0-0</sub>)

The total number of O–O interactions is given by  $\frac{N(N-1)}{2}$ . There are totally 6 water molecules hence there would be '15 O–O interactions', which are modeled using Eq. 4.28.  $q_0$  is the atomic charge on oxygen atoms = -0.6836.  $\varepsilon$  is the dielectric constant of water = 80 at 298K

Figure 4.23: O – O interaction (repulsion)

## 4.10.2 O–H interactions $(V_{0-H})$

The total number of O–H interactions is given by 2N(N - 1). There are totally 6 water molecules hence there would be '60 O–H interactions', which are modeled using Eq. 4.29.  $q_0$  is the atomic charge on oxygen atoms = -0.6836 and  $q_H$  is the atomic charge on hydrogen atoms = 0.3418.  $\varepsilon$  is the dielectric constant of water = 80 at 298K

$$V_{O-H} = \text{Sum of all } O - \text{H interactions} = \frac{q_O q_H}{\epsilon} \sum_{i=1}^{N-1} \sum_{j>i}^{N} \left[ \frac{1}{r_{ij1}} + \frac{1}{r_{ij2}} \right]$$
(4.29)



Figure 4.24: O-H interaction

# 4.10.1 H–H interactions (V<sub>H–H</sub>)

The total number of H–H interactions is given by  $\frac{4N(N-1)}{2}$ . There are totally 6 water molecules hence there would be '60 H–H interactions', which are modeled using Eq. 4.30.  $q_{\rm H}$  is the atomic charge on hydrogen atoms = 0.3418.  $\varepsilon$  is the dielectric constant of water = 80 at 298K.

 $V_{H-H} =$  Sum of all H – H interactions



Figure 4.25: H – H interaction

Hence there are a total of '135 terms' in the inter water-water potential. The new  $V_{TOT}$  potential would have extra 135 terms in addition to the 18 intra water molecular

interactions, which have been already modeled. So totally there 153 terms in the model corresponding to 6 water molecules. Eq. 4.31 depicts the final  $V_{TOT}$  potential for N water molecules.

$$V_{TOT} = \sum_{i=1}^{N} (V_2) + \sum_{i=1}^{2N} (V_{H_2O}) + \sum_{i=1}^{2N+3} V_{O-O} + \sum_{i=1}^{10N} V_{O-H} + \sum_{i=1}^{10N} V_{H-H}$$
(4.31)

After adding the inter water potentials, the model was integrated over time to check whether the law of conservation of energy still holds well. It was observed that the energy was conserved with an accuracy of 9-10 digits, as shown in Figure 4.26.



Figure 4.26: Total Energy conservation of the octahedral [Cu (H<sub>2</sub>O)  $_{6}$ ]<sup>2+</sup> system

## 4.11 Simulation of $[Cu (H_2O)_6]^{2+}$ molecule under electrode potential

In ECMP process, the  $Cu^{2+}$  ion gets liberated from the anode due to oxidation and travels towards cathode where hydrogen reduction happens. In order to observe the behavior of  $[Cu (H_2O)_6]^{2+}$  complex molecule under electrode potential, the molecule is subjected to an external electric field. The electric field is assumed to be uniform in this case.

Figure 4.27 provides a pictorial representation of the application of uniform electric field between positively charged anode and negatively charged cathode. 'V' is the applied anodic potential and R is the distance between anode and cathode in meters.



Figure 4.27: Application of external electric field

Uniform electric field ( $\epsilon$ ) can be calculated using Eq. 4.32.

$$\varepsilon = \frac{V}{R} \text{ (volts m}^{-1}\text{)} \tag{4.32}$$

External electric energy, 
$$E_e = Vq$$
 (in eV) (4.33)

$$= Vq [eV] \left[ \frac{1 \text{ hartree}}{27.2114 \text{ eV}} \right] = \frac{Vq}{27.2114} \text{ hartree}$$
$$= \frac{Vq}{27.2114 \text{ R}} \left[ \frac{\text{hartree}}{\text{m}} \right] \left[ \frac{0.529177 \times 10^{-10} \text{m}}{\text{bohr}} \right]$$
$$E_e = \frac{Vq}{R} (1.944689 \times 10^{-12}) \frac{\text{hartree}}{\text{bohr}}$$

Assuming that the conversion factor,  $\beta = 1.944689 \times 10^{-12}$ , we can calculate the force exerted on the system along X axis due to the external electric field as shown in Eq. 4.34

$$\overrightarrow{\varepsilon_{\rm X}} = \frac{\beta {\rm V} {\rm q}}{{\rm R}} \tag{4.34}$$

where,  $\overrightarrow{\varepsilon_X}$  is the Electric force due to the uniform electric field

$$\beta = 1.944689 \times 10^{-12}$$

V is the applied experimental anodic potential = 4.5 V

q is charge on Copper = +2

R is Experimental distance between anode and cathode = 0.005 m

 $\varepsilon_r$  is the dielectric constant of water = 80 at 298K

As the external electric field is applied along the X axis,  $\vec{\varepsilon}_Y = \vec{\varepsilon}_Z = 0$  and  $\vec{F}_Y = \vec{F}_Z = 0$ . V<sub>TOT</sub> has to be added with  $\vec{\varepsilon}_X$  and

$$F_{Xi}^{e} = \text{electric force on particle } i = \frac{\beta V q_{i}}{R} = \left(-\frac{\partial V_{i}}{\partial X_{i}}\right)$$

Hence,

$$\frac{\partial P_{Xi}}{\partial x} = \left(-\frac{\partial V_{TOT}}{\partial X_i}\right) + F_{Xi}^e = \left(-\frac{\partial V_{TOT}}{\partial X_i}\right) + \frac{\beta V q_i}{\epsilon_r R}$$

In the next section, Section 4.12, results obtained from the MD simulation at different condition has been presented.

#### 4.12 Simulation results and discussions:

The copper (II) hexa hydrate molecule was simulated under 3 different conditions.

- 1) No voltage condition
- 2) Experimental voltage condition (4.5 V)
- 3) High voltage condition (50 V)

Some voltages are too high to be included in the real world ECMP process, but are considered here to investigate the effects of voltage on the shape of the complex molecule and also the movement of  $Cu^{2+}$  ion.

In all the results shown below, the following conventions would apply.

Magenta sphere: Cu<sup>2+</sup> ion Red spheres: O atoms Green atoms: H atoms Blue dotted lines: Cu–O bonds Black dotted lines: O–H bonds At the end of every simulation results, the corresponding energy curves have also been reported. All the simulations were run from an initial time,  $t_0 = 10$  to final time,  $t_{\text{final}} = 50000$ , with a time step of 100.

The view angle of the molecular structure is at an azimuthal angle of 169 and an elevation angle of 9. This configuration was chosen after trying out different viewing angles and only this alignment could show the change of the molecular shape accurately.

#### 4.12.1 No voltage condition:

At no voltage condition there is no external force applied to the system. The atoms are provided with minimal thermal energy at the room temperature (298K). The following snap shots from the MD simulation, shows the step by step change of the isolated complex molecule in aqueous solution with time. Snap shots are taken at time steps of 0, 50, 150, 200, 250, 300, 350, 400, 450 and 500.



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From the above snap shots, we can clearly see that at no voltage condition the atoms reorient themselves at random positions. Though the octahedral shape of the complex molecule is distorted the H atoms stay with the O atoms indicating that the water molecules are stable. But they exhibit three modes of vibration: symmetric stretching, asymmetric stretching and bending, in correspondence to the initial energy we have provided the atoms. As the repulsive force between Cu ion and water molecules increase owing to distance between Cu and  $H_2O$ , the water molecules are repelled away from the Cu ion, as they are governed by the V<sub>2</sub> potential, which has a greater well depth.

From Figure 4.28, we can observe the movement of Cu ion in the X and Y directions. Though we have not applied any voltage, we can observe minimal movement of Cu ion. This is caused due to the repulsive and attractive forces exerted on the Cu ion by the water molecules.



Figure 4.28: Movement of Cu ion at no volt condition

With no external force applied to the system, the energy is perfectly conserved as shown in Figure 4.29





### 4.12.2 Experimental voltage condition

In order to observe the simulation under experimental voltage condition, a potential of 4.5 V was applied along the X direction. Snap shots of the system were taken at the same time steps as recorded in Section 4.12.1.









Figure 4.30: Movement of Cu ion at experimental potential condition

In this case, we can observe that the velocity of the Cu ion has increased, due to the applied voltage of 4.5V as observed in Figure 4.30.



Figure 4.31: Energy plot at experimental potential condition

Figure 4.31 depicts the variation of energy with time. Due to application of the external electric force field the energy is no longer conserved.

# 4.12.3 High voltage condition

To observe the behavior of the Cu ion at a higher voltage condition, a voltage of 50 V was applied across the system along the X direction.





-15 -10

9

-10 -15 -20

-20

-5

-5





Figure 4.32: Energy plot at higher potential condition



Figure 4.33: Movement of Cu ion at higher potential condition

From the snap shots at different time steps and from Figure 4.33, we can clearly observe that the distance moved by Cu ion with time is more than observed at 0V and 4.5V conditions. This shows that the velocity of the Cu ion has increased and it is mainly because of the increase in applied voltage. There are minimal movements in the Y direction, which is due to the repulsive and attractive forces exerted on it by water molecules. From Figure 4.32, we can see that the energy of the system has increased due to the high voltage applied to the system.

Although by running the simulation for more time, we will be able to track the movement of  $Cu^{2+}$  and water molecules in a more precise manner, the relative tolerance value of the numerical integrator was set to a very low value in order to obtain accurate results, hence the simulation would require more computation time and memory to store the results.

### 4.13 Computation of current carried by Cu<sup>2+</sup>

The current carried by the cation, in our case  $Cu^{2+}$ , can be calculated from the velocity of  $Cu^{2+}$  at varied anodic voltage conditions.

Time =  $(500 \times 100)$  time units  $\times 2.41888 \times 10^{-17} = 1.2094 \times 10^{-12}$  s

Distance = 
$$\left(\frac{1 \text{ cm}}{10^8}\right) \times (d) \times \left(\frac{0.529 \text{ A}^\circ}{\text{bohr}}\right)$$

where, d = distance travelled by  $Cu^{2+}$  in bohrs

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$$\therefore \langle v \rangle = \frac{\text{Distance}}{\text{Time}} \text{ in cm/s}$$

The current carried by a  $Cu^{2+}$  (cation),

$$I_{+} = \frac{\partial Q_{+}}{\partial t} = \frac{N_{+}z_{+}ev_{+}}{D}$$

where,  $N_{+} =$  Total number of Cu<sup>2+</sup>in the system

 $z_+$  = charge carried by  $Cu^{2+} = +2$ 

e = charge carried by electrons =  $1.602177 \times 10^{-19}$  coloumb

 $v_{+} = \text{terminal velocity of } Cu^{2+}$ 

D = Distance between cathode and anode = 0.01 m

The current carried by a single copper ion at experimental and high voltage condition has been presented in Table 4.2

Table 4.2: Current carried by  $Cu^{2+}$  at different voltage conditions

Applied voltage (V)	Velocity (m/s)	I <sub>+</sub> (amps/ion)
4.5	103.84	$3.322 \times 10^{-15}$
50	166.53	$5.335 \times 10^{-15}$

Hence, the total current carried by all the  $Cu^{2+}ions = N_{+} \times I_{+}$ 

where,  $N_{+}=\mbox{Total}$  number of  $\mbox{Cu}^{2+}\mbox{in}$  the system

#### CHAPTER V

#### SUMMARY AND FUTURE WORK

In this study, the prime emphasis was on understanding the process mechanism from an experimental and atomistic standpoint. The focus was on understanding and delineating the various electrochemical phases as stated in the literature from the ECMP process and their influence on the process. The ECMP apparatus was redesigned and modified to integrate the setup with a data acquisition (DAQ) system to monitor the current signal and track the transition of electrochemical phases with respect to the process parameters. From the experimental study of Cu ECMP using phosphoric acid based slurry, the following observations are recorded and the mechanism behind it has been elaborately discussed in chapter III.

The next significant aspect of this work is to study and understand the process further from an atomistic standpoint. Using the molecular dynamics (MD) simulation technique, the copper (II) hexa hydrate molecule, one of the chief elements of the passivation layer, was modeled and simulated under different anodic conditions. Changes in the molecular shape with increases in voltage were also observed at different time steps. Without any voltage being applied, the atoms in the molecule reorient themselves to a minimum energy configuration and energy is also conserved as there is no external force. The molecule began orienting itself, such that the copper ion moved towards the anode along the direction of application of the electric force (X direction), with application of the electric potential. As the voltage increased, the translation distance of Cu ion with time increased indicating that the velocity of the Cu ion has also increased.

Process Parameter	Conditions	Conclusions	
рН	Acidic (pH = 2)	<ul> <li>MRR increases with increase in anodic voltage</li> <li>R<sub>a</sub> increases with increase in anodic voltage</li> </ul>	
	Near Neutral (pH = 6.5)	<ul> <li>MRR increases with increase in anodic voltage but the values are way lower than the values observed at pH = 2.</li> <li>R<sub>a</sub> decreases with increase in anodic voltage</li> </ul>	
Anodic Voltage	Active	<ul> <li>R<sub>a</sub> value is high</li> <li>MRR value is low</li> </ul>	
	Passive	<ul> <li>R<sub>a</sub> is improved than active region</li> <li>MRR is slightly increased over the active region</li> </ul>	
	Transient/ Trans-passive	<ul> <li>Best R<sub>a</sub> value and high MRR are achieved</li> <li>Optimum regions to work with, in Cu ECMP</li> </ul>	

Table 5.1: Summary of the experimental results

In the future, with respect to the experimental part, better control of the pad pressure can help in achieving better control over the process. Electrochemical analysis of the slurry can provide more information about the components formed in the slurry
after the ECMP process, which can provide us more insight into the composition of the passivation layer. In terms of atomistic simulation, more water molecules can be added to the system. In addition to observing the complex molecule under different anodic potentials, the effect of variations of the pH on the complex molecule can be studied. In this work, the interaction of Cu ions with water has been studied in detail, but the work can be extended by adding other slurry components to gain a better understanding of the Cu ECMP process. Furthermore, by calculating the number of copper ions in the system we can calculate the fraction of current carried by copper ions. With addition of H<sup>+</sup> and  $PO_4^{3-}$  ions to the system we can also compute the fraction of current carried by the corresponding ions. The total current I<sub>total</sub> can then be calculated as,

$$I_{total} = i_{Cu^{2+}} + i_{H^+} + i_{PO_4^{3-}}$$

where,  $i_{Cu^{2+}} = Current carried by Cu^{2+}$  $i_{H^+} = Current carried by H^+$  $i_{PO_4^{3-}} = Current carried by PO_4^{3-}$ 

If we can create a model which can express material removal rate as a function of  $I_{total}$ , we can very well predict MRR of ECMP process. To compute the transference numbers of H<sup>+</sup> and PO<sub>4</sub><sup>3-</sup> more experimental and computational works have to be carried out. By this way we can to make this molecular dynamic simulation model work from an engineering standpoint, to predict MRR.

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### APPENDIX

## ATOMIC UNIT CONVERSIONS

### Time

1 time unit =  $2.41888 \times 10^{-17} s$ 

# Length

1 *bohr radius* =  $5.29177 \times 10^{-11}m$ 

### Energy

1 *hartree* = 2625.5 *kJ/mol* 

## Angle

1 radian = 57.295 degrees

## Weight

 $1 amu = 1.6605 \times 10^{-27} kg$ 

#### VITA

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